RENORMALIZATIONS OF THE KONTSEVICH INTEGRAL AND THEIR BEHAVIOR UNDER BAND SUM MOVES
by

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# AN ABSTRACT OF A DISSERTATION <br> submitted in partial fulfillment of the requirements for the degree DOCTOR OF PHILOSOPHY 

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KANSAS STATE UNIVERSITY
Manhattan, Kansas
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## Abstract

We generalize the definition of the framed Kontsevich integral initially presented in [LM1]. We study the behavior of the renormalized framed Kontsevich integral $\hat{Z}_{f}$ under band sum moves and show that it can be further renormalized into some invariant $\widetilde{Z}_{f}$ that is well-behaved under moves for which link components of interest are locally put on top of each other. Originally, Le, Murakami and Ohtsuki ([LM5], [LM6]) showed that another choice of normalization is better suited for moves for which link components involved in the band sum move are put side by side. We show the choice of renormalization leads to essentially the same invariant and that the use of one renormalization or the other is just a matter of preference depending on whether one decides to have a horizontal or a vertical band sum. Much of the work on $\widetilde{Z}_{f}$ relies on using the tangle chord diagrams version of $\hat{Z}_{f}([\mathrm{ChDu}])$. This leads us to introducing a matrix representation of tangle chord diagrams, where each chord is represented by a matrix, and tangle chord diagrams of degree $m$ are represented by stacks of $m$ matrices, one for each chord making up the diagram. We show matrix congruences for some appropriately chosen matrices implement on the modified Kontsevich integral $\widetilde{Z}_{f}$ the band sum move on links. We show how $\widetilde{Z}_{f}$ in matrix notation behaves under the Reidemeister moves and under orientation changes. We show that for a link $L$ in plat position, $Z_{f}(L)$ in book notation is enough to recover its expression in terms of chord diagrams. We elucidate the relation between $\check{Z}_{f}$ and $\widetilde{Z}_{f}$ and show the quotienting procedure to produce 3-manifold invariants from those as introduced in [LM5] is blind to the choice of normalization, and thus any choice of normalization leads to a 3-manifold invariant.

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Approved by:

Major Professor
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Renaud Gauthier

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## Chapter 1

## Introduction

This thesis has for aim to study 3-manifold invariants built from the framed Kontsevich integral $Z_{f}$ ([LM1]). The unframed Kontsevich integral $Z$, first introduced in $[\mathrm{K}]$, is a functional on knots that becomes an isotopy invariant of knots once it is corrected in an appropriate manner. The normalized framed Kontsevich integral defined by Le, Murakami and Ohtsuki which they denote $\hat{Z}_{f}$ [LM1] is also an isotopy invariant and will be the starting point of our study. In [LM5], a 3-manifold invariant $\Omega$ is constructed by further renormalizing $\hat{Z}_{f}$ into some other invariant $\check{Z}_{f}$. One property that $\Omega$ must satisfy for it to be a true 3 -manifold invariant is that it be invariant under band sum moves. Band sum moves between link components correspond to handle sliding on their respective 2 handles if a framed link is regarded as attaching data for 2-handles [GS]. Band sum moves on tangle chord diagrams ([ChDu]) however (tangles with chords ending on them) are not well-defined. One has to specify a window in which link components involved in the band sum are locally frozen. In [LM5], the invariance of $\Omega$ under band sum moves between link components locally put side by side relies on $\check{Z}_{f}$ behaving a certain way under such moves. In this paper, we start with $\hat{Z}_{f}$ and consider what normalization makes it well-behaved under band sum moves where link components of interest are locally on top of each other. We then show the equivalence between the normalization we get and that obtained in [LM5]. That two renormalizations come into play is not surprising since as Kontsevich initially pointed out $([\mathrm{K}])$ the Kontsevich integral depends very much on the choice of a time axis, and thus fixing two windows
where link components are rotated will lead to different results locally.

One fundamental remark is the following: from [LM1] we only use the normalization of the Kontsevich integral at local extrema which we generalize by using an arbitrary parameter $M$. Apart from enabling one to identify whether a given tangle $T$ has a local extremum by looking at the coefficients of $Z_{f}(T)$, this will allows us to locate those once $Z_{f}(T)$ is expressed in a "book notation" that we introduce in the present work. In [LM1] Le and Murakami defined a truncated invariant $Z$ obtained from a limiting procedure on the original Kontsevich integral $Z$ (same notation) on braids that enables one to compute the values of that invariant on q-tangles (or non-associative tangles). In this paper we work with geometric tangles only, that is tangles whose strands are separated from one another by a measurable distance. By the universality of $\hat{Z}_{f}$, the renormalized Kontsevich integral, the construction performed in [LM1] and the one done in the present work yield the same results. Whenever we present a result for geometric tangles initially proved by Le, Murakami and Ohtsuki for their invariant $Z_{f}$ ( or $\hat{Z}_{f}$, or even $\check{Z}_{f}$ ) using q-tangles, then we will mention the corresponding reference.

We then study the quotienting map ([LM5]) necessary for producing 3-manifold invariants from the Kontsevich integral and show that any normalization survives the quotienting process, and thus any choice of normalization yields a 3-manifold invariant. In particular this lifts the redundancy of having two normalizations $\check{Z}_{f}$ and $\widetilde{Z}_{f}$ describing the same things.

We formalize all these statements. We briefly recall the definition of knots and links, handle decompositions, the 4 dimensional 2-handlebodies determined by links, and that the diffeomorphism types of boundaries of such spaces are invariant under two operations, one of which is handle slide. Finally we will define the Kontsevich integral, its correction, and show what its expression is in the framed case. We show how it behaves under band sum
moves.

### 1.1 Links and handlebodies

We define a framed e-component link $L$ in $S^{3}$ [RK1] to be a smooth embedding $\gamma$ : $\coprod_{1 \leq i \leq e} S^{1} \rightarrow S^{3}$, for $e$ finite, along with integers $n_{1}, \ldots, n_{e}$, one for each embedding $\gamma_{i}:=\left.\gamma\right|_{S^{1}}$ of the $i$-th $S^{1}$ in $\coprod_{1 \leq i \leq e} S^{1}$. Those integers define the framing for each embedded circle, and are defined as follows. We first consider the case of knots. A knot is obtained for $e=1$. A knot $K$ in $S^{3}$ bounds an embedded compact surface called a Seifert surface [S]. The zero framing of the normal bundle of $K$ is obtained from the outward normal vector to the orientable Seifert surface $F$ of $K$. A non-zero framing given by some positive integer $n$ is obtained by twisting the zero framing $n$ times in a clockwise direction. This corresponds to associating the integer $n$ to the embedding $\gamma$ defining the knot $K$. Doing this for all embeddings of circles $\gamma_{i}, 1 \leq i \leq e$ for the link $L$ one obtains the integer numbers $n_{1}, \cdots, n_{e}$ respectively. Links will be oriented unless specified otherwise.

Now a link $L$ in $S^{3}$ determines a 4-manifold $M_{L}$ obtained by adding 2-handles to the 4-ball $D^{4}$ along the circles in $S^{3}$ defining $L$, with the gluing performed using the framing along the circles. Before constructing $M_{L}$, one therefore needs to introduce the concept of handles and that of handle decompositions. We will mainly follow [GS].

Definition 1.1.1. Let $X$ be an $n$-dimensional manifold. For $0 \leq k \leq n, D^{k} \times D^{n-k}$ is called an $n$-dimensional $k$-handle once it is attached to the boundary $\partial X$ of $X$ along $\partial D^{k} \times D^{n-k}$ via an embedding $\varphi: \partial D^{k} \times D^{n-k} \rightarrow \partial X$, called the attaching map.

If $h$ is the handle thus specified, one writes $X \cup_{\varphi} h$ to denote the resulting $n$-dimensional manifold ([GS]). One can smooth corners during the attaching map in such a manner that one views $X \cup_{\varphi} h$ as a smooth $n$-dimensional manifold. In terms of $M_{L}$ resulting from the
gluing of a 2-handle along $L$ using the framing on each circle, note that we have a deformation of $X \cup_{\varphi} h$ onto $X \cup_{\varphi \mid \partial D^{k} \times 0} D^{k} \times 0$. By the tubular neighborhood theorem, one can reconstruct $X \cup_{\varphi} h$ from the restriction of the attaching map $\varphi$ to $D^{k} \times 0$ along with a normal framing of $\left.\varphi\right|_{D^{k} \times 0}\left(\partial D^{k}\right)$. In particular, to construct $X \cup_{\varphi} h$, it suffices to have $\varphi_{c}: \partial D^{k}=S^{k-1} \rightarrow \partial X$ with a local trivialization of the normal bundle, $\varphi_{c}=\left.\varphi\right|_{D^{k} \times 0}$ where the subscript $c$ stands for core, $D^{k} \times 0$ is called the core of the $k$-handle, and one also needs a normal framing $f$ of $\varphi_{c}\left(S^{k-1}\right)$. This is giving a knot in $\partial X$ along with a framing of that knot. One can generalize this to links, and apply it to the case $n=4, k=2, X=D^{4}$. The resulting space $M_{L}$ is called a handle body since it is obtained from $D^{4}$ by attaching handles to it.

We have the following theorem:
Theorem 1.1.2 ([RK1]). $\partial M_{L}=\partial M_{L^{\prime}}$ (orientation preserving diffeomorphism) if and only if one can pass from $L$ to $L^{\prime}$ following a sequence of two operations, one being the blow up or down of a circle of framing $\pm 1$, the other being the band sum move between two link components.

These two moves are commonly referred to as Kirby I and Kirby II moves respectively. Kirby II corresponds in $M_{L}$ to doing a handle slide of one 4-dimensional 2-handle over another. There are 3-dimensional (resp. 4-dimensional) formalisms called 3-d Kirby Calculus (resp. 4-d Kirby calculus) [RK1] [RK2] [GS] meant to facilitate computations on handlebodies in 3 and 4 dimensions respectively. In both formalisms there are band sum moves between link components representing handle slides on handles, whence our interest in studying the behavior of the Kontsevich integral under such moves.

In what follows, we consider the second operation, the band sum move. It is defined as follows [RK1]: consider two link components $K_{i}$ and $K_{j}$ of $L$. Pick one knot, say $K_{j}$, push it off itself using its framing $\phi_{j}$ and in such a manner that it misses $L$. In this manner one
obtains $\tilde{K}_{j}$, in addition to having $K_{j}$. Then one connect sums $K_{i}$ with $\tilde{K}_{j}$ with the use of a band connecting $K_{i}$ to $\tilde{K}_{j}$, and the connected sum is denoted:

$$
\begin{equation*}
K_{i} \underset{b}{\#} \tilde{K}_{j}:=K_{i}^{\prime} \tag{1.1}
\end{equation*}
$$

where the $b$ underneath the pound sign indicates that one has performed a band sum ([RK1]). One insists that the band $b$ misses the rest of $L$. Formally what that means is that if $\gamma_{i}$ and $\gamma_{j}$ are two embeddings defining the two knots $K_{i}$ and $K_{j}$ in $S^{3}$ respectively, then if $b: I \times I \rightarrow S^{3}$ is an embedding such that $b(I \times I) \cap \gamma_{i}=b(\{0\} \times I)$ and $b(I \times I) \cap \gamma_{j}=b(\{1\} \times I)$, then one formally writes:

$$
\begin{equation*}
\underset{b}{\gamma_{i} \# \gamma_{j}}=\gamma_{i} \cup \gamma_{j}-b(\partial I \times I) \cup b(I \times \partial I) \tag{1.2}
\end{equation*}
$$

Observe [RK2] that the band is allowed to have any number of half twists in it, either left or right twists, and since the circles are oriented this will lead to a notion of subtracting or adding a circle to another, respectively. In the terminology of handle bodies, for an $n$-dimensional $k$-handle, $\partial D^{k} \times 0$ is called the attaching region. In our situation one has $n=4, k=2$, and the 2 -handle has $S^{1} \times 0$ as the attaching sphere, which is nothing but a knot component of the link under consideration. The band sum of one knot over another corresponds to a certain move between 2-handles called a handle slide. Following [GS], for $0 \leq k \leq n$, given two $n$-dimensional $k$-handles $h_{1}$ and $h_{2}$ attached to some $n$-dimensional manifold $X$ along its boundary $\partial X$ as done above, one can isotope the attaching sphere $S^{k-1} \times 0$ of $h_{1}$ in $\partial\left(X \cup h_{2}\right)$ and slide it along a disk $D^{k} \times\{p t\} \subset h_{2}$ until it comes back down to $\partial X$. Observe that $\partial\left(D^{k} \times\{p t\}\right)=S^{k-1} \times\{p t\}$ is the attaching sphere of $h_{2}$ pushed away from itself using the framing on that link component.

### 1.2 Isotopy invariants and Vassiliev invariants

One can define an equivalence relation on the category of knots called ambient isotopy [GS], denoted by $\sim$. An isotopy between embeddings $\gamma_{1}, \gamma_{2}: X \rightarrow Y$ is a homotopy $\gamma: X \times I \rightarrow Y$
through embeddings. An ambient isotopy between two knots $\gamma_{0}: S^{1} \rightarrow S^{3}$ and $\gamma_{1}: S^{1} \rightarrow S^{3}$ is an isotopy $\gamma: S^{1} \times I \rightarrow S^{3}$ through diffeomorphisms $\Gamma: S^{3} \times I \rightarrow S^{3}$ such that $\Gamma_{0}=i d_{S^{3}}$ and $\gamma_{t}=\Gamma_{t} \circ \gamma_{0}$ for each t . The above definition for isotopy carries over to the case of links.

In a first time, one is interested in isotopy classes of knots since original definitions for the concepts that we introduce presently where made in the case of knots. To distinguish one class from another, one needs a function on knots that takes different values on different classes, but which must of course be constant on equivalence classes, and such an object is rightfully called an isotopy invariant or invariant for short. Typically an invariant is valued in some abelian group $G$, and if we denote by $\mathbb{Z}$ Knots the abelian group generated by oriented knots, we write $\Gamma: \mathbb{Z}$ Knots $/ \sim \rightarrow G$ for an invariant $\Gamma$, it is a $G$-valued functional on the set of equivalence classes of knots. One can generalize this definition to the case of links: if $\mathbb{Z} L i n k s / \sim$ is the set of equivalence classes of oriented links, $G$ is an abelian group, then a link invariant $\Gamma$ will be a map $\Gamma: \mathbb{Z}$ Links $/ \sim \rightarrow G$.

Of those invariants, Vassiliev invariants [V] are of particular interest. Such invariants are functionals on $\mathbb{Z}$ Knots $/ \sim$ that are extended to be defined on singular knots whose singularities are transversal self-intersections. One first defines a positive crossing in the image of a knot to be:

where the arrows indicate the orientation on the portion of the knot that is being displayed. A negative crossing is represented as:


One represents a transversal self-intersection as follows:


At this point one can extend invariants of knots to also be defined on singular knots, with the use of the relation:

where $V$ is any knot invariant. By iterating this procedure, one can extend a knot invariant to be defined on knots with multiple double-crossings. If a knot invariant $V$ has its extension vanishing on knots that have more than $m$ self-intersections, one says that $V$ is a Vassiliev invariant of type $m$ :

where the argument of $V$ above has more than $m$ self-intersections. Due to the local nature of the extension of knot invariants to the case of singular knots, we can generalize the definition of Vassiliev invariant to the case of links. A Vassiliev invariant of links will be said to be of type $m$ if it evaluates to zero on any singular link with more than $m$ double-crossings (not necessarily same component intersections).

Among the Vassiliev invariants of type $m$, one is of particular interest: the degree $m$ part of the Kontsevich integral first introduced in $[\mathrm{K}]$. One should note that the Kontsevich
integral as it was initially defined is not strictly speaking a knot invariant, but once it is corrected as done below, then it becomes an honest knot invariant, and its final corrected form is a universal Vassiliev invariant in the sense that every finite type Vassiliev invariant factors through it [DBN2]. Before introducing this integral, one defines the algebra $\mathcal{A}[\mathrm{K}]$ in which it takes its values.

### 1.3 The algebra $\mathcal{A}$ of chord diagrams

For a singular oriented knot whose only singularities are transversal self-intersections, the preimage of each singular crossing under the embedding map defining the knot yields two distinct points on $S^{1}$. Each singular point in the image therefore yields a pair of points on $S^{1}$ that are conventionally connected by a chord for book keeping purposes. A knot with $m$ singular points will yield $m$ distinct chords on $S^{1}$. One refers to such a circle with $m$ chords on it as a chord diagram of degree $m$, the degree being the number of chords. The support of the graph is an oriented $S^{1}$, and it is regarded up to orientation preserving diffeomorphisms of the circle. More generally, for a singular oriented link all of whose singularities are doublecrossings, the preimage of each singular crossing under the embedding map defining the link yields a pair of distinct points on possibly different circles depending on whether the double crossing was on a same component or between different components of the link. One also connects points forming a pair by a chord. An $e$-component link with $m$ singular points will yield $m$ chords on $\coprod^{e} S^{1}$. One still calls such a graph a chord diagram. The support now is $\coprod^{e} S^{1}$ regarded up to orientation preserving diffeomorphism of each $S^{1}$. One denotes by $\mathcal{D}\left(\amalg^{e} S^{1}\right)$ the $\mathbb{C}$-vector space spanned by chord diagrams with support on $\amalg^{e} S^{1}$. One writes $\mathcal{D}$ for $\mathcal{D}\left(S^{1}\right)$. There is a grading on $\mathcal{D}\left(\amalg^{e} S^{1}\right)$ given by the number of chords featured in a diagram. If $\mathcal{D}^{(m)}\left(\amalg^{e} S^{1}\right)$ denotes the subspace of chord diagrams of degree $m$, then one writes:

$$
\begin{equation*}
\mathcal{D}\left(\amalg^{e} S^{1}\right)=\oplus_{m \geq 0} \mathcal{D}^{(m)}\left(\amalg^{e} S^{1}\right) \tag{1.4}
\end{equation*}
$$

One quotients this space by the 4-T relation which locally looks like:

where solid lines are intervals on $\amalg^{e} S^{1}$ on which a chord foot rests, and arrows indicate the orientation of each strand. One further quotients this space by the framing independence relation: if a chord diagram has a chord forming an arc on $S^{1}$ with no other chord ending in between its feet, then the chord diagram is set to zero. The resulting quotient space is the $\mathbb{C}$-vector space generated by chord diagrams mod the 4-T relation and framing independence and is denoted by $\mathcal{A}\left(\amalg^{e} S^{1}\right)$. One writes $\mathcal{A}$ for $\mathcal{A}\left(S^{1}\right)$ ([BN]). The grading of $\mathcal{D}\left(\amalg^{e} S^{1}\right)$ is preserved by the quotient, inducing a grading on $\mathcal{A}\left(\amalg^{e} S^{1}\right)$ :

$$
\begin{equation*}
\mathcal{A}\left(\amalg^{e} S^{1}\right)=\oplus_{m \geq 0} \mathcal{A}^{(m)}\left(\amalg^{e} S^{1}\right) \tag{1.5}
\end{equation*}
$$

where $\mathcal{A}^{(m)}\left(\amalg^{e} S^{1}\right)$ is obtained from $\mathcal{D}^{(m)}\left(\amalg^{e} S^{1}\right)$ by modding out by the 4-T and the framing independence relations. The connected sum of circles can be extended to chorded circles, thereby defining a product on $\mathcal{A}$ that one denotes by $\cdot([\mathrm{BN}])$, making it into an algebra that is associative and commutative [BN]. More generally $\mathcal{A}\left(\amalg^{e} S^{1}\right)$ is a module over $\otimes^{e} \mathcal{A}$. The Kontsevich integral will be valued in the graded completion $\overline{\mathcal{A}}\left(\amalg^{e} S^{1}\right)$ of the algebra $\mathcal{A}\left(\amalg^{e} S^{1}\right)$.

### 1.4 The original definition of the Kontsevich integral

As far as knots are concerned, we will work with Morse knots and geometric tangles, not q-tangles ([LM1], [LM2], [LM5], [LM6]), and for that purpose one considers the following decomposition of $\mathbb{R}^{3}$ as the product of the complex plane and the real line: $\mathbb{R}^{3}=\mathbb{R}^{2} \times \mathbb{R} \simeq$ $\mathbb{C} \times \mathbb{R}$, with local coordinates $z$ in the complex plane and $t$ on the real line for time. A Morse
knot $K$ is such that $t \circ K$ is a Morse function on $S^{1}$. If one denotes by $Z$ the Kontsevich integral functional on knots, if $K$ is a Morse knot, one defines ( $[\mathrm{K}],[\mathrm{BN}],[\mathrm{CL} 2]$ ):

$$
\begin{equation*}
Z(K):=\sum_{m \geq 0} \frac{1}{(2 \pi i)^{m}} \int_{t_{\text {min }}<t_{1}<\ldots<t_{m}<t_{\text {max }}} \sum_{P \text { applicable }}(-1)^{\varepsilon(P)} D_{P} \prod_{1 \leq i \leq m} \operatorname{dlog} \Delta z\left(t_{i}\right)\left[P_{i}\right] \tag{1.6}
\end{equation*}
$$

where $t_{\text {min }}$ and $t_{\text {max }}$ are the min and max values of $t$ on $K$ respectively, $P$ is an $m$-vector each entry of which corresponds to a pair of points on the image of the knot $K$. We write $P=\left(P_{1}, \ldots, P_{m}\right)$, where the $i$-th entry $P_{i}$ corresponds to a pair of points on the knot, and if we further situate these paired points at some height $t_{i}$, we can denote these two points by $z_{i}$ and $z_{i}^{\prime}$, so that we can write $\Delta z\left(t_{i}\right)\left[P_{i}\right]:=z_{i}-z_{i}^{\prime}$. One refers to such $P$ 's as pairings. We denote by $K_{P}$ the knot $K$ with $m$ pairs of points placed on it following the prescription given by $P$, and then connecting points at a same height by a chord. A pairing is said to be applicable if each entry corresponds to a pair of two distinct points on the knot, at the same height [BN]. For a pairing $P=\left(P_{1}, \cdots, P_{m}\right)$ giving the position of $m$ pairs of points on $K$, one denotes by $\varepsilon(P)$ the number of those points ending on portions of $K$ that are locally oriented down. For example if $P=\left(z(t), z^{\prime}(t)\right)$ and $K$ is locally oriented down at $z(t)$, then $z(t)$ will contribute 1 to $\varepsilon(P)$. We also define the length of $P=\left(P_{1}, \cdots, P_{m}\right)$ to be $|P|=m$. If we denote by $\iota_{K}$ the embedding defining the knot then $D_{P}$ is defined to be the chord diagram one obtains by taking the inverse image of $K_{P}$ under $\iota_{K}: D_{P}=\iota_{K}^{-1} K_{P}$. This generalizes immediately to the case of Morse links, and in this case the geometric coefficient will not be an element of $\overline{\mathcal{A}}$ but will be an element of $\overline{\mathcal{A}}\left(\coprod_{e} S^{1}\right)$ if the argument of $Z$ is an $e$-component link.

Now if one wants to make this integral into a true knot invariant, then one corrects it as follows. Consider the embedding in $S^{3}$ of the trivial knot as:


Consider the following correction $[\mathrm{K}]$ :

$$
\begin{equation*}
\hat{Z}:=Z(\Upsilon)^{-m} \cdot Z \tag{1.8}
\end{equation*}
$$

where the dot is the product on chord diagrams extended by linearity, and $m$ is a function that captures the number of maximal points of any knot $K$ that is used as an argument of $Z$. Defining $\nu:=Z(U)^{-1}$, this reads $\hat{Z}=\nu^{m} . Z$. Equivalently, we can define $\hat{Z}$ as being $Z$ with the provision that $\nu$ multiplies each maximal point of a given knot $K$ in the expression for $Z(K)$. As pointed out in [W], there are two possible corrections to the Kontsevich integral: the one we just presented, and the other one obtained by using $1-m$ as an exponent of the Kontsevich integral of the hump instead of just $-m$. In this manner the corrected version is multiplicative under connected sum, while using the above correction it behaves better under cabling operations. We will argue later that a modified version of $\hat{Z}$ using $1-m$ as an exponent of $\nu$ is a most convenient renormalization choice for band sum moves purposes and the construction of topological invariants of 3 -manifolds. In the case of links, there will be one such correction for each component of the link, with a power $m_{i}$ on the correction term for the $i$-th component, where $m_{i}$ is the number of maximal points of the $i$-th link component. Equivalently, $\hat{Z}(L)$ is the same as $Z(L)$ save that every $i$-th link component in the expression for $Z(L)$ is multiplied by $\nu^{m_{i}}, 1 \leq i \leq e$.

Observe that it is not enough to have the Kontsevich integral being valued in $\overline{\mathcal{D}} /$ fr.ind.. Indeed, by [DBN3], [ChDu], a Vassiliev invariant $V$ satisfies the 4 terms relation:


Using (1.3) twice on each term, we can express this relation as a linear combination:

with $\epsilon_{i}=1$ (resp. -1), $1 \leq i \leq 3$ for a positive crossing (resp. a negative crossing). In this relation all terms cancel except a collection of differences, each difference being an expression of the change of $V$ under the third Reidemeister move. $\hat{Z}$ being a Vassiliev invariant, it also satisfies the 4 terms relation. Now, for a chord diagram $D$, contracting all chords of $D$ one gets a knot $K_{D}$ to use the notations of [BN]. Then $\hat{Z}\left(K_{D}\right)=D+$ higher order terms [BN], [DBN2], [ChDu] using the fact that $\hat{Z}$ is a Vassiliev invariant. This enables us to rewrite the 4 terms relation for $\hat{Z}$ as a relation involving only chord diagrams. This relation is exactly the $4-\mathrm{T}$ relation. In other terms, the $4-\mathrm{T}$ relation must be imposed for $\hat{Z}$ to be invariant under the third Reidemeister move.

### 1.5 The Kontsevich Integral of tangles

One considers the generalization of the Kontsevich integral from knots to tangles as discussed in [BN], [LM1], [ChDu].

For this purpose, one considers a slightly more general algebra of chord diagrams as defined in [LM5]: For $X$ a compact oriented 1-dimensional manifold with labeled components, a chord diagram with support on $X$ is the manifold $X$ together with a collection of chords with feet on $X$. One represents such chord diagrams by drawing the support $X$ as solid lines, the graph consisting of dashed chords. One introduces an equivalence relation
on the set of all chord diagrams: two chord diagrams $D$ and $D^{\prime}$ with support on $X$ are equivalent if there is a homeomorphism $f: D \rightarrow D^{\prime}$ such that the restriction $\left.f\right|_{X}$ of $f$ to $X$ is a homeomorphism of $X$ that preserves components and orientation. One denotes by $\mathcal{A}(X)$ the complex vector space spanned by chord diagrams with support on $X$ modulo the 4- T and framing independence relations. $\mathcal{A}(X)$ is still graded by the number of chords as:

$$
\begin{equation*}
\mathcal{A}(X)=\oplus_{m \geq 0} \mathcal{A}^{(m)}(X) \tag{1.10}
\end{equation*}
$$

where $\mathcal{A}^{(m)}(X)$ is the complex vector space spanned by chord diagrams of degree $m$. Write $\overline{\mathcal{A}}(X)$ for the graded completion of $\mathcal{A}(X)$. We define a product on $\mathcal{A}(X)$ case by case. For example, if $X=\amalg^{e>1} S^{1}$, there is no well defined product defined on $\mathcal{A}(X)$. If $X=I^{N}$, the concatenation induces a well defined product on $\mathcal{A}\left(I^{N}\right)$. The product of two chord diagrams $D_{1}$ and $D_{2}$ in this case is defined by putting $D_{1}$ on top of $D_{2}$ and is denoted by $D_{1} \times D_{2}$. More generally, for $D_{i} \in \mathcal{A}\left(X_{i}\right), i=1,2, D_{1} \times D_{2}$ is well-defined if $X_{1}$ and $X_{2}$ can be glued strand-wise. One chord diagram of degree 1 we will use repeatedly is the following:

$$
\begin{equation*}
\Omega_{i j}=\left.\left.\left.\left.\right|_{1} \ldots\right|_{\mathrm{i}} \ldots \ldots\right|_{\mathrm{j}} \ldots\right|_{\mathrm{N}} \tag{1.11}
\end{equation*}
$$

For $T$ a tangle, one defines $Z(T) \in \overline{\mathcal{A}}(T)$ by:

$$
\begin{equation*}
Z(T):=\sum_{m \geq 0} \frac{1}{(2 \pi i)^{m}} \int_{t_{\min }<t_{1}<\ldots<t_{m}<t_{\max }} \sum_{P \text { applicable }}(-1)^{\varepsilon(P)} T_{P} \prod_{1 \leq i \leq m} \operatorname{dlog} \Delta z\left(t_{i}\right)\left[P_{i}\right] \tag{1.12}
\end{equation*}
$$

exactly as $Z(K)$ was defined in (1.6) with the difference that $T_{P}$ is the tangle $T$ with $m$ chords placed on it following the prescription given by $P$. Following [ ChDu ], one refers to $T_{P}$ as a tangle chord diagram ([ChDu]). One defines the Kontsevich integral of a tangle $T$ to be trivial if $Z(T)=T$. When working with links, we will sometimes omit the orientation on link components for convenience unless it is necessary to specify them. We will sometimes need the map $S$ ([LM5]) on chord diagrams: suppose $C$ is a component of $X$. If we reverse the orientation of $C$, one gets another oriented manifold from $X$ that we will denote by $X^{\prime}$
([LM5]). This induces a linear map:

$$
\begin{equation*}
S_{(C)}: \mathcal{A}(X) \rightarrow \mathcal{A}\left(X^{\prime}\right) \tag{1.13}
\end{equation*}
$$

defined by associating to any chord diagram $D$ in $\mathcal{A}(X)$ the element $S_{(C)}(D)$ obtained from $D$ by reversing the orientation of $C$ and multiplying the resulting chord diagram by $(-1)^{m}$ where $m$ is the number of vertices of $D$ ending on the component $C$. Suppose $Z(T)$ is known for some oriented tangle $T$, and $T^{\prime}$ is an oriented tangle with the same skeleton as $T$ 's, but with possible reversed orientations on some of its components. Then one can find $Z\left(T^{\prime}\right)$ by symply applying $S_{C}(Z(T))$ iteratively as many times as there are components $C$ of $T^{\prime}$ that have an orientation different from that of $T$.

