

AXIAL DEPENDENCE OF NUCLEAR FUEL MANAGEMENT

by

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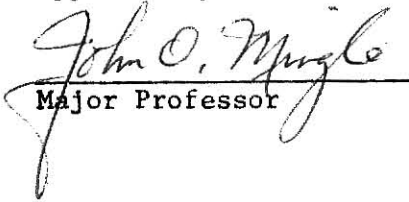
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1.0 INTRODUCTION

Nuclear reactor fuel management encompasses those aspects of fuel composition and loading, related to reactor physics, engineering, or economics, that are relevant to optimal fuel utilization within the design limits of the reactor core. Specifically, in-core management, as a subarea of the field, includes determination of the nuclear properties associated with new and partially irradiated fuel and fuel assemblies; specification of fuel loading, discharging patterns, and internal shuffling schemes; implementation of control procedures during reactor operation which are related and parallel to the objectives specified for fuel use; optimization of fuel and control strategies; and sequencing of decisions for the purpose of fuel procurement.

The interacting requirements of neutronics, materials, control and safety, availability, and economics become quite complex; thus no one simple refueling scheme is possible for all reactors, or even reactor types. However, component parts of the total picture can be studied independently. An adequate representation of the fuel's nuclear properties is available in the current literature^(2,3). It is the areas of fuel shuffling, optimal control, and economical fuel procurement that possibilities of profitable fuel management planning arise.

The objective of fuel management planning is conservation resulting in a reduction in the cost of electrical power generation. Since the costs of the various components of nuclear power generation are interrelated and difficult to predict, a non-cost type of objective, usually directly related to cost, is generally chosen. Maximization of fuel burnup, minimization of

fresh fuel consumption, minimization of local power peaking for both safety and increased thermal efficiency, and maximization of refueling cycle length are all objectives that may be used as bases for fuel management schemes.

Glasstone and Sesonske⁽⁴⁾ present a summary of early fuel management schemes. The earliest, and costliest, of these is batch loading. In batch loading the reactor is operated until it becomes subcritical and then the entire core is removed and replaced. Since the neutron flux profile and therefore burnup of fuel is uneven, the fuel near the extremes of the reactor is underused in this method. Improved burnup may be obtained through the use of a zoned core. The core is divided into radial zones of differing enrichment in an attempt to flatten the power density profile. Again, the entire core is replaced at once. Another method of improving burnup is the "in-out" shuffle. Rather than replace the entire core, a fixed fraction of the core is replaced at each refueling, with the new fuel elements being loaded in the center of the reactor and the retained portion of the old fuel being shuffled radially outward. Thus, the new fuel is placed in the position where it will experience maximum neutron irradiation and therefore maximum burnup. This scheme suffers the disadvantage that the maximum power generation is concentrated in the center of the core, causing heat removal problems and often necessitating lowered power levels. The "out-in" shuffling pattern also achieves higher average discharge burnups than the batch method and still avoids the reactivity concentration inherent in the in-out method. For the out-in method the new fuel is placed near the outer edge of the core and the older fuel is shuffled radially inward, from whence it is eventually removed. While the

out-in method is considerably better than the in-out or batch methods, this procedure can still lead to undesirable neutron flux shapes in the reactor because the reactivity difference between the new and old fuel elements in their respective zones can be fairly large. This problem can be reduced by placing elements of high burnup near those of lower burnup during shuffling. This procedure, called "scatter-loading," or "rondelay," is similar to the out-in method, but the inner regions are intermingled in a checkerboard pattern.

It can be seen that none of the above mentioned fuel management programs assures an optimal use of the fuel. Fuel management can be studied with a number of different techniques aimed at various parts of the optimality problem.

Wall and Fenech⁽⁵⁾ use a modified dynamic programming algorithm to optimize the refueling of a single-enrichment, three-zone, 1000 MW(e) pressurized water reactor, minimizing fuel cost over the life of the reactor. However, only twenty-eight combinations of replacing and reshuffling the fuel are considered since dynamic programming becomes unwieldy if too many variables are used. Wall and Fenech simulate the PWR core response with precalculated performance curves.

Because of the computational difficulties associated with dynamic programming several investigators present attempts to combine it with other techniques. Stover and Sesonske⁽⁶⁾ revise the algorithm of Wall and Fenech, reducing the number of cases which must be considered by eliminating core end states similar to those already tried and determined to be suboptimal. Stover and Sesonske apply this technique to a three-zone boiling water reactor simulated using previously fitted burnup polynomials. Fuel shuffling is again neglected because of the limitations of dynamic programming.

A third type of exhaustive search method is presented by Fagan and Sesonske⁽⁷⁾ using a quasi-equilibrium cycle in which the operating conditions of the reference 1000 MW(e) PWR are nearly identical for each cycle. A fixed fraction of the core is replaced by fuel of constant feed enrichment at each refueling. Logical selection rules guide the placement of fuel elements in a trial-and-error approach. To insure optimality, all combinations must be tried, which leads to a large computational effort. The reactor is modeled with simplified polynomial burnup relations.

Accelerated direct search methods are also in use. Naft and Sesonske⁽⁸⁾ use variational techniques to minimize the radial power peaking factor in the quasi-equilibrium cycle. A simplified semi-analytical equation is used to calculate the radial power distributions. Martinet and Santamarina⁽⁹⁾ describe an improved version of Naft and Sesonske's model. Hoshino⁽¹⁰⁾ presents an heuristic method using weighting factors applied to burnup, power peaking, and effective multiplication factor to determine rules for a one-dimensional four-zone reactor, both with and without fuel shuffling at refueling.

Because linear programming (LP) has advantages in saving computer time and storage over dynamic programming it is used by several investigators. Tabak⁽¹¹⁾ uses a simplified point reactor model with linear programming to determine the time periods between refueling which either maximized plutonium-239 production or minimized uranium-235 consumption. Fulford and Sesonske⁽¹²⁾ use LP to optimize fuel usage for a system of nuclear plants on line together. A number of researchers direct efforts to minimizing fresh fuel consumption^(13,14,15). The model of Suzuki and Kiyose⁽¹³⁾ selects new fuel and shuffles the retained fuel with it into a pattern which

minimizes fuel use. The model uses a reactor with five annular regions in which the fuel elements are identified by their exposure. The important point identified by Suzuki and Kiyose is that "the sequential nature of the refueling decision may not be as strong as generally regarded,"⁽¹⁰⁾ that is, a solution generated by stagewise optimization of fuel shuffling compared to a solution globally optimized over all refueling stages compares very closely. This implies that the simpler, less costly, stagewise optimization procedure is very nearly the best available.

The fact that cycle-by-cycle optimization is acceptable is taken advantage of by many^(17,18,19,20,21,22). Brown⁽¹⁸⁾ uses a stagewise linear programming approach to minimize the fresh fuel consumption of a PWR with three radial zones and a constant flux shape. Brown's model contains constraints on energy removal, neutronic properties, and reloading time restrictions. Chen⁽¹⁹⁾ uses LP to generate a shuffling pattern for a core which minimizes the radial power peaking factor. Chen uses the LP constraints to minimize the deviation of the fuel loading pattern from an ideal distribution described by Winterburg⁽²³⁾. Chen's program has the advantage that the individual fuel element identities are not lost.

Some investigators^(17,22,24) combine the various approaches to fuel management in multi-step techniques, using the most appropriate method for the individual parts of the problem. Melice⁽²⁴⁾ breaks the problem into two parts, first finding the number and enrichment of the new fuel assemblies and then determining the loading pattern. The reloading pattern determination is based on an optimal core k-infinity profile. The optimal profile is approximated through a trial-and-error approach using interactive graphics. This method produces good results but requires a large amount of effort, and the use of interactive graphics may in many cases be inconvenient. Sauar⁽¹⁵⁾

uses a two-step procedure minimizing the fuel cost. The first step uses LP to determine the number of fuel elements assigned to a few region core and the second step uses the interactive graphics method of Melice to fine tune the fuel element distribution. Mingle⁽¹⁷⁾ also uses a two-step program via perturbation theory. First the cycle cost is minimized in a procedure similar to that development by Brown⁽¹⁸⁾, linear programming with a zoned core and discrete burnup groups. Then perturbation theory is used to examine the effects of interchanging individual fuel assemblies in a search to minimize either the power peaking factor or an overall power maldistribution figure-of-merit.

The most sophisticated multi-step method yet proposed is a six phase technique formulated by Motoda, Herczeg, and Sesonske⁽²²⁾. In the first phase, an initial guess is made as to the optimal cycle length and control and refueling requirements for one cycle of a three-zone reactor. The guess is then iterated with a small LP program to optimize the result. For phase two, the phase one result is extrapolated to a five-zone core. Steps one and two could be combined with the added penalty of greatly increased computer time requirements. A larger, more accurate LP program is used in phase three to optimize the five-zone core. The fuel loading and shuffling requirements are first approximated in phase four using the minimum integrated k-deviation technique, and then optimized in phase five with a direct search scheme resembling that developed by Naft⁽⁸⁾. Finally, the burnup, actual cycle length, and power profile are calculated in phase six, allowing the procedure to be repeated for the next cycle.

All of the above mentioned fuel management schemes are based on obtaining the desired objectives (flat flux profile, maximum burnup, minimum cost) by either the one-dimensional (radial) or two-dimensional (X-Y) optimal

placement of fuel. A few researchers^(25,26,27,28) examine the effects of control rod programming on fuel consumption. Motoda⁽²⁵⁾ analyses a two-region, single enrichment BWR for in-out, out-in, and parallel refueling schemes for the effects of control rod programming on maximization of fuel burnup. Notably, Motoda finds that maximization of burnup is equivalent to minimization of fuel enrichment, that the rod control programming is strongly dependent on the refueling scheme used, and that the burnup is influenced much more strongly by reshuffling than by rod programming.

Since present light water reactor fuel elements are constructed so that, at least initially, the fuel properties are constant along their length, most of the previously mentioned fuel management programs are limited to two dimensions at most. According to Silvennoinen⁽¹⁾, "practically no concern has yet been given to the optimization of the axial power and burnup distributions" in light water reactors. Work has been done axially for the Canadian CANDU reactors, but since the CANDU is a heavy water moderated reactor with partial length fuel elements that are continuously replaced in a bidirectional fashion, the CANDU results are not directly applicable to PWR's. Some simplified PWR burnup calculations are described by Benedict⁽²⁹⁾ to demonstrate a two-dimensional code, but the cases described are not truly realistic because they assume a core with continuous out-in refueling in a steady state. No fuel shuffling, axial fuel variation, or optimization is attempted. Although some of the work with control rod programming is directed toward the minimization of axial power peaking, little work exists on the effects of axial fuel variation on different reactor parameters. In this paper the behavior of the fuel characteristics and reactor power profile in the axial direction

of an example PWR is investigated. In addition, the effects of varying the fuel feed enrichment in the axial direction on burnup, power profile, and other important factors are examined.

2.0 THEORY

The study of the in-core section of the fuel cycle requires a familiarity with several branches of nuclear engineering analysis. The optimal use of fuel requires an optimizing technique, such as linear programming, and a method of quantifying the function to be optimized, such as the consumption or burnup of fuel. The burnup, in turn, depends on the local power history of the reactor, which must be either known from operational experience or calculated. A brief discussion of the methods of burnup determination, fuel shuffling optimization, and reactor simulation used in this work is given below.

2.1 Burnup Equations

Nonuniform consumption of fuel in an operating reactor requires the calculation of the burnup separately for each fuel assembly. The value of the atom density, N_i , of each nuclide considered will affect the neutron flux levels in the core. However, the flux levels in turn affect the values of the nuclei density set N_i . Fortunately, the flux varies slowly in a large power reactor and the flux may be considered constant over reasonably long periods of time. Thus the initial flux can be calculated from the initially known values of N_i , and then held constant for some time increment Δt . New values of N_i can be calculated at the end of the time period using the assumed constant flux, the flux levels can be adjusted to correspond to those which would result from the new atom density distribution, and the process repeated for another time increment.

The production and consumption of the various nuclides present during irradiation is modeled with a coupled set of differential equations of the form

$$\frac{dN_i}{dt} = \text{formation rate} - \text{destruction rate} - \text{decay rate} \quad (1)$$

The equations describing the nuclides modeled in this work are simplified with the assumptions that the neutron flux for each lattice point is known and constant for each irradiation time step, that the radioactive decay of the fission products, except xenon and samarium, and the fuel is negligible, and that only one energy group is required for successful modeling. The derivation of the burnup equations can be found in references 2 and 3. For the example PWR which is modeled in this work the atom densities of the nuclides uranium-235, -236, -238, plutonium-239, -240, -241, and a composite fission-produce group are explicitly followed. The important fission products xenon-135 and samarium-151 are calculated independently of the composite group because of their special effects on reactor operation. For the model used, a time increment for which the fluxes are assumed constant on the order of two to four weeks is believed to be adequate.

2.2 Shuffling Patterns

Several parameters must be considered when new fuel is loaded into a core and the old shuffled: how much fuel, of what enrichment, and in what positions to obtain the best core performance. To reduce the number of variables in this investigation, the standard policy of one-third of the core being replaced at each reloading by fuel of constant feed enrichment is followed. The shuffling is done with the objective of minimizing the local power peaking factor by the method developed by Chen⁽¹⁹⁾.

Chen's procedure is useful in designing a loading pattern for which the fuel distribution is as close as possible to an ideal distribution. The ideal distribution is based on a fuel distribution formula developed

by Winterburg⁽²⁷⁾ for a bare cylindrical reactor which results in a constant power density distribution during operation. The differences in fuel distribution between a bare reactor and a large water reflected reactor are small. The ideal fuel distribution has the three dimensional form

$$Z(r,z) = \frac{\Sigma_a(r,z)}{\Sigma_1(r,z)} = \frac{\pi k_{eff} R}{4} \left[\sum_{M=1}^{\infty} \sum_{N=1}^{\infty} \frac{1 - (-1)^M}{M} \frac{1}{\lambda_N J_1(\lambda_N R)} \right. \\ \left. \frac{\eta - \sum_{i=1}^{\infty} [1 + \alpha_i^2 (\lambda_N^2 + \pi^2 M^2 / h^2)]}{\sum_{i=0}^{\infty} [1 + \alpha_i^2 (\lambda_N^2 + \pi^2 M^2 / h^2)]} J_0(\lambda_N r) \cos(\pi M z / h) \right]^{-1} \quad (2)$$

where

Σ_a is the macroscopic absorption cross section of all fissionable material at radius r and height above center z ,

Σ_1 is the macroscopic absorption cross section of all other non-fuel materials at r and z ,

k_{eff} is the effective multiplication factor,

R is the effective radius of the core,

η is the average number of neutrons emitted per neutron absorbed in the fissile isotopes,

α_i are the group diffusion lengths,

λ_N is the N -th zero of the equation $J_0(\lambda_N R) = 0$, and

J_0 , J_1 are Bessel functions of order zero and one, respectively.

The parameter $Z(r,z)$ represents the quality of the fuel at location (r,z) . It is an indication of the power producing ability of the fuel elements. Chen considers this fuel quality factor to be superior to enrichment, infinite multiplication factor, or burnup as a description of the state of the fuel.

The general shape of Equation 2 for a typical pressurized water reactor is given in Figure 1 as a function of radius and height. It is evident that more fissionable material is required near the periphery of the reactor to counterbalance the leakage of thermal neutrons from the core. The relative difference in shape of the radial versus axial curves is explained by the nature of the flux distribution of a uniformly enriched core, which follows a cosine shape axially and a J_0 Bessel function radially. The ideal fuel distribution must counteract these tendencies caused by leakage from the reactor. Chen's program uses linear programming to determine the shuffling pattern which minimizes deviation of the fuel element distribution for the ideal radial pattern.

While the Winterburg formula represents an ideal to which to aspire, it is not a practical solution to reactor fueling. If the local power distribution would remain constant for the life of the core, a modified batch loading method based on the ideal distribution could be used. As it is, the effects of fuel burnup decrease the fuel absorption term and increase the non-fuel absorption term in $Z(r,z)$. Thus the ideal distribution cannot be maintained for even a short period. The best that can be hoped for is a good approximation for an appreciable time.

2.3 Computer Programs

The computer simulation of a reactor power history is a complex problem. The primary program used in this study is an original two-group, two-dimensional diffusion code called BURNRZ. It calculates the local power profile and burnup profile for a cross section of a pressurized water reactor in the radial-axial plane. A supplementary code, BURNXY, is used to model the example reactor in the X-Y cross sectional plane.

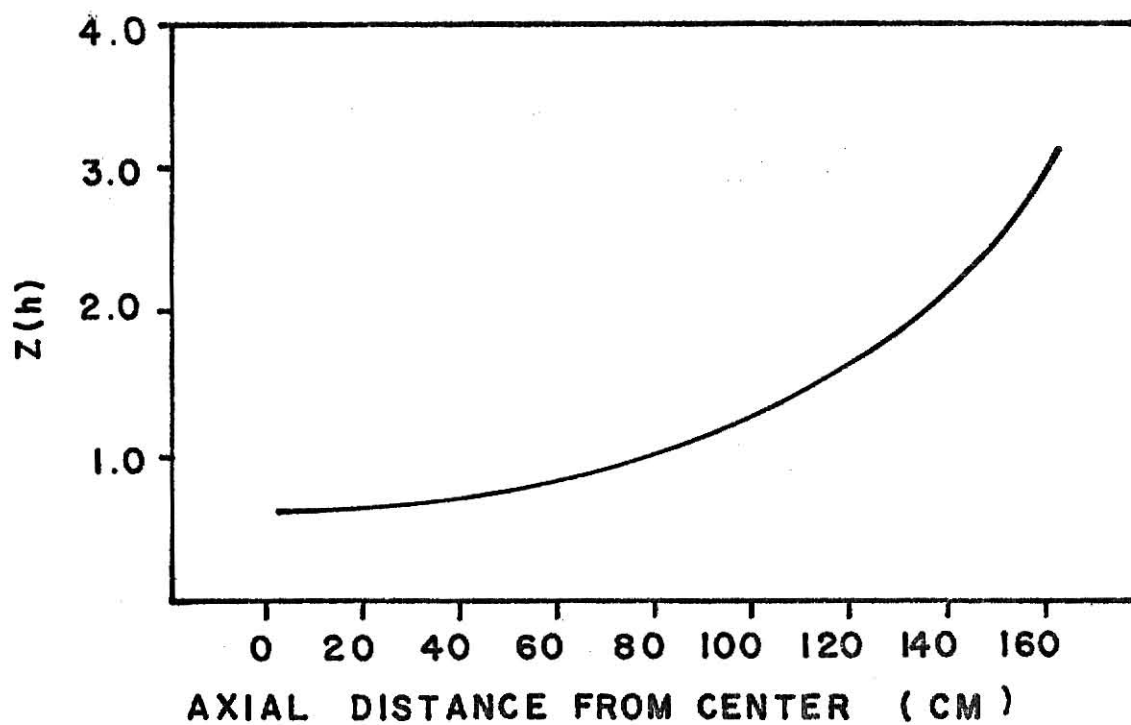
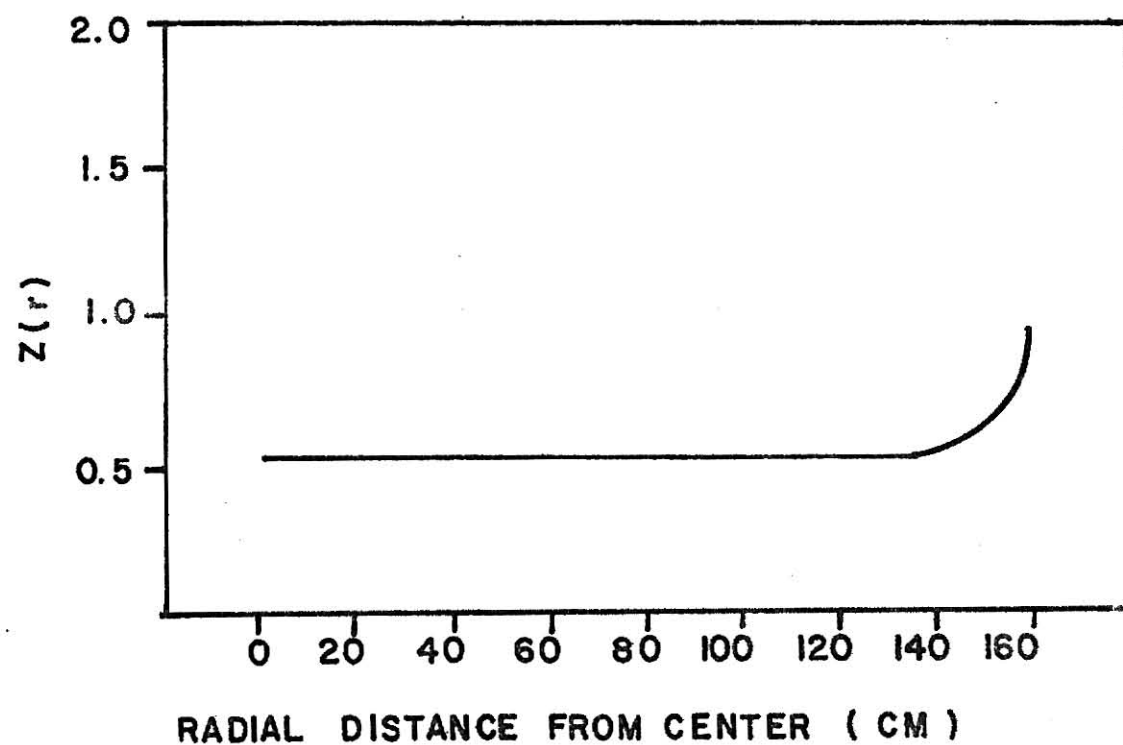