### 1.6 Integral of framed oriented links

In the framed case, in which the framing independence relation is no longer imposed, upon computing the Kontsevich integral $Z(T)$ of a tangle $T$ with local extrema, one runs into computational problems. For instance, the Kontsevich integral of the following tangle:

has possible integrands $\operatorname{dlog}\left(z_{k}-z_{k+1}\right)$ in degree 1 corresponding to the chord diagram:

where $z_{k}$ and $z_{k+1}$ are local coordinates on the strands indexed by $k$ and $k+1$ respectively, and such integrands are made to vanish by virtue of the framing independence in the unframed case:

$$
\begin{equation*}
\bigcirc=0 \tag{1.16}
\end{equation*}
$$

However, in the framed setting, without this relation, one has to somehow avoid the divergence as $z_{k}-z_{k+1} \rightarrow 0$. Le and Murakami ([LM1]) solved this problem as follows. If $\omega$ is the chord diagram defined by:

then Le and Murakami define $\varepsilon^{ \pm \omega /(2 \pi i)}$ for some $\varepsilon \in \mathbb{R}$ as a formal power series expansion of $\exp \left( \pm \frac{\omega}{2 \pi i} \log \varepsilon\right)$ where they formally write:


Then they consider the following type of tangle $T$ ([LM1]):

where $T=\left(T-T_{\varepsilon}\right) \times T_{\varepsilon}$, and the diagram above defines the two tangles $T-T_{\varepsilon}$ and $T_{\varepsilon}$. We could equally have had this tangle upside down in the event that we had a minimum. One has the following result of Le and Murakami:

Proposition 1.6.1 ([LM1]). Let T be a tangle as above. Then:

$$
\begin{equation*}
Z_{f}(T)=\lim _{\varepsilon \rightarrow 0} \varepsilon^{\omega /(2 \pi i)} Z\left(T_{\varepsilon}\right) \tag{1.17}
\end{equation*}
$$

is well-defined, while if this tangle is turned upside down so that it is a local minimum, then:

$$
\begin{equation*}
Z_{f}(T)=\lim _{\varepsilon \rightarrow 0} Z\left(T_{\varepsilon}\right) \varepsilon^{-\omega /(2 \pi i)} \tag{1.18}
\end{equation*}
$$

is well-defined as well.

Observe that it is no longer true that $Z_{f}$ is an element of the completion of $\mathcal{A}(X)=$ $\mathcal{D}(X) /(4-\mathrm{T}$, fr. ind.) since one no longer imposes the framing independence relation. Rather $Z_{f}$ becomes an element of the completion $\overline{\hat{\mathcal{A}}}(X):=\overline{\hat{\mathcal{A}}(X)}$ of the algebra $\hat{\mathcal{A}}(X)=\mathcal{D}(X) / 4$-T where we have adopted the original notation $\hat{\mathcal{A}}$ from $[\mathrm{K}]$. In Chapter 2 we will fully develop the formalism of the framed Kontsevich integral.

Let $L$ be an $e$-components framed oriented link in the blackboard framing represented by a link diagram $\mathcal{D}$. Recall that we consider Morse knots and links; one considers $\mathbb{R}^{3}$ as $\mathbb{C} \times \mathbb{R}$ and one can arrange that our knots live in $\mathbb{C} \times I$. If $t$ is the variable in $I$, then one says that a knot $K$ is a Morse knot if $t(K)$ is a Morse function $([\mathrm{K}])$. Let $m_{i}$ be the number of maximal points of the $i$-th component of $\mathcal{D}$ with respect to $t$. Then following [LM5], one defines:

$$
\begin{equation*}
\hat{Z}_{f}(L)=Z_{f}(\mathcal{D}) \cdot\left(\nu^{m_{1}} \otimes \cdots \otimes \nu^{m_{e}}\right) \in \overline{\hat{\mathcal{A}}}\left(\coprod^{e} S^{1}\right) \tag{1.19}
\end{equation*}
$$

where $\nu=Z_{f}(U)^{-1}$ and:


Further, since we regard $\hat{\mathcal{A}}\left(\amalg^{e} S^{1}\right)$ as a $\otimes^{e} \mathcal{A}$-module, each $\nu^{m_{i}}$ acts only on the $i$-th component, and it does so by connected sum [LM2]. Strictly speaking, we should write:

$$
\begin{equation*}
\nu^{m_{1}} \otimes \cdots \otimes \nu^{m_{e}} \cdot=\left(\nu^{m_{1}} \cdot\right) \otimes \cdots \otimes\left(\nu^{m_{e}} \cdot\right) \tag{1.21}
\end{equation*}
$$

It is more economical to define $\hat{Z}_{f}$ as being $Z_{f}$ with the provision that in the expression for $Z_{f}(L), \nu$ multiplies each local max of each link component. Though as defined $\hat{Z}_{f}$ is already an isotopy invariant, we will use a certain renormalization of $\hat{Z}_{f}$ that is exceptionally well-behaved under band sum moves; we will define the modified $\widetilde{Z}_{f}$ to be $\hat{Z}_{f}$ with the provision that in the expression for $\hat{Z}_{f}(L)$, each link component is multiplied by $\nu^{-1}$. If one uses $\nu$ instead as a renormalizing factor, one would get $\check{Z}_{f}$ as introduced in [LM3], [LM4], [LM5], [LM6]. It is important to note that $\check{Z}_{f}$ and $\widetilde{Z}_{f}$ are essentially the same invariant. What distinguishes them is their behavior under band sum moves.

The doubling map $\Delta$ on strands defined by:

$$
\begin{equation*}
\Delta: \nmid \mapsto \nmid \uparrow \tag{1.22}
\end{equation*}
$$

induces a map $([\mathrm{LM} 4]) \Delta: \mathcal{A}(I) \rightarrow \mathcal{A}\left(I^{2}\right)$ on chord diagrams that is defined as follows on one chord:

$$
\begin{equation*}
\Delta: \nmid \ldots \ldots \ldots \mapsto+\ldots \ldots \ldots+\uparrow \ldots \ldots \ldots \tag{1.23}
\end{equation*}
$$

and for chord diagrams of degree greater than 1, we impose that the following square be
commutative and use induction:


In [LM6], the behavior of $\check{Z}_{f}$ under band sum moves for which link components of interest are locally side by side is given by:

Theorem 1.6.2. ([LM6]) Let $L$ be a framed oriented link. Suppose $K_{1}$ and $K_{2}$ are two link components of $L$, and $K_{1}$ is band summed over $K_{2}$, which one pictorially represents as:

where $K_{1}^{\prime}$ is the result of doing a band sum move of $K_{1}$ over $K_{2}$, and one denotes by $L^{\prime}$ the link obtained from $L$ after such an operation. In the above picture, we have only symbolically displayed $K_{1}$ and $K_{2}$, and not other components that may be linked to either or both components. If one writes:

$$
\begin{equation*}
\check{Z}_{f}(L)=\sum_{\substack{\text { chord } \\ \text { diagrams } X}} c_{X} X \tag{1.26}
\end{equation*}
$$

then Le and Murakami find:

$$
\begin{equation*}
\check{Z}_{f}\left(L^{\prime}\right)=\sum_{\substack{\text { chord } \\ \text { diagrams } X}} c_{X} X^{\prime} \tag{1.27}
\end{equation*}
$$

where $X$ and its corresponding chord diagram $X^{\prime}$ after the band sum move are given below:


To be specific, the map $\Delta$ doubles strands, and the chords on those strands as well. Since one operates a band sum move here, the $\Delta$ enclosed in the box means by abuse of notation ([LM2], [LM5], [LM6]) that the doubling of strands coming with the band sum move proper has been performed and the only thing left to be done is to double chords accordingly.

For band sum moves for which link components of interest are locally on top of each other, we have the following theorem which is proved later in Chapter 4:

Theorem 1.6.3. Let $L$ be a framed oriented link. Suppose $K_{1}$ and $K_{2}$ are two link components of $L$, and $K_{1}$ is band summed over $K_{2}$, which we pictorially represent as:

where $K_{1}^{\prime}$ is the result of doing a band sum move of $K_{1}$ over $K_{2}$, and one denotes by
$L^{\prime}$ the link obtained from $L$ after such an operation. In the above picture, we have only symbolically displayed $K_{1}$ and $K_{2}$, and not other components that may be linked to either or both components. If:

$$
\widetilde{Z}_{f}(L)=\sum_{\begin{array}{c}
\text { chord }  \tag{1.29}\\
\text { diagrams } X
\end{array}} c_{X} X
$$

for coefficients $c_{X}$, then we can write:

$$
\begin{equation*}
\widetilde{Z}_{f}\left(L^{\prime}\right)=\sum_{\substack{\text { chord } \\ \text { diagrams } X}} c_{X} X^{\prime} \tag{1.30}
\end{equation*}
$$

where $X$ and its corresponding chord diagram $X^{\prime}$ after the band sum move are given below:


## Chapter 2

## The Framed Kontsevich integral

### 2.1 First definition and Preliminary lemmas

We now consider the problem of having to integrate over tangles that have local extrema in the framed case, in which one no longer imposes the framing independence relation. In that case one normalizes the integral at local extrema as was done in [LM1]. First define the following tangles and chord diagrams:


$$
\begin{equation*}
\omega=\stackrel{\underset{\mu}{\cdots}, \ldots . .}{\longrightarrow} \tag{2.1}
\end{equation*}
$$



It is convenient to define a formal tangle chord diagram consisting of a single chord stretching between two strands of opposite orientations, and to call such a graph $\Omega$. This enables one to write $T_{a}^{\mu}$ above with $m$ chords on it as $T_{a}^{\mu} \times \Omega^{m}$ or $\Omega^{m} \times T_{a}^{\mu}$. We would also have $\omega$ above written simply as $T_{\mu} \times \Omega$. Define $S: \mathcal{A} \rightarrow \hat{\mathcal{A}}$ to be the standard inclusion algebra
map that maps elements of $\mathcal{A}$ to the subspace of $\hat{\mathcal{A}}$ in which all basis chord diagrams with an isolated chord have coefficient zero as in [BNGRT], [ChDu].

Definition 2.1.1 ([LM1]). We define $Z_{f}$ of the above tangle $T_{a}$ as:

$$
\begin{equation*}
Z_{f}\left(T_{a}\right)=\lim _{\mu \rightarrow 0} \overbrace{1} \times \mu^{\Omega / 2 \pi i} \times T_{\mu}^{1} \times S Z\left(T_{a}^{\mu}\right) \tag{2.3}
\end{equation*}
$$

which Le and Murakami write in a symbolic, compact form as:

$$
\begin{equation*}
Z_{f}\left(T_{a}\right):=\lim _{\mu \rightarrow 0} \mu^{\omega /(2 \pi i)} \times S Z\left(T_{a}^{\mu}\right) \tag{2.4}
\end{equation*}
$$

where in [LM1], $\omega^{m}$ is defined by:

$$
\begin{align*}
& \left.\omega^{m}=\underset{\underset{\mu}{\longrightarrow}}{\substack{\ldots \\
\hdashline \cdots \\
\hdashline}}\right\} m \tag{2.5}
\end{align*}
$$

$$
\begin{align*}
& =\overparen{\longrightarrow} \times T_{\mu}^{1} \times \Omega^{m} \tag{2.7}
\end{align*}
$$

For a local minimum, if we take our tangle $T^{a}$ to be a single local minimum, then one uses the normalization $Z_{f}\left(T^{a}\right)=\lim _{\mu \rightarrow 0} S Z\left(T_{\mu}^{a}\right) \times T_{1}^{\mu} \times \mu^{-\Omega / 2 \pi i} \times \stackrel{\leftrightarrow}{\longrightarrow}$ also written $Z_{f}\left(T^{a}\right)=\lim _{\mu \rightarrow 0} S Z\left(T_{\mu}^{a}\right) \times \mu^{-\omega /(2 \pi i)}$. In that expression $\omega$ and $T_{\mu}^{a}$ are our old $\omega$ and $T_{a}^{\mu}$ respectively, flipped upside down.

Lemma 2.1.2. For $a, b>0$, we have:

with a plus sign for strands with same orientations, and a minus sign for opposite orientations.

Proof. The proof is a simple computation:

$$
\begin{align*}
Z\left(T_{a}^{b}\right) & =\sum_{n \geq 0} \frac{1}{(2 \pi i)^{n}} \int_{0<t_{1}<\cdots<t_{n}<1}( \pm 1)^{n} \Omega^{n} \times T_{a}^{b} \times \operatorname{dlog} \Delta z\left(t_{1}\right) \cdots \operatorname{dlog} \Delta z\left(t_{n}\right)  \tag{2.8}\\
& =\sum_{n \geq 0} \frac{1}{(2 \pi i)^{n}}( \pm 1)^{n} \Omega^{n} \times T_{a}^{b} \times I_{n} \tag{2.9}
\end{align*}
$$

with:

$$
\begin{align*}
I_{n} & =\int_{0<t_{1}<\cdots<t_{n}<1} \operatorname{dlog} \Delta z\left(t_{1}\right) \cdots \operatorname{dlog} \Delta z\left(t_{n}\right)  \tag{2.10}\\
& =\int_{0<t_{2}<\cdots<t_{n}<1} \operatorname{dlog} \Delta z\left(t_{2}\right) \cdots \operatorname{dlog} \Delta z\left(t_{n}\right) \cdot \log \frac{\Delta z\left(t_{2}\right)}{a}  \tag{2.11}\\
& =\int_{0<t_{3}<\cdots<t_{n}<1} \operatorname{dlog} \Delta z\left(t_{3}\right) \cdots \operatorname{dlog} \Delta z\left(t_{n}\right) \int_{t_{2}<t_{3}} \operatorname{dlog} \Delta z\left(t_{2}\right) \log \frac{\Delta z\left(t_{2}\right)}{a}  \tag{2.12}\\
& =\int_{0<t_{3}<\cdots<t_{n}<1} \operatorname{dlog} \Delta z\left(t_{3}\right) \cdots \operatorname{dlog} \Delta z\left(t_{n}\right) \int_{t_{2}<t_{3}} \operatorname{dlog} \frac{\Delta z\left(t_{2}\right)}{a} \log \frac{\Delta z\left(t_{2}\right)}{a}  \tag{2.13}\\
& =\int_{0<t_{3}<\cdots<t_{n}<1} \operatorname{dlog} \Delta z\left(t_{3}\right) \cdots \operatorname{dlog} \Delta z\left(t_{n}\right) \cdot \frac{1}{2} \log ^{2} \frac{\Delta z\left(t_{3}\right)}{a}  \tag{2.14}\\
& =\int_{0<t_{4}<\cdots<t_{n}<1} \operatorname{dlog} \Delta z\left(t_{4}\right) \cdots \operatorname{dlog} \Delta z\left(t_{n}\right) \int_{t_{3}<t_{4}} \operatorname{dlog} \Delta z\left(t_{3}\right) \cdot \frac{1}{2} \log ^{2} \frac{\Delta z\left(t_{3}\right)}{a}  \tag{2.15}\\
& =\int_{0<t_{4}<\cdots<t_{n}<1} \operatorname{dlog} \Delta z\left(t_{4}\right) \cdots \operatorname{dlog} \Delta z\left(t_{n}\right) \cdot \frac{1}{3} \cdot \frac{1}{2} \log ^{3} \frac{\Delta z\left(t_{4}\right)}{a}  \tag{2.16}\\
& =\cdots=\frac{1}{n!} \log ^{n} \frac{\Delta z\left(t_{n}=1\right)}{a}=\frac{1}{n!} \log ^{n} \frac{b}{a} \tag{2.17}
\end{align*}
$$

so that:

$$
\begin{align*}
Z\left(T_{a}^{b}\right) & =\sum_{n \geq 0} \frac{1}{(2 \pi i)^{n}}( \pm 1)^{n} \Omega^{n} \times T_{a}^{b} \times I_{n}  \tag{2.18}\\
& =\sum_{n \geq 0} \frac{1}{(2 \pi i)^{n}}( \pm 1)^{n} \Omega^{n} \times T_{a}^{b} \times \frac{1}{n!} \log ^{n} \frac{b}{a}  \tag{2.19}\\
& =\sum_{n \geq 0} \frac{1}{n!}\left(\left( \pm \frac{\Omega}{2 \pi i}\right) \log \frac{b}{a}\right)^{n} \times T_{a}^{b}  \tag{2.20}\\
& =\sum_{n \geq 0} \frac{1}{n!}\left(\log (b / a)^{ \pm \frac{\Omega}{2 \pi i}}\right)^{n} \times T_{a}^{b}  \tag{2.21}\\
& =e^{\log \left((b / a)^{ \pm \Omega / 2 \pi i}\right)} \times T_{a}^{b}=(b / a)^{ \pm \Omega / 2 \pi i} \times T_{a}^{b} \tag{2.22}
\end{align*}
$$

Corollary 2.1.3. For $\epsilon>0$ :
with a plus sign for same orientations, a minus sign for opposite orientations.

Corollary 2.1.4.

$$
\begin{equation*}
Z\left(\stackrel{\leftrightarrow}{\rightarrow}_{\rightarrow}^{\rightarrow}\right)=\epsilon^{\mp \Omega / 2 \pi i} \times \stackrel{\leftrightarrow}{\rightarrow} \epsilon_{\stackrel{1}{\rightarrow}} \tag{2.24}
\end{equation*}
$$

with a plus sign for same orientations and a minus sign for opposite orientations.

## Lemma 2.1.5.

$$
\begin{equation*}
Z_{f}\left(T_{a}\right)=\overleftrightarrow{\leftrightarrow}^{\infty} \times S Z\left(T_{a}^{1}\right)=: \overleftrightarrow{\leftrightarrow}^{\infty} \times S Z\left(T_{a}^{1 \text {-resolved }}\right) \tag{2.25}
\end{equation*}
$$

where $T_{a}^{1 \text {-resolved }}$ is identical with the tangle $T_{a}$ except that it is seen as being analytically probed by isolated chords near the local maximum that deform it into a spout of opening width 1 , whence the name 1-resolved.

Proof. Using Definition 2.1.1 for the same tangle $T_{a}$ :

$$
\begin{align*}
& Z_{f}\left(T_{a}\right)=\lim _{\mu \rightarrow 0} \xrightarrow[\leftrightarrow]{\leftrightarrow} \times \mu^{\Omega / 2 \pi i} \times T_{\mu}^{1} \times S Z\left(T_{a}^{\mu}\right)  \tag{2.26}\\
& =\lim _{\mu \rightarrow 0} \xrightarrow[\rightarrow]{\rightarrow} \times\left(\sum_{m \geq 0} \frac{1}{m!} \frac{1}{(2 \pi i)^{m}} \Omega^{m} \log ^{m} \mu\right) \times T_{\mu}^{1} \times S Z\left(T_{a}^{\mu}\right)  \tag{2.27}\\
& =\lim _{\mu \rightarrow 0} \xrightarrow[\rightarrow]{\rightarrow} \times \sum_{m \geq 0}\left(\frac{1}{m!} \frac{1}{(2 \pi i)^{m}} \Omega^{m} \log ^{m} \mu \times T_{\mu}^{1}\right) \times S Z\left(T_{a}^{\mu}\right)  \tag{2.28}\\
& =\lim _{\mu \rightarrow 0} \xrightarrow[\rightarrow]{\leftrightarrow} \times \sum_{m \geq 0} \frac{1}{m!} \frac{1}{(2 \pi i)^{m}} T_{\mu}^{1} \times \Omega^{m} \log ^{m} \mu \times S Z\left(T_{a}^{\mu}\right)  \tag{2.29}\\
& =\lim _{\mu \rightarrow 0} \xrightarrow[\rightarrow]{\rightarrow} \times S Z(\underbrace{\stackrel{\text { rir }}{\gtrless}}_{\rightarrow \mu^{-}}) \times S Z\left(T_{a}^{\mu}\right) \tag{2.30}
\end{align*}
$$

$$
\begin{align*}
& =\underset{\rightarrow}{\leftrightarrow} \times S Z()_{\underset{a}{\rightarrow}}^{\stackrel{1}{\rightarrow}}()  \tag{2.32}\\
& =\overleftrightarrow{\rightarrow} \times S Z\left(T_{a}^{1}\right) \tag{2.33}
\end{align*}
$$

where in (2.30) we used Lemma 2.1.4, and in going from (2.31) to (2.32) we have used the invariance of the Kontsevich integral under horizontal deformations $[\mathrm{BN}],[\mathrm{ChDu}]$.

Lemma 2.1.6.

$$
\begin{equation*}
Z_{f}\left(T^{a}\right)=S Z\left(T_{1}^{a}\right) \times \stackrel{\leftrightarrow}{\overleftrightarrow{\longrightarrow}}=S Z\left(T_{1-\text { resolved }}^{a}\right) \times \stackrel{\leftrightarrow}{\overleftrightarrow{\longrightarrow}} \tag{2.34}
\end{equation*}
$$

Proof. Same as for the preceding lemma.
Remark 2.1.7. From Lemmas 2.1.5 and 2.1.6 we might want to generalize the definition of $Z_{f}$ near local extrema to:
and:

$$
\begin{equation*}
Z_{f}[\mathrm{M}]\left(T^{a}\right)=S Z()_{\underset{M}{ }}^{\stackrel{a}{\rightarrow}}() \times \bigcup=: S Z\left(T^{a \text { M-resolved }}\right) \tag{2.36}
\end{equation*}
$$

thereby defining a family of framed Kontsevich integrals $Z_{f}[\mathrm{M}]$ parametrized by $M>0$. Without loss of generality, we can focus on a local maximum. An equivalent definition to (2.35) would be:

$$
\begin{align*}
Z_{f}[\mathrm{M}]\left(T_{a}\right) & =\bigcap \times \lim _{(\mu / M) \rightarrow 0}(\mu / M)^{\Omega / 2 \pi i} \times T_{\mu}^{M} \times S Z\left(T_{a}^{\mu}\right)  \tag{2.37}\\
& =\Omega \times \lim _{\xi \rightarrow 0} \xi^{\Omega / 2 \pi i} \times T_{M \xi}^{M} \times S Z\left(T_{a}^{M \xi}\right) \tag{2.38}
\end{align*}
$$

Definition 2.1.1 being a special case thereof for which $M=1$. To see that (2.37) is welldefined, it suffices to write:

$$
\begin{aligned}
& \times \lim _{(\mu / M) \rightarrow 0}(\mu / M)^{\Omega / 2 \pi i} \times T_{\mu}^{M} \times S Z\left(T_{a}^{\mu}\right) \\
& \quad=\lim _{(\mu / M) \rightarrow 0}(M / \mu)^{-\Omega / 2 \pi i} \times T_{\mu}^{M} \times S Z\left(T_{a}^{\mu}\right) \\
& \quad=\lim _{(\mu / M) \rightarrow 0} S Z\left(T_{\mu}^{M}\right) \times S Z\left(T_{a}^{\mu}\right)
\end{aligned}
$$

where in the last step we have used Lemma 2.1.2 with $a=\mu, b=M$ and mixed orientations on the strands since we work at a local maximum. We have a concatenation of tangles $T_{\mu}^{M} \times T_{a}^{\mu}=T_{a}^{M}$ leading to the equality $S Z\left(T_{\mu}^{M}\right) \times S Z\left(T_{a}^{\mu}\right)=S Z\left(T_{a}^{M}\right)$. Using it in the last equality above, we get:

$$
\begin{align*}
& \bigcap \times \lim _{(\mu / M) \rightarrow 0}(\mu / M)^{\Omega / 2 \pi i} \times T_{\mu}^{M} \times S Z\left(T_{a}^{\mu}\right) \\
& =\int_{(\mu / M) \rightarrow 0} S Z\left(T_{a}^{M}\right)  \tag{2.39}\\
& =\bigcap \times S Z()_{\underset{a}{\rightarrow}}^{\rightarrow \rightarrow}() \tag{2.40}
\end{align*}
$$

This leads to generalizing the definition of the normalization for the value of $Z_{f}$ of local extrema.

Definition 2.1.8. For $M>0$ :

$$
\begin{equation*}
Z_{f}[\mathrm{M}]\left(T_{a}\right)=\bigcap \times S Z()_{\overleftrightarrow{a}}^{-\boldsymbol{\rightarrow}}() \tag{2.41}
\end{equation*}
$$

and:

$$
\begin{equation*}
Z_{f}[\mathrm{M}]\left(T^{a}\right)=S Z()_{\leftrightarrow}^{\rightarrow}() \times \bigcup \tag{2.42}
\end{equation*}
$$

Remark 2.1.9. By invariance of $Z$ under horizontal deformations:


Since $Z$ is defined for Morse knots, there is no ambiguity as to what:

mean; it is really

$$
\lim _{\epsilon \rightarrow 0} \rightarrow \mathfrak{F} \epsilon \text { and } \lim _{\epsilon \rightarrow 0} \rightarrow \underbrace{\epsilon}
$$

respectively. The first equality in (2.43) is justified as $Z\left(T_{a}^{b}\right)$ is well-defined. To make sense of the second equality, one would have to make sense of the values of $Z$ on those pinched tangles. The problem is the same as the one encountered for local extrema. If we formally define:

$$
\mathcal{A}=\lim _{\epsilon \rightarrow 0} \rightarrow \notin \epsilon \text { and }=\lim _{\epsilon \rightarrow 0} \rightarrow \leftarrow \epsilon
$$

then we can define:

## Definition 2.1.10.

$$
Z_{f}[\mathrm{M}](\overbrace{\underset{a}{\longrightarrow}}^{\sim})=\lambda \times S Z()_{\underset{a}{\rightarrow}}^{\rightarrow \rightarrow}()
$$

and:

$$
Z_{f}[\mathrm{M}](\underbrace{\stackrel{a}{\rightarrow}})=S Z()_{M}^{\stackrel{a}{\rightarrow}}() \times \text { ( }
$$

Remark 2.1.11. Long after this work was finished, it was pointed to the author ([DM]) that a similar construction has already been performed ([ZD]). In the present work we just generalize the work of Le, Murakami et al ([LM1]) concerning the normalization of the Kontsevich integral at local extrema. In [ZD], the problem of divergences at local extrema is circumvented by invoking a QFT technique without naming it, but which from the looks of it should be multiplicative renormalization. A scale $\mu$ is fixed (the equivalent of $M$ for us). Without loss of generality one works with a local maximum, which is opened to some width $\mu$, the opening bottleneck having a length $\epsilon$, as in:


Then:


While performing such a computation, only chords both of whose feet are on the bottleneck are allowed. This is imposed while in our picture this follows precisely because we have a bottleneck such that long chords vanish on it. Further it is understood that the bottleneck
has a width $\epsilon$. Thus in the limit the bottleneck is cut off. What we do however is retain the top of the tangle with a bottleneck and formally close it off.

One last remark is in order here. Though the following Lemma is true in the unframed case, it fails in the unframed case as the note that follows points out.

Lemma 2.1.12. ([ChDu]) For a tangle $T, Z(T)$ is invariant under a dilation of $T$ by a factor of $\Lambda$ with respect to the time axis.

Proof. We have:

$$
\begin{aligned}
& Z(T)=\sum_{m \geq 0} \frac{1}{(2 \pi i)^{m}} \int_{0<t_{1}<\cdots<t_{m}<1} \sum_{|P|=m}(-1)^{\epsilon(P)} T_{P} \cdot \prod_{1 \leq i \leq m} \operatorname{dlog} \Delta z\left(t_{i}\right)\left[P_{i}\right] \\
& =\sum_{m \geq 0} \frac{1}{(2 \pi i)^{m}} \sum_{|P|=m}(-1)^{\epsilon(P)} T_{P} \int_{0<t_{1}<\cdots<t_{m}<1} \prod_{1 \leq i \leq m} \operatorname{dlog} \Delta z\left(t_{i}\right)\left[P_{i}\right]
\end{aligned}
$$

If we squeeze the tangle $T$ uniformally with respect to the time axis by a factor of $\Lambda$ then each coordinate $z$ along $T$ is being mapped to $z / \Lambda$. Thus $\operatorname{dlog} \Delta z \mapsto \operatorname{d} \log (\Delta z / \Lambda)=\operatorname{dlog} \Delta z$ making the above multivariate integral invariant, as well as $Z(T)$ which is a sum of such integrals.

Note 2.1.13. $Z_{f}[\mathrm{M}]$ is not invariant under resizing by a factor of $\Lambda$ the way $Z$ was. For example:

is non trivial if $A \neq M$, while if we resize the tangle on the left hand side of the above equation by a factor of $\Lambda=M / A$, then we have:

$$
\begin{equation*}
Z_{f}[\mathrm{M}](\overbrace{M})=\Omega \times S Z(|\underset{M}{\leftrightarrow}|) \tag{2.46}
\end{equation*}
$$

which is trivial.

### 2.2 Properties of the framed Kontsevich integral

It is well-known that $Z$ is invariant under horizontal deformations ([BN]), which was further generalized for the parameter-free, framed Kontsevich integral $Z_{f}$ ([LM1]). We show this persists in the presence of a parameter:

Lemma 2.2.1. For $M>0, Z_{f}[\mathrm{M}]$ is invariant under horizontal deformations.
Proof. $Z_{f}[\mathrm{M}]$ is defined locally. In particular, away from local extrema $Z_{f}[\mathrm{M}]=Z$, which is invariant under horizontal deformations.

The following is the framed generalization of a well-known result in the unframed case ([BN]), also proved in the framed case for Le and Murakami's parameter-free $Z_{f}$ ([LM1]).

Proposition 2.2.2. For $M>0, Z_{f}[\mathrm{M}]$ is multiplicative.
Proof. Multiplicativity means that if $T=T_{1} \times \cdots \times T_{n}$ where local extrema and pinched extremities are located at times other than those at which a concatenation is performed, then $Z_{f}[\mathrm{M}](T)=\times_{1 \leq i \leq n} Z_{f}[\mathrm{M}]\left(T_{i}\right)$. By definition of $Z_{f}$, away from those extrema $Z_{f}=S Z$, which is multiplicative.

The following Theorem was initially presented in [LM1] for the parameter-free $Z_{f}$ defined on q-tangles, which we generalize as follows:

Theorem 2.2.3. $p Z_{f}[\mathrm{M}]=Z$ for all $M>0$, where $p: \hat{\mathcal{A}} \rightarrow \mathcal{A}$.

Proof. Fix $M>0$. By definition of $Z_{f}(L)$ for $L$ a framed oriented link with possible pinched extremities, for a given tangle chord diagram and its corresponding coefficient, the contribution of long chords to that coefficient is computed from $L$ punctured (c.f. (2.37) and (2.38)) and is thus independent of $M$. It reproduces the value one would obtain from $Z(L)$ in the unframed case. Isolated chords near an extremum are projected out by $p$. The result follows.

The following two Theorems were originally presented in [LM1] for the parameter-free $\hat{Z}_{f}$.

Theorem 2.2.4. For $M>0, \hat{Z}_{f}[\mathrm{M}]$ is multiplicative.
Proof. This follows from $Z_{f}[\mathrm{M}]$ being multiplicative (Proposition 2.2.2) and the definition of $\hat{Z}_{f}[\mathrm{M}]$.

Theorem 2.2.5. For $M>0, \hat{Z}_{f}[\mathrm{M}]$ is an isotopy invariant.
Proof. $\hat{Z}_{f}[\mathrm{M}]$ is invariant under all the Reidemeister moves since $Z$ is, except the straightening of humps (henceforth referred to as the cancellation of critical points move) discussed in (2.51) and the following move:


We have to show $Z_{f}[\mathrm{M}]$ is invariant under this move, and by definition of $\hat{Z}_{f}[\mathrm{M}]$ it will follow that so is this latter. By definition of $Z_{f}[\mathrm{M}]$ :


As far as non isolated chords are concerned, this amounts to computing:

which is invariant under the unframed counterpart to the move depicted in (2.47) and is
therefore invariant under that move as well. For isolated chords and mixed chords made up of isolated and long chords, we use:


Under the move (2.47) the strands are rotated $180^{\circ}$ with respect to the time axis, and complex coordinates are switched accordingly to leave the computation of the above limit invariant by virtue of the fact that $\operatorname{dlog}\left(z-z^{\prime}\right)=\operatorname{dlog}\left(z^{\prime}-z\right)$.

We now consider the cancellation of critical points move. Since $Z_{f}[\mathrm{M}]$ is invariant under all moves except this particular move, we can drag a hump along a component until it reaches a local max. Then squeezing this local max into a needle along with the hump on
it, we essentially have to consider the following:


Thus multiplying by $\nu^{2}$ on both sides, we have the invariance of $\hat{Z}_{f}$ under the cancellation of critical points move.

Proposition 2.2.6. For $L$ a link, $Z_{f}[\mathrm{M}](L)$ is independent of $M>0$ and we denote it by $Z_{f}(L)$.

Proof. A change in scale in $\mathbb{R}^{3}$ by a factor of $\lambda \neq 0$ induces a dilation map $d_{\lambda}$ on tangles,
which in turn induces a map $d_{\lambda}^{*}$ on $Z_{f}[\mathrm{M}]$. We write:

$$
\begin{align*}
d_{\lambda}^{*} Z_{f}[\mathrm{M}]\left(T_{a}\right)=Z_{f}[\mathrm{M}]\left(d_{\lambda} T_{a}\right) & =Z_{f}[\mathrm{M}]\left(T_{\lambda a}\right)  \tag{2.56}\\
& =\overbrace{\text { M }} \times S Z\left(T_{\lambda a}^{M}\right)  \tag{2.57}\\
& =S Z\left(d_{\lambda} T_{a}^{M / \lambda}\right)  \tag{2.58}\\
& =\overbrace{\text { M/ }} \times S Z\left(T_{a}^{M / \lambda}\right)  \tag{2.59}\\
& =\overbrace{\sim} \times S\left(T_{a}^{M / \lambda}\right)  \tag{2.60}\\
& =Z_{f}[M / \lambda]\left(T_{a}\right) \tag{2.61}
\end{align*}
$$

where in the next to last step we have used the invariance of the original Kontsevich integral under dilations. It is not difficult to see that for a tangle $T$ with only one local maximum, the above computation generalizes to:

$$
\begin{align*}
d_{\lambda}^{*} Z_{f}[\mathrm{M}](T) & =Z_{f}[\mathrm{M}]\left(d_{\lambda} T\right)  \tag{2.62}\\
& =\bigcap_{\times S Z}\left(\left(d_{\lambda} T\right)^{\mathrm{M}-\text { resolved }}\right)  \tag{2.63}\\
& =\bigcap \times S Z\left(d_{\lambda}\left(T^{M / \lambda-\text { resolved }}\right)\right)  \tag{2.64}\\
& =\bigcap \times S Z\left(T^{M / \lambda \text {-resolved }}\right)  \tag{2.65}\\
& =Z_{f}[M / \lambda](T) \tag{2.66}
\end{align*}
$$

Likewise, one would show that for a tangle $T$ with a unique local minimum, one has $d_{\lambda}^{*} Z_{f}[\mathrm{M}](T)=Z_{f}[M / \lambda](T)$. The same computations essentially carry over to the case of tangles with pinched extremities. On associators and crossings, for any $M>0, Z_{f}[\mathrm{M}]=Z$. Thus on elementary tangles, as well as for tangles with pinched extremities, for $M>0$ we have:

$$
\begin{equation*}
d_{\lambda}^{*} Z_{f}[\mathrm{M}]=Z_{f}[M / \lambda] \tag{2.67}
\end{equation*}
$$

Let $K$ be a knot that we decompose as a concatenation of $l$ elementary tangles as $K=\times_{1 \leq i \leq l} T_{i}$. Then:

$$
\begin{align*}
d_{\lambda}^{*} Z_{f}[\mathrm{M}](K) & =Z_{f}[\mathrm{M}]\left(d_{\lambda} K\right)  \tag{2.68}\\
& =Z_{f}[\mathrm{M}]\left(d_{\lambda}\left(\times_{i} T_{i}\right)\right)  \tag{2.69}\\
& =Z_{f}[\mathrm{M}]\left(\times_{i}\left(d_{\lambda} T_{i}\right)\right)  \tag{2.70}\\
& =\times_{1 \leq i \leq l} Z_{f}[\mathrm{M}]\left(d_{\lambda} T_{i}\right)  \tag{2.71}\\
& =\times_{i} Z_{f}[M / \lambda]\left(T_{i}\right)  \tag{2.72}\\
& =Z_{f}[M / \lambda]\left(\times_{i} T_{i}\right)  \tag{2.73}\\
& =Z_{f}[M / \lambda](K) \tag{2.74}
\end{align*}
$$

Thus $d_{\lambda}^{*} Z_{f}[\mathrm{M}]=Z_{f}[M / \lambda]$ on knots as well, and this generalizes easily to the case of links. Now observe that by isotopy invariance of $\hat{Z}_{f}[\mathrm{M}]$ we have:

$$
\begin{equation*}
\hat{Z}_{f}[\mathrm{M}]\left(d_{\lambda} K\right)=\hat{Z}_{f}[\mathrm{M}](K) \tag{2.75}
\end{equation*}
$$

from which it follows that:

$$
\begin{equation*}
Z_{f}[\mathrm{M}]\left(d_{\lambda} K\right)=Z_{f}[\mathrm{M}](K) \tag{2.76}
\end{equation*}
$$

and coupled with (2.74), we get $Z_{f}[\mathrm{M}](K)=Z_{f}[M / \lambda](K)$ for all $M>0, \lambda>0$, thus $Z_{f}[\mathrm{M}](K)$ is independent of $M$, and we denote it by $Z_{f}(K)$. Likewise, for a link $L$, one would show following the same line of reasoning that $Z_{f}[\mathrm{M}](L)$ is independent of $M$, and we denote it by $Z_{f}(L)$.

Corollary 2.2.7. For $L$ a link, $\hat{Z}_{f}[\mathrm{M}](L)$ is independent of $M>0$ and we denote it by $\hat{Z}_{f}(L)$, as it coincides with Le and Murakami's $\hat{Z}_{f}(L)$.

Proof. This follows readily from the definition of $\hat{Z}_{f}[\mathrm{M}]$ and the previous Proposition.

## Chapter 3

## Fundamental results

In this section we essentially discuss two very important results. The first, implicit in much of the work that has been done on the Kontsevich integral, is the long chords Lemma. A proof in the unframed case can be found in $[\mathrm{ChDu}]$. We show it is also true in the framed case, and not only for knots but for links as well. The second result is the statement $\Delta \hat{Z}_{f}=\hat{Z}_{f}(\Delta)([\mathrm{LM} 2])$ which will require a few lemmas before we proceed to proving it.

### 3.1 The long chords lemma

To discuss the long chord lemma, we will follow the notations of $[\mathrm{ChDu}]$. Suppose $L$ is a link, with a distinguished tangle $T$ as in:

located between times $t_{\min }$ and $t_{\max }$. The idea is that if one shrinks the tangle $T$, in the limit when it is shrunk to zero size the chords between it and the rest of the link $L$ in the
expression for $Z_{f}(L)$ go away. The exact statement is that if one denotes by $Z_{f, T}(L)$ that part of $Z_{f}(L)$ where in the horizontal strip between the times $t_{\min }$ and $t_{\max }$ one only has chords both feet of which are either on $T$ or not at all, then $Z_{f}(L)=Z_{f, T}(L)$.

### 3.1.1 The Kontsevich integral of $n$ unlinked circles

The case of $n$ unlinked circles is a simple instance where computations can be done by hand. In this case, the long chords Lemma is true because all tangle chord diagrams have zero coefficients. The case of $n$ unlinked circles corresponds to the simple case where both $T$ and $T^{\prime}$ are trivial. By invariance of $Z$ under horizontal deformations, we essentially have to show that $n$ unlinked circles as in:

have trivial Kontsevich integral. We first show the result in the unframed case and then prove it in the framed case.

Lemma 3.1.1.1. The Kontsevich integral of $n$ unlinked circles as above is trivial.

Proof. We first consider the types of chords we can have and their corresponding log differentials. It is convenient to present the $n$ strands as circles equally spaced by a distance $\delta$ as
in:

with a similar setting for bottom strands, which number $n$ as well. For chords with support on those strands, we have long chords stretching from a $q$-th strand on one side to a $p$-th strand on the other side of the knot, $0 \leq p, q \leq n-1$, short chords stretching between the $q$-th and $p$-th strands, $0 \leq q, p \leq n-1$, both on a same side of the knot, and self chords. We first consider a long chord stretching from the $q$-th strand on the left to the $p$-th one on the right as in:


Such a chord diagram we denote by $C_{(q, p)}^{(n)+}$, where the plus sign indicates that we consider the upper portion of $n$ unlinked circles, and $(q, p)$ indicates that we consider one chord only stretching from the $q$-th strand on the left to the $p$-th strand on the right. At height $t$, the separation from one chord foot to the other is $\Delta z(t)=\sqrt{(a+p \delta)^{2}-t^{2}}+\sqrt{(a+q \delta)^{2}-t^{2}}$ from which we get:

$$
\begin{equation*}
\operatorname{dlog} \Delta z(t)=-\frac{t d t}{\sqrt{(a+q \delta)^{2}-t^{2}} \cdot \sqrt{(a+p \delta)^{2}-t^{2}}} \tag{3.4}
\end{equation*}
$$

For a chord stretching from the $p$-th strand to the $q$-th strand on the right, that is for $C_{(p, q)}^{(n)+}$ :

we get the same log differential. We also have short chords between strands on the left as in the following tangle chord diagram that we denote by $C_{(q p,-)}^{(n)+}$ :

with chords of length $\sqrt{(a+q \delta)^{2}-t^{2}}-\sqrt{(a+p \delta)^{2}-t^{2}}$, leading to a $\log$ differential:

$$
\begin{equation*}
\operatorname{dlog} \Delta z(t)=\frac{t d t}{\sqrt{(a+q \delta)^{2}-t^{2}} \cdot \sqrt{(a+p \delta)^{2}-t^{2}}} \tag{3.7}
\end{equation*}
$$

and short chords on strands on the right side of the knot as in the chord diagram $C_{(-, p q)}^{(n)+}$ :

with the same log differential. We also have to consider self-chords that do not cancel if they are not isolated, such as in the chord diagram $C_{(q, q)}^{(n)+}$ :


These chords have length $2 \sqrt{(a+q \delta)^{2}-t^{2}}$ and thus yield a log differential:

$$
\begin{equation*}
\mathrm{d} \log \Delta z(t)=-\frac{t d t}{(a+q \delta)^{2}-t^{2}} \tag{3.10}
\end{equation*}
$$

We obtain similar chord lengths and corresponding log differentials for chords stretching between strands at the bottom of the knot. What we have at this point is that the log differentials for each chord are odd functions times the differential $d t$.

We see that we essentially have three possibilities for chords: long chords, short chords, and non-isolated self-chords. For long chords and short chords we assume that $q>p$ for computational purposes. Chord diagrams are graded by their number of chords, which are numbered from the bottom up. The $k$-th chord at time $t_{k}$ is stretching between the $q_{k}$-th strand and the $p_{k}$-th one. The Kontsevich integral of $n$ unlinked trivial knots is a sum over $m \geq 0$ of terms proportional to integrals of the form:

$$
\begin{equation*}
\int_{\substack{-a-p_{m-1} \delta<t_{m-1}<\min \left(a+p_{m-1} \delta, t_{m}\right) \\-a-p_{1} \delta<t_{1}<\min \left(a+p_{1} \delta, t_{2}\right)}}^{-a-p_{m} \delta<t_{1}<a+p_{m} \delta_{m}} \prod_{1 \leq i \leq m} \mathrm{~d} \log \Delta z\left(t_{i}\right) \tag{3.11}
\end{equation*}
$$

times some tangle chord diagrams. For convenience, we write this integral as:

$$
\begin{equation*}
\int_{-a-p_{m} \delta<t_{m}<a+p_{m} \delta_{m}} \mathrm{~d} \log \triangle z\left(t_{m}\right) \Psi\left(t_{m}\right) \tag{3.12}
\end{equation*}
$$

where $\Psi\left(t_{m}\right)$ comes from the integration of $m-1 \log$ differentials. As we have shown above, the $\log$ differential is odd, the region of integration is symmetric with respect to zero, so if we can show that $\Psi\left(t_{m}\right)$ is an even function of $t_{m}$ then we would have this $m$-th term in the Kontsevich integral to be vanishing, and this for all such terms and for all $m \geq 0$, so this would show the result. Thus it suffices to show that $\Psi$ is an even function. We have:

$$
\begin{equation*}
\Psi\left(t_{2}\right)=\int_{-a-p_{1} \delta<t_{1}<\min \left(a+p_{1} \delta, t_{2}\right)} d \log \Delta z\left(t_{1}\right) \tag{3.13}
\end{equation*}
$$

If $t_{2}>a+p_{1} \delta$, then $\Psi\left(t_{2}\right)=\int_{-a-p_{1} \delta}^{a+p_{1} \delta} \mathrm{~d} \log \Delta z\left(t_{1}\right)=0=\Psi\left(-t_{2}\right)$. If $t_{2}<a+p_{1} \delta$ then $\Psi\left(t_{2}\right)=\log \Delta z\left(t_{2}\right)-\log \Delta z\left(-a-p_{1} \delta\right)$ while:

$$
\begin{align*}
\Psi\left(-t_{2}\right) & =\int_{-a-p_{1} \delta<t_{1}<-t_{2}} \operatorname{dlog} \Delta z\left(t_{1}\right)  \tag{3.14}\\
& =\log \Delta z\left(-t_{2}\right)-\log \Delta z\left(-a-p_{1} \delta\right)  \tag{3.15}\\
& =\log \Delta z\left(t_{2}\right)-\log \Delta z\left(-a-p_{1} \delta\right)=\Psi\left(t_{2}\right) \tag{3.16}
\end{align*}
$$

Thus $\Psi\left(t_{2}\right)$ is an even function of $t_{2}$. Suppose $\Psi\left(t_{m-1}\right)$ is an even function of $t_{m-1}$. Then we have:

$$
\begin{equation*}
\Psi\left(t_{m}\right)=\int_{-a-p_{m-1} \delta<t_{m-1}<\min \left(a+p_{m-1} \delta, t_{m}\right)} d \log \Delta z\left(t_{m-1}\right) \Psi\left(t_{m-1}\right) \tag{3.17}
\end{equation*}
$$

an integral of an odd integrand, which we denote by $\Gamma$. Let $A=a+p_{m-1} \delta$. If $t_{m}>A$ then $\Psi\left(t_{m}\right)=\int_{-A}^{A} \Gamma(t) d t=0=\Psi\left(-t_{m}\right)$. If $t_{m}<A$ then $\Psi\left(-t_{m}\right)=\int_{-A}^{-t_{m}} \Gamma(t) d t=\int_{A}^{t_{m}} \Gamma(u) d u$, while:

$$
\begin{align*}
\Psi\left(t_{m}\right) & =\int_{-A<t<t_{m}} \Gamma(t) d t  \tag{3.18}\\
& =\int_{A}^{-t_{m}} \Gamma(u) d u  \tag{3.19}\\
& =\int_{A}^{t_{m}} \Gamma(u) d u+\int_{t_{m}}^{-t_{m}} \Gamma(u) d u  \tag{3.20}\\
& =\int_{A}^{t_{m}} \Gamma(u) d u=\Psi\left(-t_{m}\right) \tag{3.21}
\end{align*}
$$

and $\Psi\left(t_{m}\right)$ is therefore an even function. This being true for any $m \geq 0$ it is true for all $m \geq 0$, so $\Psi$ is an even function and we have shown the result. This proof is using basic
properties of the Kontsevich integral. We can give a second proof that uses some invariance properties of the Kontsevich integral per se. It suffices to invoke the invariance of $Z$ under moves other than the cancellation of critical points move as well as its multiplicative property to write:

$$
\begin{align*}
Z(\underbrace{S^{1}+\cdots+S^{1}}_{n}) & =n Z\left(S^{1}\right)  \tag{3.22}\\
& =n Z(\bigcap) \times Z(\longrightarrow)  \tag{3.23}\\
& =n \bigcap \times \longrightarrow=n S^{1} \tag{3.24}
\end{align*}
$$

Lemma 3.1.1.2. The framed Kontsevich integral of $n$ unlinked circles is trivial.

Proof. As for the preceding lemma, we will give two proofs. The first one uses the definition of the framed Kontsevich integral, the second is much shorter and just uses some invariance properties of the framed Kontsevich integral. For the first proof, we consider a segmentation of $n$ circles at each local extremum and use the multiplicativity of the framed Kontsevich
integral.

$$
\begin{align*}
& Z_{f}\left(n S^{1}\right)=Z_{f}(\xlongequal{\Longrightarrow(0)})  \tag{3.25}\\
& =Z_{f}[\mathrm{M}](\sim) \times Z_{f}[\mathrm{M}](/ \sim) \times \\
& \cdots \times Z_{f}[\mathrm{M}]((\cdots(\bigcirc) \cdots)) \times \cdots \\
& \times Z_{f}[\mathrm{M}](\downarrow /) \times Z_{f}[\mathrm{M}](>)  \tag{3.26}\\
& =\overbrace{\infty}^{n} \times \lim _{\epsilon_{1}, \cdots, \epsilon_{2 n} \rightarrow 0} S Z(\overbrace{-}^{\infty} \\
& \cdots \times S Z(\text { ? } \\
& S Z(\text { C }  \tag{3.27}\\
& =S Z\left(n S^{1}\right)+\text { mixed chords } \tag{3.28}
\end{align*}
$$

The first term is equal to $n S Z\left(S^{1}\right)=n S^{1}$ by the preceding lemma, and mixed chords here means that this is the part of $Z_{f}\left(n S^{1}\right)$ with at least one integrated chord that would have been isolated in the unframed case. One new problem arises however from the fact that
those chords can now escape through a spout. Supposing we consider such a chord at a local maximum with other long chords below that will prevent the isolated chord to go all the way down through the spout at the local minimum, in this instance we can no longer invoke a symmetry argument. A paradigm example is the following:

whose coefficient modulo sign and a power of $1 / 2 \pi i$ is:

$$
\begin{align*}
& \int \begin{array}{c}
-a-p_{m+1} \delta<t_{m+1}<a+p_{m+1} \delta+A \\
-a-p_{m+1} \delta<t_{m}<\min \left(a+p_{m+1} \delta, t_{m+1}\right) \\
-a-p_{m+1} \delta<t_{1}<\min \left(a+p_{m+1} \delta, t_{2}\right)
\end{array} \prod_{1 \leq i \leq m+1} d \log \Delta z\left(t_{i}\right)  \tag{3.30}\\
& -\quad .
\end{align*}
$$

where $\Delta z\left(t_{m+1}\right)$ corresponds to the top chord which is escaping through the puncture, until it reaches the top of the open tangle at the top in a time $A$. All the other chords stop at $a+p_{m+1} \delta$. Using the notations of the previous proof, we consider:

$$
\begin{equation*}
\int_{-a-p_{m+1} \delta<t_{m}<\min \left(a+p_{m+1} \delta, t_{m+1}\right)}^{-a-p_{m+1} \delta<t_{m+1}<a+p_{m+1} \delta+A} \operatorname{dlog} \Delta z^{\prime}\left(t_{m+1}\right) \operatorname{dlog} \Delta z\left(t_{m}\right) \Psi\left(t_{m}\right) \tag{3.31}
\end{equation*}
$$

with $\Delta z^{\prime}\left(t_{m+1}\right)$ the length of the isolated chord. We write this integral as:

$$
\begin{align*}
\int_{-a-p_{m+1} \delta}^{a+p_{m+1} \delta} & d \log \Delta z^{\prime}\left(t_{m+1}\right) \int_{-a-p_{m+1} \delta}^{t_{m+1}} \mathrm{~d} \log \Delta z\left(t_{m}\right) \Psi\left(t_{m}\right) \\
& +\int_{a+p_{m+1} \delta}^{a+p_{m+1} \delta+A} \mathrm{~d} \log \Delta z^{\prime}\left(t_{m+1}\right) \int_{-a-p_{m+1} \delta}^{a+p_{m+1} \delta} \mathrm{~d} \log \Delta z\left(t_{m}\right) \Psi\left(t_{m}\right) \tag{3.32}
\end{align*}
$$

The second integral vanishes because the region of integration in $t_{m}$ is symmetric, dlog $\Delta$ $z\left(t_{m}\right)$ is odd, and $\Psi\left(t_{m}\right)$ is an even function from the preceding proof. The first integral can
be rewritten as:

$$
\begin{equation*}
\int_{-a-p_{m+1} \delta}^{a+p_{m+1} \delta} \mathrm{~d} \log \Delta z^{\prime}\left(t_{m+1}\right) \Psi\left(t_{m+1}\right) \tag{3.33}
\end{equation*}
$$

with $\Psi\left(t_{m+1}\right)$ even and $\operatorname{dlog} \Delta z\left(t_{m+1}\right)$ odd since the domain of integration is the circle. The region of integration being symmetric, the integral vanishes. For the more general case of having more than one isolated chord on top, say $n$ of them, we would get $2+n-1$ integrals of either of the forms discussed above, which have been showed to vanish, therefore giving a vanishing contribution to such more complicated mixed tangle chord diagrams. Thus any mixed tangle chord diagram with what would have been isolated chords in the unframed case on top has a vanishing coefficient. We would similarly show that the same result holds for isolated chords at the bottom. Thus all tangle chord diagrams with mixed chords vanish. We are left with $Z_{f}\left(n S^{1}\right)=n S^{1}$. The second proof uses the invariance of $Z_{f}$ under moves other than the cancellation of critical points move. Using this property, we can write the framed Kontsevich integral of $n$ unlinked circles as the framed Kontsevich integral of $n$ parallel circles of same radius $a>M$ separated by a distance $\Lambda \gg a$. For $\Lambda$ very large, chords stretching from one circle to another yield a negligible contribution to the Kontsevich integral. Equivalently, one could have put all $n$ circles on top of each other and we would not have had chords between them in the expression for $Z_{f}\left(n S^{1}\right)$, thereby yielding the same result. We are left with:

$$
\begin{align*}
& Z_{f}\left(n S^{1}\right)=Z_{f}[\mathrm{M}]\left(n S^{1}\right)=\overbrace{}^{n} \times\left(n S Z\left({\underset{\sim}{M}}_{\stackrel{M}{\rightarrow}}^{)}\right)\right) \times \checkmark^{n}  \tag{3.34}\\
& =\overbrace{}^{n} \times n \overbrace{M}^{\stackrel{M}{\rightarrow}}) \times \cup^{n}=n S^{1} \tag{3.35}
\end{align*}
$$

### 3.1.2 $Z_{f}(L)$ : General case

We now tackle the general case: $T$ is non trivial.

Theorem 3.1.2.1 ([ChDu], [CL]). For a knot $K$ and a distinguished tangle $T: Z(K)=$ $Z_{T}(K)$

Proposition 3.1.2.2. For $L$ a link and $T$ a distinguished tangle, then $Z_{f}(L)=Z_{f, T}(L)$.
Proof. The preceding Theorem is valid for links as well. Indeed, the proof consists in squeezing the tangle $T$ in a window of width $\epsilon$ and to let $\epsilon \rightarrow 0$. In the limit of very small values of $\epsilon$, chords between $T$ and the rest of $L$ have a negligible contribution to $Z_{f}(L)$. More generally, tangle chord diagrams with at least one such long chord have a negligible contribution. In the framed case, for $L$ with $2 n$ local extrema, it suffices to write:

$$
\begin{equation*}
Z_{f}(L)=Z_{f}[\mathrm{M}](L)=\Omega^{n} \times S Z\left(L^{M \text {-resolved }}\right) \times \circlearrowleft^{n} \tag{3.36}
\end{equation*}
$$

If we consider $L^{M \text {-resolved }}$, then the part of $L$ that constitutes $T$ is also $M$-resolved. Thus by Theorem 3.1.2.1, we have:

$$
\begin{equation*}
Z_{f}(L)=\overbrace{}^{n} \times S Z_{T^{\mathrm{M} \text {-resolved }}}\left(L^{M \text {-resolved }}\right) \times \circlearrowleft^{n} \tag{3.37}
\end{equation*}
$$

which is equal to $Z_{f, T}(L)$ by definition of this latter.