Figure 1. Radial and Axial Behavior of the Ideal Fuel Distribution for Constant Power Density in a Square Cylindrical Unreflected Reactor

The basic flow pattern for both BURNRZ and BURNXY is given in Figure 2. More detailed flow diagrams for the two codes are given in the appendices. The significant difference in the two programs lies in the different geometry of the two systems they solve. BURNXY solves for the power profile of a core with one-quarter symmetry in the X-Y grid where each fuel element is individually represented by a finite-difference mesh point. The quarter core is represented in Figure 3. The unique part of the program BURNRZ is Subroutine XY-TO-R, which converts the information input in the standard X-Y form into an equivalent series of concentric annular rings. The X-Y core pattern is averaged as depicted in Figure 3. The rings are separated by a distance of one fuel element width. Which fuel elements and fractions of fuel elements are present in each annular zone is first determined through an approximate procedure involving straight line segments to represent the arc of each zone boundary through the X-Y grid cells. The resultant concentric rings represent coaxial cylinders in three dimensions. Each of these cylinders is further broken into segments axially, so that the XY-TO-R process results in the configuration shown in Figure 4. The atom densities of the nuclides present in the fuel elements contained within any annular ring are averaged, with the resultant average being used to describe the ring as a whole. The procedure is repeated for each axial division of each cylinder. The net result is a radial-axial, R-Z, distribution of isotopes. Each element of this distribution is used as a mesh point in the flux, power, and burnup calculations.

The neutron flux calculations for BURNXY and BURNRZ are similar. Both are based on numerical solution of the two-energy-group, two-dimensional neutron diffusion equation set up for the appropriate geometry. The solution to the diffusion equation is obtained by use of the central-difference finite

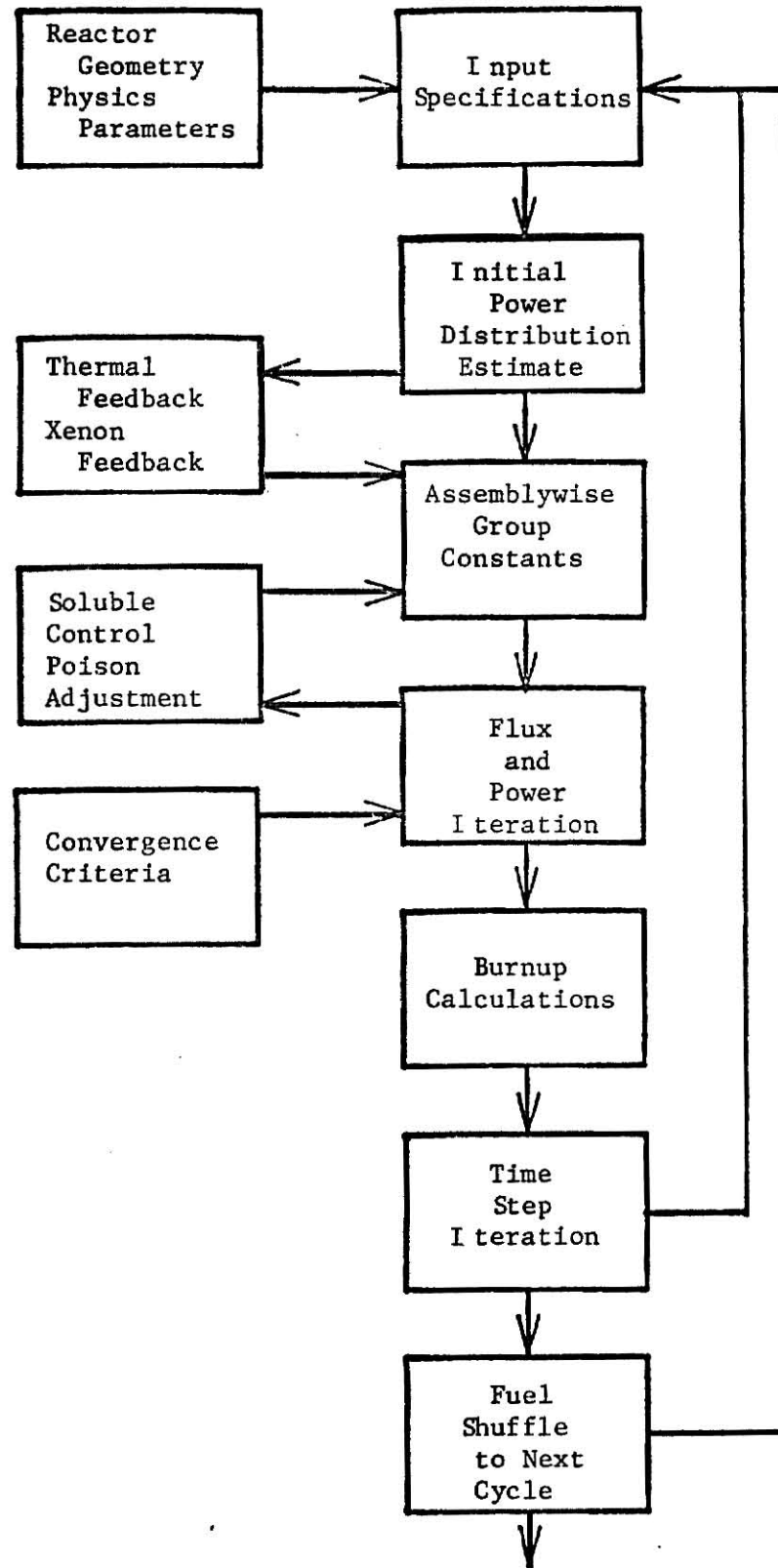


Figure 2. Basic Program Flow for Computer Simulation of Reactor Performance

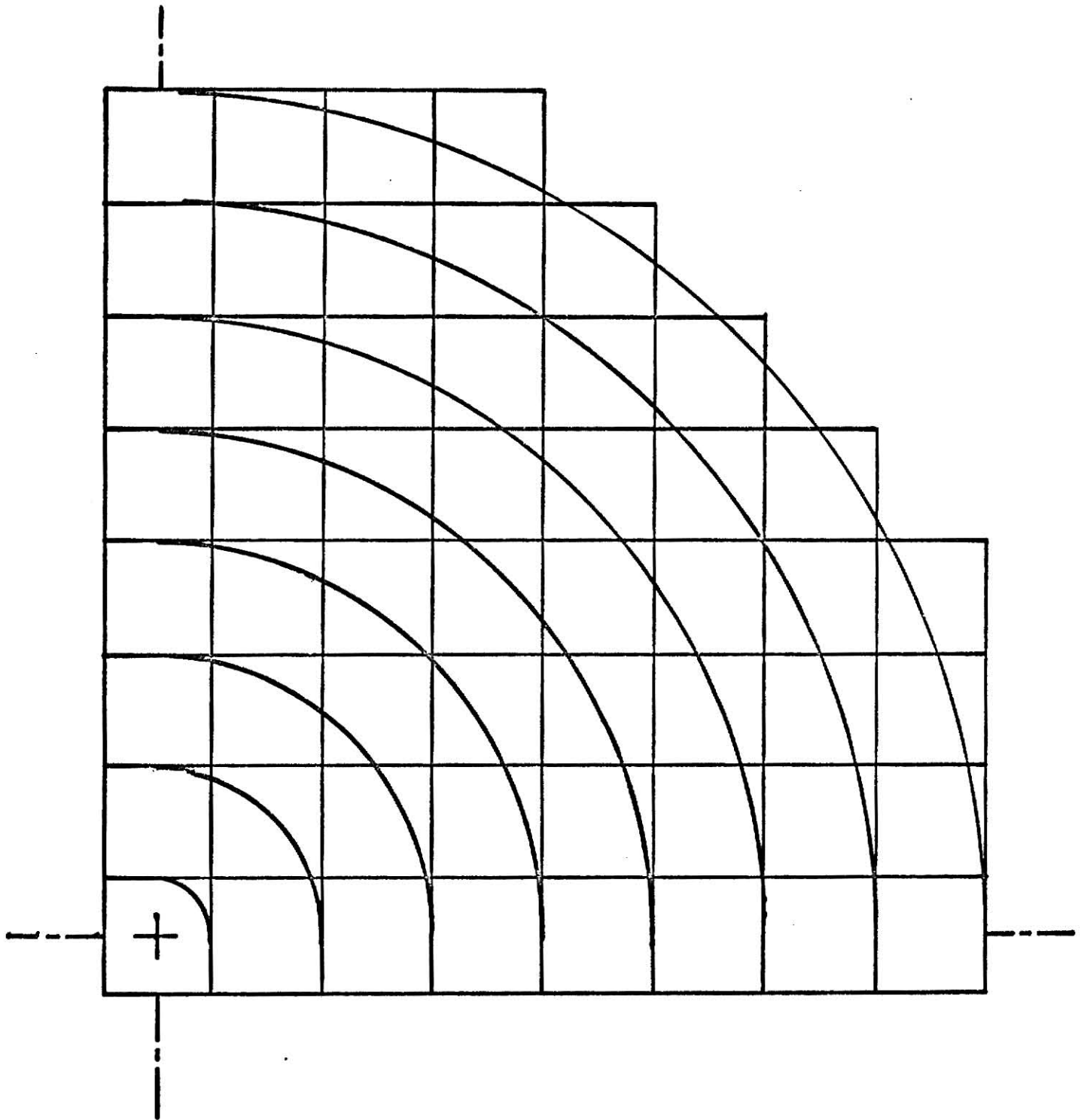


Figure 3. Illustration of Radial Averaging Procedure for Quarter-Symmetric Reactor Core Used in Program BURNRZ

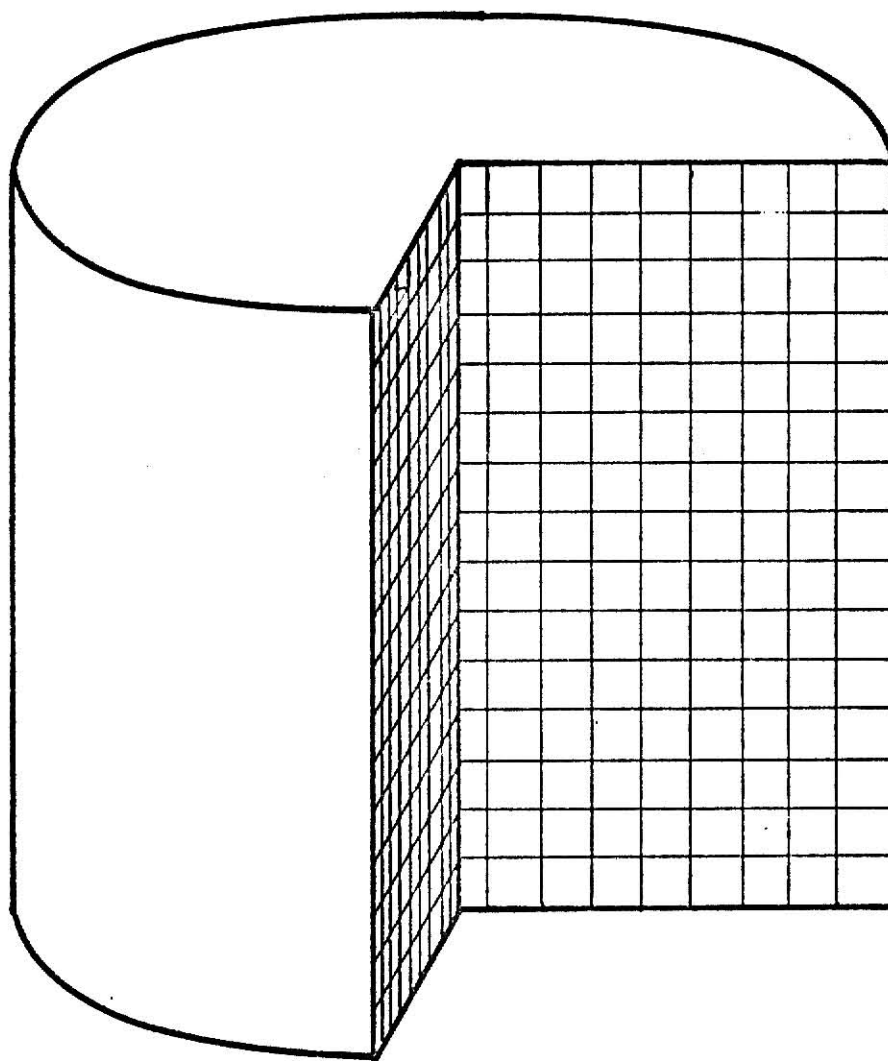


Figure 4. Illustration of Radial-Axial Averaged Fuel Distribution
Used in Program BURNRZ

difference approximation applied to each mesh point. A point successive overrelaxation iterative technique is used to determine the solution of the resultant set of coupled linear equations. For the program BURNRZ, the properties of the mesh points are further complicated by the change in the temperature of the upward flowing cooling water. The changes in the moderating properties of the heated coolant/moderator are accounted for in Subroutine LIQHOH and the material properties of each mesh point are suitably corrected. This procedure is of course unnecessary for the code BURNXY, for which the axially averaged water properties are used. The soluble boron concentration required to maintain criticality may be found if desired, and the flux levels modified to compensate for it.

The known fluxes imply known power densities. A useful parameter is the local relative power, defined as

$$P_i = \text{Local Power Density/Core Average Power Density} \quad (3)$$

The core maximum value of relative power is known as the core power peaking factor. While the power peaking factor is an important parameter in reactor operations, alone it is not particularly descriptive of core power distribution. Mingle⁽¹⁷⁾ defines a core "Figure of Merit," FOM, as

$$\text{FOM} = \left\{ \frac{1}{N} \sum_{i=1}^N |P_i - 1| \right\}^{-1} \quad (4)$$

The Figure of Merit is thus a measure of the deviation of core local power from average. Since the reciprocal of the average deviation is used, a large FOM indicates a flat core power distribution. From an ideal point of view, an infinite Figure of Merit would be desirable.

With the known neutron fluxes the nuclide concentrations for each mesh point are calculated as described in Section 2.1 for a time increment

for which the flux is assumed constant. To save computer time, between major flux calculations the flux is normalized once to compensate for burnup and maintain a constant power output, and atom densities are updated. Every second flux update, then, the new atom densities are used to find new properties for each mesh point. The flux determination and burnup calculations may be repeated as many times as required.

If the BURNRZ simulation is for a cycle other than that corresponding to cycle 1, new fuel, fuel may be shuffled before other calculations are performed. This is done in a procedure that is essentially the reverse of the X-Y to R-Z transformation. The R-Z distribution is returned to an X-Y distribution, the old fuel is subtracted, new added, and the result shuffled in the X-Y mode, and the core is reaveraged to the R-Z mode once again. This is done since all present PWR fuel elements are manufactured as full-core-length units and partial shuffle is physically impossible. Since the individual fuel assembly identity is lost in the initial R-Z averaging the return to X-Y geometry is not exact, and the accuracy of the fuel shuffling is thereby limited. However, it is believed that it is of sufficient accuracy to provide a strong indication of the reactor trends. As a check on the accuracy of the invert-shuffle-convert, the process uses as a comparison the X-Y atom densities of the companion code BURNXY.

The refueling shuffling patterns are derived by the linear programming code LPCHEN, developed by Chen⁽¹⁹⁾ to approximate the ideal fuel distribution described in section 2.4. The code performs the calculations on a model of the core broken into radial and azimuthal zones. There are constraints on radial uniformity, azimuthal uniformity, and local uniformity, as well as a normalization constraint. The input to LPCHEN is from another

of Chen's codes, ZVALUE, which functions identically to BURNXY but which also provides the values of $Z(r,z)$ used in the LP optimization.

Thus four codes are used in this work. The main one is BURNRZ with its companion BURNXY. The shuffling is done with the combination of the somewhat redundant ZVALUE and the LP program LPCHEN.

3.0 THE REFERENCE REACTOR

The example reactor used in this analysis is a typical pressurized water reactor (PWR). It is based on the Westinghouse PWR to be used at Wolf Creek Generating Station near Burlington, Kansas. Pertinent data are taken from the Westinghouse RESAR-3⁽³⁰⁾.

The reactor is designed to produce 3411 MW thermal and to operate at a nominal pressure of 2250 psia and temperature averaging 590 degrees Fahrenheit. It contains 193 fuel assemblies each composed of a 17 x 17 array of fuel pins. The core is twelve feet high and approximately eleven feet in diameter. The standard cycle one core contains 111 tons of uranium dioxide fuel in three enrichments, 3.1%, 2.6%, and 2.1%. The initial enrichment distribution is as shown in Figure 5. Some other pertinent information is given in Table 1.