### 3.2 The statement $\Delta \hat{Z}_{f}=\hat{Z}_{f} \Delta$

### 3.2.1 The doubling operator $\Delta$

Before setting to prove the relation $\hat{Z}_{f} \Delta=\Delta \hat{Z}_{f}([L M 2],[L M 3],[L M 4],[L M 5],[L M 6])$, we remind the reader what $\Delta$ means. We first recall the well-known result $Z(\Delta T)=\Delta Z(T)$ if $T$ is a tangle without local extrema [LM2]. We illustrate this equality with the simple tangle $T$ :


We now double the right strand of $T$ to get another tangle $\Delta T$ whose Kontsevich integral
we would like to compute. Consider:


We will display only the degree 1 terms of $Z(\Delta T)$ :

$$
\left.Z(\Delta T)=\cdots \pm \frac{1}{2 \pi i} \log \frac{a}{b} / \cdots \cdots \cdots\right)\left\langle\frac{1}{2 \pi i} \log \frac{a+\delta}{b+\delta} / \cdots \cdots \cdots+\cdots+\cdots\right.
$$

while:

$$
\begin{align*}
& \Delta Z(T)=\cdots+\Delta\left( \pm \frac{1}{2 \pi i} \log \frac{a}{b} / \cdots \cdots \cdots\right)+\cdots  \tag{3.40}\\
& =\cdots \pm \frac{1}{2 \pi i} \log \frac{a}{b} \cdot \Delta / \cdots \cdots+\cdots+\cdots  \tag{3.41}\\
& =\cdots \pm \frac{1}{2 \pi i} \log \frac{a}{b} / \cdots \cdots \cdots+\cdots+\frac{1}{2 \pi i} \log \frac{a}{b} / \cdots \cdots \cdots \cdot{ }^{\prime} / \cdots \tag{3.42}
\end{align*}
$$

which shows that $Z(\Delta T)=\Delta Z(T)$ is true only in the limit $\delta \rightarrow 0$. Thus in general writing $\Delta T$ means the replica of the tangle $T$ is geometrically infinitesimally close to $T$ itself and thus can be considered to be analytically coincident with the domain of $T$. We adopt the notation $\Delta(\delta)$ to denote the map on tangles that creates a replica of a tangle a distance $\delta$ off of it. Some ambiguity arises concerning the position of the replica, whether it should be to the right or to the left of an original tangle without local extrema, and what convention is to be used when there are local extrema. Typically it will be clear from the context what $\Delta(\delta)$ does. Irrespective of the particular $\Delta(\delta)$, one has:

$$
\begin{equation*}
Z(\Delta T)=Z\left(\lim _{\delta \rightarrow 0} \Delta(\delta) T\right) \tag{3.43}
\end{equation*}
$$

which is independent of the distance between the original tangle and its copy. We will call $\Delta T$ a doubled tangle and $\Delta(\delta) T$ a $\delta$-doubled tangle for $\delta \neq 0$.

Lemma 3.2.1.1. ([LM2]) $Z(\Delta T)=\Delta Z(T)$ if $T$ has no local extrema, nor does it have pinched extremities.

Proof. It suffices to write:

$$
\begin{equation*}
Z(\Delta T)=Z\left(\lim _{\delta \rightarrow 0} \Delta(\delta) T\right)=\Delta Z(T) \tag{3.44}
\end{equation*}
$$

Warning 3.2.1.2. The statement $Z(\Delta T) \neq \Delta Z(T)$ if $T$ has a local extremum is properly written $Z(\Delta(\delta) T) \neq \Delta Z(T)$ if $\delta \neq 0$. Pinched extremities have not been considered in the literature so far, and a same remark would hold in that case. In words, the vanishing of the coefficient of chords with feet on parallel $\delta$-doubled strands is no longer true in case we double a local extremum or a pinched extremity. Without loss of generality, consider a tangle consisting of a single local maximum, its feet being separated by a distance $2 a$. By invariance of $Z$ under horizontal deformations including tweaking local extrema into needles, it suffices to consider $Z$ of a semi-circle of radius $a$. At the first order already, the coefficient of:

is:

$$
\begin{align*}
\frac{1}{(2 \pi i)} & \log \frac{\sqrt{(a+\delta)^{2}-a^{2}}-\sqrt{a^{2}-a^{2}}}{\delta} \\
& =\frac{1}{(4 \pi i)} \log \frac{2 a+\delta}{\delta} \tag{3.46}
\end{align*}
$$

The chord diagram with a same short chord on the right hand side has the same coefficient. For a $\delta$-doubled local max, we also have long chords of the form:

with the following coefficient in the Kontsevich integral, using the same orientations as above for consistency:

$$
\begin{align*}
-\frac{1}{(2 \pi i)} & \log \frac{\sqrt{(a+\delta)^{2}-a^{2}}+\sqrt{a^{2}-a^{2}}}{2 a+\delta} \\
& =-\frac{1}{(2 \pi i)} \log \frac{\sqrt{\delta} \sqrt{2 a+\delta}}{2 a+\delta} \\
& =-\frac{1}{(4 \pi i)} \log \frac{\delta}{2 a+\delta} \\
& =\frac{1}{(4 \pi i)} \log \frac{2 a+\delta}{\delta} \tag{3.48}
\end{align*}
$$

The chord diagram with a long chord:

has the same coefficient. What we have so far is that long chords and short chords have
the same coefficient, and thus the first order term of the Kontsevich integral of a $\delta$-doubled local max is equal to $(1 /(\pi i)) \log (2 a+\delta) / \delta$, which is non zero. What we have shown is that the Kontsevich integral of a $\delta$-doubled local max is non trivial, and likewise the Kontsevich integral of a $\delta$-doubled local min is non trivial either. This means in particular that $\Delta Z(T)=Z(\Delta(\delta) T)$ is not true of tangles $T$ with local extrema since the statement is not even true in first order.

### 3.2.2 The relation between $\Delta \hat{Z}_{f}$ and $\hat{Z}_{f} \Delta$

By definition of $\Delta$, we can now make sense of $Z_{f}(\Delta K)$. The following Lemma is a statement about $Z_{f}[\mathrm{M}](\Delta(\delta) T)$. The Proposition that follows proves that $Z_{f}(\Delta K)=\Delta Z_{f}(K)$.

Lemma 3.2.2.1. ([LM2]) If $T$ has a local extremum, $M>0, \delta>0$, then $Z_{f}[\mathrm{M}](\Delta(\delta) T) \neq$ $\Delta Z_{f}[\mathrm{M}](T)$.

Proof. It suffices to consider the tangle:

$$
\begin{equation*}
T_{a}=\overleftrightarrow{\longrightarrow} \tag{3.50}
\end{equation*}
$$

a semi-circle of radius $a$. We have:

$$
\begin{align*}
Z_{f}[\mathrm{M}]\left(\Delta(\delta) T_{a}\right) & =\overbrace{}^{2} \times S Z\left(\left(\Delta(\delta) T_{a}\right)^{M-\text {-resolved }}\right)  \tag{3.51}\\
& =S Z\left(\Delta(\delta) T_{a}\right)+\text { other terms } \tag{3.52}
\end{align*}
$$

where "other terms" are tangle chord diagrams with at least one chord that would have been isolated on $\Delta(\delta) T_{a}$ in the unframed case. Thus in degree 1 , the only tangle chord diagram with a chord between $T_{a}$ and its replica is coming from $S Z\left(\Delta(\delta) T_{a}\right)$ where we have 4 occurrences of such a term:

with overall coefficient

$$
\begin{equation*}
\frac{1}{\pi i} \log ((2 a+\delta) / \delta) \tag{3.54}
\end{equation*}
$$

On the other hand:

$$
\begin{align*}
& \Delta Z_{f}[\mathrm{M}]\left(T_{a}\right)=\Delta\left(\bigcap \times S Z\left(T_{a}{ }^{M \text {-resolved }}\right)\right)  \tag{3.55}\\
& =\Delta\left(\Omega \times S Z()_{\underset{a}{\rightarrow}}^{\rightarrow \rightarrow}()\right)  \tag{3.56}\\
& =\Delta \Omega \times \Delta S Z()_{\vec{a}}^{\rightarrow \rightarrow}()  \tag{3.57}\\
& =\Delta \bigcap \times S Z(\Delta)_{a}^{\rightarrow \rightarrow}()  \tag{3.58}\\
& =\Delta \bigcap \times\left(\alpha \lim _{\epsilon \rightarrow 0} \rightarrow\right)+\cdots((+\alpha)) \cdots((+\cdots) \tag{3.59}
\end{align*}
$$

with $\alpha=\lim _{\epsilon \rightarrow 0}-1 / 2 \pi i \log ((M+\epsilon) /(a+\epsilon))=-1 / 2 \pi i \log M / a$ (note that this is one operation we wouldn't have been able to do at local extrema), yielding an overall coefficient of $-1 / \pi i \log (M / a)$ for the degree one tangle chord diagram with one chord between $T_{a}$ and its replica. This is finite, whereas (3.54) is large for small values of $\delta$. Not only are the coefficients of $Z_{f}\left(\Delta(\delta) T_{a}\right)$ and $\Delta Z_{f}\left(T_{a}\right)$ different, they have unlike behaviors as $\delta$ becomes very small.

In the preceding Lemma we studied the limiting behavior of $Z_{f}[\mathrm{M}](\Delta(\delta) T)$ as $\delta \rightarrow 0$ and saw that it was different from $\Delta Z_{f}[\mathrm{M}](T)$. However this does not say anything about $Z_{f}[\mathrm{M}](\Delta T)$ as it equals $Z_{f}[\mathrm{M}]\left(\lim _{\delta \rightarrow 0} \Delta(\delta) T\right)$.

Note that the result that follows was initially proved in [LM2] for the parameter-free $\hat{Z}_{f}$ and what we call the $\delta$-doubling map in this work.

Proposition 3.2.2.2. $\Delta Z_{f}(K)=Z_{f}(\Delta K)$.
Proof. On knots, $Z_{f}=Z_{f}[\mathrm{M}]$ for all $M>0$. Fix $M>0$. By multiplicativity of $Z_{f}[\mathrm{M}]$, it suffices to consider the doubling operation on one local extremum and one pinched extremity. Without loss of generality we can consider a local maximum. The proof for a roof-shaped extremity is exactly the same. We have:


Thus if we double the local max, we get:

$$
\begin{align*}
& =\overbrace{}^{2} \times S Z(\Delta)_{\vec{a}}^{\boldsymbol{H}_{\rightarrow}^{\rightarrow}}()  \tag{3.62}\\
& =\Delta \Omega \times \Delta S Z()_{\vec{a}}^{\rightarrow \rightarrow}()  \tag{3.63}\\
& =\Delta Z_{f}(\xrightarrow[a]{\longleftrightarrow}) \tag{3.64}
\end{align*}
$$

Corollary 3.2.2.3. If $L$ is an e-components framed oriented link, $K_{i}$ a component of $L$, $1 \leq i \leq e, \Delta_{i}$ the doubling map on $K_{i}$, then

$$
\begin{equation*}
\Delta_{i} Z_{f}(L)=Z_{f}\left(\Delta_{i} L\right) \tag{3.65}
\end{equation*}
$$

Proof. There are three types of chords to consider: chords none of whose feet are on $K_{i}$, for which the contribution to $Z_{f}(L)$ make (3.65) true, chords that have only one foot on $K_{i}$, with a corresponding contribution to $Z_{f}(L)$ that make (3.65) hold by definition of $\Delta_{i}$, and chords both of whose feet are on $K_{i}$, for which we invoke Proposition 3.2.2.2 above.

## Corollary 3.2.2.4.

$$
\begin{equation*}
\Delta \nu^{-1}=\nu^{-1} \otimes \nu^{-1} \tag{3.66}
\end{equation*}
$$

Proof. It follows from the previous proposition applied to $K=U$ and the fact that $\nu=$ $Z_{f}(U)^{-1}$ that:

$$
\begin{equation*}
\nu^{-1} \otimes \nu^{-1}=Z_{f}(\Delta U)=\Delta Z_{f}(U)=\Delta \nu^{-1} \tag{3.67}
\end{equation*}
$$

Proposition 3.2.2.5. For $K$ and $K^{\prime}$ two knots, $\alpha=\sum_{P} c_{P} K_{P}$ and $\beta=\sum_{Q} d_{Q} K_{Q}^{\prime}$ in $\hat{\mathcal{A}}(K)$ and $\hat{\mathcal{A}}\left(K^{\prime}\right)$ respectively, we have:

$$
\begin{equation*}
\Delta \alpha \cdot \Delta \beta=\Delta(\alpha \cdot \beta) \tag{3.68}
\end{equation*}
$$

where the product $\Delta \alpha \cdot \Delta \beta$ is defined by:


Proof. On the one hand:

$$
\begin{equation*}
\Delta \alpha \cdot \Delta \beta=\sum c_{P} d_{Q} \Delta K_{P} \cdot \Delta K_{Q}^{\prime} \tag{3.70}
\end{equation*}
$$

On the other hand:

$$
\begin{equation*}
\Delta(\alpha \cdot \beta)=\sum c_{P} d_{Q} \Delta\left(K_{P} \cdot K_{Q}^{\prime}\right) \tag{3.71}
\end{equation*}
$$

We show $\Delta K_{P} \cdot \Delta K_{Q}^{\prime}=\Delta\left(K_{P} \cdot K_{Q}^{\prime}\right)$ from which the result will follow. If we isolate the local minimum of $K$ that will be used in the product, removing this local minimum the
remaining tangle is denoted by $K^{+}$. Likewise if we isolate the local maximum of $K^{\prime}$ that will be used in the product, removing this local maximum the remaining tangle is denoted by $K^{\prime-}$. Then we consider:


$$
\begin{equation*}
=\Delta\left(K_{P} \cdot K_{Q}\right) \tag{3.73}
\end{equation*}
$$

Lemma 3.2.2.6. $\left(\Delta \nu^{-1}\right)^{-1}=\Delta \nu$
Proof. It follows from the previous Proposition that:

$$
\begin{equation*}
\Delta \nu^{-1} \cdot \Delta \nu=\Delta\left(\nu^{-1} \cdot \nu\right)=\Delta 1=1 \tag{3.74}
\end{equation*}
$$

Corollary 3.2.2.7. ([LM4]) $\Delta \nu=\nu \otimes \nu$.
Proof. From the previous Lemma $\Delta \nu=\left(\Delta \nu^{-1}\right)^{-1}$. By Corollary 3.2.2.4 this equals $\left(\nu^{-1} \otimes\right.$ $\left.\nu^{-1}\right)^{-1}$, which equals $\nu \otimes \nu$.

Corollary 3.2.2.8. $\Delta\left(\nu^{p}\right)=\nu^{p} \otimes \nu^{p}$ for all $p \in \mathbb{Z}$.
Proof. Fix $p \in \mathbb{Z}$. By Proposition 3.2.2.5, we have $\Delta\left(\nu^{p}\right)=(\Delta \nu)^{p}$. By Lemmas 3.2.2.4 and 3.2.2.7, we have that this equals $(\nu \otimes \nu)^{p}$, which is further equal to $\nu^{p} \otimes \nu^{p}$.

Theorem 3.2.2.9. ([LM2]) For an e-component link $L, \Delta_{i}$ the doubling map on the $i$-th component, $1 \leq i \leq e$ we have $\hat{Z}_{f}\left(\Delta_{i} L\right)=\Delta_{i} \hat{Z}_{f}(L)$.

Proof. By definition, one has:

$$
\begin{equation*}
\hat{Z}_{f}(L)=\nu^{m_{1}} \otimes \cdots \otimes \nu^{m_{e}} \cdot Z_{f}(L) \tag{3.75}
\end{equation*}
$$

Thus if one doubles the $i$-th component, $1 \leq i \leq e$ :

$$
\begin{align*}
\Delta_{i} \hat{Z}_{f}(L) & =\Delta_{i}\left(\nu^{m_{1}} \otimes \cdots \otimes \nu^{m_{e}} \cdot Z_{f}(L)\right)  \tag{3.76}\\
& =\nu^{m_{1}} \otimes \cdots \otimes \Delta_{i}\left(\nu^{m_{i}}\right) \otimes \cdots \otimes \nu^{m_{e}} \cdot \Delta_{i} Z_{f}(L)  \tag{3.77}\\
& =\nu^{m_{1}} \otimes \cdots \otimes\left(\Delta_{i} \nu\right)^{m_{i}} \otimes \cdots \otimes \nu^{m_{e}} \cdot \Delta_{i} Z_{f}(L)  \tag{3.78}\\
& =\nu^{m_{1}} \otimes \cdots \otimes 1 \otimes \cdots \otimes \nu^{m_{e}} \cdot\left(\left(\Delta_{i} \nu\right)^{m_{i}} \cdot \Delta_{i} Z_{f}(L)\right)  \tag{3.79}\\
& =\nu^{m_{1}} \otimes \cdots \otimes 1 \otimes \cdots \otimes \nu^{m_{e}} \cdot\left((\nu \otimes \nu)^{m_{i}} \cdot \Delta_{i} Z_{f}(L)\right)  \tag{3.80}\\
& =\nu^{m_{1}} \otimes \cdots \otimes(\nu \otimes \nu)^{m_{i}} \otimes \cdots \otimes \nu^{m_{e}} \cdot \Delta_{i} Z_{f}(L)  \tag{3.81}\\
& =\nu^{m_{1}} \otimes \cdots \otimes\left(\nu^{m_{i}} \otimes \nu^{m_{i}}\right) \otimes \cdots \otimes \nu^{m_{e}} \cdot \Delta_{i} Z_{f}(L)  \tag{3.82}\\
& =\nu^{m_{1}} \otimes \cdots \otimes\left(\nu^{m_{i}} \otimes \nu^{m_{i}}\right) \otimes \cdots \otimes \nu^{m_{e}} \cdot Z_{f}\left(\Delta_{i} L\right)  \tag{3.83}\\
& =\hat{Z}_{f}\left(\Delta_{i} L\right) \tag{3.84}
\end{align*}
$$

where in (3.78) we have used Lemma 3.2.2.7, and in (3.82) we used $\Delta Z_{f}=Z_{f} \Delta$ (Proposition 3.2.2.2). This being true for any $1 \leq i \leq e$, the result follows: $\hat{Z}_{f}(\Delta) \equiv \Delta \hat{Z}_{f}$.

## Chapter 4

## Behavior of the Kontsevich integral under band sum moves

In [LM3], [LM5], [LM6], Le, Murakami and Ohtsuki considered the band sum move of one link component over another by locally considering the two link components of interest to be side by side. The band sum thus looks horizontal as in:

$$
\begin{equation*}
|\mapsto \underset{\sim}{\smile}| \tag{4.1}
\end{equation*}
$$

Kontsevich introduced a modified Kontsevich integral in his seminal paper $[\mathrm{K}]$ that turned out to be multiplicative under connected sum. The band sum move being somewhat similar in nature to a connected sum, one could investigate whether this modified integral is also well-behaved under band sum moves for which link components of interest are locally on top of each other. In this case the band sum looks like:


We claim that once we generalize Kontsevich's modified integral for knots to links this is indeed the case. A band sum move of some component $K_{i}$ over some other component $K_{j}$ of a link $L$ is effected by doubling $K_{j}$ into itself and a copy that we will denote by $K_{j}^{(2)}$ and
then by connect summing $K_{i}$ with this copy $K_{j}^{(2)}$.

### 4.1 Unlinked case

We assume both $K_{i}$ and $K_{j}$ have orientations such that the band sum move of $K_{i}$ over $K_{j}$ results in the addition of $K_{i}$ to $K_{j}$ in the terminology of [RK1]. The case of reverse orientations will be treated in Section 8. We want to compute the Kontsevich integral of:

$$
K_{i} \# \Delta K_{j}\left\{\begin{array}{l}
K_{i} K_{i} \# K_{j}^{(2)}  \tag{4.3}\\
K_{j}
\end{array}\right.
$$

where $K_{i} \# \Delta K_{j}$ is a short-hand for $\left(K_{i} \# K_{j}^{(2)}\right) \amalg K_{j}$. We have:
Proposition 4.1.1. If $K_{i}$ and $K_{j}$ are unlinked:

$$
\begin{equation*}
\hat{Z}_{f}\left(K_{i} \# \Delta K_{j}\right)=\nu^{-1} \otimes 1 \cdot \hat{Z}_{f}\left(K_{i}\right) \cdot \Delta \hat{Z}_{f}\left(K_{j}\right) \tag{4.4}
\end{equation*}
$$

Proof. We compute:

where in the second line:

$$
\begin{align*}
& \hat{Z}_{f}[\mathrm{M}](\square):=\hat{Z}_{f}[\mathrm{M}]\left(K_{i}^{+}\right)  \tag{4.7}\\
& \hat{Z}_{f}[\mathrm{M}](\square):=\hat{Z}_{f}[\mathrm{M}]\left(K_{j} \amalg K_{j}^{(2)-}\right) \tag{4.8}
\end{align*}
$$

We seek to rewrite (4.6). It suffices to consider:

$$
\begin{align*}
\hat{Z}_{f}\left(K_{i}\right) & =\hat{Z}_{f}[\mathrm{M}]\left(K_{i}\right) \\
& =\hat{Z}_{f}[\mathrm{M}]\left(K_{i}^{+}\right) \times \hat{Z}_{f}[\mathrm{M}](\longrightarrow)  \tag{4.9}\\
& =\hat{Z}_{f}[\mathrm{M}]\left(K_{i}^{+}\right) \times Z_{f}[\mathrm{M}](\square) \tag{4.10}
\end{align*}
$$

as well as:

$$
\begin{align*}
\hat{Z}_{f}\left(\Delta K_{j}\right) & =\hat{Z}_{f}[\mathrm{M}]\left(\Delta K_{j}\right)=\hat{Z}_{f}[\mathrm{M}](\uparrow) \times \hat{Z}_{f}[\mathrm{M}]\left(K_{j} \amalg K_{j}^{(2)-}\right)  \tag{4.11}\\
& =\nu \otimes 1 \cdot Z_{f}[\mathrm{M}](\square) \times \hat{Z}_{f}[\mathrm{M}]\left(K_{j} \amalg K_{j}^{(2)-}\right) \tag{4.12}
\end{align*}
$$

where the first component in the tensor product refers to $K_{j}^{(2)}$ and the second to $K_{j}$ proper. Then:

$$
\begin{align*}
Z_{f}[\mathrm{M}](\uparrow) \times \hat{Z}_{f}[\mathrm{M}]\left(K_{j} \amalg K_{j}^{(2)-}\right) & =\nu^{-1} \otimes 1 \cdot \hat{Z}_{f}\left(\Delta K_{j}\right)  \tag{4.13}\\
& =\nu^{-1} \otimes 1 \cdot \Delta \hat{Z}_{f}\left(K_{j}\right) \tag{4.14}
\end{align*}
$$

by Theorem 3.2.2.9 which states that $\Delta \hat{Z}_{f}=\hat{Z}_{f}(\Delta)$. Then we can rewrite (4.6) as:

$$
\begin{align*}
& \hat{Z}_{f}[\mathrm{M}](\square) \times \hat{Z}_{f}[\mathrm{M}](\square) \\
& =\hat{Z}_{f}[\mathrm{M}](\uparrow) \times \hat{Z}_{f}[\mathrm{M}](|\quad|) \times \hat{Z}_{f}[\mathrm{M}](\square)  \tag{4.15}\\
& =\hat{Z}_{f}[\mathrm{M}](\uparrow) \times S Z(| |) \times \hat{Z}_{f}[\mathrm{M}](\square) \tag{4.16}
\end{align*}
$$

$$
\begin{align*}
& \times \hat{Z}_{f}[\mathrm{M}](\square)  \tag{4.17}\\
& =\left(\hat{Z}_{f}[\mathrm{M}]\left(K_{i}^{+}\right) \times Z_{f}[\mathrm{M}](\circlearrowright)\right) \#\left(Z_{f}[\mathrm{M}](\curvearrowright) \times \hat{Z}_{f}[\mathrm{M}]\left(K_{j} \amalg K_{j}^{(2)-}\right)\right)  \tag{4.18}\\
& =\left(\hat{Z}_{f}[\mathrm{M}]\left(K_{i}^{+}\right) \times Z_{f}[\mathrm{M}](\circlearrowright)\right) \#\left(\nu^{-1} \otimes 1 \cdot \Delta \hat{Z}_{f}\left(K_{j}\right)\right)  \tag{4.19}\\
& =\hat{Z}_{f}\left(K_{i}\right) \cdot\left(\nu^{-1} \otimes 1 \cdot \Delta \hat{Z}_{f}\left(K_{j}\right)\right)  \tag{4.20}\\
& =\nu^{-1} \otimes 1 \cdot \hat{Z}_{f}\left(K_{i}\right) \cdot \Delta \hat{Z}_{f}\left(K_{j}\right) \tag{4.21}
\end{align*}
$$

so we have essentially proved that we have:

$$
\begin{equation*}
\hat{Z}_{f}\left(K_{i} \# \Delta K_{j}\right)=\nu^{-1} \otimes 1 \cdot \hat{Z}_{f}\left(K_{i}\right) \cdot \Delta \hat{Z}_{f}\left(K_{j}\right) \tag{4.22}
\end{equation*}
$$

Remark 4.1.2. Observe that (4.22) written in the form:

$$
\begin{equation*}
\hat{Z}_{f}\left(K_{i} \# \Delta K_{j}\right)=\nu^{-1} \otimes 1 \cdot \hat{Z}_{f}\left(K_{i}\right) \cdot \hat{Z}_{f}\left(\Delta K_{j}\right) \tag{4.23}
\end{equation*}
$$

is surprisingly similar to the statement $[\mathrm{K}]$ :

$$
\begin{equation*}
\hat{Z}_{f}\left(K_{i} \# K_{j}\right)=\nu^{-1} \otimes 1 \cdot \hat{Z}_{f}\left(K_{i}\right) \cdot \hat{Z}_{f}\left(K_{j}\right) \tag{4.24}
\end{equation*}
$$

One now defines, as Kontsevich did, a modified version of the hatted Kontsevich integral.

Definition 4.1.3. We define $\widetilde{Z}_{f}$ to be $\hat{Z}_{f}$ with the provision that every link component in its argument has a factor of $\nu^{-1}$ multiplying it. Equivalently:

$$
\begin{equation*}
\widetilde{Z}_{f}(L)=\nu^{-1} \otimes \cdots \otimes \nu^{-1} \cdot \hat{Z}_{f}(L) \tag{4.25}
\end{equation*}
$$

Proposition 4.1.4. ([K]) If $K_{i}$ and $K_{j}$ are unlinked:

$$
\begin{equation*}
\widetilde{Z}_{f}\left(K_{i} \# \Delta K_{j}\right)=\widetilde{Z}_{f}\left(K_{i}\right) \cdot \Delta \widetilde{Z}_{f}\left(K_{j}\right) \tag{4.26}
\end{equation*}
$$

Proof. Multiplying both sides of (4.22) by $\nu^{-1} \otimes \nu^{-1}$ where the first $\nu^{-1}$ factor acts on $K_{i} \# K_{j}^{(2)}$, the second on $K_{j}$, we get:

$$
\begin{equation*}
\nu^{-1} \otimes \nu^{-1} \cdot \hat{Z}_{f}\left(K_{i} \# \Delta K_{j}\right)=\nu^{-1} \otimes 1 \cdot \hat{Z}_{f}\left(K_{i}\right) \cdot \nu^{-1} \otimes \nu^{-1} \cdot \Delta \hat{Z}_{f}\left(K_{j}\right) \tag{4.27}
\end{equation*}
$$

Now using:

$$
\begin{align*}
\nu^{-1} \otimes \nu^{-1} \cdot \Delta \hat{Z}_{f}\left(K_{j}\right) & =\Delta\left(\nu^{-1}\right) \cdot \Delta \hat{Z}_{f}\left(K_{j}\right)  \tag{4.28}\\
& =\Delta\left(\nu^{-1} \cdot \hat{Z}_{f}\left(K_{j}\right)\right)  \tag{4.29}\\
& =\Delta \widetilde{Z}_{f}\left(K_{j}\right) \tag{4.30}
\end{align*}
$$

we get:

$$
\begin{equation*}
\widetilde{Z}_{f}\left(K_{i} \# \Delta K_{j}\right)=\widetilde{Z}_{f}\left(K_{i}\right) \cdot \Delta \widetilde{Z}_{f}\left(K_{j}\right) \tag{4.31}
\end{equation*}
$$

On the other hand for two unlinked knots:

Proposition 4.1.5. If $K_{i}$ and $K_{j}$ are unlinked:

$$
\begin{equation*}
\widetilde{Z}_{f}\left(K_{i} \amalg K_{j}\right)=\widetilde{Z}_{f}\left(K_{i}\right)+\widetilde{Z}_{f}\left(K_{j}\right) \tag{4.32}
\end{equation*}
$$

Proof. For two unlinked knots, we have:

$$
\begin{equation*}
\hat{Z}_{f}\left(K_{i} \amalg K_{j}\right)=\hat{Z}_{f}\left(K_{i}\right)+\hat{Z}_{f}\left(K_{j}\right) \tag{4.33}
\end{equation*}
$$

That this is true is immediate. It suffices to compute $\hat{Z}_{f}\left(K_{i} \amalg K_{j}\right)$ in two ways. The first is performed by putting the two knots side by side, so that:

$$
\begin{equation*}
\hat{Z}_{f}\left(K_{i} \amalg K_{j}\right)=\hat{Z}_{f}\left(K_{i}\right)+\hat{Z}_{f}\left(K_{j}\right)+\text { terms with chords between } K_{i} \text { and } K_{j} \tag{4.34}
\end{equation*}
$$

The second way to perform this computation is by putting one knot above the other, in which case we do not get chords between one component and the other, and by isotopy invariance of $\hat{Z}_{f}$ it follows that in the above sum the terms with chords between the two components vanish and we are left with (4.33). From there, by modifying the hatted Kontsevich integral for each component, we get:

$$
\begin{align*}
\nu^{-1} \otimes \nu^{-1} & \cdot \hat{Z}_{f}\left(K_{i} \amalg K_{j}\right) \\
& =\left(\nu^{-1} \cdot\right) \otimes\left(\nu^{-1} \cdot\right) \hat{Z}_{f}\left(K_{i} \amalg K_{j}\right)  \tag{4.35}\\
= & \left(\nu^{-1} \cdot\right) \otimes\left(\nu^{-1} \cdot\right)\left(\hat{Z}_{f}\left(K_{i}\right)+\hat{Z}_{f}\left(K_{j}\right)\right)  \tag{4.36}\\
= & \nu^{-1} \cdot \hat{Z}_{f}\left(K_{i}\right)+\nu^{-1} \cdot \hat{Z}_{f}\left(K_{j}\right) \tag{4.37}
\end{align*}
$$

or:

$$
\begin{equation*}
\widetilde{Z}_{f}\left(K_{i} \amalg K_{j}\right)=\widetilde{Z}_{f}\left(K_{i}\right)+\widetilde{Z}_{f}\left(K_{j}\right) \tag{4.38}
\end{equation*}
$$

Remark 4.1.6. Before the band sum move, we have:

$$
\begin{equation*}
\widetilde{Z}_{f}\left(K_{i} \amalg K_{j}\right)=\widetilde{Z}_{f}\left(K_{i}\right)+\widetilde{Z}_{f}\left(K_{j}\right) \tag{4.39}
\end{equation*}
$$

and after the band sum move:

$$
\begin{equation*}
\widetilde{Z}_{f}\left(K_{i} \# \Delta K_{j}\right)=\widetilde{Z}_{f}\left(K_{i}\right) \cdot \Delta \widetilde{Z}_{f}\left(K_{j}\right) \tag{4.40}
\end{equation*}
$$

The essence of these two statements is that whatever chords we have on the $j$-th component in the expression for $\widetilde{Z}_{f}$ before band sum move are doubled after the move.