Table 1. Reference Reactor Parameters

Average Operating Temperatures (Degrees Fahrenheit)	
Inlet	557.3
Outlet	618.9
Overall average	590.8
Coolant Flow Rate (Pounds/hr)	1.358×10^9
Physics Parameters	
Fast Fission Factor	1.0055
Fission to Resonance Non-leakage Probability	0.995
Resonance Escape Probability	0.7088
ν -235	2.07
ν -239	1.87
Fermi Age (cm^2)	36.0
Diffusion Length (cm^2)	21.4
Total Buckling (cm^{-2})	2.77×10^{-4}
Mesh Size (cm)	21.4
Average Plant Capacity Factor	0.75

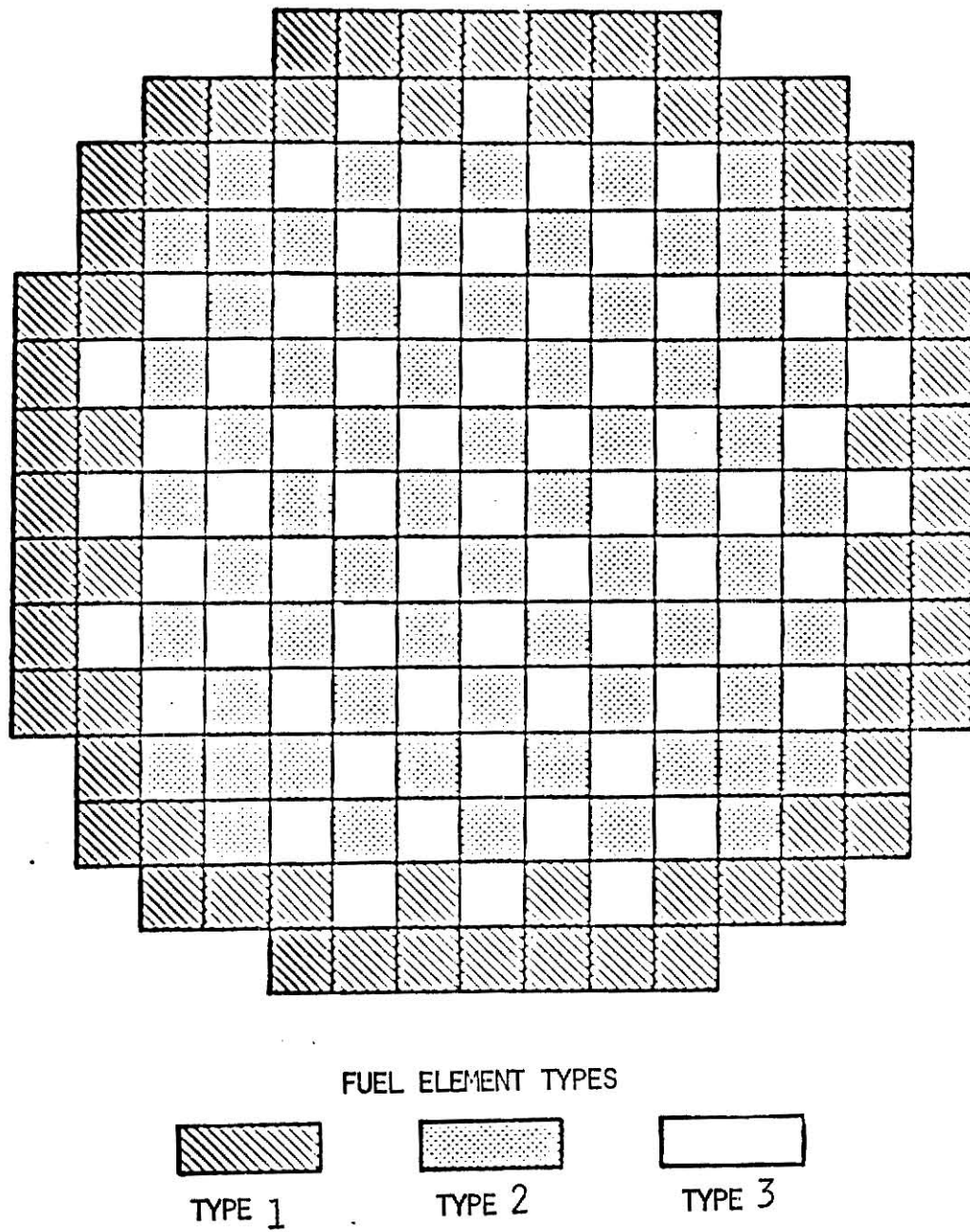


Figure 5. Initial Fuel Element Distribution for Example PWR

4.0 RESULTS

Several cases are investigated. A reference standard: a series of three consecutive burnup cycles of length one year, ten months, and one year, respectively, is calculated for a reference reactor core using three types of fuel elements of axially constant fuel enrichment. This reference case closely follows procedures presently used for PWR's. Three experimental cases, the fuel elements of which have axial enrichment variations, are then presented. Each experimental case is a different attempt to approximate the Winterburg ideal fuel distribution described in Section 2.2. The fuel assembly axial enrichment distributions are summarized in Table 2. The loading patterns for the first cycle of each case examined are the same as shown previously in Figure 5. The new fuel elements inserted in the cores at the beginning of cycles 2 and 3 are always of the type 1 element shown in Table 2 for the particular case. So that fuel burnup among the four separate cases investigated may best be compared, all are calculated for the same burnup period lengths as the reference case.

The code BURNRZ has provisions for the calculation of the equilibrium xenon concentration. The xenon concentration is considered to be directly proportional to the neutron flux levels. Unfortunately, this procedure for determination of xenon concentration is only approximate. The procedure leads to large calculated flux oscillations from top to bottom in the core with time. It is believed that the xenon calculation procedure does not adequately model the xenon oscillations. Therefore, the equilibrium xenon is not calculated in this work. Since the xenon concentration is proportional to the power level, it tends to depress the local power in areas of high reactivity below where it would normally be. The critical concentration of soluble control boron is also not calculated in this work.

TABLE 2

Axial Fuel Enrichment Distribution in Fuel Assemblies, Each Number Represents the U-235 Fraction in a 9.6 inch Segment of a Fuel Assembly.

Type 1	<u>Reference</u> Type 2	Type 3	Type 1	<u>Case 1</u> Type 2	Type 3
0.031	0.026	0.021	0.031	0.031	0.031
0.031	0.026	0.021	0.031	0.026	0.021
0.031	0.026	0.021	0.031	0.026	0.021
0.031	0.026	0.021	0.031	0.026	0.021
0.031	0.026	0.021	0.031	0.026	0.021
0.031	0.026	0.021	0.026	0.026	0.021
0.031	0.026	0.021	0.026	0.026	0.021
0.031	0.026	0.021	0.026	0.026	0.021
0.031	0.026	0.021	0.026	0.026	0.021
0.031	0.026	0.021	0.026	0.026	0.021
0.031	0.026	0.021	0.026	0.026	0.021
0.031	0.026	0.021	0.031	0.026	0.021
0.031	0.026	0.021	0.031	0.026	0.021
0.031	0.026	0.021	0.031	0.026	0.021
0.031	0.026	0.021	0.031	0.026	0.021
0.031	0.026	0.021	0.031	0.031	0.031
Type 1	<u>Case 2</u> Type 2	Type 3	Type 1	<u>Case 3</u> Type 2	Type 3
0.032	0.032	0.032	0.032	0.032	0.032
0.032	0.026	0.021	0.031	0.026	0.021
0.032	0.026	0.021	0.031	0.026	0.021
0.032	0.026	0.021	0.031	0.026	0.021
0.032	0.026	0.021	0.031	0.026	0.021
0.032	0.026	0.021	0.031	0.026	0.021
0.032	0.026	0.021	0.031	0.026	0.021
0.032	0.026	0.021	0.031	0.026	0.021
0.032	0.026	0.021	0.031	0.026	0.021
0.032	0.026	0.021	0.031	0.026	0.021
0.032	0.026	0.021	0.031	0.026	0.021
0.032	0.026	0.021	0.031	0.026	0.021
0.032	0.026	0.021	0.031	0.026	0.021
0.032	0.032	0.032	0.032	0.032	0.032

since it requires excessive computational effort. The boron tends to increase the power peaking factor, and the xenon to decrease it, so by ignoring both factors it is believed that the ultimate result is not greatly changed. Also, none of the X-Y programs has a xenon correction, so direct comparison is facilitated.

Figures 6, 7, 8, 9, 10, and 11 are example X-Y and R-Z power profiles put in a form for ready comprehension. Figures 12 and 13 are X-Y power profile data. Tables 3, 4, and 5 are X-Y discharge burnup data. Figures 14 and 15 are R-Z power profile data, and Figures 16, 17, 18, 19, 20, 21, and 22 are R-Z burnup data.

4.1 Discussion of Reference Case

The reference case models the performance of pressurized water reactors as they are presently constructed. Three types of fuel elements of constant axial enrichment of 3.1%, 2.6%, and 2.1% are used. For the initial burnup cycle, the power density of the core is highest in the central part of the core at Beginning-of-Life (BOL) and gradually shifts outward over time both radially and axially, until at the core End-of-Life (EOL) the preponderance of the power output comes from the upper and lower outer edges of the core. The power density shift follows from a thermal neutron flux density shift as the inner regions of the core burn. This basic pattern is followed from BOL to EOL of every case studied, but is most pronounced in the reference case. To illustrate the nature of the power production shift, Figure 6 contains example data on BOL and EOL local relative power for the symmetric quarter core in the X-Y plane as it is provided by the program BURNXY. To aid comprehension of these data, Figure 7 presents a graphical representation of the BOL numbers for the

0.922	0.931	0.782	0.587				
1.018	1.037	0.953	0.707				
0.958	1.300	0.874	1.082	0.854	0.588		
0.931	1.356	0.930	1.295	1.031	0.721		
1.157	0.993	1.099	0.908	1.017	1.037	0.589	
1.101	0.927	1.120	0.949	1.154	1.278	0.721	
0.995	1.160	0.966	1.101	0.909	1.019	0.857	
0.888	1.077	0.903	1.112	0.942	1.154	1.031	
1.173	0.998	1.154	0.967	1.104	0.912	1.089	0.583
1.053	0.876	1.067	0.899	1.112	0.950	1.296	0.707
1.009	1.179	0.999	1.157	0.971	1.107	0.881	0.789
0.865	1.046	0.873	1.067	0.903	1.120	0.930	0.953
1.188	1.013	1.182	1.003	1.169	1.002	1.312	0.941
1.039	0.862	1.046	0.876	1.077	0.927	1.365	1.038
1.016	1.190	1.013	1.181	1.003	1.169	0.967	0.931
0.860	1.039	0.865	1.053	0.888	1.101	0.931	1.018

0.922	BOL Local Power Density
1.018	EOL Local Power Density

Figure 6. Example Cycle 1 Local Relative Power Density Data from Program BURNXY

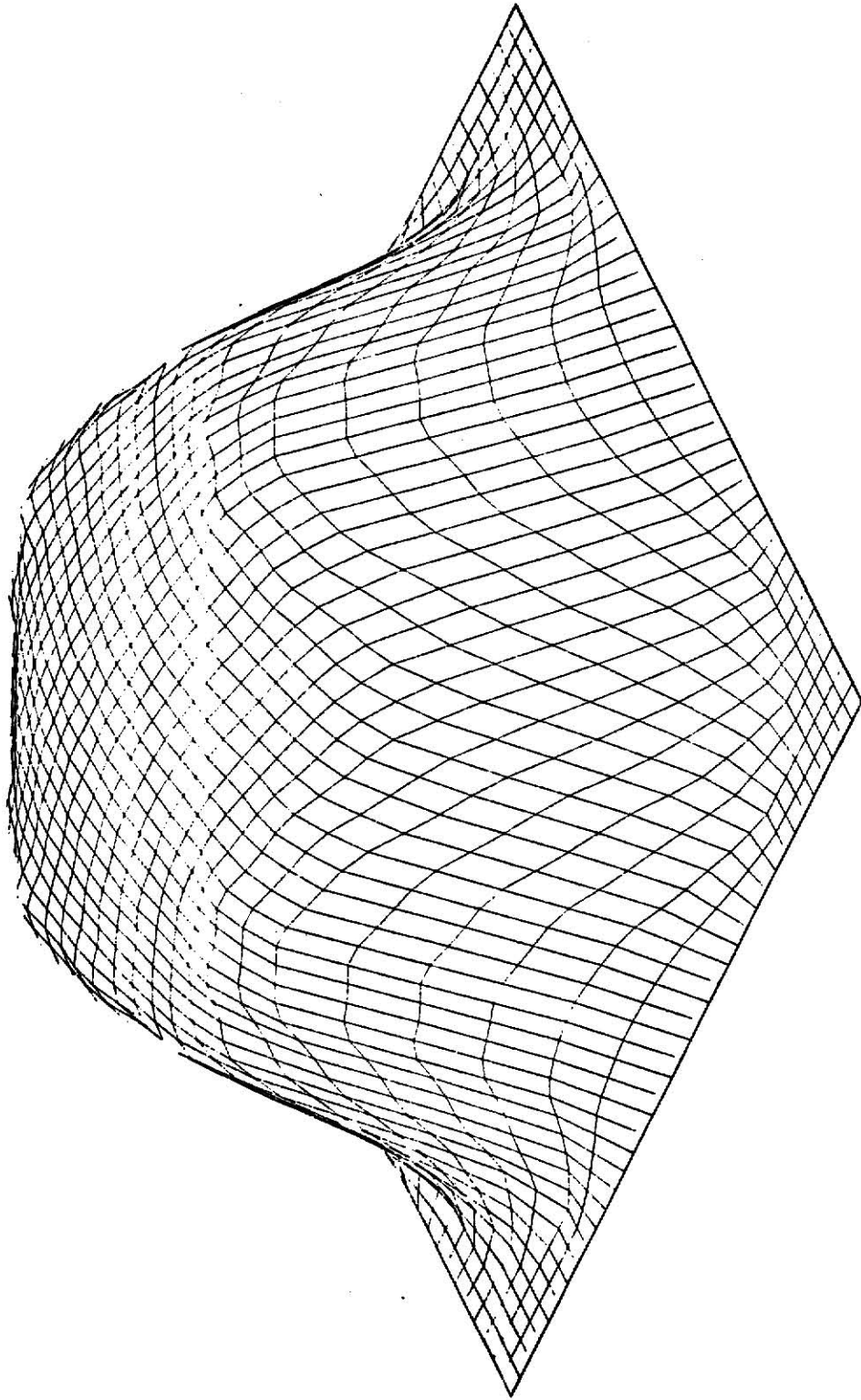


Figure 7. Example Cycle 1 BOL X-Y Local Relative Power Density Profile Map. The Distance Above the Base Plane is Proportional to Power Density.

entire core. The height above the base plane is proportional to the value of relative power at that position. Figure 8 is a similar representation for the EOL case. The shift in power density from the center to the edges of the core is evident. Radially, the cycle 1 reference case power peaking factor, the maximum value of relative power, declines with increased burnup and the position of maximum local power density moves from near the core center line toward the outer edge. The axial peaking factor behaves in a similar manner, but the area of maximum power generation splits into two parts which migrate slowly to the reactor extremes. Figure 9 is similar to Figure 6 but contains reference data for one-half of the reactor R-Z cross sectional plane. Figure 10 is a graphical representation of the BOL local power values. The curves enclose areas in which the local power is equal to or greater than that indicated on the curve. Figure 11 represents the R-Z plane at cycle 1 EOL. The decline in the maximum value of power peaking due to the diffusion of power production to the edges of the core results in a higher reactor Figure of Merit of EOL than at BOL. The total core burnup at EOL is fairly even, concentrated somewhat in the inner two-thirds of the core. This is to be expected, since the neutron leakage is negligible in the inner core regions. It is also desirable since the lower enrichment fuel scattered in the inner regions is generally that which is first discharged at refueling.

The second cycle of the reference case is shorter than the first, since the burnup of the first cycle is so even. This causes the remaining shuffled fuel to be less neutronically potent and the cycle 2 BOL core effective multiplication factor to be lower. The local power peaking factor at BOL is less than that for the cycle 1 BOL, and actually goes lower than at any time in cycle 1 about halfway through the second cycle,

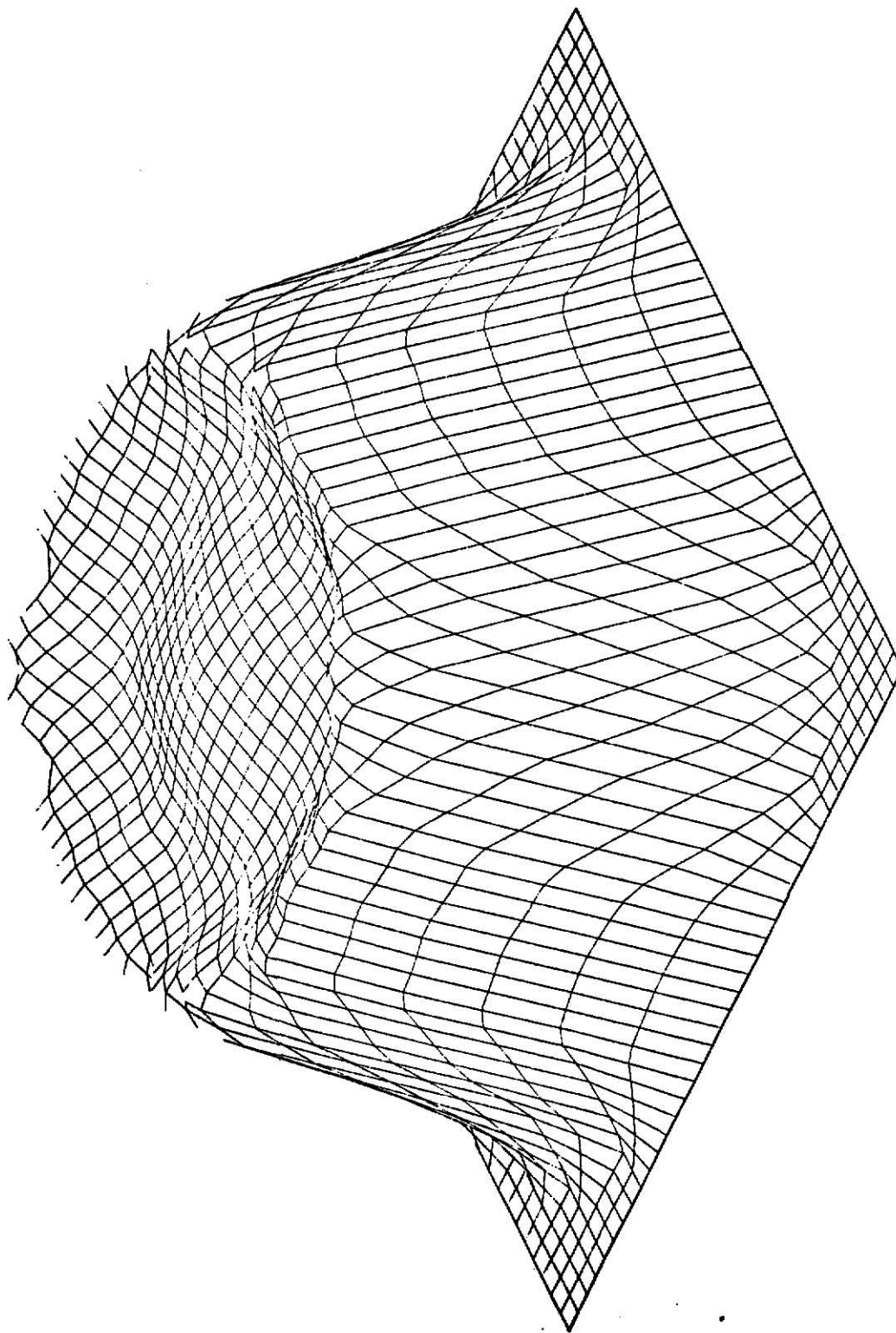


Figure 8. Example Cycle 1 EOL X-Y Local Power Density Profile Map. The Distance Above the Base Plane is Proportional to Power Density.

0.419 0.664	0.472 0.774	0.441 0.738	0.435 0.774	0.407 0.778	0.386 0.806	0.338 0.767	0.270 0.648	0.087 0.175
0.764 0.988	0.863 1.164	0.804 1.103	0.795 1.165	0.744 1.177	0.706 1.240	0.619 1.201	0.496 1.028	0.160 0.254
1.071 1.016	1.213 1.203	1.129 1.137	1.118 1.206	1.048 1.226	0.997 1.313	0.875 1.299	0.702 1.139	0.224 0.267
1.323 0.938	1.507 1.114	1.400 1.051	1.389 1.119	1.303 1.143	1.244 1.239	1.094 1.246	0.880 1.115	0.287 0.254
1.511 0.862	1.732 1.026	1.606 0.967	1.597 1.031	1.500 1.056	1.438 1.153	1.267 1.172	1.022 1.062	0.319 0.238
1.627 0.816	1.879 0.973	1.738 0.915	1.733 0.975	1.630 0.999	1.570 1.095	1.388 1.120	1.123 1.022	0.346 0.226
1.671 0.795	1.945 0.950	1.794 0.891	1.794 0.949	1.691 0.971	1.637 1.065	1.453 1.092	1.180 1.002	0.358 0.219
1.646 0.792	1.930 0.974	1.775 0.886	1.782 0.942	1.683 0.962	1.638 1.055	1.459 1.083	1.190 0.998	0.355 0.216
1.558 0.801	1.841 0.958	1.690 0.895	1.701 0.949	1.610 0.965	1.576 1.058	1.410 1.088	1.155 1.006	0.340 0.215
1.419 0.821	1.688 0.983	1.545 0.915	1.561 0.969	1.481 0.983	1.457 1.076	1.309 1.107	1.078 1.028	0.312 0.217
1.237 0.861	1.481 1.032	1.354 0.958	1.372 1.012	1.304 1.023	1.290 1.117	1.164 1.149	0.938 1.071	0.275 0.223
1.026 0.926	1.234 1.112	1.127 1.029	1.144 1.086	1.090 1.094	1.083 1.190	0.982 1.221	0.815 1.138	0.230 0.234
0.792 1.001	0.958 1.207	0.873 1.115	0.889 1.175	0.848 1.180	0.847 1.278	0.770 1.302	0.642 1.204	0.180 0.246
0.546 0.977	0.662 1.211	0.603 1.117	0.615 1.178	0.588 1.180	0.588 1.273	0.536 1.280	0.448 1.161	0.125 0.241
0.293 0.705	0.356 0.860	0.324 0.794	0.311 0.838	0.316 0.839	0.317 0.901	0.290 0.894	0.242 0.794	0.067 0.175

0.419 0.664

BOL Local Power Density
EOL Local Power Density

Figure 9. Example Cycle 1 Local Relative Power Density Data from Program BURNRZ

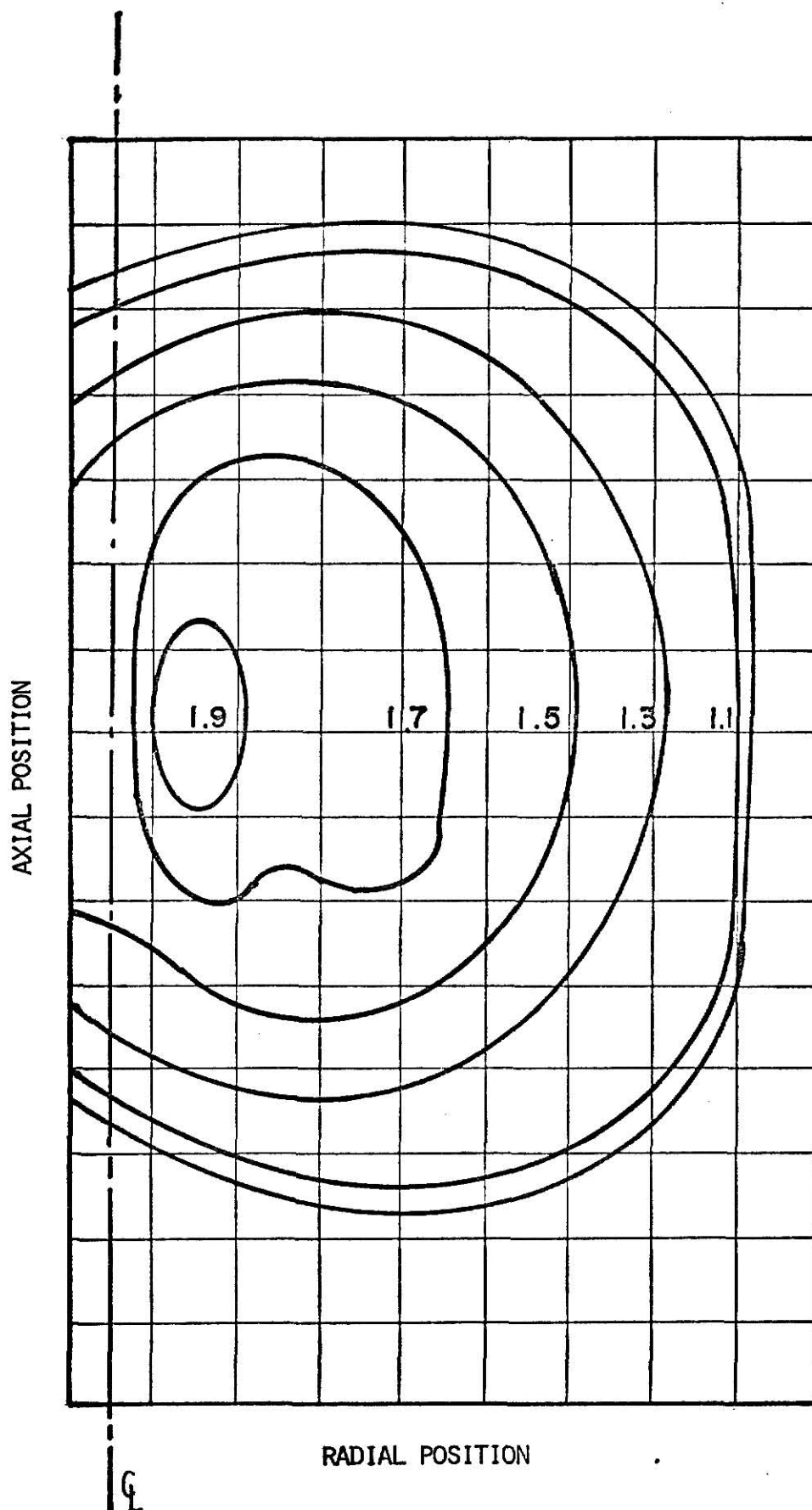


Figure 10. Example Cycle 1 BOL R-Z Local Relative Power Density Profile Map. The Curves Enclose Volumes in which the Relative Power is Greater than Indicated on the Curve.

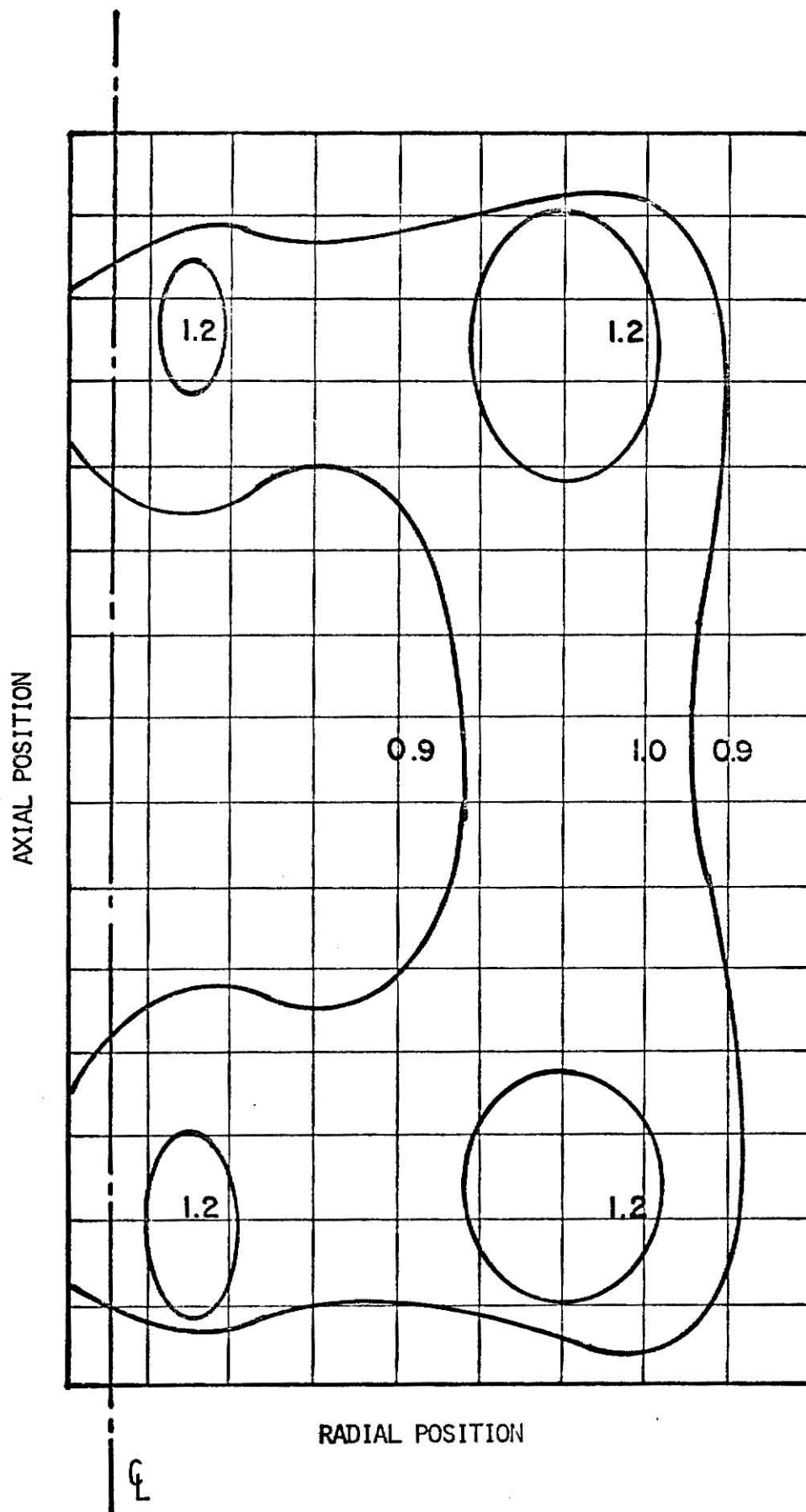


Figure 11. Example Cycle 1 EOL R-Z Local Relative Power Density Profile Map. The Curves Enclose Volumes in which the Relative Power is Greater than Indicated on Curve.

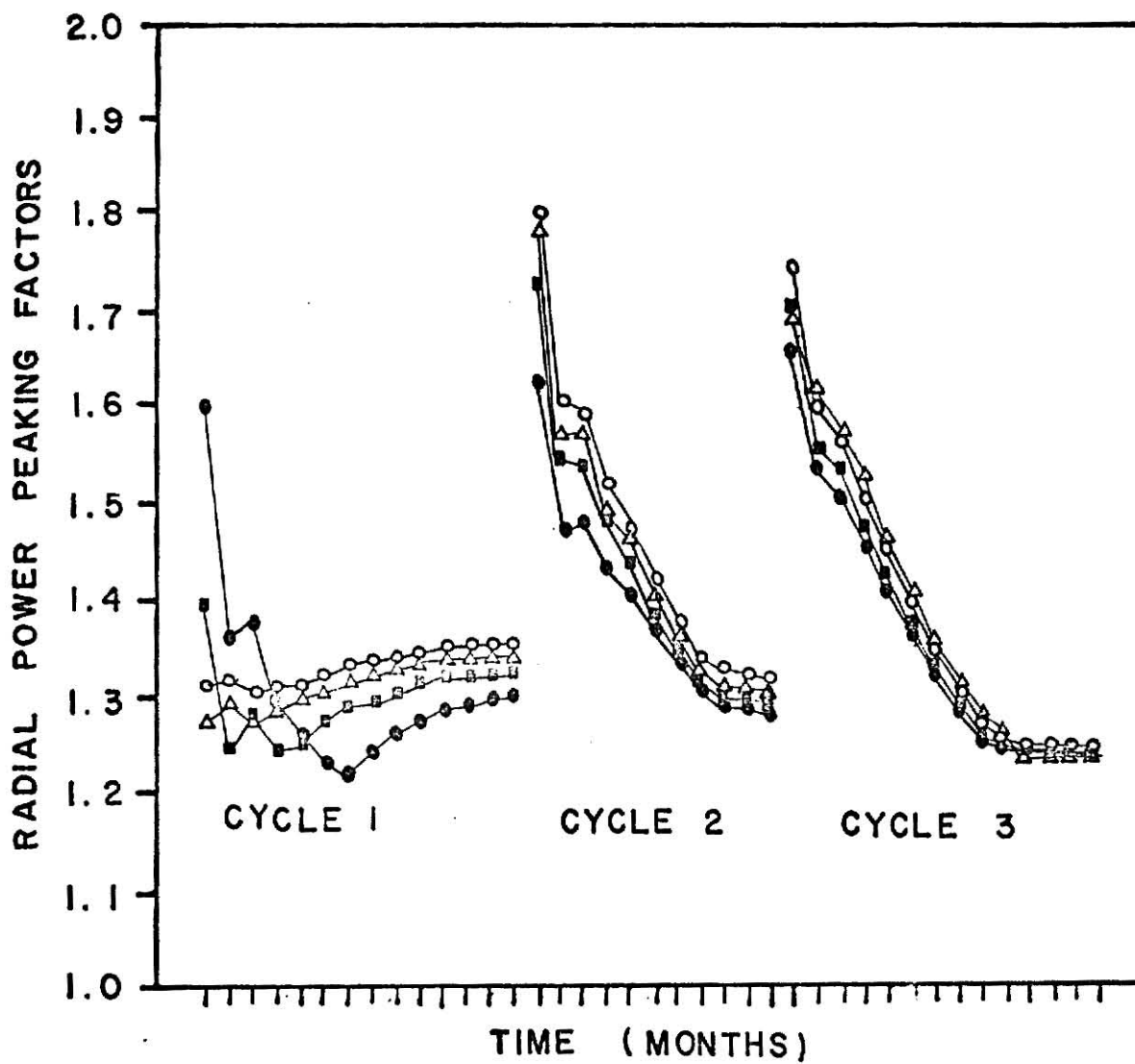
but it then increases again near EOL as the flux maxima move to the reactor core extremes seeking to escape the highly burned up central regions.

Radially the power profile begins and ends much higher at the outer edge of the core. This implies that much of the cycle two power production is occurring in the new fuel.

Figures 12 and 14 are depictions of the power peaking factor behavior with time for the X-Y and R-Z conditions, respectively. Figures 13 and 15 display the actions of the core Figures of Merit respectively for X-Y and R-Z cases. On all graphs the open circle represents the data for the reference case.

The third cycle of the reference case lasts as long as does the first. The last of the low enrichment fuel elements are shuffled out between cycles two and three and the whole core is made up of fuel that, at least originally, is of 3.1% enrichment. This fact combines with the short life of cycle two, thus inflicting less burnup on the elements retained from that core, implying that the cycle three core should last longer than that of the second cycle. As a result of this higher reactivity, the BOL flux and power profiles are rather highly peaked towards the center of the core. The power density quickly becomes more uniform as the center burns out and the flux maxima glide out toward the reactor edges. Finally, as this process proceeds, the maxima begin again to increase slightly as the power generation concentrates at the outer edges.

In all three cycles of the reference case the core burnup is quite uneven axially. In cycle one, the center parts of some fuel elements suffer as much as 143% more burnup than the extreme end sections. All the elements are at least twice as burned near the center as near the edges. Cycles two and three are not quite as severe, but still average 80% more and 60% more, respectively, in the center than the edges.