### 4.2 Linked case

We now treat the general case wherein $K_{i}$ and $K_{j}$ are linked to other components of the given link they are part of, and are linked between themselves. In this case the desired map on $\hat{Z}_{f}$ that should be induced by the band sum move on link components turns out to be non-existent. A way around this difficulty amounts to fixing a small window in which only strands from the link components involved in the band sum show, and are disconnected from the rest of the link as it were, or "frozen", and then a map on $\hat{Z}_{f}$ induced by the band sum on link components with a fixed window can be worked out ([DM]). In [O], [LM3], [LM5], [LM6] the window is fixed so that the two link components involved in the band sum are side by side. We chose a window in which they are locally on top of each other.

If $L$ is the link under consideration, we can arrange that by isotopy invariance of $\hat{Z}_{f}$ we just have to compute:

where the tangle $T$ represents the link $L$ minus one local max from $K_{i}$ and one local min from $K_{j}$, both local extrema having been moved to a side using the isotopy invariance of $\hat{Z}_{f}$. Using the same trick, it suffices to compute:


Proposition 4.2.1. If $L$ is an e-components framed oriented link, $L^{\prime}$ is the link resulting from the band sum of a link component $K_{i}$ over some other link component $K_{j}, 1 \leq i, j \leq e$, then:

$$
\begin{equation*}
\hat{Z}_{f}\left(L^{\prime}\right)=\underset{M}{|\stackrel{\nu}{\nu}|} \cdot \Delta_{j} \hat{Z}_{f}[\mathrm{M}](L \backslash \underset{\sim}{\varkappa}) \tag{4.43}
\end{equation*}
$$

where • is the product on $Z$ induced by the long chords lemma ([ChDu]) and $L \backslash \underset{\curvearrowleft}{\hookrightarrow}$ is the link $L$ from which its portion involving $K_{i}$ and $K_{j}$ to the left that looks like $\underset{\sim}{\bigodot}$ has been severed.

Proof.

where in the second step we have used the long chords lemma ([ChDu]). In the same paper it is shown that the long chords lemma induces a product on tangle chord diagrams written •, that in this case would read:


Note that this product is not the product on $\hat{\mathcal{A}}$ induced by the connected sum of $S^{1}$,s, nor is it a concatenation per se. All that is left is to compute the first term in the product as the second exactly equals:

$$
\begin{equation*}
\hat{Z}_{f}[\mathrm{M}]\left(\Delta_{j}(L \backslash \underset{\prec}{\smile})\right)=\Delta_{j} \hat{Z}_{f}[\mathrm{M}](L \backslash \underset{\sim}{\smile}) \tag{4.46}
\end{equation*}
$$

## Lemma 4.2.2.

$$
\begin{equation*}
\hat{Z}_{f}[\mathrm{M}](\underset{\leftrightarrow}{\stackrel{\leftrightarrow}{\longleftrightarrow}})=\mid \underset{M}{|\stackrel{\rightharpoonup}{\nu}|} \tag{4.47}
\end{equation*}
$$

Proof. We present two proofs. The first proof consists in computing $Z_{f}(\curvearrowright)$ exactly. It suffices to write:

$$
\begin{align*}
& \xrightarrow[\mathrm{M}]{\longrightarrow}=Z_{f}[\mathrm{M}](\Omega)  \tag{4.48}\\
& =Z_{f}[\mathrm{M}]\left(\Re_{)}\right.  \tag{4.49}\\
& =Z_{f}[\mathrm{M}](\curvearrowright) \times Z_{f}[\mathrm{M}](\curvearrowright)  \tag{4.50}\\
& =\bigcap \times Z_{f}[\mathrm{M}](\Re) \tag{4.51}
\end{align*}
$$

where in the first step we used the fact that the two strands are analytically coincident and thus are $M$-resolved together. Moreover, since the base is chosen to be $M, Z_{f}[\mathrm{M}](\Omega)$ is trivial. In the second step we used the invariance of $Z_{f}[\mathrm{M}]$ under the stretching move thus displayed. It follows that:

$$
\begin{equation*}
z_{f}(\Re)=\Re \tag{4.52}
\end{equation*}
$$

and:

$$
\begin{equation*}
\hat{Z}_{f}[\mathrm{M}](\curvearrowleft)=1 \otimes \nu \cdot Z_{f}[\mathrm{M}](\curvearrowright)=|\sqrt{\nu}| \tag{4.53}
\end{equation*}
$$

Second proof:

$$
\begin{align*}
& \hat{Z}_{f}[\mathrm{M}](\underset{\sim}{\stackrel{M}{\longleftrightarrow}})=| | \times \hat{Z}_{f}[\mathrm{M}](\curvearrowleft)  \tag{4.54}\\
& =\circlearrowleft_{\#} \overbrace{Z_{f}[\mathrm{M}]}(\curvearrowleft)  \tag{4.55}\\
& =\left(\hat{Z}_{f}[\mathrm{M}](\triangleleft) \# Z_{f}[\mathrm{M}](\curvearrowright)\right) \times \hat{Z}_{f}[\mathrm{M}](\curvearrowleft)  \tag{4.56}\\
& =\hat{Z}_{f}[\mathrm{M}](\circlearrowright) \# Z_{f}[\mathrm{M}](\curvearrowright) \times \hat{Z}_{f}[\mathrm{M}](\curvearrowleft)  \tag{4.57}\\
& =Z_{f}[\mathrm{M}](\checkmark) \# \nu^{-1} \otimes 1 \cdot \hat{Z}_{f}[\mathrm{M}](\curvearrowright) \times \hat{Z}_{f}[\mathrm{M}](\curvearrowright)  \tag{4.58}\\
& =\hat{Z}_{f}[\mathrm{M}](\longrightarrow) \# \nu^{-1} \otimes 1 \cdot \hat{Z}_{f}[\mathrm{M}] \text { (~) }  \tag{4.59}\\
& =\#_{K_{j}^{(2)}}^{K_{i}} \nu^{-1} \otimes 1 \cdot \hat{Z}_{f}[\mathrm{M}]\left(\bigcap_{\bigcap}^{\bigcup}\right) \tag{4.60}
\end{align*}
$$

where in the last relation we denoted by $\#_{K_{j}^{(2)}}^{K_{i}}$ a connected sum between $K_{i}$ and $K_{j}^{(2)}$ and we used the following lemma:

## Lemma 4.2.3.

$$
\begin{equation*}
\hat{Z}_{f}[\mathrm{M}](\smile) \# \hat{Z}_{f}[\mathrm{M}](\text { ( ) })=\#_{K_{j}^{(2)}}^{K_{i}} \hat{Z}_{f}[\mathrm{M}]\left(\bigcap_{\bigcap}^{\bigcup}\right) \tag{4.61}
\end{equation*}
$$

Proof. Consider the following tangle:

$$
\begin{equation*}
\mid \bigcup_{\bigcap} \tag{4.62}
\end{equation*}
$$

Then:

$$
\begin{equation*}
\hat{Z}_{f}[\mathrm{M}](\underbrace{}_{\text {గी }})=\hat{Z}_{f}[\mathrm{M}](\| \circlearrowright) \times \hat{Z}_{f}[\mathrm{M}](\| / \text { ) } \tag{4.63}
\end{equation*}
$$

so that:

$$
\begin{equation*}
\#_{K_{j}^{(2)}}^{K_{i}} \hat{Z}_{f}[\mathrm{M}]\left(\mid \bigcap_{\bigcap}^{\longrightarrow}\right)=\#_{K_{j}^{(2)}}^{K_{i}}\left(\hat{Z}_{f}[\mathrm{M}](\mid \circlearrowleft) \times \hat{Z}_{f}[\mathrm{M}]\left(\|^{(\Omega)}\right)\right. \tag{4.64}
\end{equation*}
$$

If in either expression we consider only those tangle chord diagrams none of whose chords have any foot on the straight strand, then:

$$
\begin{equation*}
\#_{K_{j}^{(2)}}^{K_{i}} \hat{Z}_{f}[\mathrm{M}]\left(\bigcap_{\mathrm{n}}\right)=\hat{Z}_{f}[\mathrm{M}]\left(\text { ) } \# \hat{Z}_{f}[\mathrm{M}](\text { (®) }\right. \tag{4.65}
\end{equation*}
$$

So far we have:

$$
\begin{equation*}
\hat{Z}_{f}[\mathrm{M}](\curvearrowright)=\#_{K_{j}^{(2)}}^{K_{i}} \nu^{-1} \otimes 1 \cdot \hat{Z}_{f}[\mathrm{M}](\overbrace{)}^{( }) \tag{4.66}
\end{equation*}
$$

We have:

$$
\begin{equation*}
\hat{Z}_{f}[\mathrm{M}](\overparen{\Omega})=\nu \otimes \nu \cdot Z_{f}[\mathrm{M}](\Delta \underset{\mathrm{M}}{\Omega}) \tag{4.67}
\end{equation*}
$$

In $Z_{f}(\Delta \Omega)$ there are no non-self chords between one local max and its replica as they would be integrated from $M$ at the top down to the base which is of width $M$. Thus:

$$
\begin{equation*}
Z_{f}[\mathrm{M}](\Delta \underset{\mathrm{M}}{\Omega})=\bigcap \tag{4.68}
\end{equation*}
$$

From which it follows that:

and:

so that:


From which we get:


This completes the proof of the Lemma.

Back to the proof of the Proposition:

$$
\begin{align*}
\hat{Z}_{f}[\mathrm{M}]\left(L^{\prime}\right) & =\hat{Z}_{f}[\mathrm{M}](\curvearrowleft) \cdot \Delta_{j} \hat{Z}_{f}[\mathrm{M}](L \backslash \underset{\sim}{\varkappa})  \tag{4.73}\\
& =|\stackrel{\nu}{ }| \cdot \Delta_{j} \hat{Z}_{f}[\mathrm{M}](L \backslash \underset{\sim}{\sim}) \tag{4.74}
\end{align*}
$$

Proposition 4.2.4. For $L$ an $e$-components framed oriented link, $L^{\prime}$ the link resulting from the band sum move of one component $K_{i}$ over some other component $K_{j}, 1 \leq i, j \leq e$, then
the resulting induced map on $\widetilde{Z}_{f}$ is given by:

$$
\begin{align*}
& \checkmark \\
& \widetilde{Z}_{f}(L)=\widetilde{Z}_{f}[\mathrm{M}](L)=\stackrel{\nu}{\sim} \cdot \widetilde{Z}_{f}[\mathrm{M}](L \backslash \backsim)  \tag{4.75}\\
& \longmapsto \widetilde{Z}_{f}\left(L^{\prime}\right)=\widetilde{Z}_{f}[\mathrm{M}]\left(L^{\prime}\right)=\Delta_{\|}^{\|} \quad \Delta_{j} \tilde{Z}_{f}[\mathrm{M}](L \backslash \backsim) \tag{4.76}
\end{align*}
$$

which we can symbolically write as:

$$
\begin{equation*}
\tilde{z}_{f}(L)=\sum \lambda \stackrel{\square}{\square} \quad \mapsto \tilde{z}_{f}\left(L^{\prime}\right)=\sum \lambda \tag{4.77}
\end{equation*}
$$

Proof. The previous proposition showed that:

$$
\begin{align*}
& \left.\hat{Z}_{f}(L)=\hat{Z}_{f}[\mathrm{M}](L)={ }^{\nu}\right) \cdot \hat{Z}_{f}[\mathrm{M}](L \backslash \bigcup)  \tag{4.78}\\
& \longmapsto \hat{Z}_{f}\left(L^{\prime}\right)=\hat{Z}_{f}[\mathrm{M}]\left(L^{\prime}\right)=\underset{M}{\stackrel{\nu}{\Delta \mid} \mid} \cdot \Delta_{j} \hat{Z}_{f}[\mathrm{M}](L \backslash \underset{\sim}{\rightsquigarrow}) \tag{4.79}
\end{align*}
$$

We will use this statement and determine what normalization $\widetilde{Z}_{f}$ leads to a statement of the form:

$$
\widetilde{Z}_{f}(L)=\sum \lambda \stackrel{\widetilde{Z}_{f}\left(L^{\prime}\right)=\sum \lambda, ~ \longmapsto \square}{\square}
$$

Let $\widetilde{Z}_{f}$ be $\hat{Z}_{f}$ with the provision that every link component in the expression of $\hat{Z}_{f}(L)$
is multiplied by $\nu^{p}$ for some $p$ yet to be determined. Without loss of generality we can work with a two component link. The proof for other link components in addition to those involved in the band sum move is the same but is just a clutter of tensor products. Before band sum move:

$$
\begin{align*}
& \widetilde{Z}_{f}[M](L)=\nu^{p} \otimes \nu^{p} \cdot \hat{Z}_{f}[\mathrm{M}](L)  \tag{4.80}\\
& \left.=\nu^{p} \otimes \nu^{p} \cdot \sqrt{\nu}\right) \cdot \hat{Z}_{f}[\mathrm{M}](L \backslash \bigcup)  \tag{4.81}\\
& \text { 四 } \\
& =\underbrace{\sqrt[\nu^{1+}]{u}} \cdot \hat{Z}_{f}[\mathrm{M}](L \backslash \underset{\sim}{\cup})  \tag{4.82}\\
& =\sqrt{\nu}) \cdot \widetilde{Z}_{f}[M](L \backslash \underset{\sim}{\sim}) \tag{4.83}
\end{align*}
$$

After the band sum move:

$$
\begin{align*}
\widetilde{Z}_{f}[M]\left(L^{\prime}\right) & =\nu^{p} \otimes \nu^{p} \cdot \mid\langle\sqrt{\nu}| \cdot \Delta \hat{Z}_{f}[\mathrm{M}](L \backslash \bigcup)  \tag{4.84}\\
& =\nu^{p} \otimes \nu^{p} \cdot \mid\langle\nu| \cdot \Delta\left(\nu^{-p} \otimes \nu^{-p} \cdot \widetilde{Z}_{f}[M](L \backslash \backsim)\right) \tag{4.85}
\end{align*}
$$

$$
\begin{equation*}
=\nu^{p} \otimes \nu^{p} \cdot|\sqrt{\nu}| \cdot \nu^{-p} \otimes \nu^{-p} \otimes \nu^{-p} \cdot \Delta \widetilde{Z}_{f}[M](L \backslash \underset{\sim}{\smile}) \tag{4.86}
\end{equation*}
$$

$$
\begin{equation*}
=\nu^{p} \otimes \nu^{p} \cdot \nu^{-2 p} \otimes \nu^{1-p} \cdot \curvearrowright \cdot \Delta \widetilde{Z}_{f}[M](L \backslash \bigcup) \tag{4.87}
\end{equation*}
$$

$$
\begin{equation*}
=\nu^{-p} \otimes \nu \cdot \curvearrowright \cdot \Delta \widetilde{Z}_{f}[M](L \backslash \underset{\curvearrowleft}{\smile}) \tag{4.88}
\end{equation*}
$$

$$
\begin{equation*}
=\nu^{-p / 2} \cdot \nu^{-p / 2} \otimes \nu \cdot \curvearrowright \cdot \Delta \widetilde{Z}_{f}[M](L \backslash \asymp) \tag{4.89}
\end{equation*}
$$

$$
\begin{equation*}
=\overbrace{\frac{\Delta \nu}{\|} \cdot \Delta \widetilde{Z}_{f}[M](L \backslash \backsim)}^{\frac{1}{\nu^{-\frac{p+1}{2}}} \overbrace{}^{\nu^{-\frac{p+1}{2}}}} \tag{4.91}
\end{equation*}
$$

Thus if we want a statement of the form:

then using (4.83) it follows that we must have $p=-1$, whence the result.

## Chapter 5

## The book notation

### 5.1 Linking matrix and degree 1 Kontsevich integral

It was originally Kontsevich's idea ([K]) to generalize Gauss' definition of the linking number via an integral, to the algebraic completion of the sum of generalized such integrals valued in tangle chord diagrams. The old observation of Yetter [Y2] that the degree 1 part of the framed Kontsevich integral of a link with labeled components behaves like the linking matrix under band sum move is thus not very surprising, though it has been generally overlooked in most of the literature that tends to focus on knots rather than links. He sees that as a motivation for introducing the book notation that we will cover in the next section. For now we show his claim in the simple case of a two component link $L=K_{i} \cup K_{j}$ where the two knots $K_{i}$ and $K_{j}$ are trivial and unlinked for the simplicity of exposition. We take the following basis for $\mathcal{A}\left(S^{1} \amalg S^{1}\right)$ :


We write:

$$
\begin{equation*}
Z_{1}(L)=c_{i}\left(\frac{1}{2} \circlearrowleft\right)+c_{j}\left(\frac{1}{2} \circlearrowleft\right)+c_{i j}(\square) \tag{5.4}
\end{equation*}
$$

After the band sum move of $K_{i}$ over $K_{j}$, the resulting link is denoted by $L^{\prime}$ and we have:

$$
\begin{align*}
Z_{1}\left(L^{\prime}\right) & =c_{i}\left(\frac{1}{2}\right.  \tag{5.5}\\
& =c_{i}\left(\frac{1}{2}\right. \\
& +c_{j}\left(\frac{1}{2}\right. \\
& +c_{i j}( \pm)^{\prime}+c_{j}\left(\frac{1}{2}\right.  \tag{5.6}\\
& =\left[c_{i}+c_{j} \pm 2 c_{i j}\right] \\
& +\left[ \pm c_{j}+c_{i j}\right] \tag{5.7}
\end{align*}
$$

Now observe that in the basis for $\mathcal{A}\left(S^{1} \amalg S^{1}\right)$ we chose, the coefficients of $Z_{1}$ transform like the linking matrix entries under band sum move:

$$
\begin{align*}
& L_{i} \mapsto L_{i}^{\prime}  \tag{5.8}\\
&=L_{i}+L_{j} \pm 2 L_{i j}  \tag{5.9}\\
& L_{j} \mapsto L_{j}^{\prime}=L_{j}  \tag{5.10}\\
& L_{i j} \mapsto L_{i j}^{\prime}=L_{i j} \pm L_{j}
\end{align*}
$$

The matrix congruence that implements the band sum move on both the coefficients of $Z_{1}(L)$ and the linking matrix is given by:

$$
M=\left[\begin{array}{rr}
1 & 0  \tag{5.11}\\
\pm 1 & 1
\end{array}\right]
$$

We check that if:

$$
A=\left(\begin{array}{ll}
a_{i i} & a_{i j}  \tag{5.12}\\
a_{j i} & a_{j j}
\end{array}\right)
$$

stands for either of

$$
\left[\begin{array}{ll}
L_{i} & L_{i j}  \tag{5.13}\\
L_{j i} & L_{j}
\end{array}\right] \text { or }\left[\begin{array}{ll}
c_{i} & c_{i j} \\
c_{j i} & c_{j}
\end{array}\right]
$$

with $L_{i j}=L_{j i}$ and $c_{i j}=c_{j i}$, then:

$$
\begin{align*}
M^{T} A M & =\left[\begin{array}{ll}
1 & \pm 1 \\
0 & 1
\end{array}\right]\left[\begin{array}{ll}
a_{i i} & a_{i j} \\
a_{j i} & a_{j j}
\end{array}\right]\left[\begin{array}{ll}
1 & 0 \\
\pm 1 & 1
\end{array}\right] \\
& =\left[\begin{array}{ll}
a_{i i}+a_{j j} \pm 2 a_{i j} & a_{i j} \pm a_{j j} \\
a_{j i} \pm a_{j j} & a_{j j}
\end{array}\right] \tag{5.14}
\end{align*}
$$

In the following, we seek to generalize such a transformation.

### 5.2 Books of matrices

Note that in the previous section, as is customary whenever we compute the Kontsevich integral of tangles, the diagrams are tangle chord diagrams and not chord diagrams as elements of $\hat{\mathcal{A}}\left(\amalg S^{1}\right)$ as initially defined by Kontsevich. Indeed, for a link $L$ sliced into $n$ horizontal strips in each of which we have a tangle $T_{i}, 1 \leq i \leq n$ such that $L=T_{1} \times \cdots \times T_{n}$, it is generally understood that when we compute $\hat{Z}_{f}(L)$ by using the multiplicativity property of the Kontsevich integral $\hat{Z}_{f}$ :

$$
\begin{equation*}
\hat{Z}_{f}(L)=\hat{Z}_{f}\left(T_{1}\right) \times \cdots \times \hat{Z}_{f}\left(T_{n}\right) \tag{5.15}
\end{equation*}
$$

the resulting object $\hat{Z}_{f}(L)$ is a sum of tangle chord diagrams with coefficients in front of each diagram being obtained from the Kontsevich integral itself. By definition, $\widetilde{Z}_{f}$ is also multiplicative and $\widetilde{Z}_{f}(L)$ is also a sum of tangle chord diagrams with complex coefficients. The Kontsevich integral $\widetilde{Z}_{f}$ of a link $L$ can be written:

$$
\begin{equation*}
\widetilde{Z}_{f}(L)=\sum_{\text {chord diagr. } \mathrm{x}} c_{X} X=\sum_{m \geq 0} \sum_{|X|=m} c_{X} X \tag{5.16}
\end{equation*}
$$

where by $|X|=m$ we mean that the tangle chord diagram $X$ has chord degree $m$, and we sum over all such tangle chord diagrams, for all $m \geq 0$. In what follows, we fix $m \geq 1$. We will be working with an $e$-component link $L=\coprod_{1 \leq l \leq e} K_{l}$.

### 5.2.1 Vertical slicing of tangles

Before slicing links, we have to fully determine where local extrema will be located on any given link. The definitions of $\hat{Z}_{f}$ and $\widetilde{Z}_{f}$ each introduce factors of $\nu=Z_{f}(U)^{-1}$ which will yield additional local extrema on the link $L$ upon acting on $Z_{f}(L)$. Since those factors of $\nu$ are acting on local maxima, it suffices to consider products of the form:

$$
\nu \cdot Z_{f}[\mathrm{Q}](\overbrace{a}^{\longrightarrow})
$$

for $Q>0$. Observe that if we write:

$$
\begin{equation*}
Z_{f}\left(\text { ) }=\sum_{|P| \geq 0} c_{P} \cdot(\text { ) }\right. \tag{5.17}
\end{equation*}
$$

then:

$$
\begin{equation*}
\nu=Z_{f}(\sqrt{\sim})^{-1}=\sum_{|P| \geq 0} d_{P} \cdot(\sqrt{\mid} \tag{5.18}
\end{equation*}
$$

with coefficients $d_{P}$ such that

$$
\begin{equation*}
\nu \cdot z_{f}(\sqrt{ })=\sqrt{ } \tag{5.19}
\end{equation*}
$$

It follows that:

$$
\begin{align*}
& \nu \cdot Z_{f}[\mathrm{Q}](\underset{a}{\longrightarrow})=\nu \cdot \bigcap \times S Z(\overbrace{a}^{\rightarrow \rightarrow}()  \tag{5.20}\\
& =\sum_{|P| \geq 0} d_{P} \cdot(\sqrt{\longrightarrow})_{P} \# \sum_{|R| \geq 0} e_{R} \cdot()_{\vec{a}}^{Q}()_{R} \\
& =\sum_{|P|,|R| \geq 0} d_{P} e_{R} \cdot(\square)_{(P, R)} \tag{5.21}
\end{align*}
$$

This simple computation shows that the skeleton $L$ will be modified in the expression for $Z_{f}(L)$ by the introduction of a hook of the above form at each local max but for one local $\max$ on each component by definition of $\widetilde{Z}_{f}$.

Definition 5.2.1.1. Let $\widetilde{L}$ be the link $L$ with all but one local max on each component being tweaked into a left pointing hook as above.

We consider the band sum move of $K_{i}$ over $K_{j}$. The result of such a band sum move is denoted by $L^{\prime}$. However we work with $\widetilde{L}$ and $\widetilde{L^{\prime}}$ as the rest of the section will make clear. We slice $\widetilde{L}$ into N vertical strips as follows. For a local maximum on the $s$-th component $K_{s}$ such as:

the slicing is performed as follows:


We want each local max to be enclosed within two vertical slices to distinguish neighboring local extrema of a same component. If we call the vertical slices on either side of a local max dividing slices, it follows in practice that consecutive local extrema share a dividing slice is sufficient as we will see later. We do this at each local max of each component $K_{s}$ of the link $\widetilde{L}$. We slice each local min of each component in like manner, keeping in mind that consecutive local extrema can share dividing slices. We number those vertical strips formed from this slicing procedure starting from the left.

In the above situation, we would have the strips with the following labels:


We now discuss the labeling of the chords. For each time, we have a chord. Thus it is natural to number the chords from the bottom up. If $t_{1}<\cdots t_{m}$ are the different times corresponding to $m$ different chords, then those corresponding chords will be labeled $1, \cdots, m$ respectively.

### 5.2.2 Matrix representation of tangle chord diagrams

Consider the generic situation of one chord stretching between 2 components of $L$ indexed by $s$ and $t$ :

where we have displayed only portions of the $s$ and $t$ components on which the $a$-th chord is ending, $1 \leq a \leq m$ and $1 \leq n, p \leq N$ are strip indices. We will index such a chord by $(s, t)$ and $(n, p)$. Chords can be moved up or down along a link and thus their feet may end up in vertical strips with different indices. To avoid such ambiguity, we require that chords be moved up along the link until at least one of their feet reaches a local max.

We represent each such chord by a $e N \times e N$ matrix in the basis given by the ordering of the components, and of the strips, as $(1,2, \cdots, N, 1,2, \cdots, N, \cdots, N)$ the first $N$ vectors $1,2, \cdots, N$ corresponding to the first component of $L$, followed by those for the second component, and so on, until the $e$-th component. The st block of that matrix will carry information about chords between the $s$ and $t$ components of the link. In the above situation, the $s t$ block has all its entries zero, except the $n p$ entry which is one. Note that the matrix is symmetric, so all blocks are empty for that particular chord, except the $t s$-th one, whose $p n$ entry is one as well. We represent such a matrix as follows where without loss of generality
we have chosen $s<t$ and $n<p$ :


We will refer to such a matrix as a page, and we denote it by $A_{s, t, n, p}$ with obvious notations. We do this for all chords of a given tangle chord diagram $X$ of degree $m$. The information about its chords will therefore be given by $m$ ordered pages from the bottom up, the collection of which will be referred to as a book. We denote a book as follows:

$$
\begin{equation*}
A_{I, J, U, V}:=\underset{1 \leq a \leq m}{\times} A_{i_{a}, j_{a}, u_{a}, v_{a}} \tag{5.27}
\end{equation*}
$$

where $I, J, U$ and $V$ are multi-indices defined by:

$$
\begin{align*}
I & =\left(i_{1}, \cdots, i_{m}\right)  \tag{5.28}\\
J & =\left(j_{1}, \cdots, j_{m}\right)  \tag{5.29}\\
U & =\left(u_{1}, \cdots, u_{m}\right)  \tag{5.30}\\
V & =\left(v_{1}, \cdots, v_{m}\right) \tag{5.31}
\end{align*}
$$

with $1 \leq i_{l}, j_{l} \leq e$ are component indices, $1 \leq u_{l}, v_{l} \leq N$ are strip indices for $1 \leq l \leq m$. We denote the size of such multi-indices by $|I|=|J|=|U|=|V|=m$. In the above example, for the $a$-th chord, we have $i_{a}=s, j_{a}=t, u_{a}=n$ and $v_{a}=p$. In the product of pages defining a book, pages are ordered from the bottom up. Now instead of using the notation $X$ for a tangle chord diagram, we use the book notation $A_{I, J, U, V}$ which incorporates the
information about the chords on the tangle. We have:

$$
\begin{equation*}
Z_{f}(L)=\sum_{\substack{\text { chord } \\ \text { diagr. X }}} a_{X} X=\sum_{m \geq 0} \sum_{\substack{|I|=|J|=m \\|U|=|V|=m}} a_{I, J, U, V} A_{I, J, U, V}(L) \tag{5.32}
\end{equation*}
$$

However $\widetilde{Z}_{f}(L)$ uses powers of

$$
\begin{equation*}
\nu=Z_{f}(U)^{-1}=Z_{f}(\sqrt{ }) \tag{5.33}
\end{equation*}
$$

and with the above slicing performed on $\widetilde{L}$ we are able to use books $A_{I J U V}(\widetilde{L})=A_{\text {IJUV }}$. We can therefore write:

$$
\begin{equation*}
\widetilde{Z}_{f}(L)=\sum_{\substack{\text { chord } \\ \text { diagr. } \mathrm{x}}} c_{X} X=\sum_{m \geq 0} \sum_{\substack{|I|=|J|=m \\|U|=|V|=m}} c_{I, J, U, V} A_{I, J, U, V} \tag{5.34}
\end{equation*}
$$

where we have $c_{I, J, U, V}:=c_{A_{I J U V}}$.