- REFERENCE
- CASE 1
- △ CASE 2
- CASE 3

Figure 12. Behavior of Radial Power Peaking Factors with Increased Burnup

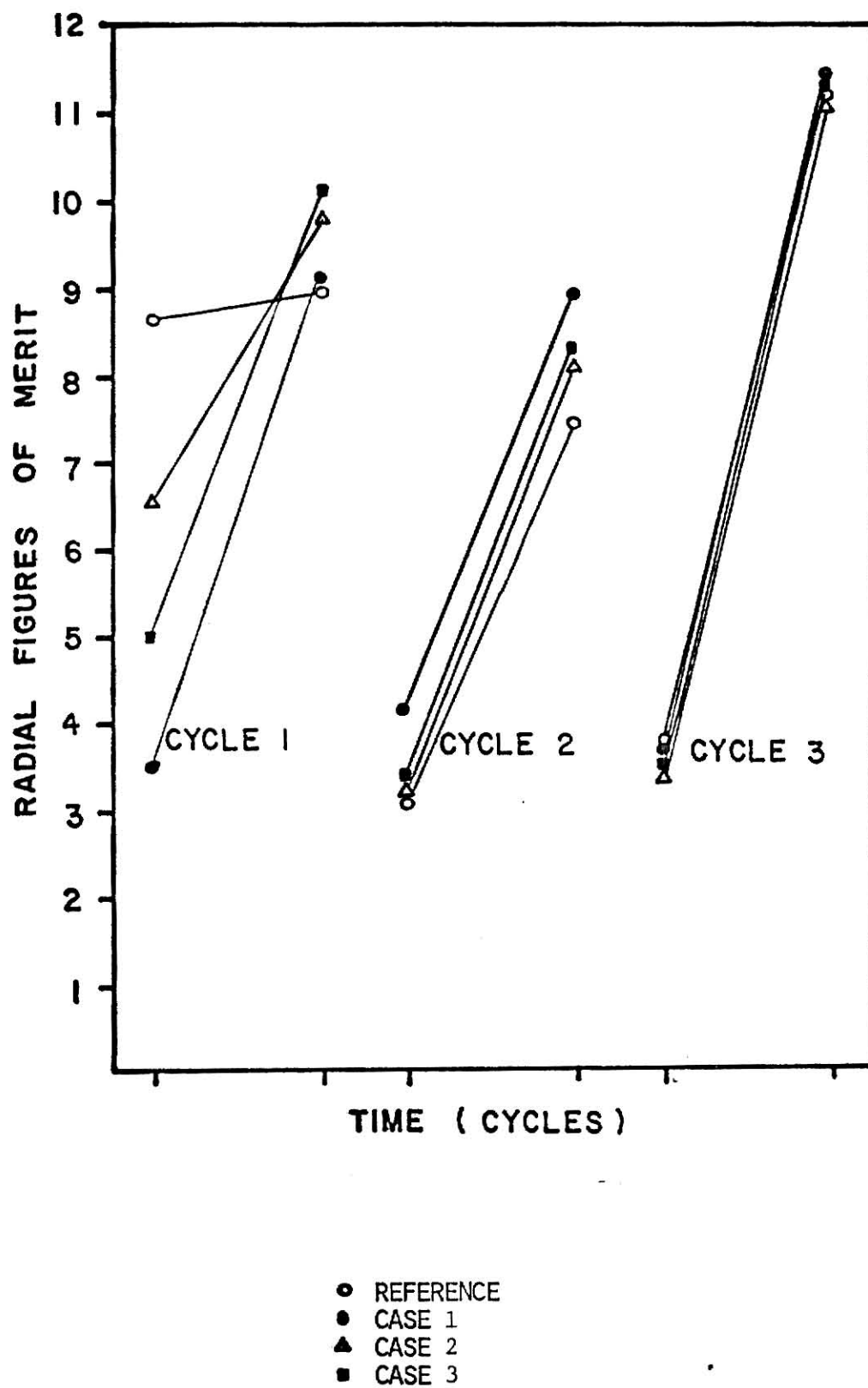


Figure 13. Behavior of Radial Figures of Merit with Increased Burnup

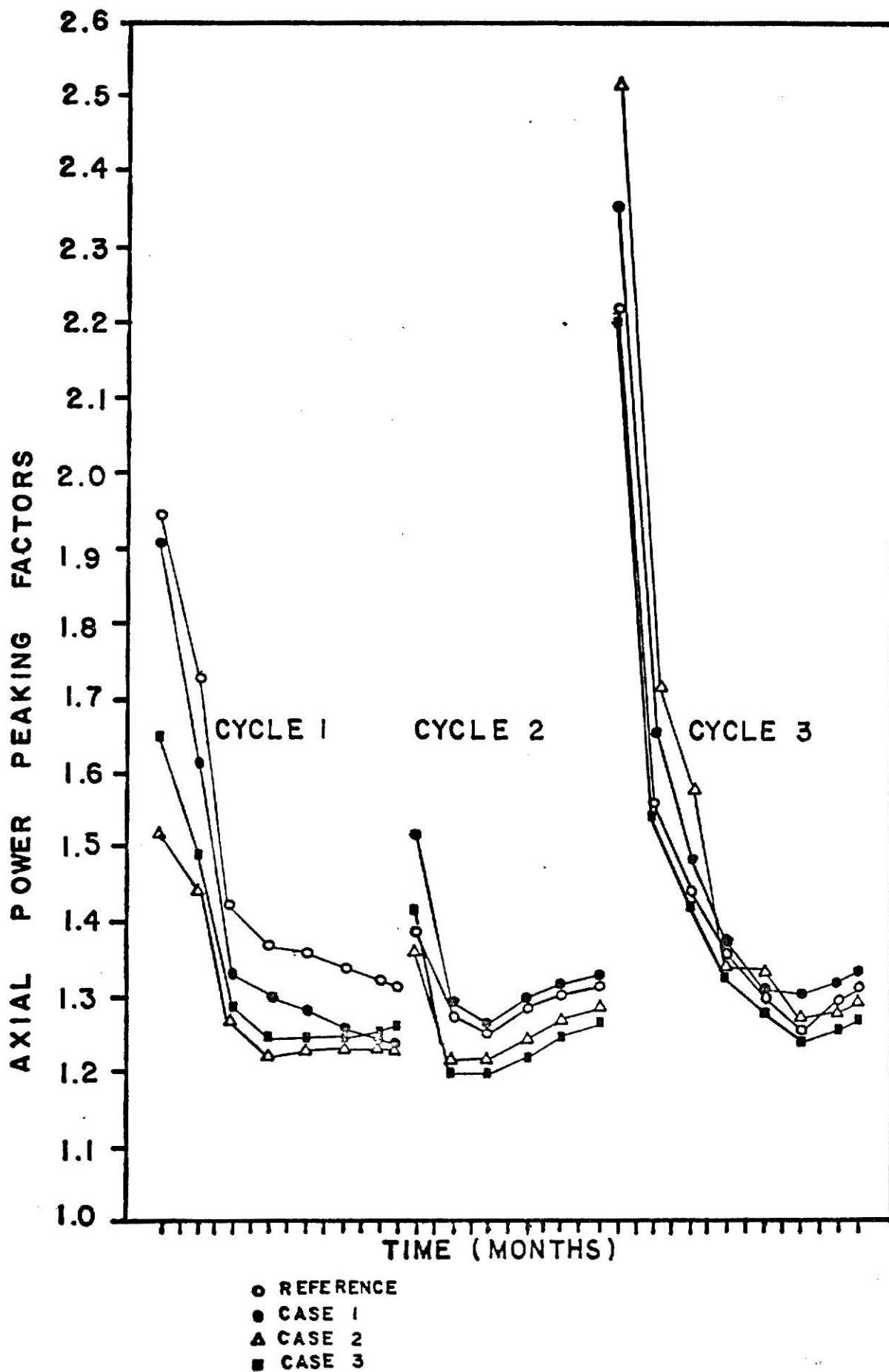


Figure 14. Behavior of Axial Power Peaking Factors with Increased Burnup

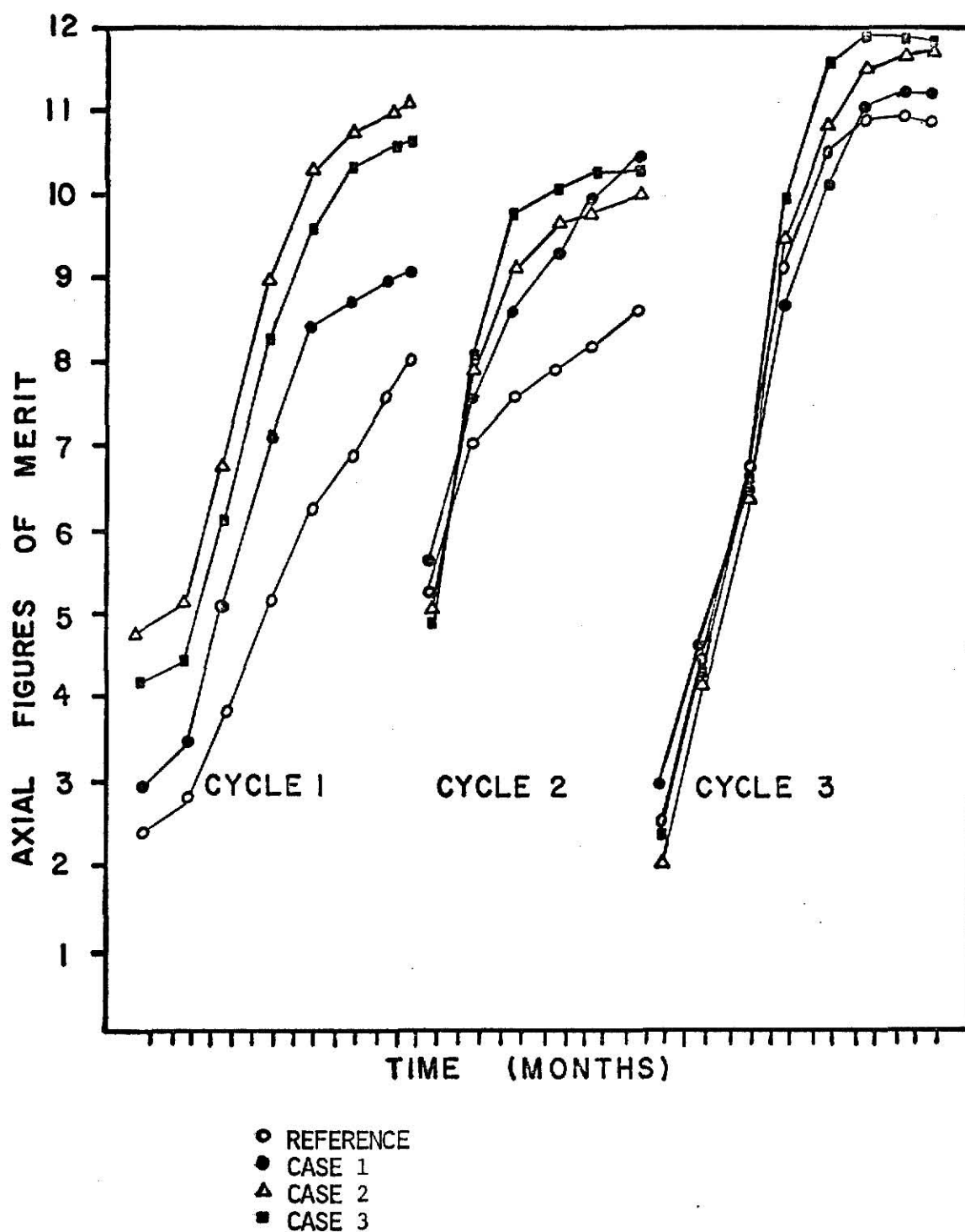


Figure 15. Behavior of Axial Figures of Merit with Increased Burnup

Figures 16, 20, 21 and 22 illustrate this problem. Figure 16 shows the column average burnup for all three cycles of the reference case. Figures 20, 21, and 22 show the percent deviation from the column average, that is, the standard deviation of the fifteen mesh points along the fuel element from the mean value divided by the mean. The higher the coefficient of variation expressed as a percent, the greater the unevenness of the burnup. An ideal case, perfectly even burnup, would have zero deviation. This display of the underutilization of the fuel in the reactor shows the need for some means of evening the fuel burnup throughout the core.

4.2 Case 1 Discussion

The fuel elements used in Case 1 are designed to lower both the axial and radial power peaking factors and increase the burnup of the fuel at the reactor top and bottom edges by more closely approximating the ideal Winterburg fuel distribution. The fuel in the top and bottom segments of all fuel elements in the core is replaced by fuel of 3.1% enrichment, and that in the elements of the outer ring lowered to 2.6% in the middle third of each element. Element compositions are given in Table 2.

The radial behavior of all cycles of Case 1 is very similar to that of the reference case cycles, except at the very start of cycle 1, when the lowered enrichment of the Type 1 fuel elements causes a greater central peaking tendency. The similarity in X-Y power peaking factors can be seen in Figure 12 and the similarity in radial Figures of Merit in Figure 13. That the variations in axial enrichment do not greatly effect the radial power distribution can also be seen in Tables 3, 4, and 5 which give the X-Y discharge burnup for each cycle of each case. The discharged fuel for each case is very near the burnup of the reference case.

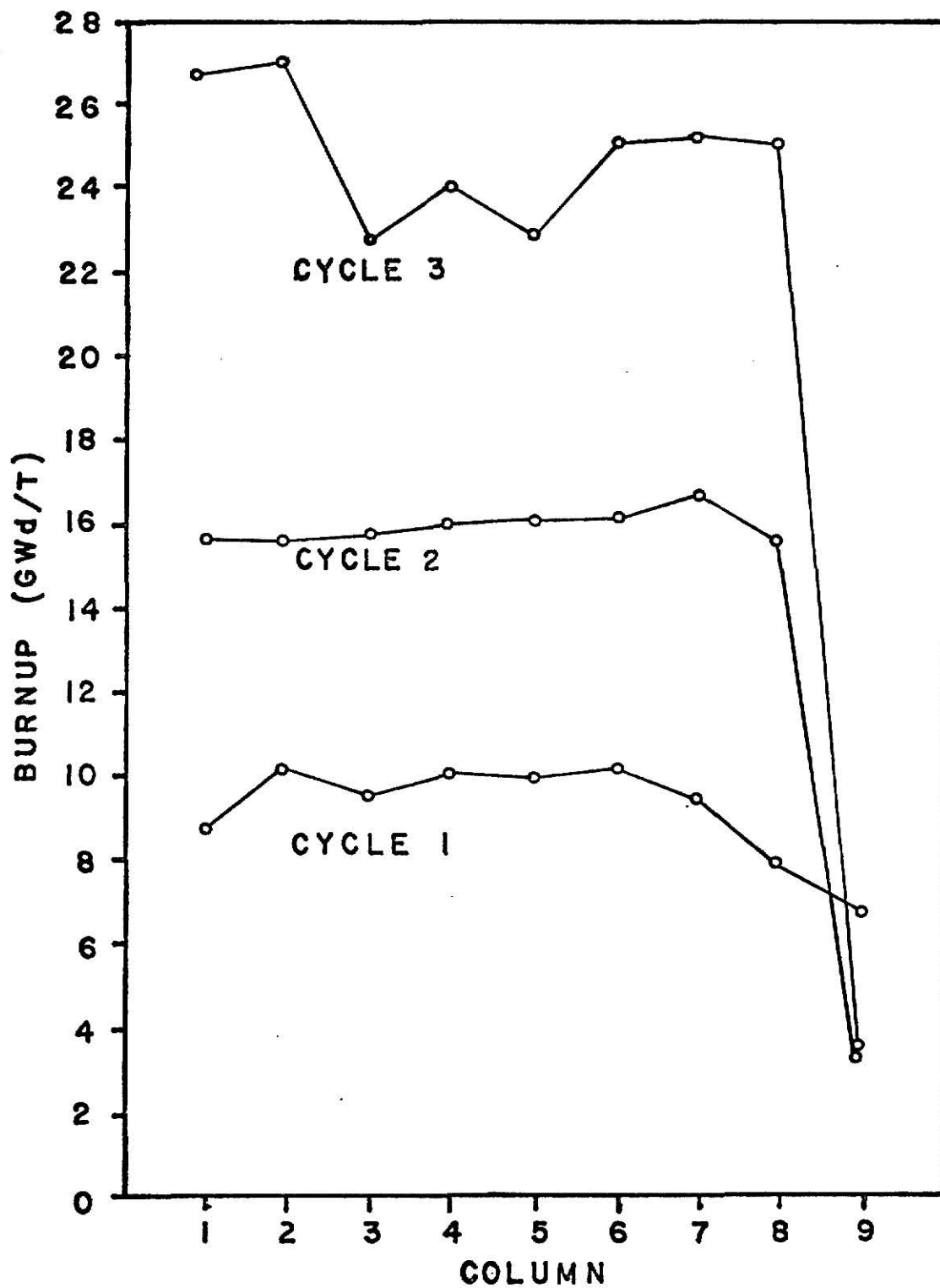


Figure 16. EOL Column Average Burnup for the Reference Case

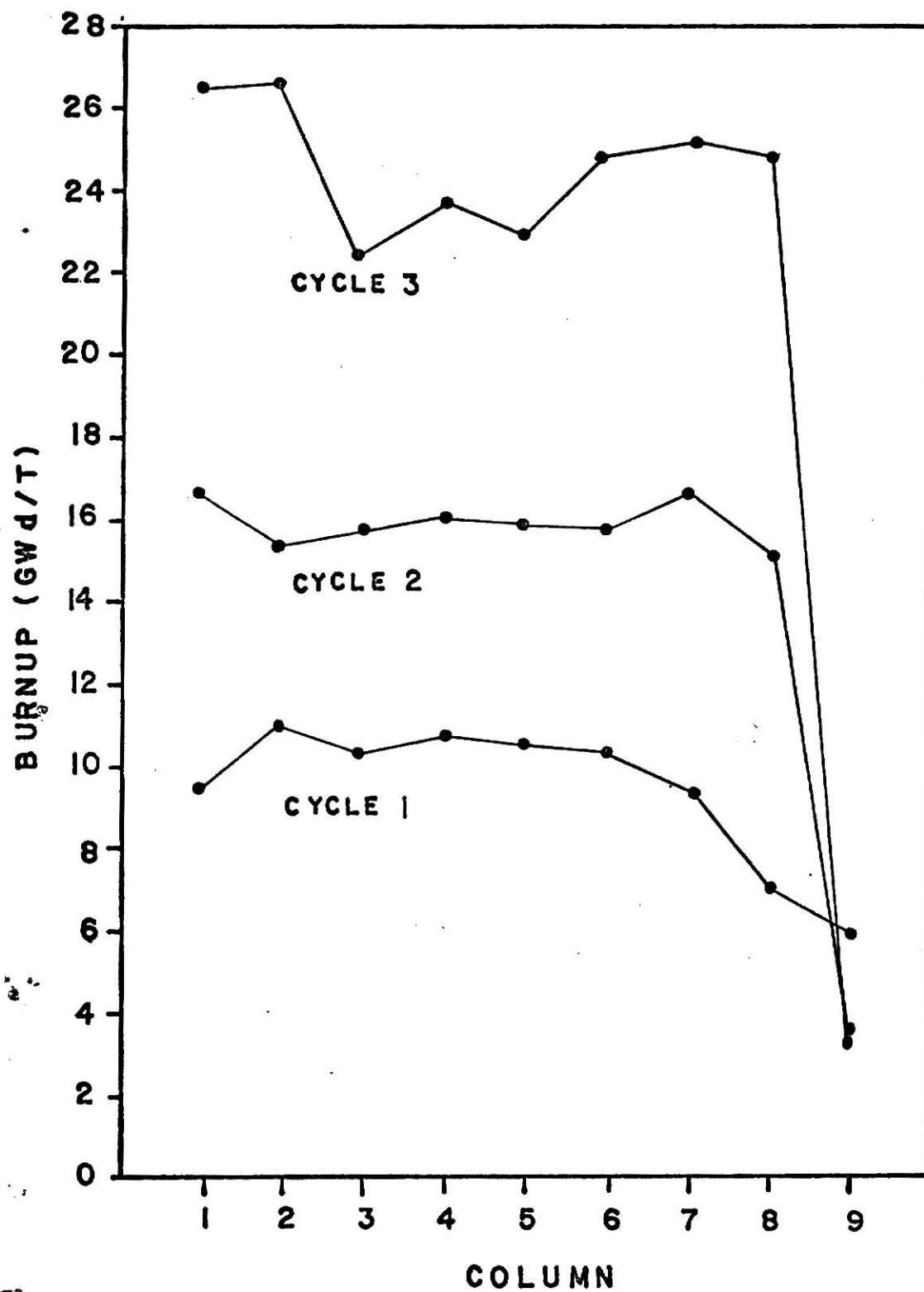


Figure 17. EOL Column Average Burnup for Case 1

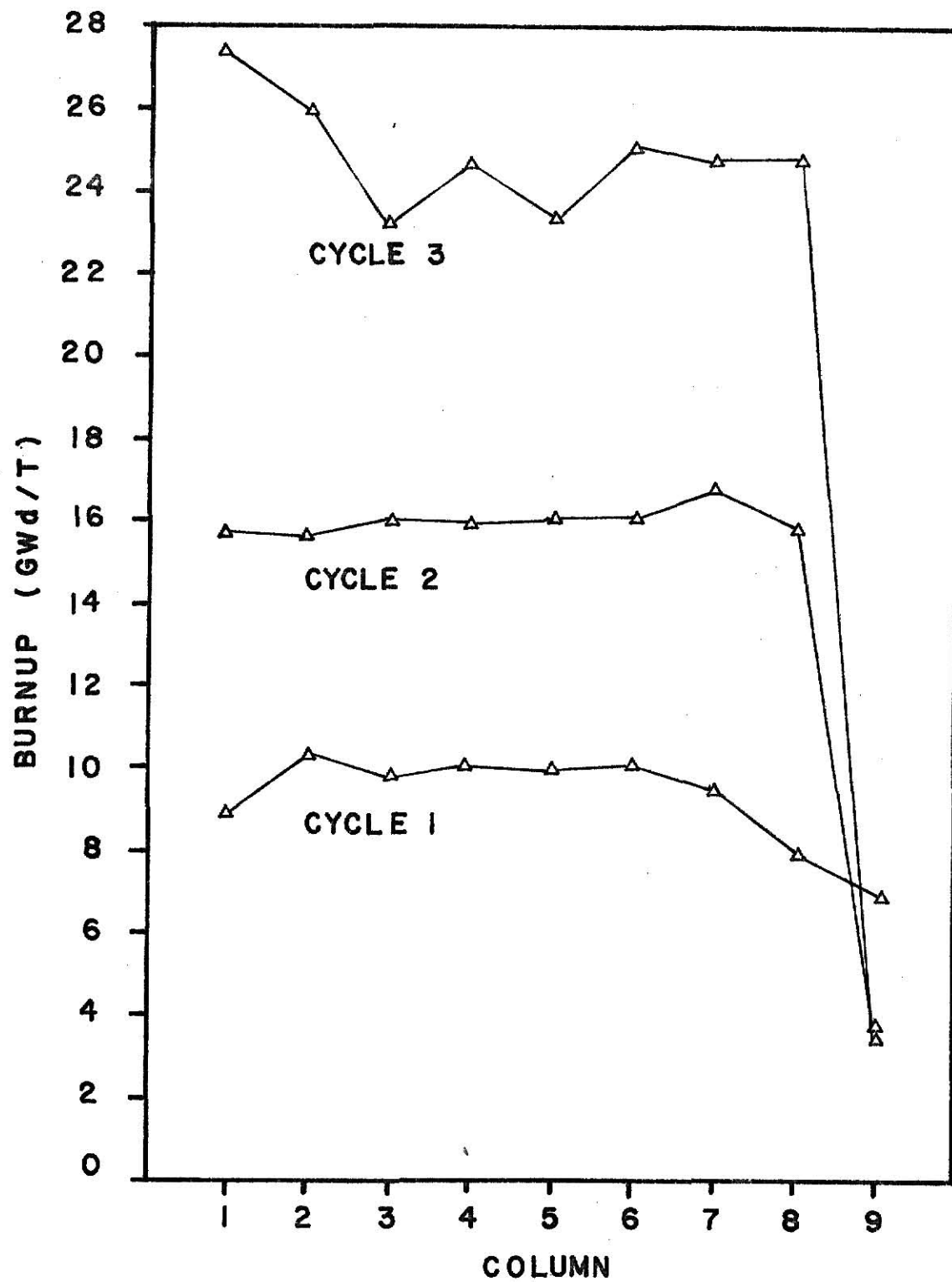


Figure 18. EOL Column Average Burnup for Case 2

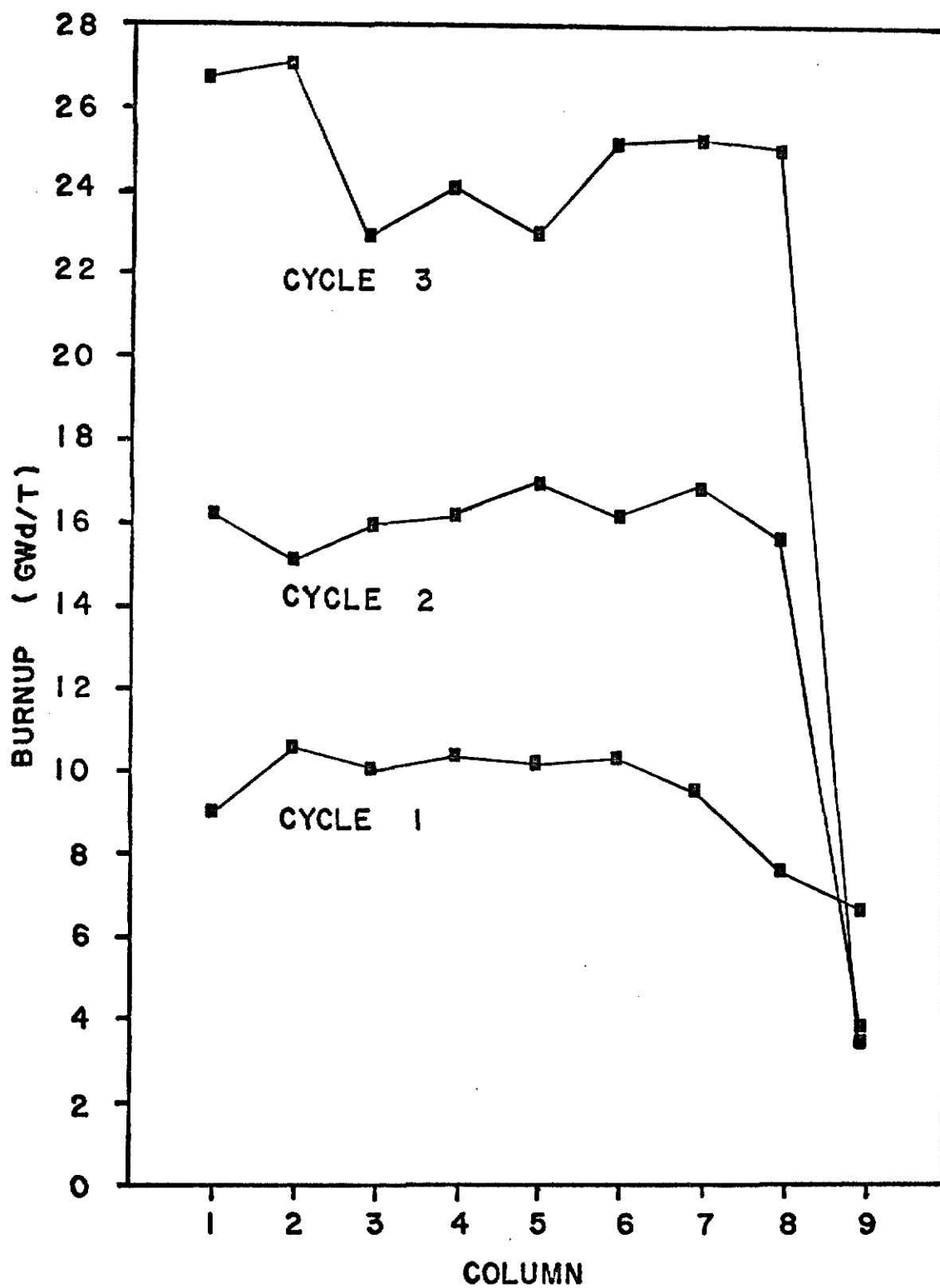


Figure 19. EOL Column Average Burnup for Case 3

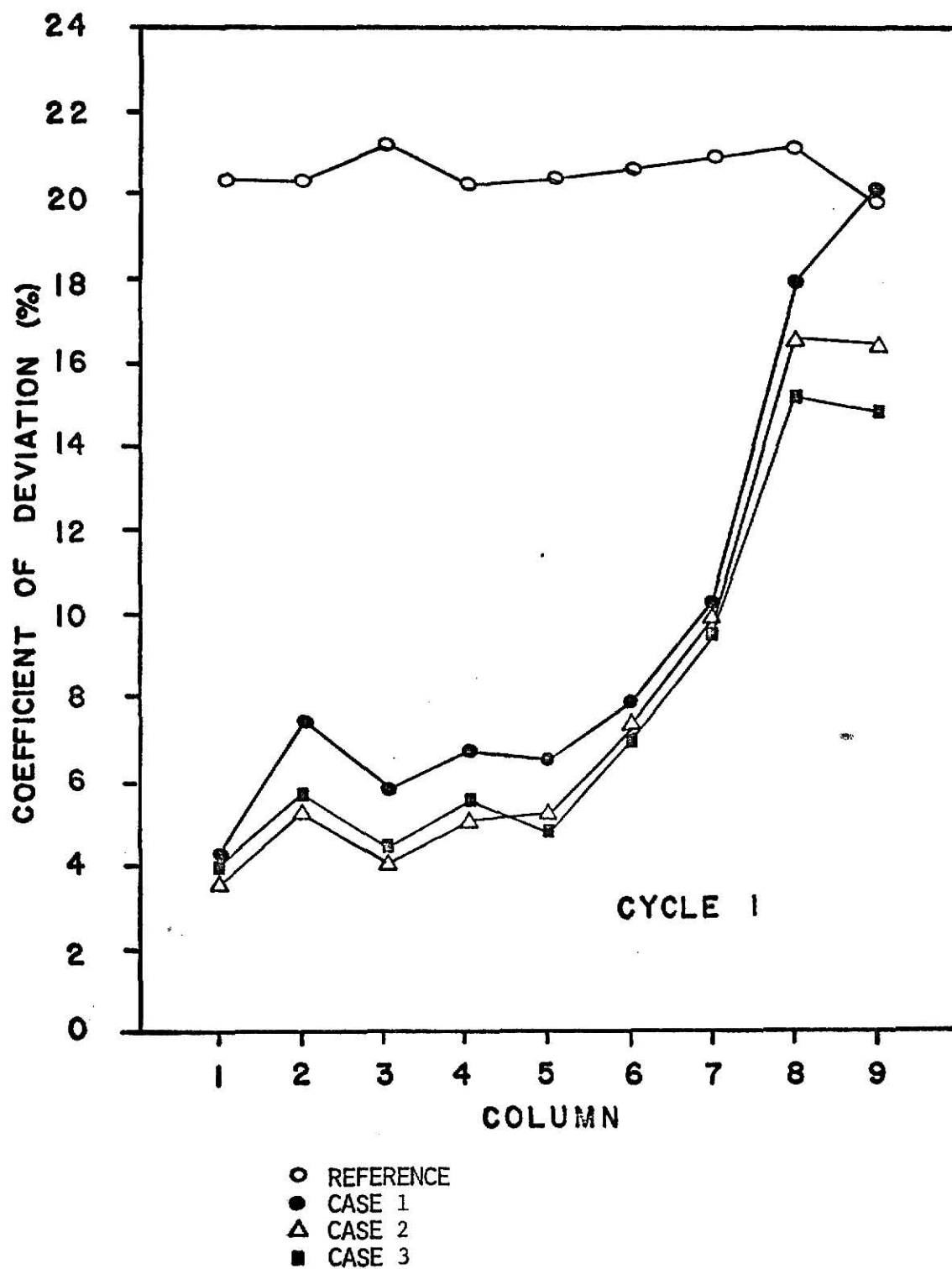


Figure 20. Coefficient of Deviation from Column Average Burnup for Cycle 1 of all Cases Examined

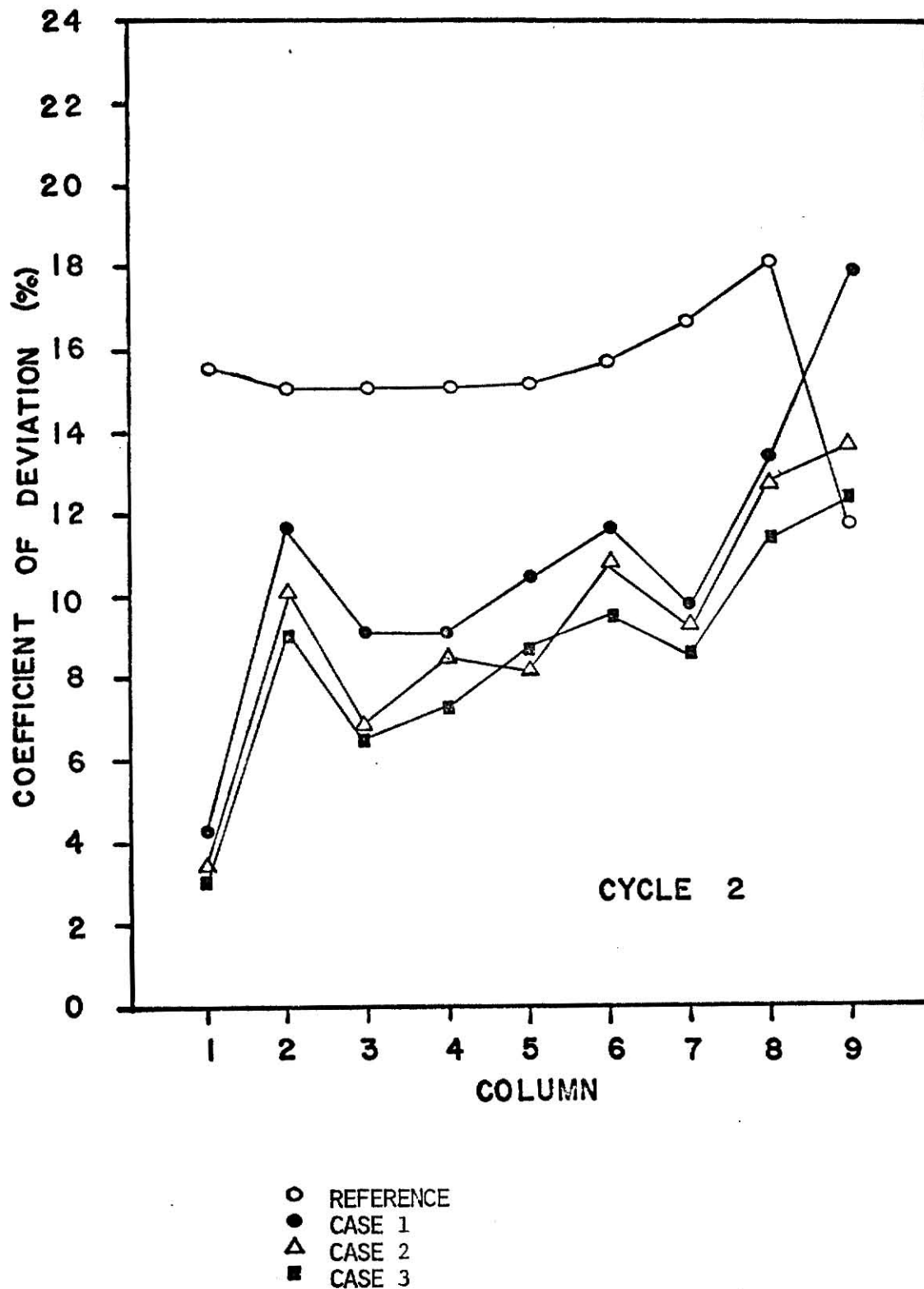


Figure 21. Coefficient of Deviation from Column Average Burnup for Cycle 2 of all Cases Examined

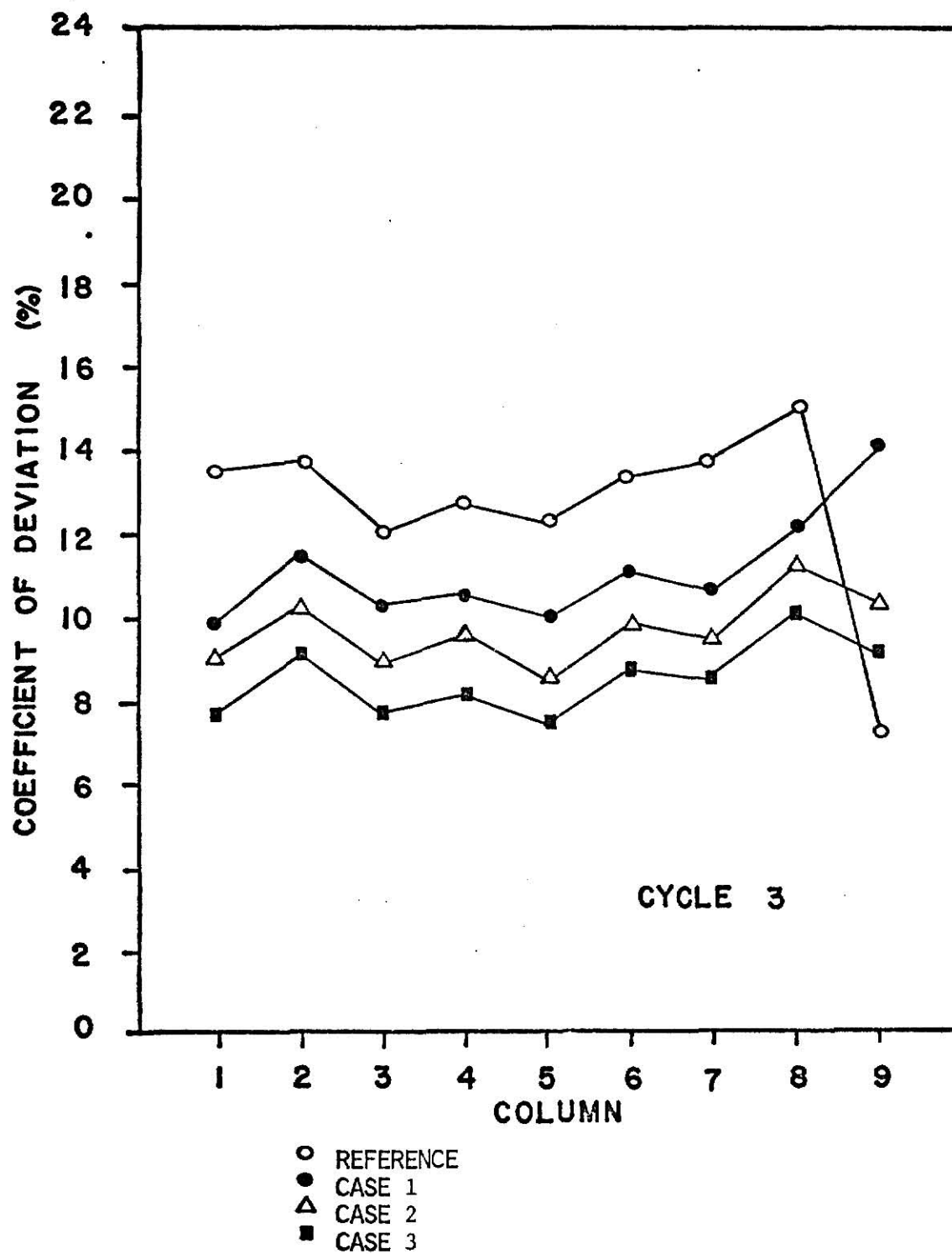


Figure 22. Coefficient of Deviation from Column Average Burnup for Cycle 3 of All Cases Examined

TABLE 3

Burnup of Discharged Fuel Assemblies, X-Y Determination
(MWd/T)

Assembly	Reference	Cycle 1		
		Case 1	Case 2	Case 3
1	8,718	8,767	9,365	8,852
2	9,163	9,258	9,390	9,334
3	9,176	9,630	9,399	9,474
4	9,184	9,642	9,528	9,484
5	9,231	10,084	9,551	9,721
6	9,250	10,109	9,565	9,749
7	9,260	10,123	9,608	9,767
8	9,297	10,132	9,634	9,817
9	9,312	10,149	9,632	9,828
10	9,324	10,202	9,637	9,833
11	9,327	10,216	9,638	9,840
12	9,331	10,216	9,643	9,847
13	9,337	10,229	9,666	9,847
14	9,355	10,237	9,782	9,863
15	9,500	10,245	9,800	9,897
16	9,516	10,246	11,063	9,916
Average	9,268	9,968	9,681	9,692
Std. Dev.	172	415	377	272
Coeff. of Dev.	1.85%	4.16%	3.89%	2.80%

TABLE 4

Burnup of Discharged Fuel Assemblies, X-Y Determination
(MWd/T)

Cycle 2				
<u>Assembly</u>	<u>Reference</u>	<u>Case 1</u>	<u>Case 2</u>	<u>Case 3</u>
1	13,215	14,043	13,589	13,798
2	13,515	14,880	13,916	14,360
3	16,253	17,611	14,229	16,802
4	16,543	17,974	15,934	17,119
5	16,582	18,006	16,581	17,161
6	16,676	18,057	16,656	17,237
7	16,874	18,127	16,755	17,394
8	16,923	18,151	16,907	17,413
9	17,041	18,191	16,923	17,473
10	17,064	18,224	17,002	17,622
11	17,788	18,443	17,121	17,806
12	18,872	19,678	17,760	19,137
13	18,898	19,720	18,916	19,174
14	20,147	20,987	20,134	20,391
15	20,148	20,990	20,140	20,392
16	20,469	21,302	20,498	20,707
Average	17,313	18,399	17,071	17,746
Std. Dev.	2,024	1,902	2,032	1,877
Coeff. of Dev.	11.69%	10.34%	11.90%	10.58%

TABLE 5

Burnup of Discharged Fuel Assemblies, X-Y Determination
(MWd/T)

Cycle 3				
<u>Assembly</u>	<u>Reference</u>	<u>Case 1</u>	<u>Case 2</u>	<u>Case 3</u>
1	20,713	20,164	20,820	20,544
2	23,251	22,411	23,163	22,970
3	24,109	23,199	23,755	23,715
4	24,129	23,218	23,995	23,736
5	24,151	23,295	24,837	23,824
6	24,870	24,253	25,623	24,916
7	25,284	24,715	25,637	25,218
8	25,957	24,811	25,811	25,517
9	25,999	25,122	25,869	25,529
10	26,230	25,481	26,067	25,796
11	26,534	25,938	26,380	26,203
12	28,320	27,240	27,968	28,017
13	29,268	27,994	29,272	28,834
14	30,114	28,843	29,437	29,628
15	30,674	29,178	30,153	30,044
16	31,279	29,494	30,493	30,531
Average	26,299	25,335	26,206	25,939
Std. Dev.	2,837	2,572	2,608	2,721
Coeff. of Dev.	10.79%	10.15%	9.95%	10.49%

The general cycle 1 axial behavior of Case 1 follows that of the reference case. The power peaking factors are uniformly about 5 to 10 percent lower and the Figures of Merit are roughly 30 to 40 percent higher as seen in Figures 14 and 15. The standard deviations of the axial fuel burnup are much smaller, as much as 75 percent lower, averaging about 60 percent lower, indicating that the fuel is being consumed much more evenly.