## Chapter 6

## Behavior of the Kontsevich integral $\widetilde{Z}_{f}$ under band sum moves using the book notation

What we have in (1.30) is the following: if $h_{i j}$ denotes the band sum move map on links corresponding to the band sum move of the $i$-th component of $L$ over its $j$-th component, then we can write $L^{\prime}=h_{i j}(L)$, so that $\widetilde{Z}_{f}\left(L^{\prime}\right)=\widetilde{Z}_{f}\left(h_{i j} L\right)$. What we would like however is to find a map $\mathbf{h}_{i j}$ induced from $h_{i j}$ that acts on the Kontsevich integral expressed in book notation to yield its corresponding values after band sum moves. In so doing, we use a different notation for the Kontsevich integral in book notation as once it is written as a linear combination of books, the resulting object is no longer a link invariant. We write ${ }^{b} Z$ for $Z$ written in book notation. We seek a map $\mathbf{h}_{i j}$ that implements the band sum move on ${ }^{b} \widetilde{Z}_{f}$. We claim that there is such a map, and that moreover for any link $L$, we have $\mathbf{h}_{i j}\left({ }^{b} \widetilde{Z}_{f}(L)\right)={ }^{b} \widetilde{Z}_{f}\left(h_{i j} L\right)$. In other terms, the following diagram is commutative:


It is important to remember that using the isotopy invariance of $\widetilde{Z}_{f}$, we can arrange that the band sum does not introduce new local extrema, and is far away enough from the rest of the link $L$ that we can use the Long Chords Lemma.

We have the following fact about books; since no two chords can be positioned at the same height $t$ on a tangle, pages, which represent chords, can allow for the possibility to hold many other chords by virtue of the non-simultaneity of chords. In this manner there is no ambiguity as to what chord in a page is represented by which entry. Further if a page holds information about more than one chord, we can split the matrix in as many matrices as there are chords represented in the original page. We illustrate this situation presently:


This represents a page carrying information about two chords, one between the $g$-th and $k$-th components, the other between the $h$-th and $l$-th components. For the first chord, the foot on $K_{g}$ is in the $n$-th strip, and the foot on $K_{k}$ is in the $u$-th strip. For the second chord, the foot on $K_{h}$ is in the $p$-th strip, and the foot on $K_{l}$ is in the $v$-th strip. We have indicated to the left and above the matrix the block indices $g, h, k$ and $l$, and in small letters the strips within the blocks where a foot is ending, and those are $n, p, u$ and $v$. Such a matrix splits as follows:


Since a page holds some information about one chord only by non-simultaneity of chords, if a page displays the information about more than one chord, we can isolate the information about each chord as a direct sum of pages each of which carries information about a unique chord. If we call the original matrix $A$ and the two spin-off matrices $B$ and $C$, then inserting $A$ in a book of $m$ pages as follows:

$$
\begin{equation*}
A^{-} \times A \times A^{+} \tag{6.4}
\end{equation*}
$$

where $A^{-}$are the first $m^{-}$pages of the book, and $A^{+}$are the last $m^{+}$pages, with $m^{-}+1+$ $m^{+}=m$, then we can write:

$$
\begin{equation*}
A^{-} \times A \times A^{+}=A^{-} \times B \times A^{+}+A^{-} \times C \times A^{+} \tag{6.5}
\end{equation*}
$$

representing two chord diagrams, one of which has its $\left(m^{-}+1\right)$-st chord represented by the page $B$, the other has its $\left(m^{-}+1\right)$-st chord represented by $C$. We have done this for a page $A$ containing the information about two chords. We generalize (6.4) by iterating this process for pages that contain information about more than 2 chords and generalize (6.5)
by iteration for books with two or more pages written as a sum of matrices. All this can be easily summarized by saying that

$$
\bigoplus_{m \geq 0} \underset{1 \leq i \leq m}{\times} \operatorname{Symm}_{e N \times e N}
$$

forms a tensor algebra.

We now discuss the band sum move proper. By virtue of the fact that we have:
or by abuse of notation:

then studying the doubling of chords during the band sum move can be done one chord after another. It suffices to work with one page at a time. During the band sum move of the handle corresponding to the $i$-th component $K_{i}$ over the $j$-th component $K_{j}$ of the link $L$, we encounter two different situations:

where $n<p$. In this first case, a given chord starts and ends on the $j$-th component, with matrix representation given by:

where without loss of generality we have chosen $i<j$ and the case $i>j$ is dealt with by a simple change of basis. Under a band sum move we obtain:

which equals:

a sum of chord diagrams that correspond, in this order, to the following sum of matrices:

which combines into:


The second situation we can have is the case where the chord starts on the $j$-th component but ends on some other $l$-th component. Without loss of generality we can pick $i<j<l$. For other arrangements of these indices we modify the basis for our matrices accordingly. We discuss the case $l=i$ right after since it doesn't consist in a basis change. Pictorially we have:

with the following matrix representation:


After a band sum move we get the following chord diagram:

which equals:

represented, in this order, by the sum of matrices:

combining into:


For the case $l=i$, we pictorially have:

with the following matrix representation:


After a band sum move, we get the following diagram:

which equals:

represented in this order by the sum of matrices:

combining into:


We present now our main result:

Theorem 6.1. For an e-component link $L$ for which the $i$-th component $K_{i}$ is being band summed over the $j$-th component $K_{j}$, the induced map on ${ }^{b} \widetilde{Z}_{f}$ is denoted by $\mathbf{h}_{i j}$ and is defined by:

$$
\begin{equation*}
\mathbf{h}_{i j}{ }^{b} \widetilde{Z}_{f}(L)=\sum_{m \geq 0} \sum_{\substack{|I|=|J|=m \\|U|=|V|=m}} c_{I J U V}{\underset{1}{1 \leq a \leq m}}_{\times} M_{i j}^{T} A_{i_{a}, j_{a}, u_{a}, v_{a}} M_{i j} \tag{6.21}
\end{equation*}
$$

where $M_{i j}$ is an $e N \times e N$ matrix with ones on its diagonal and the $j i$ block is the $N \times N$ identity matrix $I_{N}$. For $i<j$, we write such a matrix as:


If we define:

$$
\begin{equation*}
M_{i j}^{m}:=\underset{1 \leq a \leq m}{\times} M_{i j} \tag{6.23}
\end{equation*}
$$

acting on books of $m$ pages chord-wise, then we can rewrite the above formula in compact form as:

$$
\begin{equation*}
\mathbf{h}_{i j}^{b} \widetilde{Z}_{f}(L)=\sum_{m \geq 0} \sum_{\substack{|I|| | J|=m\\| U|=|V|=m}} c_{I J U V} M_{i j}^{m T} A_{I J U V} M_{i j}^{m} \tag{6.24}
\end{equation*}
$$

We may even generalize this further by defining $M_{i j}$ to be the product of as many matrices $M_{i j}$ as the book they operate on have pages, which leads to having the even simpler formula:

$$
\begin{equation*}
\mathbf{h}_{i j}{ }^{b} \widetilde{Z}_{f}(L)=M^{\cdot T}\left({ }^{b} \widetilde{Z}_{f}(L)\right) M \tag{6.25}
\end{equation*}
$$

Proof. The only cases that need to be studied are those where a chord has a foot on the $j$-th component, since all other tangle chord diagrams are left invariant by $\mathbf{h}_{i j}$, and of those there are only two kinds that we presented just before the statement of the Theorem. By elementary matrix multiplication, we have in the first case, for one page:

which is what we expected from the above considerations. In the second case of interest, if $l \neq i$ :

which is what was expected as well. In case $l=i$ :


This completes the proof.

## Chapter 7

## Isotopy invariance of $\widetilde{Z}_{f}$ in book notation

From Theorem 2.2.5 $\hat{Z}_{f}$ is an isotopy invariant. By definition, so is $\widetilde{Z}_{f}$. If we consider $\widetilde{Z}_{f}(L)$ as just a formal sum of books with complex coefficients however, we have a representation of tangle chord diagrams by books that is no longer isotopy invariant, whence the introduction in the previous chapter of the notation ${ }^{b} Z$. We use the fact that ambient isotopic links are in the same class if and only if they are related by the Reidemeister moves [Rei] [Y]. Instead of studying the behavior of ${ }^{b} \widetilde{Z}_{f}(L)$ under any arbitrary isotopy, we study Reidemeister moves of tangles. They are pictured as follows:



which is the framed version of the original $\Omega .1$ Reidemeister move.



Proposition 7.1. The book representation of $\widetilde{Z}_{f}(L)$ of any link $L$ is invariant under the Reidemeister moves $\Omega .2$ and $\Omega .3$.

Proof. In both cases we have an equivalence of two tangle diagrams which yield the same page for any chord ending on them after possibly moving the chord up the tangle.

Those are the only moves that leave ${ }^{b} \widetilde{Z}_{f}(L)$ invariant. For the other Reidemeister moves, there are a few changes to be performed on each page.

### 7.1 Behavior of ${ }^{b} \widetilde{Z}_{f}$ under $\Delta . \pi .1$

In a slice presentation of knots, and using the isotopy invariance of $\widetilde{Z}_{f}$, it suffices to consider strands other than those involved in the move $\Delta . \pi .1$ to be vertical, and the straightened out strand involved in $\Delta . \pi .1$ to be straight as well. For instance, one of the moves involved in $\Delta . \pi .1$ would look like:

In the expression for $\hat{Z}_{f}$, powers of $\nu$ cancel self-chords from the Kontsevich integral of the hump. Chords with one foot on the hump however are not canceled, whereas in the expression for $\hat{Z}_{f}$ of the right hand side of (7.6), all of whose strands are vertical, there are no such chords. This is in particular true of $\widetilde{Z}_{f}$. Moreover mixed chord diagrams involving self chords on the hump and long chords with one foot on the hump are not canceled by $\nu$. They are canceled by squeezing the hump into a window of vanishing width. For each long chord with a foot on an ascending (resp. descending) part of the hump, there is a long chord with a foot on a descending (resp. ascending) part of the hump, each tangle chord diagram with $\log$ differentials that differ in sign and therefore cancel each other off. This is how we have all long chords ending on the hump go. Thus we seek a transformation on the expression for ${ }^{b} \widetilde{Z}_{f}$ of the left hand side of (7.6) that effectively gets rid of chords with a foot on the hump.

Proposition 7.1.1. The $e N \times e(N-8)$ matrix $\operatorname{diag}\left(M_{\text {hump } \mapsto \mid}\right)$ with:

$$
\begin{equation*}
\mathrm{M}_{\mathrm{hump} \mapsto 1}=\left(\frac{I_{n} \mid}{\frac{O_{8 \times(N-8)}}{\sqrt{I_{N-n-8}}}}\right) \tag{7.7}
\end{equation*}
$$

implements the move (7.6). In case the hump is on the first strand from the left:

$$
\begin{equation*}
\mathrm{M}_{\text {hump } \mapsto 1}=\binom{\frac{O_{8 \times(N-8)}}{}}{I_{N-8}} \tag{7.8}
\end{equation*}
$$

If we let $I_{0}$ be the empty matrix then (7.7) implements the move (7.6) for all $n \geq 0$.
Proof. First observe that since we have a local max on the hump, the skeleton for both $\hat{Z}_{f}$
and $\widetilde{Z}_{f}$ is as follows, with added vertical slices for numbering purposes:


The straightened out strand on the right hand side of $(7.6)$ is in the $(n+1)$-st strip. The matrix that takes care of relabeling all the strips under the map depicted in (7.6) and kills all the chords on the hump is, for the block corresponding to the component on which the hump is located:

$$
\begin{equation*}
\mathrm{M}_{\text {hump } \mapsto 1}=\left(\frac{I_{n} \mid}{\frac{O_{8 \times(N-8)}}{\sqrt{I_{N-n-8}}}}\right) \tag{7.10}
\end{equation*}
$$

It is a $N \times(N-8)$ matrix. In other horizontal slices, chords ending on strands other than the one with the hump are moved up to be either in strips indexed 1 through $n$, or $n+9$ through $N$. Therefore the above matrix also takes care of the transformation of those chords under the move (7.6). We conclude that the $e N \times e(N-8)$ matrix $\operatorname{diag}\left(M_{\text {hump } \mapsto \mid}\right)$ with $M_{\text {hump } \mapsto \mid}$ given above implements the move (7.6).

One would similarly show that the move:
with a local slicing of the hump given by:

is taken care of by the same $e N \times e(N-8)$ matrix $\operatorname{diag}\left(M_{\text {hump }}\right)$ defined above.
Remark 7.1.2. We could equally have taken the following slicing and skeleton:

and the matrix would have been the same.

We have shown:

Proposition 7.1.3. The map $\mathbf{h}_{\Delta . \pi .1 \text {, hump } \mapsto \mid}$ that acts on $\widetilde{Z}_{f}(L)$ in book notation to effect two of the $\Delta . \pi .1$ Reidemeister moves as depicted in (7.6) and (7.11) is given on a page by:
where $M_{\Delta \cdot \pi .1, \text { hump } \mapsto \mid}$ is the $e N \times e(N-8)$ matrix $\operatorname{diag}\left(M_{\text {hump } \mapsto \mid}\right)$ with $M_{\text {hump } \mapsto \mid}$ given by:

$$
\begin{equation*}
\mathrm{M}_{\text {hump } \mapsto 1}=\left(\frac{I_{n} \mid}{\frac{O_{8 \times(N-8)}}{\sqrt{I_{N-n-8}}}}\right) \tag{7.15}
\end{equation*}
$$

for $n \geq 0$.

Proposition 7.1.4. The matrix that implements the move:

$$
\begin{equation*}
|\ldots| \quad|\quad| \rightarrow|\ldots| \bigcap^{\prime}|\ldots| \tag{7.16}
\end{equation*}
$$

with the same slicing convention as in (7.9) is given by the $N \times(N+8)$ matrix:

$$
\mathrm{M}_{\mid \mapsto \text { hump }}=\left(\begin{array}{c||c|c}
I_{n} & &  \tag{7.17}\\
& O_{N \times 8} & I_{N-n}
\end{array}\right)
$$

for $n \geq 0$.

Proof. In other horizontal slices, chords ending on strands other than the one on which the hump is located are moved up to end in strips indexed 1 through $n$ or $n+9$ through $N$. Therefore the above matrix also takes care of the transformation of those chords under the move (7.16). We conclude that the $e N \times e(N+8)$ matrix $\operatorname{diag}\left(M_{\mid \mapsto \text { hump }}\right)$ with $M_{\mid \mapsto \text { hump }}$ given above implements the move (7.16).

Likewise, one would show that the same matrix implements the move:

$$
\begin{array}{|l|}
\cdots  \tag{7.18}\\
\cdots
\end{array}|\quad| \cdots|\cdots| \quad \bigcup|\ldots|
$$

with the slicing on the hump given as in (7.12). We have shown:

Proposition 7.1.5. The map $\mathbf{h}_{\Delta . \pi .1, \mid \mapsto \text { hump }}$ that acts on $\widetilde{Z}_{f}(L)$ in book notation to effect two of the $\Delta . \pi .1$ Reidemeister moves as depicted in (7.16) and (7.18) is given on a page by:

$$
\begin{equation*}
\underset{\mid \mapsto h i n m p}{\mathbf{h}_{\mid . \pi .1}} A:=\underset{\mid \mapsto \text { hump }}{M_{\Delta . \pi .1}^{T}} A M_{\mid \mapsto \text { hump }}^{\Delta . \pi .1}, \tag{7.19}
\end{equation*}
$$

where $M_{\Delta . \pi .1, \mid \mapsto \text { hump }}$ is the $e N \times e(N+8) \operatorname{diag}\left(M_{\mid \mapsto \text { hump }}\right)$ matrix with $M_{\mid \mapsto \text { hump }}$ given by:

$$
\mathrm{M}_{\mid \mapsto \text { hump }}=\left(\begin{array}{l||l|l}
I_{n} & &  \tag{7.20}\\
& O_{N \times 8} & I_{N-n}
\end{array}\right]
$$

Finally we consider the map that implements the following move on strands:

$$
\begin{equation*}
|\cdots| \int_{\Omega}|\cdots| \rightarrow|\cdots| \bigcup_{n}|\cdots| \tag{7.21}
\end{equation*}
$$

With slicing conventions as in (7.9) and (7.12), the corresponding transformation matrix is given by the $N \times N$ matrix:

$$
\mathrm{M}_{\text {hump } \mapsto \text { hump }}=\left[\begin{array}{ccc}
I_{n} & &  \tag{7.22}\\
& O_{8 \times 8} & \\
& & \begin{array}{|c}
I_{N-n-8}
\end{array}
\end{array}\right]=\mathrm{M}_{\text {hump } \mapsto \mid} \cdot \mathrm{M}_{\mid \mapsto \text { hump }}
$$

In other horizontal slices, chords ending on strands other than the one on which the hump is located are moved up to end in strips indexed 1 through $n$ or $n+9$ through $N$. Therefore the above matrix also takes care of the transformation of those chords under the move (7.21). We conclude that the $e N \times e N$ matrix $\operatorname{diag}\left(M_{\text {hump } \mapsto \text { hump }}\right)$ with $M_{\text {hump } \mapsto \text { hump }}$ given above implements the move (7.21). We have shown:

Proposition 7.1.6. The map $\mathbf{h}_{\Delta . \pi .1 \text {, hump } \mapsto \text { hump }}$ that acts on $\widetilde{Z}_{f}(L)$ in book notation to effect the $\Delta . \pi .1$ Reidemeister move as depicted in (7.21) is given on a page by:

$$
\begin{equation*}
\underset{\text { hump } \mapsto \text { hump }}{\Delta \cdot \pi \cdot 1} A:=\underset{\text { hump } \uparrow \text { hump }}{\Delta} A M \underset{\text { hump } \mapsto \text { hump }}{\Delta \cdot \pi \cdot 1} \tag{7.23}
\end{equation*}
$$

where $M_{\Delta . \pi .1, \text { hump } \mapsto \text { hump }}$ is the $e N \times e N$ matrix $\operatorname{diag}\left(M_{\text {hump } \mapsto \text { hump }}\right)$ with $M_{\text {hump } \mapsto \text { hump }}$ given by:

$$
\mathrm{M}_{\text {hump } \rightarrow \text { hump }}=\left[\begin{array}{ccc}
I_{n} & &  \tag{7.24}\\
& & \\
& & O_{8 \times 8} \\
& & \boxed{I_{N-n-8}}
\end{array}\right]
$$

### 7.2 Behavior of ${ }^{b} \widetilde{Z}_{f}$ under $\Delta . \pi .2$

The two tangle diagrams involved in the statement for the $\Delta . \pi .2$ Reidemeister move have the following presentation with vertical strips for book purposes:

and:


Proposition 7.2.1. With tangle diagrams as above, the map $\mathbf{h}_{\Delta . \pi .2}$ that acts on $\widetilde{Z}_{f}(L)$ in book notation to effect the Reidemeister move $\Delta . \pi .2$, is given on a page by:

$$
\begin{equation*}
\mathbf{h}_{\Delta . \pi .2} A:=M_{\Delta . \pi .2}^{T} A M_{\Delta . \pi .2} \tag{7.27}
\end{equation*}
$$

where $M_{\Delta . \pi .2}$ is the $e N \times e N$ identity matrix save for the $t t$ block which is of the form:


Proof. Recall that chords on a given tangle are moved until they reach a local max. Keeping this in mind, under the move $\Delta . \pi .2$, only the $t$-th component moves, and thus only chords with a foot on this component need be considered. Suppose a chord is stretching between that component and the $s$-th component. Moving the chord up the $s$-th component so that it's localized near its local max, the foot of that chord on the $s$-th component is either in the $n+1$-st or $n+2$-nd strip, and its other foot on the $t$-th component is then in the $n$-th or $n+3$-rd strip depending on which tangle we are looking at. Under $\Delta . \pi .2$, chords with one foot in the $n$-th strip (resp. the $n+3$-rd strip) are moved to end up in the $n+3$-rd strip (resp. the $n$-th strip). Suppose now a chord is stretching between the $t$-th component and another $l$-th component, $l \neq s$. For a needle shaped local max on the $s$-th component, the corresponding vertical strips indexed by $n+1$ and $n+2$ are sufficiently narrow that chords from some $l$-th component, $l \neq s$ end on the $t$-th component in some strip other than the $n+1$-st or $n+2$-nd strips, that is in the $n$-th or $n+3$-rd strips, which are interchanged under $\Delta . \pi .2$. We conclude that $\Delta . \pi .2$ is implemented by a $e N \times e N$ matrix with a $t t$ block of the form above which effectively switches the $n$-th and $n+3$-rd strips for $n \geq 1$, taking $I_{0}$ to be the empty matrix in the event that $n=1$ indexes the first vertical strip.

### 7.3 Behavior of ${ }^{b} \widetilde{Z}_{f}$ under $\Omega$.1. $f$

The two equivalent tangle diagrams under this move are represented as follows along with the strips for book representation purposes:

and:


Proposition 7.3.1. With tangle diagrams as above, the map $\mathbf{h}_{\Omega .1 . f}$ that acts on $\widetilde{Z}_{f}(L)$ in book notation to effect the Reidemeister move $\Omega$.1.f on the $l$-th component of $L$, is given on a page by:

$$
\begin{equation*}
\mathbf{h}_{\Omega .1 . f} A:=M_{\Omega .1 . f}^{T} A M_{\Omega .1 . f} \tag{7.31}
\end{equation*}
$$

where $M_{\Omega .1 . f}$ is the $e N \times e N$ identity matrix save for the $l l$ block which is of the form:

with ones on the transverse diagonal.

Proof. With the vertical slicing as in (7.29) and (7.30), it is immediate that the above matrix implements the move $\Omega$.1.f on the $l$-th component. Chords with feet on other strands are in strips indexed 1 through $n$ or $n+10$ through $N$ and stay put under $\Omega$.1.f, so for those we just use an identity matrix. We conclude that the above matrix implements the move $\Omega .1 . f$.

### 7.4 Adding and deleting strips

If $M_{n}^{+}$denotes the matrix that inserts a strip between the $n$-th and $n+1$-st strip, then the transformation on pages that effects such a change is the $e N \times e(N+1)$ matrix $\operatorname{diag}\left(M_{n}^{+}\right)$ with:

$$
\mathrm{M}_{n}^{+}={ }^{\mathrm{n}+1}\left(\begin{array}{l}
\left.I_{n}\right|^{0} \mathrm{n}+1  \tag{7.33}\\
\vdots \\
0
\end{array}\right)
$$

and if $M_{n}^{-}$denotes the matrix that deletes the $n$-th strip, then the transformation on pages that effects such a change is the $e(N+1) \times e N$ matrix $\operatorname{diag}\left(M_{n}^{-}\right)$with:

$$
\begin{equation*}
\mathrm{M}_{n}^{-}=\left(\frac{I_{n-1}^{\mathrm{n}-1} \mid}{I_{N-n+2}}\right) \tag{7.34}
\end{equation*}
$$

## Chapter 8

## Behavior of $\widetilde{Z}_{f}(L)$ in book notation under a change of orientation

### 8.1 Behavior of ${ }^{b} \widetilde{Z}_{f}$ under orientation change

Recall that the links we deal with are oriented. Under a change of orientation on a component of a link $L$, any chord diagram with a foot on that component has its coefficient being multiplied by -1 . The map on chord diagrams in $\hat{\mathcal{A}}\left(\amalg S^{1}\right)$ that effects this change is denoted by $S_{r}$ for the $r$-th component of the link $L$ whose orientation is being changed, $1 \leq$ $r \leq e$ [LM3]. We generalize this to tangle chord diagrams: $S_{r}$ is the map $\hat{A}\left(\coprod_{1 \leq l \leq e} K_{l}\right) \rightarrow$ $\hat{A}\left(\coprod_{1 \leq l \neq r \leq e} K_{l}+S_{r} K_{r}\right)$ induced from a change of orientation $S_{r} K_{r}$ on $K_{r}$. The book notation should reflect this change. Pictorially, the following local tangle chord diagram:

is represented in book notation by:

where without loss of generality we have chosen $r<t$. The same tangle chord diagram with the reverse orientation on the $r$-th component:

has book representation:


Proposition 8.1. With tangle diagrams as above, the map $\mathbf{h}_{S, r}$ that acts on $\widetilde{Z}_{f}(L)$ in book notation to effect the orientation change $S_{r}$ on the $r$-th component of $L$, is given on a page by:

$$
\begin{equation*}
\mathbf{h}_{S, r} A:=M_{S, r}^{T} A M_{S, r} \tag{8.5}
\end{equation*}
$$

where $M_{S, r}$ is the $e N \times e N$ identity matrix save for the $r r$ block which is the negative of the $N \times N$ identity matrix, $-I_{N}$ :


There are as many modified blocks such as the $r r$-th above as there are components whose
orientation is changed. For a chord with feet on the $r$-th and $t$-th components both of whose orientations are reversed, the matrix $M_{S, r, t}$ with two blocks $r r$ and $t t$ as $-I_{N}$ will effect the orientation change while leaving the coefficient in front of corresponding chord diagrams unchanged since $(-1) \times(-1)=1$.

Proof. Matrix multiplication.

### 8.2 Behavior of $\widetilde{Z}_{f}(L)$ as $K_{i}$ is being subtracted from $K_{j}$

We now consider the effect of having a subtraction of $K_{i}$ from $K_{j}$ as a result of operating a band sum move of $K_{i}$ over $K_{j}$. As defined in [RK1], this corresponds to having a band such that upon doing the band sum move the components $K_{i}$ and $K_{j}$ end up having opposite orientations. If we locally represent the pieces of those two components with a chord stretching between them as:

then this chord behaves as follows under a band sum move of $K_{i}$ over $K_{j}$ :


If the latter tangle chord diagram has a coefficient of $c_{X}$ in the expression for $\widetilde{Z}_{f}$, the former has a coefficient $-c_{X}$ however due to the orientation change resulting in the Kontsevich integral picking up an overall minus sign. It follows that the statement of Theorem 1.6.3 is no longer true in this case. We can nevertheless remedy this as follows.

Proposition 8.2.1. ([LM4], [LM5]) $Z_{f}\left(S_{K_{i}} L\right)=S_{K_{i}} Z_{f}(L)$ for $1 \leq i \leq e$.
Proof. This follows from the definition of $Z_{f}$ and the fact that $Z\left(S_{K_{i}} L\right)=S_{K_{i}} Z(L)$.
Theorem 8.2.2. $([\mathrm{LM} 4],[\mathrm{LM} 5]) \hat{Z}_{f}\left(S_{K_{i}} L\right)=S_{K_{i}} \hat{Z}_{f}(L)$ for $1 \leq i \leq e$.

Proof. It suffices to write:

$$
\begin{align*}
\hat{Z}_{f}\left(S_{K_{i}} L\right) & =\nu^{m_{1}} \otimes \cdots \otimes \nu^{m_{i}} \otimes \cdots \otimes \nu^{m_{e}} \cdot Z_{f}\left(S_{K_{i}} L\right)  \tag{8.9}\\
& =\nu^{m_{1}} \otimes \cdots \otimes S_{K_{i}} S_{K_{i}} \nu^{m_{i}} \otimes \cdots \otimes \nu^{m_{e}} \cdot S_{K_{i}} Z_{f}(L)  \tag{8.10}\\
& =S_{K_{i}}\left(\nu^{m_{1}} \otimes \cdots \otimes\left(S_{K_{i}} \nu\right)^{m_{i}} \otimes \cdots \otimes \nu^{m_{e}} \cdot Z_{f}(L)\right)  \tag{8.11}\\
& =S_{K_{i}} \hat{Z}_{f}(L) \tag{8.12}
\end{align*}
$$

Theorem 8.2.3. $\widetilde{Z}_{f}\left(S_{K_{i}} L\right)=S_{K_{i}} \widetilde{Z}_{f}(L)$ for $1 \leq i \leq e$.

Proof. Same as for the previous Theorem.

Then we consider the following diagram, where without loss of generality we have reversed the orientation of the $i$-th component:


First applying the map $S_{K_{i}}$ to $\widetilde{Z}_{f}(L)$ we obtain $\widetilde{Z}_{f}\left(S_{K_{i}} L\right)$ by Theorem 8.2.3. We write this quantity as $\sum c_{X} \cdot X$. We can apply Theorem 1.6.3 for $S_{K_{i}} L$ since then the band sum move results in $K_{i}$ being added to $K_{j}$. We obtain $\sum c_{X} \cdot X^{\prime}$ under the band sum move. This equals $\widetilde{Z}_{f}\left(\left[S_{K_{i}} L\right]^{\prime}\right)$. By further reversing the orientation of $K_{i}$, we get, using Theorem 8.2.3 again, $\widetilde{Z}_{f}\left(S_{K_{i}}\left[S_{K_{i}} L\right]^{\prime}\right)$. Since $S_{K_{i}}\left[S_{K_{i}} L\right]^{\prime}=L^{\prime}$, this last quantity is $\widetilde{Z}_{f}\left(L^{\prime}\right)$. We have that the band sum move in the event of a subtraction is given by closing the above diagram, and this corresponds to the composition of $S_{K_{i}}$, a band sum, and $S_{K_{i}}$. Each of those maps can be represented in book notation by three respective matrices whose action is by congruence. This composition is then easily represented by a congruence with a matrix obtained by multiplying those three matrices.