The BOL power profile of cycle 2 begins to show the difficulty inherent in using a multi-enrichment replacement fuel element. The BOL axial power peaking factor is about 10 percent higher than in the reference case. Although the axial peaking factor temporarily dips below the reference case about halfway through the cycle, the EOL axial peaking factor is also higher. This happens because the power generation is primarily coming from the top and bottom thirds of the reactor. The in-shuffling of the initial Type 1 fuel element from cycle 1 has produced a concentration of fuel near the top and bottom, and left the center underenriched.

The third cycle of Case 1 is much like the second, except that at BOL the axial power profile is already split into top and bottom parts due to the enrichment variation. The problem is compounded at EOL. The local power peaks exceed those of the reference case through most of the third core's life, even though overall the core burnup is more even than the reference.

The final burnup profile of Case 1 indicates that, while the deviations from average burnup are lower than the reference case (Figure 22) the elements are now being underused in the center because of the enrichment variation of the replacement elements.

4.3 Discussion of Case 2

Case 2 is again an attempt to achieve the ideal distribution. The results of case 1 indicate that the Type 1 fuel elements in cycle 1 and the replacement elements should be of more uniform enrichment and that the changed top and bottom segments of the other elements should be of higher enrichment. Therefore Case 2 core interior fuel elements have a 3.2% enriched segment top and bottom and the edge elements are entirely of 3.2% enriched fuel.

As in Case 1, the radial behavior of all cycles of Case 2 differs little from the reference case. The power peaking factors, radial Figure of Merit, and average discharge burnup are not significantly changed.

In all three cycles of Case 2 the general behavior of the axial power profile again follows that of the reference case. The peak power is considerably lower at all times except the beginning of the third cycle, when the in-shuffling of the higher enriched, 3.2%, fuel elements causes a temporary overconcentration near the core center. The axial Figures of Merit are higher than in the reference case, and in the first cycle, higher than any other case as well, indicating more even power profiles and therefore axial burnup, except again at the start of cycle 3, when it is slightly lower. The deviations from column average burnup are noticeable lower for all three cycles (Figures 20, 21, and 22) than either the reference case or Case 1, indicating better axial fuel consumption.

4.4 Discussion of Case 3

Case 2, while it has noticeable advantages over Case 1 and the reference case, has a disadvantage in its higher power peaking factor at the beginning of cycle 3. The fuel elements of Case 3 are designed to rectify this problem.

The fuel elements are like those of the reference case but with the enrichments of top and bottom segments of all elements raised to 3.2% as in Case 2.

The general radial and axial behavior of the power and flux distribution follows the general trend dictated by the reference case. The maximum axial power peak is always lower than the reference and the axial Figures of Merit are always higher. The radial averaged behavior is almost identical with the reference. As a result of having such a smooth power profile the Case 3 burnup distributions are also proportionately smoother. The columnwise burnup deviation from average for Case 3 is approximately equal to that for Case 2 in cycle 1 (Figure 20) but in cycles 2 and 3 is lower in the inner regions of the core, from which the discharged fuel comes, than for any of the other cases (Figures 21 and 22). Thus the fairly simple step of increasing the enrichment of the tips of the fuel elements results in a considerable improvement in fuel use.

5.0 CONCLUSIONS

The axial variations in fuel enrichment considered in this work are relatively small, and thus cause only small deviations from the reference in most factors. Figures 12 and 13 show that there is little difference in the radial power peaking factors or radial Figures of Merit among the four cases. Tables 3, 4, and 5, giving the X-Y discharge burnups, verify the observation that the X-Y results are little changed. The X-Y discharge burnup averages for the three cycles of all four cases follow similar patterns and vary by only a few percent. A most striking verification is given by Figures 16, 17, 18, and 19, the axial column average burnups. The four figures are almost identical, indicating the radial average burnup is unchanged by small variations in the axial enrichment.

While the axial enrichment variations have little effect on the radially averaged behavior of the reactor, they do have a significant impact on axial reactor response. The closer approximation of the Winterburg ideal fuel distribution has an immediate effect on all the pertinent parameters. Figure 14 shows the decrease of axial power peaking factor for the three experimental cases in cycle 1, for Cases 2 and 3 in cycle 2, and for Case 3 in cycle 3. Figure 15 shows the increase in axial Figure of Merit for all cycles of all the experimental cases, indicating a more even core power density profile. Finally, Figures 20, 21, and 22 demonstrate the leveling of the burnup profile along the length of the fuel elements, showing improved fuel usage.

The benefits of small axial variations are plain. Fuel burnup can be increased and power maldistribution decreased. Unfortunately, the fuel

distribution factor $Z(r,z)$ is highly sensitive to fuel enrichment. A deviation in uranium-235 enrichment of only one-tenth of one percent is quite significant. The difficulty in quality control in the manufacture of fuel elements will probably prohibit the implementation of axial variation in the near future. However, the idea is sound and is deserving of further study.

6.0 SUGGESTIONS FOR FURTHER STUDY

A logical extension of this work would be the investigation of a number of burnup cycles beyond the third to determine if an equilibrium cycle is reached, and what the effects of axial variation of fuel enrichment are on the equilibrium should one be reached.

The short length of cycle 2 in all cases indicates that it should be examined. The fuel enrichments or distributions in cycle 1 could be modified so that the shuffle into cycle 2 results in a core with more desirable characteristics.

While the advantages of axial variation of operating parameters are clear it is not obvious that compensation for the added cost in enrichment and manufacture of the axially graded elements is obtained. A study of the costs and benefits would be in order.

Finally, all the enrichment variations studied have been derived from an empirical fit to the ideal distribution. It would be advantageous if the fuel element configuration and shuffling could be devised by a linear programming model like that of Chen's. Conditions on radial and axial conformity to the ideal distribution could be formulated, as well as constraints to insure that fuel elements are shuffled as continuous units. This would ensure a three-dimensional adherence to the ideal distribution and thus give a reasonable base from which to launch further burnup studies.

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APPENDIX A: THE CODE BURNRZ

A.1 BURNRZ Functions

The code BURNRZ calculates the core burnup, power and flux profiles, and reactivity changes of a load of PWR fuel from start to end of a single burnup cycle. The processes or governing equations used in the dominant steps of the program are discussed below. The discussion follows the progression of the code flow diagram presented in the next section (A.2).

The code inputs are, first, the microscopic cross sections and nuclear physics constants of the elements used in the fuel and structure. Read in next are some of the basic lattice cell parameters; the diffusion coefficients, Fermi Age, diffusion length, moderator/fuel/cladding volume ratios, disadvantage factors, fast fission factor, and resonance escape probabilities for the fuel assemblies. Finally, the basic core size and geometry variables and reactor power level and operating history are input.

Subroutine XYTOR does the radial averaging. The square grid array of fuel assemblies is broken into an equivalent set of concentric annular rings. The amount of each fuel element present in each ring is found by a straight-line-segment technique, in which the arc of the zone boundary through the X-Y grid cell is represented by a straight line. The area of the cell within the zone is then approximated with various triangle or trapezoid area summations. When the fraction of each fuel element in a zone is known, the atom densities of the nuclides in that element are multiplied by the fraction and added to the zone, resulting ultimately in an overall zone average atom density.

The subroutine LIQHOH is used as a tabular storage of thermodynamic data for the PWR coolant water. The data is read in for water at the operating pressure of the PWR and for a wide range of temperatures. Since

the pressure is constant, a specification of any other property of the coolant defines the state of water. Thus, if the subroutine is called, giving one property, any of several other properties can be returned. The variables considered, besides pressure, are temperature, enthalpy, and density. The returned values are found by interpolation or extrapolation from the tabulated data.

The burnup calculation is based on the atom densities of the fission products present in the fuel element. Expressed in terms of megawatt-days per metric ton of uranium, the burnup is found from the equation

$$\text{B.U.} = \frac{1}{2} \frac{(\text{atoms of fission products/cm}^3)}{(\text{atoms of fuel originally present/cm}^3)} \frac{(\text{MWd})}{(1.052 \text{ g U-235})} \frac{(10^6 \text{ g})}{(\text{ton})} \quad (\text{A.1})$$

since to a good approximation of fissioning of 1.052 grams of U-235 is required for the release of one megawatt-day of energy. Use of atoms of fission products allows the contribution of plutonium to be included, with the assumption that the energy release is roughly equivalent for Pu-238 and U-235.

Subroutine CPROP is used to determine the group constants for each mesh point in the calculations. For points outside the reactor or reflector, a vacuum boundary is assumed, with zero properties, which implies that any neutron which gets that far is immediately lost to the system. For the water reflector the properties are set to those input for the moderator, except for those cases for which soluble boron control poison is used, for which the diffusion length is modified to account for the increased absorptions. The regions containing fuel assemblies are treated as homogenized regions in standard unit cell practice. The fuel macroscopic thermal fission cross sections are based on the product of atom densities and

microscopic cross sections. The fast group removal cross sections are based on the number of neutrons per fission and the fast fission factor, as well as the fission cross sections, as

$$\Sigma_R = \left(\sum_{i=1}^6 v_i \Sigma_i \right) (\epsilon - 1) / (v_{238}) \quad (A.2)$$

The thermal utilization is found as

$$f = \frac{N_{\text{fuel}} \sigma_{f \text{ fuel}}}{N_{\text{fuel}} \sigma_{f \text{ fuel}} + N_{\text{poison}} \sigma_{a \text{ poison}} + \Sigma_a \text{ material} (\phi_m / \phi_f) (V_m / V_f)} \quad (A.3)$$

The diffusion length is approximated as

$$L_{\text{fuel}}^2 = L_{\text{moderator}}^2 (1 - f) \quad (A.4)$$

The infinite multiplication factor is

$$k_{\infty} = p f \frac{\Sigma(v_i \Sigma_i)}{\Sigma(\Sigma_i)} \quad (A.5)$$

Subroutine RZFLUX solves the two-dimensional two-group neutron diffusion equations

$$-\nabla \cdot D_f \nabla \phi_f + \Sigma_R \phi_f = v \Sigma_f \epsilon \phi_{th} / k_{\infty} \quad (A.6)$$

$$-\nabla \cdot D_{th} \nabla \phi_{th} + \Sigma_{tot} \phi_{th} = p \Sigma_R \phi_f \quad (A.7)$$

The radial-axial two-dimensional Laplacian operator gives, with the group notation suppressed and the source term represented by S,

$$-\frac{1}{r} \frac{\partial}{\partial r} D r \frac{\partial \phi}{\partial r} - \frac{\partial}{\partial z} D \frac{\partial \phi}{\partial z} + \Sigma \phi = S \quad (A.8)$$

This equation is integrated over the radial-axial area around the mesh point as

$$\int_{r-h/2}^{r+h/2} 2\pi r \, dr \int_{z-l/2}^{z+l/2} dz \left[-\frac{1}{r} \frac{\partial}{\partial r} D r \frac{\partial \phi}{\partial r} - \frac{\partial}{\partial z} D \frac{\partial \phi}{\partial z} + \Sigma \phi \right] = f \quad (A.9)$$

where

$$f = \int_{r-h/2}^{r+h/2} 2\pi r \, dr \int_{z-\ell/2}^{z+\ell/2} dz \, S \quad (\text{A.10})$$

Equation (A.9) integrates to the form

$$\begin{aligned} -2\pi\ell \left[D r \frac{\partial \phi}{\partial r} \Big|_{r+h/2} - D r \frac{\partial \phi}{\partial r} \Big|_{r-h/2} \right] - \pi r^2 \Big|_{r-h/2}^{r+h/2} \left[D \frac{\partial \phi}{\partial z} \Big|_{z+\ell/2} - D \frac{\partial \phi}{\partial z} \Big|_{z-\ell/2} \right] \\ + \pi r^2 \Big|_{r-h/2}^{r+h/2} \ell \Sigma \phi = f \end{aligned} \quad (\text{A.11})$$

The last term on the left hand side of (A.11) is evaluated as

$$r^2 \Big|_{r-h/2}^{r+h/2} = 2rh \quad (\text{A.12})$$

The derivatives in (A.11) may be approximated, assuming the diffusion coefficient to be constant, with a central difference approximation as

$$D r \frac{\partial \phi}{\partial r} \Big|_{r+h/2} = D(r+h/2) [\phi_{r+h} - \phi_r]/h \quad (\text{A.13})$$

Substituting these results into Equation (A.11) and assuming a square box around the mesh point so that $h=\ell$, the resultant equation is

$$\begin{aligned} -(1 + h/2r)\phi_{r+h,z} - (1 - h/2r)\phi_{r-h,z} - \phi_{r,z+h} - \phi_{r,z-h} \\ + [4 + h^2 \Sigma/D]\phi_{r,z} = f/2\pi D r \end{aligned} \quad (\text{A.14})$$

Equation (A.14) is the general form of the diffusion equation solved by the code BURNRZ.

A good derivation of the finite difference equations is given in Bell and Glasstone⁽⁴⁾ for X-Y geometry. The procedure is equivalent to the R-Z case with only small variations in the Laplacian operator.

The resultant set of coupled linear equations for the flux density at each mesh point is solved using point simultaneous overrelaxation. In symbolic form the matrix equations can be written as

$$A\phi = S \quad (\text{A.15})$$

The matrix A can be partitioned into a diagonal matrix D , an upper triangular matrix E , and a lower triangular matrix F , such that

$$A = D - E - F . \quad (\text{A.16})$$

Substitution into equation (A.15) gives

$$(D - E - F)\phi = S \quad (\text{A.17})$$

which may be rewritten as

$$D\phi = E\phi + F\phi + S . \quad (\text{A.18})$$

Since D is a diagonal matrix, it has an easily determined inverse, so

$$\phi^{(m+1)} = D^{-1} [E\phi^{(m+1)} + F\phi^{(m)} + S] \quad (\text{A.19})$$

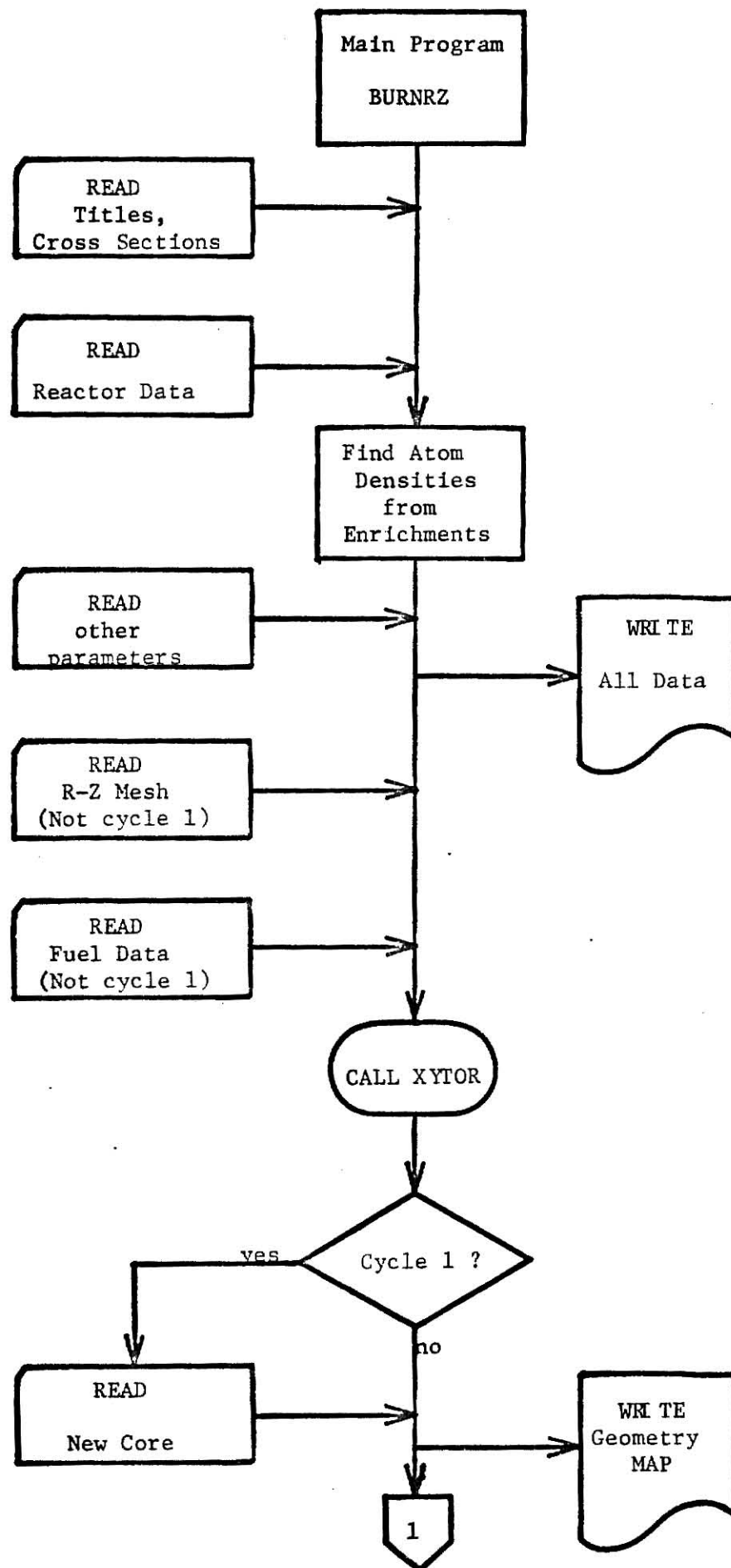
where the superscripts on the solution vector ϕ indicate that in the iterative solution the components of ϕ for the $m+1$ iteration are calculated from the values of the m -th iteration and that the new values are used as soon as they are calculated. Equation (A.19) can be written as