## Chapter 9

## Recovering $Z_{f}(L) \in \overline{\mathcal{A}}\left(\amalg S^{1}\right)$ from ${ }^{b} Z_{f}(L)$

Proposition 9.1. For an unknown $e$-component link $L$ represented as a plat, from ${ }^{b} Z_{f}(L)$ we can recover its expression in $\hat{\mathcal{A}}\left(\amalg^{e} S^{1}\right)$, without knowing the skeleton on which tangle chord diagrams represented by the books are based.

Proof. It suffices to know the algebraic number of crossings between strands. Indeed, chord diagrams are blind to the winding of the skeleta of tangle chord diagrams. Now, the degree one coefficient of the following tangle (where a crossing is either positive or negative, $\lambda>0$, $\epsilon= \pm 1, n$ the number of half-twists):

in $Z_{f, 1}(L)$ is:

$$
\begin{equation*}
\pm \frac{1}{2 \pi i} \log \frac{\lambda e^{i \epsilon n \pi} \Delta z}{\Delta z}= \pm\left(\frac{1}{2 \pi i} \log \lambda\right) \pm \epsilon \frac{n}{2} \tag{9.2}
\end{equation*}
$$

Thus if we have an odd number of half-twists between 2 strands, $n / 2$ is fractional, an integer
otherwise. It follows that from ${ }^{b} Z_{f, 1}(L)$ we can determine for a fixed strand the collection of strands to its right near the top of $L$ it has an odd number of half-twists with, or equivalently what are those strands to its right near the top of $L$ it passes on the right at the bottom of $L$. Doing this for all strands we can determine the permutation that to any strand near the top of $L$ associates a strand near the bottom of the link. This is sufficient to associate to any book the element of $\mathcal{A}\left(\amalg^{e} S^{1}\right)$ it corresponds to.

## Chapter 10

## Normalizations in LMO Invariant Theory

Le, Murakami and Ohtsuki [LM5], [LM6] defined a quotienting map that enables one to construct a 3-manifold invariant from a renormalized version $\check{Z}_{f}$ of the hatted framed Kontsevich integral $\hat{Z}_{f}$. In Section 4.2, we showed using only the long chords lemma [ ChDu ], the multiplicativity of the framed Kontsevich integral (Proposition 2.2.2), as well as the fact that $\hat{Z}_{f}$ is an isotopy invariant (Theorem 2.2.5), that the normalization $\widetilde{Z}_{f}=\nu^{-1} \hat{Z}_{f}$ leads to a link invariant that is well-behaved under band sum moves for which in a small window where such moves are performed, link components of interest are locally on top of each other, and the resulting band is vertical. Independently, Le, Murakami and Ohtsuki showed ([LM3], [LM5], [LM6]) that if one uses the formalism of q-tangle diagrams [LM1] and carefully considers associators, one finds that another normalization $\check{Z}_{f}=\nu \hat{Z}_{f}$ is well-behaved under band sum moves for which in a small window where such moves are performed, link components of interest are locally put side by side, and the resulting band looks horizontal. In what follows we discuss these normalizations and show how they relate to each other. We then discuss the quotienting map [LM5] necessary for producing 3-manifold invariants from the Kontsevich integral and show that any normalization survives the quotienting process, and thus any choice of normalization yields a 3-manifold invariant.

### 10.1 Relation between $\widetilde{Z}_{f}$ and LMO's $\check{Z}_{f}$

In Chapter 4 we showed that under band sum moves,

$$
\begin{equation*}
\widetilde{Z}_{f}(L)=\sum_{\substack{\text { chord } \\ \text { diagrams } X}} c_{X} X \tag{10.1}
\end{equation*}
$$

maps to:

$$
\begin{equation*}
\widetilde{Z}_{f}\left(L^{\prime}\right)=\sum_{\substack{\text { chord } \\ \text { diagrams } X}} c_{X} X^{\prime} \tag{10.2}
\end{equation*}
$$

where $L$ is a given link, $L^{\prime}$ is the same link except that one of its components is band summed over another of its components by a vertical band that is a band whose midline is parallel to the time axis, and $c_{X}$ are coefficients arising from computing the renormalized framed Kontsevich integrals of the appropriate framed links. In those expressions, $X$ and $X^{\prime}$ are related as follows:


We proved such a statement by stretching one strand from each of the two components involved in the band sum to one side as in:

and by performing a band sum:

where in this particular instance, the $j$-th component is being band summed over. In [LM2], [LM3], [LM5], [LM6], [O] however, computations are performed by considering strands involved in the band sum side by side, resulting in a horizontal band that is a band whose midline is transverse to the time axis:

$$
\begin{equation*}
\left.\hat{Z}_{f}() \square\right) \quad \mapsto \quad \hat{Z}_{f}(\cap \square) \tag{10.5}
\end{equation*}
$$

where the open strand is the portion of the link component operating the band sum and the closed component represents the link component being band summed over. This pictorial representation is symbolic and does not display how complex the link components are and whether they are linked to other link components. In order to compare both computations, we will use the equivalent normalization procedure whereby each local extremum is multiplied by $\nu^{1 / 2}$ instead of having only local maxima being multiplied by $\nu$. This we do only in this section.

We consider an e-component link $L$. We focus on two of its components $K_{i}$ and $K_{j}$, $1 \leq i, j \leq e$ and consider the band sum of $K_{i}$ over $K_{j}$. The resulting link after band sum we call L'. Using the isotopy invariance of $\hat{Z}_{f}$ we can pull up one local max from each of $K_{i}$
and $K_{j}$ so that we compute:

where $T$ is what remains of $L$, and we have displayed the local maximum from $K_{i}$ as the outer strand. Le Murakami et al ([LM3], [LM5], [LM6]) would compute $\hat{Z}_{f}$ after band sum move as:


The protruding local maxima at the top are not necessary for Le and Murakami's computation, but are convenient to emphasize the relation with the computation that we suggest in this work:


An intermediate stage both computations should be equal to by isotopy invariance of $\hat{Z}_{f}$ is the following:


In the next section, we show (10.7) evaluates to (10.9). In the section that follows, we show (10.8) evaluates to this intermediate value, and this will show that whether we compute $\hat{Z}_{f}\left(L^{\prime}\right)$ using LMO's procedure of locally putting link components involved in the band sum side by side or putting them on top of each other as done in the present work, one really gets the same thing. In the last section, we take stock and draw some conclusions.

### 10.1.1 LMO's computation

In [LM3], [LM5], [LM6], [O], a window is fixed inside of which one has the following tangle:


Then one has:

$$
\begin{equation*}
\hat{Z}_{f}\left(L^{\prime}\right)=\hat{Z}_{f}(\stackrel{\bigcup}{\bigcap}) \cdot \Delta_{j} \hat{Z}_{f}(L \backslash \mid) \tag{10.11}
\end{equation*}
$$

In [O], using q-tangles, it is shown that:


Outside of the window one has $\Delta_{j} \hat{Z}_{f}[\mathrm{M}](L \backslash| |)$. This comes from computing $\hat{Z}_{f}[\mathrm{M}]\left(\Delta_{j}(L \backslash\right.$ $\|\|)$. We have equality of these two quantities in the case of geometric tangles only when the duplicate $K_{j}^{(2)}$ of $K_{j}$ is analytically coincident with $K_{j}$, that is in the window one really works with:


Then:

$$
\begin{equation*}
\hat{Z}_{f}\left(L^{\prime}\right)=\hat{Z}_{f}[\mathrm{M}]\left(L^{\prime}\right)=\hat{Z}_{f}[\mathrm{M}]\left(\bigcap_{\Omega}\right) \cdot \Delta_{j} \hat{Z}_{f}[\mathrm{M}](L \backslash \mid) \tag{10.14}
\end{equation*}
$$

where the value of $\left.\hat{Z}_{f}[\mathrm{M}]()_{\text {) }}\right)$ is given by the following lemma:

Lemma 10.1.1.1. For $M>0$,

$$
\begin{equation*}
\hat{Z}_{f}[\mathrm{M}]\left(\bigcap_{M}\right)={ }^{\nu^{\frac{1}{2}}} \tag{10.15}
\end{equation*}
$$

and:

$$
\begin{equation*}
\hat{Z}_{f}[\mathrm{M}](\stackrel{\square}{\rightrightarrows})=\stackrel{{ }_{\nu^{\frac{1}{2}}}}{ } \tag{10.16}
\end{equation*}
$$

Proof. It suffices to consider the case of a local maximum, the local minimum computation being the same. We have, by isotopy invariance of $\hat{Z}_{f}$ that:

$$
\begin{equation*}
\hat{Z}_{f}[\mathrm{M}](\bigcap)=\hat{Z}_{f}[\mathrm{M}]\left(\bigcap_{\xrightarrow[M]{ }}^{\stackrel{M}{\longrightarrow}}\right)=\nu^{1 / 2} \otimes 1 \cdot Z_{f}[\mathrm{M}](\bigcap) \tag{10.17}
\end{equation*}
$$

where the first factor in the tensor product multiplies the local max and the identity multiplies the straight strand. We have:

$$
\begin{equation*}
Z_{f}[\mathrm{M}](\bigcap)=\bigcap \tag{10.18}
\end{equation*}
$$

from which we get:

Using this lemma it follows:

$$
\begin{equation*}
\hat{Z}_{f}[\mathrm{M}]\left(\bigcup_{)}\right)=\underbrace{\nu^{\nu^{\frac{1}{2}}}} \tag{10.20}
\end{equation*}
$$

$\hat{Z}_{f}\left(L^{\prime}\right)$ as computed by Le, Murakami et al and rewritten for geometric tangles as done above does equal (10.9) as the following Proposition shows:

## Proposition 10.1.1.2.

$$
\begin{equation*}
\hat{Z}_{f}[\mathrm{M}]\left(\bigcup_{\cap}\right) \cdot \Delta_{j} \hat{Z}_{f}[\mathrm{M}](L \backslash| |)=\hat{Z}_{f}(\cap) \tag{10.21}
\end{equation*}
$$

Proof. We slice $L \backslash \mid$ into three horizontal strips, the middle one containing the window that was initially frozen. To the far right of that strip, strands see those in the window as straight strands as $C \gg M$. It follows that we can replace in the above expression $\hat{Z}_{f}[\mathrm{M}]\left(\bigcup_{\text {, }}\right)$ by $\hat{Z}_{f}[\mathrm{M}]$ of the middle strip. By multiplicativity of the Kontsevich integral this equals

which by isotopy invariance of $\hat{Z}_{f}$ is equal to (10.9) that we reproduce below for convenience:


### 10.1.2 Alternate computation

In the present paper we considered doing:


Slice $L^{\prime}$ into three horizontal slices so that the window containing the tangle:

$$
\begin{equation*}
r \tag{10.25}
\end{equation*}
$$

is exactly enclosed within the middle strip. Since $A \gg M$ and $B \gg M$ strands other than those in the window see those from the window as being straight, and thus we can write:

$$
\begin{align*}
& \hat{Z}_{f}[\mathrm{M}](\curvearrowleft) \cdot \Delta_{j} \hat{Z}_{f}[\mathrm{M}](L \backslash \bigcup)  \tag{10.26}\\
& =\hat{Z}_{f}[\mathrm{M}](\sim) \times \hat{Z}_{f}[\mathrm{M}](\text { ) }  \tag{10.27}\\
& =\hat{Z}_{f}[\mathrm{M}](\cap) \tag{10.28}
\end{align*}
$$

where in the last step we have used the isotopy invariance of $\hat{Z}_{f}$.

### 10.1.3 Taking stock: conclusions

In [O], [LM3], [LM5], [LM6], the behavior of $\hat{Z}_{f}$ under band sum moves on link components should be more properly called map on $\hat{Z}_{f}$ induced by band sum moves on link components with windows. The following statement holds at the level of q-tangles:

In other terms:

where the pack of chords displayed above comes from computing $\hat{Z}_{f}(L \backslash \mid)$. In other terms those chords are very much dependent on the window we have chosen. Multiplying both expressions by a power of $\nu$ on each link component we obtain the following statement about $\check{Z}_{f}$ :


That is for $\check{Z}_{f}$ one has the simple statement:


In the present paper we prefer to work with geometric tangles, for which a computation such as the one above can make use of windows, provided we deal with analytically coincident
tangles. In this case we obtain:

$$
\begin{align*}
\hat{Z}_{f}(L) & =\hat{Z}_{f}[\mathrm{M}](\smile) \cdot \hat{Z}_{f}[\mathrm{M}](L \backslash \smile)  \tag{10.37}\\
& \longmapsto \hat{Z}_{f}[\mathrm{M}]\left(L^{\prime}\right)=\hat{Z}_{f}[\mathrm{M}](\curvearrowleft) \cdot \Delta_{j} \hat{Z}_{f}[\mathrm{M}](L \backslash \underset{\curvearrowleft}{\curvearrowleft}) \tag{10.38}
\end{align*}
$$

which we can explicity write as:

$$
\begin{align*}
& \underbrace{}_{\nu^{\frac{1}{2}}} \\
& =\underbrace{\sqrt[\nu^{\frac{1}{2}}]{ }} \cdot \hat{Z}_{f}[\mathrm{M}](L \backslash \bigcup) \mapsto\left|\sqrt{\nu^{\frac{1}{2}}}\right| \cdot \Delta_{j} \hat{Z}_{f}[\mathrm{M}](L \backslash \bigcup) \tag{10.39}
\end{align*}
$$

In other terms:

where the pack of chords come from $\hat{Z}_{f}[\mathrm{M}](L \backslash \underset{\bigcap}{\cong})$ as well as $\nu^{1 / 2}$ and is thus dependent on the choice of window. Multiplying this mapping by a factor of $\nu^{-1}$ on each link component
we get:

that is we have a statement of the form:

$$
\begin{equation*}
\widetilde{Z}_{f}(L)=\sum \rho \stackrel{\square}{\square} \quad \longmapsto \widetilde{Z}_{f}\left(L^{\prime}\right)=\sum \rho \tag{10.45}
\end{equation*}
$$

Thus the above statement, as well as:

are true in their own right, but are not related by a simple renormalization as the packs of chords displayed in each statement are very much dependent on the window selected in $\hat{Z}_{f}(L)$ and $\hat{Z}_{f}\left(L^{\prime}\right)$, from which one normalization or the other led to symbolic, window dependent statements about $\check{Z}_{f}$ or $\widetilde{Z}_{f}$. However if one keeps track of what those chords really
represent and one concatenates elements properly, one truly gets the following commutative diagram:

$$
\begin{align*}
& \check{Z}_{f}(L)=\sum \mu  \tag{10.47}\\
& \hat{Z}_{f}(| |) \cdot \hat{Z}_{f}\left(L \backslash| | \longrightarrow \ldots \hat{Z}_{f}\left(\text { Un }_{n}^{\prime \prime}\right) \cdot \Delta_{j} \hat{Z}_{f}(L \backslash| |)\right. \\
& \hat{Z}_{f}^{\prime \prime}(L) \longrightarrow \hat{Z}_{f}^{\prime \prime}
\end{align*}
$$

where horizontal maps denote a band sum move, and both diagonal and vertical maps are renormalizations. To be specific, ascending diagonal maps are multiplications of every link component in the expressions for $\hat{Z}_{f}(L)$ or $\hat{Z}_{f}\left(L^{\prime}\right)$ by $\nu$, thus defining the corresponding $\check{Z}_{f}$ values, and descending diagonal maps are multiplications by $\nu^{-1}$ as prescribed by the definition of $\widetilde{Z}_{f}$. The vertical maps are multiplications of each link component of $\check{Z}_{f}(L)$ and $\check{Z}_{f}\left(L^{\prime}\right)$ by $\nu^{-2}$ to get the corresponding $\widetilde{Z}_{f}$ values.

### 10.2 3-Manifold Invariants constructed from the Kontsevich Integral

In [LM5] a quotienting map is defined that allows one to construct a 3-manifold invariant from $\check{Z}_{f}$. The proof of invariance of the equivalence class of such an object under Kirby move II leaves room for other renormalizations to also be invariant by inserting a fixed power of $\nu$
on each link component. We first show that this is so and introduce a family of 3-manifold invariants built from renormalizations of $\hat{Z}_{f}$. In the section that follows we generalize this further by arguing that renormalization by elements of $\hat{\mathcal{A}}\left(\amalg^{e} S^{1}\right)$ that verify certain more general conditions also lead to a family of 3-manifold invariants.

### 10.2.1 Renormalizations of $\hat{Z}_{f}$ and the 3-manifold invariants they generate

We first review the requisite formalism and definitions as covered in [LM5]. If $\mathcal{A}(X)$ denotes the algebra of chord diagrams with support on $X \bmod$ the STU relation ([BN]):

then one defines $\mathcal{A}(X)$ to be obtained by including in $\mathcal{A}(X)$ dashed trivial circles as well, with a graded completion denoted by $\mathcal{A}(X)^{\wedge}$. For $2 n$ points arranged in a circle one denotes by $P_{n}$ the relation equating the sum of all pairings of those $2 n$ points by dotted arcs to zero. Denote by $L_{<2 n}$ the equivalence relation between elements of $\mathcal{A}\left(\amalg S^{1}\right)$ defined as follows: for $D_{1}, D_{2} \in \mathscr{A}\left(\amalg S^{1}\right), D_{1} \sim D_{2}$ if and only if $D_{1}-D_{2}$ is a sum of chord diagrams with less than $2 n$ chord feet on each summand. $O_{n}$ is the equivalence relation such that a dashed circle is equivalent to $-2 n$. In this section we use tangle chord diagrams that explicitly display chords ending on them. We build our reasoning on the statement (4.77):

$$
\begin{equation*}
\widetilde{Z}_{f}(L)=\sum \lambda \widetilde{Z}_{f}\left(L^{\prime}\right)=\sum \lambda \quad \square \tag{10.49}
\end{equation*}
$$

We could equally have started from:

as both renormalizations are equivalent by the following commutative diagram:


However chords displayed in (10.49) and (10.50) are different as this diagram shows. To avoid confusion, we fix chords that are shown to be complementary to the choice of window for which a convenient normalization is $\widetilde{Z}_{f}$. Thus we work with (10.49).

We now define a renormalization of such an invariant indexed by the number $k+1$ of
powers of $\nu$ by which it is renormalized as:

$$
\begin{align*}
& \stackrel{\vee}{ }(k)  \tag{10.52}\\
& \mathrm{Z}_{\mathrm{f}}
\end{align*}=\nu^{k+1} \otimes \cdots \otimes \nu^{k+1} \cdot \widetilde{Z}_{f}
$$

in such a manner that $\mathrm{Z}_{\mathrm{f}}{ }^{\vee(0)}=\hat{Z}_{f}, \mathrm{Z}_{\mathrm{f}}{ }^{\vee(-1)}=\widetilde{Z}_{f}$, and $\mathrm{Z}_{\mathrm{f}}{ }^{\vee(1)}=\check{Z}_{f}$ as defined in [LM5]. If we multiply each link component of (10.49) by a $k+1$-st power of $\nu$, then we get:


The following result was originally given in [LM5] in the case $k=1$.
Proposition 10.2.1.1. ([LM5]) For $k \in \mathbb{Z}, L$ an $e$-components framed oriented link and $n$ an integer, the class $\left[\mathrm{Z}_{\mathrm{f}}{ }^{\vee}(k)(L)\right]$ in $\mathcal{A}\left(\coprod^{e} S^{1}\right)^{\wedge} / L_{<2 n}, P_{n+1}, O_{n}$ is independent of the orientation of the link $L$ and is invariant under band sum moves.

Proof. We first show the invariance of $\left[\mathrm{Z}_{\mathrm{f}}{ }^{\vee}(k)(L)\right]$ under orientation change. For $K$ a link component, $S_{(K)}$ the map on links that implements the orientation reversal on $K$, we have:

$$
\begin{align*}
\vee_{\mathrm{f}} \mathrm{Z}_{\mathrm{f}}\left(S_{(K)} L\right) & =\nu^{k} \otimes \cdots \otimes \nu^{k} \cdot \hat{Z}_{f}\left(S_{(K)} L\right)  \tag{10.54}\\
& =\nu^{k} \otimes \cdots \otimes \nu^{k} \cdot S_{(K)} \hat{Z}_{f}(L)  \tag{10.55}\\
& =S_{(K)}\left(\nu^{k} \otimes \cdots \otimes S_{(K)} \nu^{k} \otimes \cdots \otimes \nu^{k} \cdot \hat{Z}_{f}(L)\right)  \tag{10.56}\\
& =S_{(K)} \quad{ }^{\vee(k)} \mathrm{Z}_{\mathrm{f}}(L) \tag{10.57}
\end{align*}
$$

Thus if we write $\mathrm{Z}_{\mathrm{f}}{ }^{\vee(k)}(L)=\sum c_{D} \cdot D$ where $D$ represent tangle chord diagrams supported on $L$ and $c_{D}$ are the corresponding coefficients in the expansion of $\mathrm{Z}_{\mathrm{f}}{ }^{\vee}(k)(L)$, then we want
to show $[D]=\left[S_{(K)} D\right]$. This was proved in [LM5]. Thus $\left[\mathrm{Z}_{\mathrm{f}}{ }^{\vee}(k)(L)\right]=\left[\mathrm{Z}_{\mathrm{f}}{ }^{\vee}(k)\left(S_{(K)} L\right)\right]$, and this for any link component $K$.

We now show the invariance under band sum moves. We have to prove that:
or equivalently that:

$$
\begin{equation*}
\left[\otimes^{e} \nu^{k+1} \cdot \widetilde{Z}_{f}(L)\right]=\left[\otimes^{e} \nu^{k+1} \cdot \widetilde{Z}_{f}\left(L^{\prime}\right)\right] \tag{10.59}
\end{equation*}
$$

We follow exactly the same line of reasoning as in [LM5]. We write $\widetilde{Z}_{f}(L)=\sum c_{D} \cdot D$, so that:

$$
\begin{align*}
& \mathrm{Z}_{\mathrm{f}}(k)  \tag{10.60}\\
& \mathrm{Z}_{\mathrm{f}}(L)=\otimes^{e} \nu^{k+1} \cdot \sum c_{D} \cdot D  \tag{10.61}\\
&=\sum c_{D} \otimes^{e} \nu^{k+1} \cdot D  \tag{10.62}\\
&=\sum c_{D^{\prime}, D} D^{\prime} \cdot D
\end{align*}
$$

where the chord diagrams $D^{\prime}$ come from powers of $\nu^{k+1}$ and the coefficients $c_{D^{\prime}, D}$ take account of the expansion of such powers of $\nu$. We have $\widetilde{Z}_{f}\left(L^{\prime}\right)=\sum c_{D} \cdot \Delta D$, so that:

$$
\begin{align*}
& \mathrm{Z}_{\mathrm{f}}(k)  \tag{10.63}\\
&\left.=\sum L^{\prime}\right)  \tag{10.64}\\
&=\otimes^{e} \nu^{k+1} \cdot \sum c_{D} \cdot \Delta D  \tag{10.65}\\
&=\sum c_{D} \otimes^{e} \nu^{k+1} \cdot \Delta D \\
& D^{\prime} \cdot \Delta D
\end{align*}
$$

We want to show that $\left[D^{\prime} \cdot D\right]=\left[D^{\prime} \cdot \Delta D\right]$. If there are less than $2 n$ chord feet on $K_{2}$, then both classes vanish by $L_{<2 n}$. If there are exactly $2 n$ chord feet on $K_{2}$, then all terms but one in $\left[D^{\prime} \cdot \Delta D\right]$ have less than $2 n$ chord feet on $K_{2}$, and therefore die by $L_{<2 n}$. The surviving term is exactly $\left[D^{\prime} \cdot D\right]$. If there are more than $2 n$ chord feet on $K_{2}$, we use Lemma 3.1 of [LM5] that allows us to write a chord diagram with more than $2 n$ chord feet on $K_{2}$ as a linear combination of chord diagrams with less chord feet on it. Applying this lemma repeatedly, we can reduce the number of chord feet on $K_{2}$. The proof involves only the use
of the STU relation and $P_{n+1}$. If we denote by $\cdot \Delta$ the map $D^{\prime} \cdot D \mapsto D^{\prime} \cdot \Delta D$, then we want to show that it commutes with the STU relation and $P_{n+1}$. That this commutes with $P_{n+1}$ is immediate since this relation does not involve the support of chord diagrams. Thus we turn to the commutation with the STU relation. In decreasing the degree of $D^{\prime} \cdot D$, we can decrease the number of chords coming from $D^{\prime}$ or the number of chords from $D$. We start by considering the case where we decrease the number of chords from $D$. If $D^{\prime}$ were trivial, then we would just use the commutative diagram showing the commutativity of $\Delta$ and STU given in [LM5]. Adding $D^{\prime}$ decomposed as $D_{1}^{\prime}$ on one strand and $D_{2}^{\prime}$ on the second to the previous commutative diagram and replacing $\Delta$ by $\cdot \Delta$, we still have commutativity of the resulting diagram. If we now decrease the number of chords in $D^{\prime}$, then we wish to apply the STU relation to chords from $D^{\prime}$, and this commutes with applying the map $\cdot \Delta$ as well. Thus $\left[D^{\prime} \cdot D\right]=\left[D^{\prime} \cdot \Delta D\right]$, and this for any chord diagram $D^{\prime}$ coming from a $k+1$-st power of $\nu$, and for all summands $D$ of $\widetilde{Z}_{f}(L)$. Thus $\left[\otimes^{e} \nu^{k+1} \cdot \widetilde{Z}_{f}(L)\right]=\left[\otimes^{e} \nu^{k+1} \cdot \widetilde{Z}_{f}\left(L^{\prime}\right)\right]$, or equivalently $\left[\mathrm{Z}_{\mathrm{f}}{ }^{\vee(k)}(L)\right]=\left[\mathrm{Z}_{\mathrm{f}}^{\vee(k)}\left(L^{\prime}\right)\right]$.

In order to have an invariant under Kirby move I, one moves the above equivalence class into another quotient space. Before doing so we make a few definitions, all of which once again can be found in the original paper [LM5]. If $X$ is a set with $m$ ordered elements labeled $0,1, \cdots, m-1$, write $\mathcal{A}(m)=\mathcal{A}(X)$. For $\sigma \in S_{m-2}$, define $T_{\sigma}$ to be the graph:


Define:

$$
\begin{equation*}
T_{m}=\sum_{\sigma \in S_{m-2}} \frac{(-1)^{r(\sigma)}}{(m-1) \cdot \frac{(m-2)!}{r(\sigma)!(m-2-r(\sigma))!}} T_{\sigma} \tag{10.67}
\end{equation*}
$$

where $r(\sigma)=\operatorname{Card}\{i \mid \sigma(i)>\sigma(i+1)\}$. For $m \geq 2$, define $T_{m}^{n} \in \mathcal{A}(m)$ by:

$$
\begin{equation*}
T_{m}^{n}=\sum_{\substack{m_{1}+m_{2}+\cdots+m_{n}=m \\ m_{1} \geq m_{2} \geq \cdots \geq m_{n} \geq 2}} \coprod_{1 \leq i \leq n} T_{m_{i}} \tag{10.68}
\end{equation*}
$$

where for each choice of a configuration $\left(m_{1}, \cdots, m_{n}\right)$ the cyclic order of the $T_{m_{i}}$ 's is preserved. If $m<2 m$, we set $T_{m}^{n}=0$. Define:

$$
\begin{equation*}
\iota_{n}: \mathcal{A}\left(\amalg^{e} S^{1}\right) \rightarrow \mathcal{A}(\varnothing) \tag{10.69}
\end{equation*}
$$

by sending a solid circle with $m$ chords ending on it to $T_{m}^{n}$ with those $m$ chords grafted onto the $m$ points defining the support of $T_{m}^{n}$. We now consider $\left[\mathrm{Z}_{\mathrm{f}}{ }^{\vee(k)}(L)\right] \in \mathcal{A}\left(\amalg^{e} S^{1}\right)^{\wedge} / L_{<2 n}, P_{n+1}, O_{n}$, map it into $\mathcal{A}(\varnothing)^{\wedge} / P_{n+1}, O_{n}$, which we further project to $\mathcal{A}(\varnothing) / D_{>n}, P_{n+1}, O_{n}$. The resulting element we write $\left[\iota_{n}\left(\mathrm{Z}_{\mathrm{f}} \vee(k)(L)\right)\right]$. The relation $D_{>n}$ sets to zero chord diagrams of degree greater than $n$. This last quotient space is isomorphic to $\mathcal{A}(\varnothing) / D_{>n}$ by Lemma 3.4 of [LM5].

The following result was originally given in [LM5] in the case $k=1$.
Proposition 10.2.1.2. $\left[\iota_{n}\left(\mathrm{Z}_{\mathrm{f}}{ }^{\vee}(k)(L)\right)\right] \in \mathcal{A}(\phi) / D_{>n}$ is independent of the orientation of $L$ and is invariant under band sum moves.

Proof. This follows from the previous Proposition and by following the proof of the same statement for $\check{Z}_{f}$, Lemma 3.5 of [LM5].

Recall the following definitions originally introduced in [LM5]: $U_{+}$(resp. $U_{-}$) is the trivial knot with +1 (resp. -1 ) framing, and one puts an algebra structure on $\mathcal{A}(\phi) / D_{>n}$ whereby the disjoint union in the argument of the Kontsevich integral results in a product of equivalence classes. One lets $\sigma_{+}$(resp. $\sigma_{-}$) be the number of positive (resp. negative) eigenvalues of the linking matrix for $L$. If $M$ denotes the 3-manifold obtained from a surgery on the link $L$, we would like to define a family of invariants:

$$
\begin{equation*}
\left.\left.\left.\Omega_{k, n}(M)=\left[\iota_{n}\left(\stackrel{\vee(k)}{\mathrm{Z}_{\mathrm{f}}}\left(U_{+}\right)\right)\right]^{-\sigma_{+}} \stackrel{\vee\left(\iota_{n}\right.}{ } \stackrel{\vee(k)}{\mathrm{Z}_{\mathrm{f}}}\left(U_{-}\right)\right)\right]^{-\sigma_{-}}\left[\iota_{n}\left(\mathrm{Z}_{\mathrm{f}}(L) \mathrm{L}\right)\right)\right] \in \mathcal{A}(\varnothing) / D_{>n} \tag{10.70}
\end{equation*}
$$

with $\Omega_{1, n}(M)=\Omega_{n}(M)$ in [LM5]. If $\alpha^{(d)}$ denotes the degree $d$ part of a sum of chord diagrams $\alpha$ as in [LM5], then we define:

$$
\begin{equation*}
\Omega_{k}(M)=1+\sum_{d \geq 1} \Omega_{k, d}(M)^{(d)} \tag{10.71}
\end{equation*}
$$

with $\Omega_{1}(M)=\Omega(M)$ in [LM5]. The rest of this paper is devoted to showing that $\left[\iota_{n}\left(\mathrm{Z}_{\mathrm{f}}{ }^{\vee}(k)\left(U_{ \pm}\right)\right)\right]$are invertible, and we will mainly follow the argument given in [LM5] to emphasize that their proof can be easily generalized to the present case.

One first defines a linear map ([LM5]):

$$
\begin{equation*}
p: \mathcal{A}^{\wedge}(X \amalg X) \rightarrow \mathcal{A}^{\wedge}(X) \otimes \mathcal{A}^{\wedge}(X) \tag{10.72}
\end{equation*}
$$

that sends any chord diagram $D$ with support on $X \amalg X$ with at least one chord stretching between one $X$ and the other copy of $X$ to zero, $\left.\left.D\right|_{X_{L}} \otimes D\right|_{X_{R}}$ otherwise, where $L$ and $R$ refer to left and right copies respectively. If $X=\cup_{1 \leq i \leq e} X_{i}$, then by successively applying the doubling map $\Delta$ on each of the $e$ components of $X$, one defines a map ([LM5]):

$$
\begin{equation*}
\Delta_{\left(X_{1}, \cdots, X_{e}\right)}: \mathcal{A}^{\wedge}(X) \rightarrow \mathcal{A}^{\wedge}(X \amalg X) \tag{10.73}
\end{equation*}
$$

Define ([LM5]):

$$
\begin{equation*}
\hat{\Delta}:=p \circ \Delta_{\left(X_{1}, \cdots, X_{e}\right)} \tag{10.74}
\end{equation*}
$$

In other terms, $\hat{\Delta}$ doubles all components and kills those chord diagrams that have at least one chord stretching between one copy of $X$ and the other. We have the following fact:

$$
\begin{equation*}
\hat{\Delta}\left(\vee_{\mathrm{Z}}^{\vee(k)}(L)\right)=\stackrel{\vee(k)}{\mathrm{Z}_{\mathrm{f}}}(L) \otimes{\stackrel{\vee(k)}{\mathrm{Z}_{\mathrm{f}}}(L)}^{\mathrm{Z}} \tag{10.75}
\end{equation*}
$$

Indeed:

$$
\begin{align*}
\hat{\Delta}\left(\stackrel{\vee(k)}{\left.\mathrm{Z}_{\mathrm{f}}(L)\right)}\right. & =p \circ \Delta_{\left(X_{1}, \cdots, X_{e}\right)}\left(\mathrm{Z}_{\mathrm{f}}^{(k)}(L)\right)  \tag{10.76}\\
& =p \circ\left(\Delta_{\left(X_{1}\right)}\left(\nu^{k+1}\right) \otimes \cdots \otimes \Delta_{\left(X_{e}\right)}\left(\nu^{k+1}\right) \cdot \Delta_{\left(X_{1}, \cdots, X_{e}\right)} \widetilde{Z}_{f}(L)\right)  \tag{10.77}\\
& =p \circ\left((\Delta \nu)^{k+1} \otimes \cdots \otimes(\Delta \nu)^{k+1} \cdot \widetilde{Z}_{f}(\Delta L)\right)  \tag{10.78}\\
& =p \circ(\underbrace{\nu^{k+1} \otimes \cdots \otimes \nu^{k+1}}_{2 e} \cdot \widetilde{Z}_{f}(\Delta L))  \tag{10.79}\\
& =p{ }^{\vee(k)}  \tag{10.80}\\
& \mathrm{Z}_{\mathrm{f}}(\Delta L)  \tag{10.81}\\
& \quad \vee(k) \\
& \mathrm{Z}_{\mathrm{f}}(L) \otimes \mathrm{Z}_{\mathrm{f}}(L)
\end{align*}
$$

The following result was originally given in [LM5] in the case $k=1$.

Lemma 10.2.1.3. $\left[\iota_{n}\left(\mathrm{Z}_{\mathrm{f}}{ }^{\vee}(k)\left(U_{+}\right)\right)\right]$and $\left[\iota_{n}\left(\mathrm{Z}_{\mathrm{f}}{ }^{\vee(k)}\left(U_{-}\right)\right)\right]$are invertible in $\mathcal{A}(\varnothing) / D_{>n}$.
Proof. The proof is the same as in [LM5]. Successive applications of $\hat{\Delta}$ lead to defining:

$$
\begin{align*}
\hat{\Delta}^{(1)} & =\hat{\Delta}  \tag{10.82}\\
\hat{\Delta}^{(2)} & =(\hat{\Delta} \otimes 1) \circ \hat{\Delta}  \tag{10.83}\\
\hat{\Delta}^{(3)} & =(\hat{\Delta} \otimes 1 \otimes 1) \circ \hat{\Delta}^{(2)}  \tag{10.84}\\
\quad & \\
\hat{\Delta}^{(k)} & =(\hat{\Delta} \otimes \underbrace{1 \otimes \cdots \otimes 1}_{k-1}) \circ \hat{\Delta}^{(k-1)} \tag{10.85}
\end{align*}
$$

Using the fact that $\hat{\Delta}\left(\mathrm{Z}_{\mathrm{f}}{ }^{\vee}(k)(L)\right)=\mathrm{Z}_{\mathrm{f}}{ }^{\vee(k)}(L) \otimes \mathrm{Z}_{\mathrm{f}}{ }^{\vee}(k)(L)$, we readily find:

$$
\begin{equation*}
\hat{\Delta}^{(n-1)} \stackrel{\vee(k)}{\mathrm{Z}_{\mathrm{f}}}\left(U_{ \pm}\right)=\left(\stackrel{\vee(k)}{\mathrm{Z}_{\mathrm{f}}}\left(U_{ \pm}\right)\right)^{\otimes n} \tag{10.86}
\end{equation*}
$$

One defines two maps induced by $\hat{\Delta}$. The first is denoted $\hat{\Delta}_{n, m}([\operatorname{LM} 5]):$

$$
\begin{align*}
& \left.\hat{\Delta}_{n, m}: \mathcal{A}\left(\amalg^{e} S^{1}\right) \wedge\right) / L_{<2(n+m)}, P_{n+m+1}, O_{n+m} \\
& \rightarrow \stackrel{\mathcal{A}}{\left.\left(\amalg^{e} S^{1}\right)^{\wedge} / L_{<2 n}, P_{n+1}, O_{n} \otimes \stackrel{\mathcal{A}}{ }\left(\amalg^{e} S^{1}\right)^{\wedge} / L_{<2 m}, P_{m+1}, O_{m},{ }^{2}\right)} \tag{10.87}
\end{align*}
$$

One uses in what follows the less cumbersome notation $P O_{n}$ for $P_{n+1} O_{n}$ as employed in [LM5]. The second map is denoted $\hat{\Delta}_{1, \ldots, 1}^{(l)}$ :

$$
\begin{equation*}
\hat{\Delta}_{1, \cdots, 1}^{(l)}: \mathcal{A}(\varnothing) / D_{>l+1} \rightarrow\left(\mathcal{A}(\varnothing) / D_{>1}\right)^{\otimes l+1} \tag{10.88}
\end{equation*}
$$

Now we invoke Lemma 4.2 of [LM5] which states that the following diagram is commutative (shown for three integers $n=n_{1}+n_{2}$ ):

$$
\begin{array}{cc}
\mathcal{A}\left(\amalg^{e} S^{1}\right) / L_{<2 n}, P O_{n} & \xrightarrow[\iota_{n}]{ } \\
\hat{\Delta}_{n_{1}, n_{2}} \downarrow & \\
 \tag{10.89}\\
\dot{\mathcal{A}}\left(\amalg^{e} S^{1}\right) / L_{<2 n_{1}}, P O_{n_{1}} \otimes \mathcal{A}\left(\amalg^{e} S^{1}\right) / L_{<2 n_{2}}, P O_{n_{2}} \xrightarrow{\iota_{n_{1}} \otimes \iota_{n_{2}}} / P \mathcal{A}(\varnothing) / P O_{n_{1}} \otimes \mathcal{A}(\varnothing) / P O_{n_{2}}
\end{array}
$$

We apply this lemma $n-1$ times. The composition of the $n-1$ maps on the left is:

$$
\begin{equation*}
(\hat{\Delta}_{1,1} \otimes \underbrace{1 \otimes \cdots \otimes 1}_{n-2}) \circ(\hat{\Delta}_{2,1} \otimes \underbrace{1 \otimes \cdots \otimes 1}_{n-3}) \circ \cdots \circ\left(\hat{\Delta}_{n-2,1} \otimes 1\right) \circ \hat{\Delta}_{n-1,1} \tag{10.90}
\end{equation*}
$$

We denote by $\left(\hat{\Delta}^{(n-1)} \mathrm{Z}_{\mathrm{f}}^{\vee(k)}\left(U_{ \pm}\right)\right)_{1, \cdots, 1}$ the image of $\mathrm{Z}_{\mathrm{f}}^{\vee(k)}\left(U_{ \pm}\right)$under this map. The commutativity of the $n-1$ squares yields:

$$
\begin{equation*}
\hat{\Delta}_{1, \cdots, 1}^{(n-1)}\left(\left[\iota_{n} \stackrel{\vee}{ } \mathrm{Z}_{\mathrm{f}}\left(U_{ \pm}\right)\right]\right)=\left[\iota_{1}^{\otimes n}\left(\hat{\Delta}^{(n-1)} \stackrel{\vee(k)}{\mathrm{Z}_{\mathrm{f}}}\left(U_{ \pm}\right)\right)_{1, \cdots, 1}\right] \tag{10.91}
\end{equation*}
$$

Now observe that:

From which it follows that:

$$
\begin{equation*}
\left[\iota_{1}^{\otimes n}\left(\hat{\Delta}^{(n-1)} \stackrel{\vee(k)}{\mathrm{Z}_{\mathrm{f}}}\left(U_{ \pm}\right)\right)_{1, \cdots, 1}\right]=\left[\iota_{1} \mathrm{Z}_{\mathrm{f}}\left(U_{ \pm}\right)\right]^{\otimes n} \tag{10.93}
\end{equation*}
$$

We wish to show that $\left[\iota_{1} \mathrm{Z}_{\mathrm{f}}{ }^{\vee(k)}\left(U_{ \pm}\right)\right]$is non-trivial. This will show that the constant part of $\left[\iota_{n} \mathrm{Z}_{\mathrm{f}} \vee(k)\left(U_{ \pm}\right)\right]$is non-zero, implying that it is invertible. Without loss of generality, we can work with $U_{-}$:

$$
\begin{equation*}
U_{-}= \tag{10.94}
\end{equation*}
$$

We have, using the fact that $Z_{f}(K)=Z_{f}[\mathrm{M}](K)$ for any $M>0$ :

$=\bigcap \times e^{-\Omega / 2} \times \bigcup$
$=\bigcirc-\frac{1}{2} \bigcirc+\cdots$
On the other hand:

$$
\begin{equation*}
\nu=\bigcirc+\text { terms of even order } \tag{10.99}
\end{equation*}
$$

Thus:

$$
\begin{equation*}
\stackrel{\vee}{\vee} \mathrm{Z}_{\mathrm{f}}\left(U_{-}\right)=\nu^{k+1} \cdot \nu^{-1} \cdot \nu^{1} \cdot Z_{f}\left(U_{-}\right)=\bigcirc-\frac{1}{2} \bigcirc+\text { higher order terms } \tag{10.100}
\end{equation*}
$$

In $\mathcal{A}\left(\amalg^{e} S^{1}\right)^{\wedge} / L_{<2}, P_{2}, O_{1}$, the relation $L_{<2}$ gets rid of the first term. Under the mapping $\iota_{1}$, the second term maps to $T_{2}^{1}$ joined to an isolated chord, which is a trivial dashed circle, which according to the relation $O_{1}$ is equivalent to $-2 \cdot 1$. Thus $\left[\iota_{1} \mathrm{Z}_{\mathrm{f}}{ }^{\vee}(k)\left(U_{-}\right)\right]=-1 / 2 \cdot(-2)+$ higher order terms $=1+$ higher order terms, and therefore $\left[\iota_{n} Z_{f}{ }^{\vee}(k)\left(U_{-}\right)\right]$is invertible. One can show that the same result holds for $\left[\iota_{n} \mathrm{Z}_{\mathrm{f}}{ }^{\vee}(k)\left(U_{+}\right)\right]$.

This result shows that $\Omega_{k, n}(M)$ is well-defined. We have the following result originally given in [LM5] in the case $k=1$.

Theorem 10.2.1.4. For $k \in \mathbb{Z}, L$ a framed oriented link, $M$ the 3-manifold obtained from a surgery on $S^{3}$ along $L$, the element $\Omega_{k, n}(M)$ defined by:

$$
\begin{equation*}
\left.\left.\left[\iota_{n}\left(\stackrel{\vee(k)}{\mathrm{Z}_{\mathrm{f}}}\left(U_{+}\right)\right)\right]^{-\sigma_{+}}\left[\iota_{n}\left(\stackrel{\vee(k)}{\mathrm{Z}_{\mathrm{f}}}\left(U_{-}\right)\right)\right]^{-\sigma_{-}} \stackrel{\vee\left(\iota _ { n } \left(\mathrm{Z}_{\mathrm{f}}\right.\right.}{\mathrm{Z}}(L)\right)\right] \in \mathcal{A}(\phi) / D_{>n} \tag{10.101}
\end{equation*}
$$

is a topological invariant of $M$ for any integer $n$.

Proof. $\Omega_{k, n}$ is well-defined by the previous Lemma. Invariance under orientation change and band sum moves follows from Proposition 10.2.1.2. We work out the invariance under the first Kirby move, which consists of adding to or subtracting from a link a circle with $\pm 1$ framing. If one adds a circle with framing +1 , the framing of the component it is added to increases by 1 , and likewise if one adds a circle with framing -1 , the framing of the component it is added to decreases by 1 . Then it suffices to write:

$$
\begin{equation*}
\left.\left[\iota_{n} \stackrel{\vee(k)}{\mathrm{Z}_{\mathrm{f}}}\left(L \amalg U_{+}\right)\right)\right]=\left[\stackrel{\vee(k)}{\left.\left.\iota_{n}\left(\mathrm{Z}_{\mathrm{f}}(L)\right)\right]\left[\iota_{n} \stackrel{\vee(k)}{\mathrm{Z}_{\mathrm{f}}}\left(U_{+}\right)\right)\right]}\right. \tag{10.102}
\end{equation*}
$$

or equivalently:

$$
\begin{equation*}
\left[\stackrel{\vee(k)}{\left.\left.\iota_{n}\left(\mathrm{Z}_{\mathrm{f}}\left(U_{+}\right)\right)\right]^{-1}\left[\iota_{n}\left(\vee_{\mathrm{f}}^{\vee(k)}\left(L \amalg U_{+}\right)\right)\right]=\left[\iota_{n} \stackrel{\vee(k)}{\mathrm{Z}_{\mathrm{f}}}(L)\right)\right] .}\right. \tag{10.103}
\end{equation*}
$$

so that:

$$
\begin{align*}
& \left.\left.\quad \stackrel{\vee(k)}{\left[\iota_{n}\left(\mathrm{Z}_{\mathrm{f}}\left(U_{+}\right)\right)\right]^{-\left(\sigma_{+}+1\right)}\left[\iota _ { n } \left(\mathrm{Z}_{\mathrm{f}}(k)\right.\right.}\left(U_{-}\right)\right)\right]^{-\sigma_{-}}\left[\iota_{n}\left(\mathrm{Z}_{\mathrm{f}}\left(L \amalg U_{+}\right)\right)\right]= \\
& \stackrel{\vee(k)}{\vee} \stackrel{\vee(k)}{ }^{\stackrel{\vee}{ }\left(\iota_{n}\left(\mathrm{Z}_{\mathrm{f}}\left(U_{+}\right)\right)\right]^{-\sigma_{+}}\left[\iota_{n}\left(\mathrm{Z}_{\mathrm{f}}\left(U_{-}\right)\right)\right]^{-\sigma_{-}}\left[\iota_{n}\left(\mathrm{Z}_{\mathrm{f}}(L)\right)\right]} \tag{10.104}
\end{align*}
$$

where on the left hand side of that equation $\sigma_{+}+1=\sigma_{+}\left(L \amalg U_{+}\right)$. We have a similar computation for $L \amalg U_{-}$.

This shows that the following quantity:

$$
\begin{equation*}
\Omega_{k}(M)=1+\sum_{d \geq 1} \Omega_{k, d}(M)^{(d)} \tag{10.105}
\end{equation*}
$$

is well-behaved and defines a family of topological invariants of 3-manifolds $M$, indexed by a natural number $k$.

### 10.2.2 General renormalizations of $\hat{Z}_{f}$ and the invariants they generate

In this section, we consider more general renormalizations of $\widetilde{Z}_{f}$ and study the conditions put on them to yield 3 -manifold invariants. We consider an element $\gamma \in \hat{\mathcal{A}}\left(S^{1}\right)$ whose expansion we write $\sum_{D} c_{D} \cdot D$ for chord diagrams $D$ with support on $S^{1}$ and coefficients $c_{D}$. We consider the general renormalization:

$$
\begin{equation*}
\widetilde{\mathcal{Z}}_{f}[\gamma](L)=\underbrace{\gamma \otimes \cdots \otimes \gamma}_{e} \cdot \widetilde{Z}_{f}(L) \tag{10.106}
\end{equation*}
$$

for $L$ an $e$-component link. We have:

$$
\begin{equation*}
\widetilde{Z}_{f}\left(U_{ \pm}\right)=\nu^{-1} \cdot \nu \cdot Z_{f}\left(U_{ \pm}\right) \tag{10.107}
\end{equation*}
$$

with:

$$
\begin{equation*}
Z_{f}\left(U_{ \pm}\right)=\bigcap \times e^{ \pm \Omega / 2} \times \circlearrowleft \tag{10.108}
\end{equation*}
$$

We use the notation $\theta$ to denote the unique chord diagram of degree one on the circle, and $\theta^{p}$ the solid circle with $p$ parallel chords on it. Then we can write:

$$
\begin{equation*}
\overparen{\Omega} \Omega^{p} \times \circlearrowleft=\theta^{p} \tag{10.109}
\end{equation*}
$$

so that:

$$
\begin{equation*}
\widetilde{Z}_{f}\left(U_{ \pm}\right)=Z_{f}\left(U_{ \pm}\right)=\sum_{p \geq 0} \frac{( \pm 1)^{p}}{p!} \frac{1}{2^{p}} \theta^{p} \tag{10.110}
\end{equation*}
$$

We consider:

$$
\begin{align*}
\widetilde{\mathcal{Z}}_{f}[\gamma]\left(U_{ \pm}\right) & =\gamma \cdot \widetilde{Z}_{f}\left(U_{ \pm}\right)=\gamma \cdot Z_{f}\left(U_{ \pm}\right)  \tag{10.111}\\
& =\sum_{D} c_{D} \cdot D \cdot \sum_{p \geq 0} \frac{( \pm 1)^{p}}{p!} \frac{1}{2^{p}} \theta^{p}  \tag{10.112}\\
& =\sum_{D, p \geq 0} c_{D} \frac{( \pm 1)^{p}}{p!} \frac{1}{2^{p}} D \cdot \theta^{p} \tag{10.113}
\end{align*}
$$

We simplify this expression by letting:

$$
\begin{equation*}
\lambda_{p, \epsilon, D}=c_{D} \frac{\epsilon^{p}}{p!} \frac{1}{2^{p}} \tag{10.114}
\end{equation*}
$$

with $\epsilon= \pm 1$, so that:

$$
\begin{equation*}
\widetilde{\mathcal{Z}}_{f}[\gamma]\left(U_{ \pm}\right)=\sum_{D, p \geq 0} \lambda_{p, \epsilon, D} D \cdot \theta^{p} \tag{10.115}
\end{equation*}
$$

Recall that the chord degree is half the number of vertices, trivalent vertices included. Thus for $D$ a chord diagram, if $|D|$ denotes the chord degree, we decompose that as $|D|^{e}+|D|^{i}$, where $|D|^{e}$ is half the number of external (or univalent) vertices, on solid lines, and $|D|^{i}$ is half the number of internal (or trivalent) vertices. We first regard $\widetilde{\mathcal{Z}}_{f}[\gamma]\left(U_{ \pm}\right)$as an element of $\mathcal{A}\left(S^{1}\right)^{\wedge} / L_{<2 n}, P O_{n}$ that we map by $\iota_{n}$ into $\mathcal{A}(\varnothing)^{\wedge} / P O_{n}$. We write:

$$
\begin{align*}
\iota_{n}\left(\widetilde{\mathcal{Z}}_{f}[\gamma]\left(U_{ \pm}\right)\right) & =\iota_{n}\left(\sum_{D, p \geq 0} \lambda_{p, \epsilon, D} D \cdot \theta^{p}\right)  \tag{10.116}\\
& =\sum_{D, p \geq 0} \lambda_{p, \epsilon, D} \iota_{n}\left(D \cdot \theta^{p}\right)  \tag{10.117}\\
& =\sum_{m \geq 2 n} \sum_{|D|^{e}+p=m / 2} \lambda_{p, \epsilon, D} \iota_{n}\left(D \cdot \theta^{p}\right) \tag{10.118}
\end{align*}
$$

One can see $\iota_{n}\left(D \cdot \theta^{p}\right)$ in the above sum by pulling all the $m$ chords from $D \cdot \theta^{p}$ out of the solid circle, and replacing this latter with $T_{m}^{n}$ where its support of $m$ points is grafted onto the $m$ feet of chords from $D \cdot \theta^{p}$. We write the resulting object $D \cdot \theta^{p}\left(T_{m}^{n}\right)$. We will be interested in the constant part of such objects. Depending on $m, n$, and the choice of $D$ and $p$, and using $P_{n+1}$ in conjunction with $O_{n}$, we can write:

$$
\begin{equation*}
\iota_{n}\left(D \cdot \theta^{p}\right)=D \cdot \theta^{p}\left(T_{m}^{n}\right)=\alpha_{n, m, D}+\text { chord diagrams } \tag{10.119}
\end{equation*}
$$

for some numbers $\alpha_{n, m, D}$, so that:

$$
\begin{align*}
\iota_{n}\left(\widetilde{\mathcal{Z}}_{f}[\gamma]\left(U_{ \pm}\right)\right) & =\sum_{m \geq 2 n} \sum_{|D|^{e}+p=m / 2} \lambda_{p, \epsilon, D} D \cdot \theta^{p}\left(T_{m}^{n}\right)  \tag{10.120}\\
& =\sum_{m \geq 2 n} \sum_{|D|^{e}+p=m / 2} \lambda_{p, \epsilon, D} \alpha_{n, m, D}+\text { chord diagrams of degree } \geq 1 \tag{10.121}
\end{align*}
$$

Once we have this quantity, we project it to $\mathcal{A}\left(S^{1}\right) / D_{>n}, P O_{n} \simeq \mathcal{A}(\varnothing) / D_{>_{n}}$. In this latter space we define:

$$
\begin{equation*}
\Omega_{\gamma, n}(M)=\left[\iota_{n}\left(\widetilde{\mathcal{Z}}_{f}[\gamma]\left(U_{+}\right)\right)\right]^{-\sigma_{+}}\left[\iota_{n}\left(\widetilde{\mathcal{Z}}_{f}[\gamma]\left(U_{-}\right)\right)\right]^{-\sigma_{-}}\left[\iota_{n}\left(\widetilde{\mathcal{Z}}_{f}[\gamma](L)\right)\right] \tag{10.122}
\end{equation*}
$$

and:

$$
\begin{equation*}
\Omega_{\gamma}(M)=1+\sum_{d \geq 1} \Omega_{\gamma, d}(M)^{(d)} \tag{10.123}
\end{equation*}
$$

We have the following result:

Theorem 10.2.2.1. For $\gamma=\sum_{D} c_{D} \cdot D \in \hat{\mathcal{A}}\left(S^{1}\right), L$ a framed oriented link, $M$ the 3manifold obtained from a surgery on $S^{3}$ along $L$, then $\Omega_{\gamma}(M)$ is a topological invariant of $M$ provided the coefficients of $\gamma$ satisfy for all integers $n$ :

$$
\begin{equation*}
\sum_{m \geq 2 n} \sum_{|D|^{e}+p=m / 2} \lambda_{p, \epsilon, D} \alpha_{n, m, D} \neq 0 \tag{10.124}
\end{equation*}
$$

where $\epsilon= \pm 1$ in:

$$
\begin{equation*}
\lambda_{p, \epsilon, D}=c_{D} \frac{\epsilon^{p}}{p!} \frac{1}{2^{p}} \tag{10.125}
\end{equation*}
$$

and:

$$
\begin{equation*}
D \cdot \theta^{p}\left(T_{m}^{n}\right)=\alpha_{n, m, D}+\text { chord diagrams of degree } \geq 1 \tag{10.126}
\end{equation*}
$$

Proof. The proof of Proposition 10.2.1.1 still holds for $\widetilde{\mathcal{Z}}_{f}[\gamma](L) \in \mathcal{A}\left(\amalg^{e} S^{1}\right)^{\wedge} / L_{2 n}, P O_{n}$ and shows that this quantity is independent of orientation change and band sum moves. In order to have $\Omega_{\gamma, n}(M)$ well-defined we need that $\left[\iota_{n}\left(\widetilde{\mathcal{Z}}_{f}[\gamma]\left(U_{ \pm}\right)\right)\right]$be invertible. In the case where $\gamma=\nu^{k+1}$ we used the fact that $\Delta \nu=\nu \otimes \nu$ which we can no longer assume is true for a general $\gamma$. Thus instead of following the same proof as for $\nu^{k+1}$, we show directly that for all $n$, the elements $\left[\iota_{n}\left(\widetilde{\mathcal{Z}}_{f}[\gamma]\left(U_{ \pm}\right)\right)\right]$are invertible. We have computed in (10.121) that:

$$
\begin{equation*}
\iota_{n}\left(\widetilde{\mathcal{Z}}_{f}[\gamma]\left(U_{ \pm}\right)\right)=\sum_{m \geq 2 n} \sum_{|D|^{e}+p=m / 2} \lambda_{p, \epsilon, D} \alpha_{n, m, D}+\text { chord diagrams of degree } \geq 1 \tag{10.127}
\end{equation*}
$$

showing that the constant part of $\iota_{n}\left(\widetilde{\mathcal{Z}}_{f}[\gamma]\left(U_{ \pm}\right)\right)$is:

$$
\begin{equation*}
\sum_{m \geq 2 n} \sum_{2\left(|D|^{e}+p\right)=m} \lambda_{p, \epsilon, D} \alpha_{n, m, D} \tag{10.128}
\end{equation*}
$$

and thus we need this to be non-zero for $\left[\iota_{n}\left(\widetilde{\mathcal{Z}}_{f}[\gamma]\left(U_{ \pm}\right)\right)\right]$to be invertible. The proof of invariance of $\Omega_{\gamma, n}(M)$ under Kirby move I is the same as that in Theorem 10.2.1.4 and
carries over in our case to show that $\Omega_{\gamma, n}$ is invariant under such a move. This makes this quantity a topological invariant for any integer $n$, and thus $\Omega_{\gamma}(M)$ becomes a topological invariant of $M$.

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