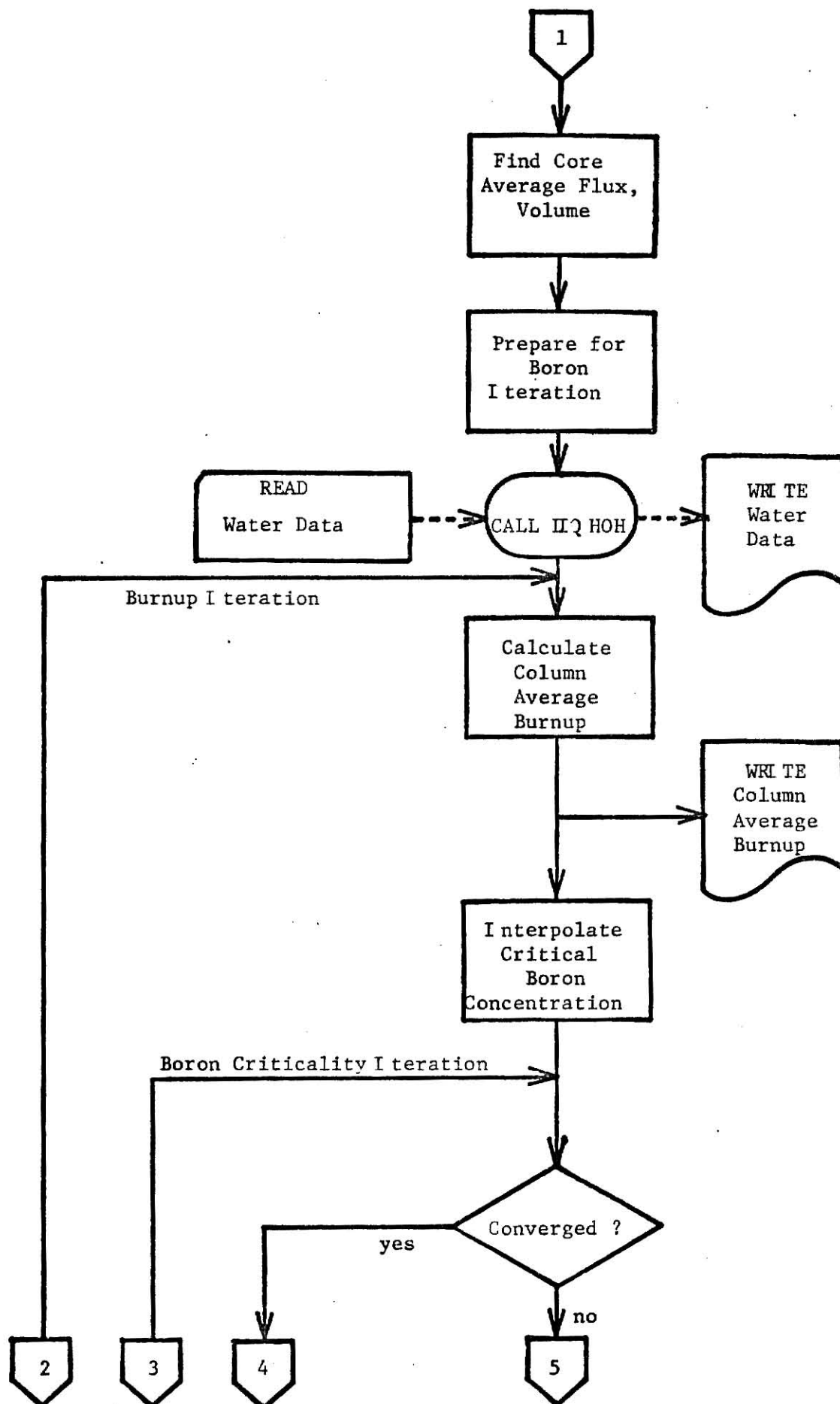
$$\phi^{(m+1)} = \phi^{(m)} + [D^{-1}(E\phi^{(m+1)} + F\phi^{(m)} + S) - \phi^{(m)}]_{\omega} \quad (\text{A.20})$$

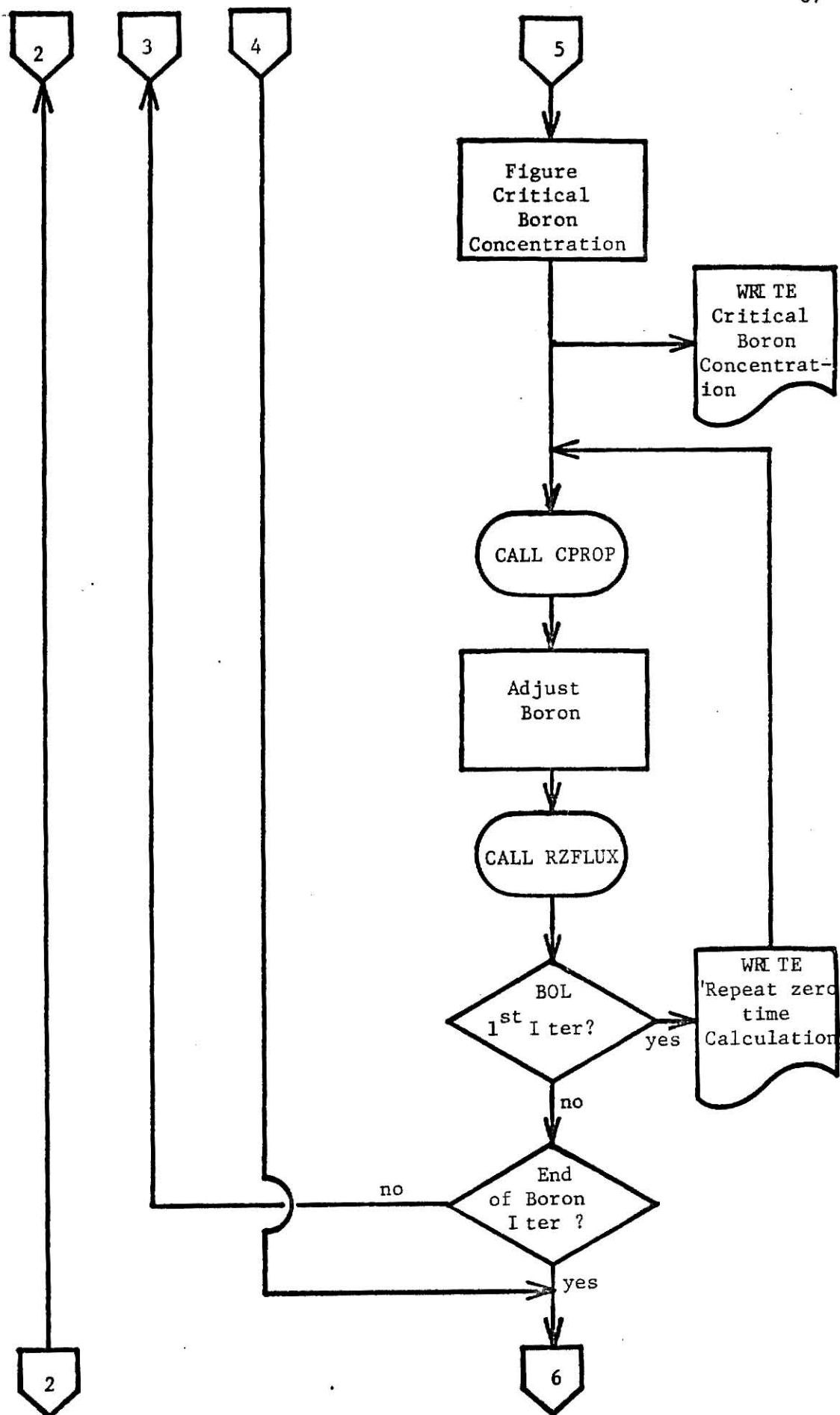
The term in brackets is the correction term added to the m -th iteration. An acceleration parameter ω , usually with a value between one and two, may be added to give the scheme an increased rate of convergence. A good discussion of successive overrelaxation, with methods of choosing the optimum relaxation parameter ω , is given by Varga⁽³¹⁾.

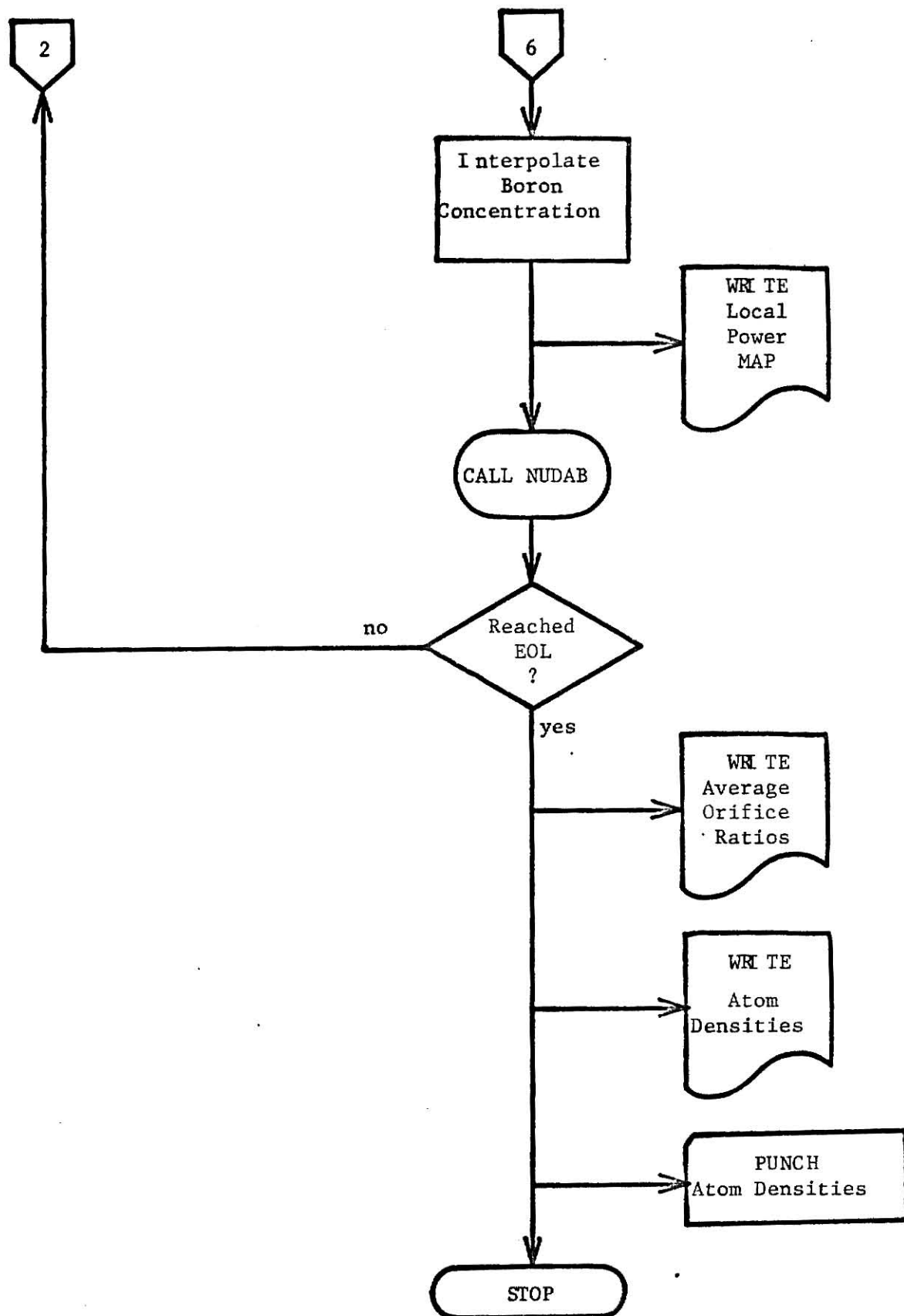
The subroutine NUDAB is used to find the atom densities if the various nuclides followed in the program. The nuclide densities are modeled as a set of coupled differential equations. The derivation of these equations and a complete listing thereof can be found in Hawk⁽³⁾ or in Chen⁽¹⁹⁾.

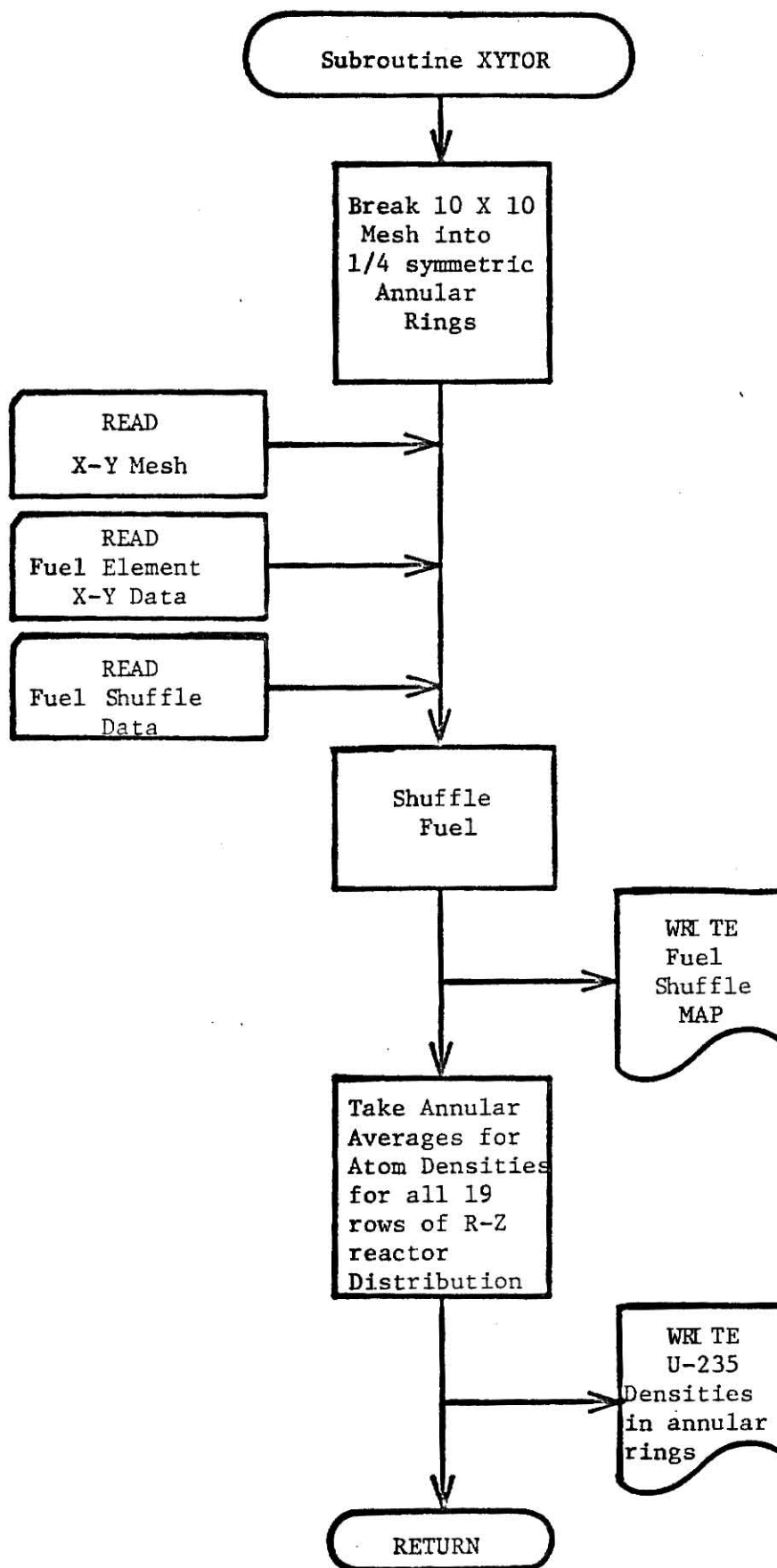
A.2 Flow Diagram of the Code BURNRZ

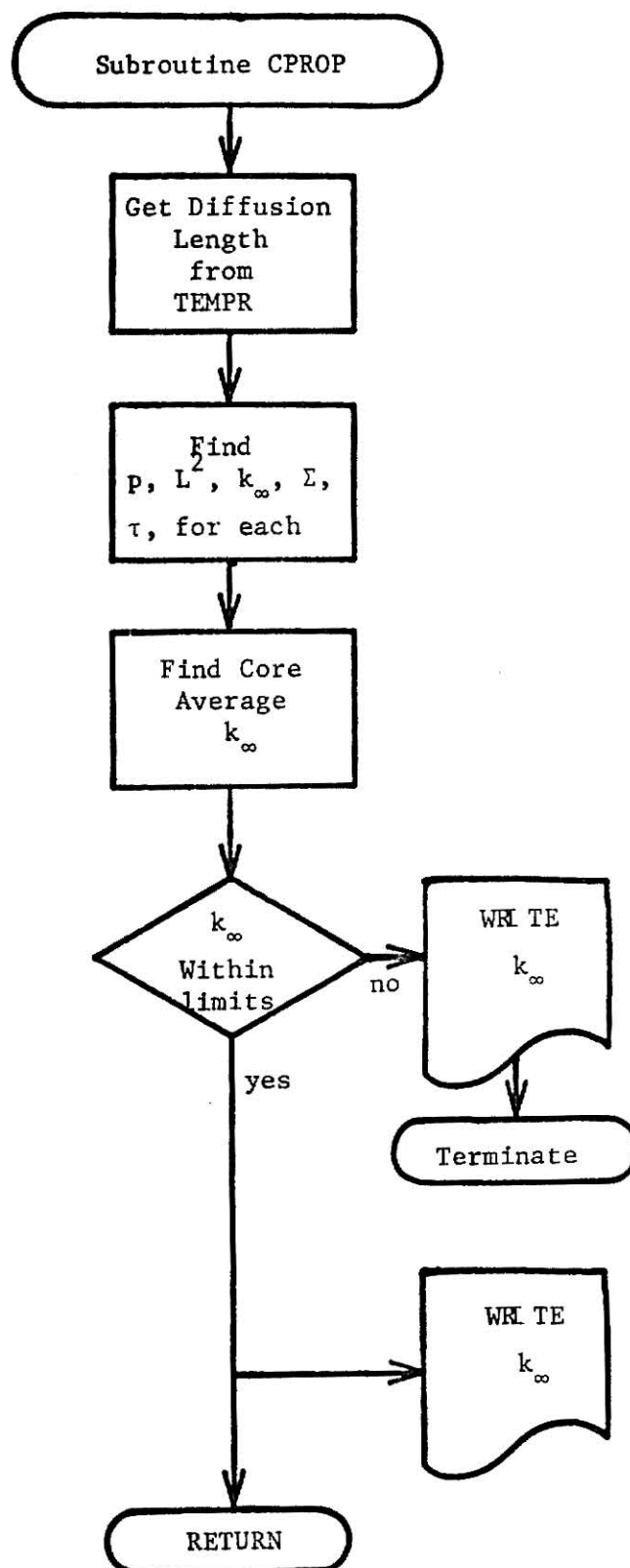


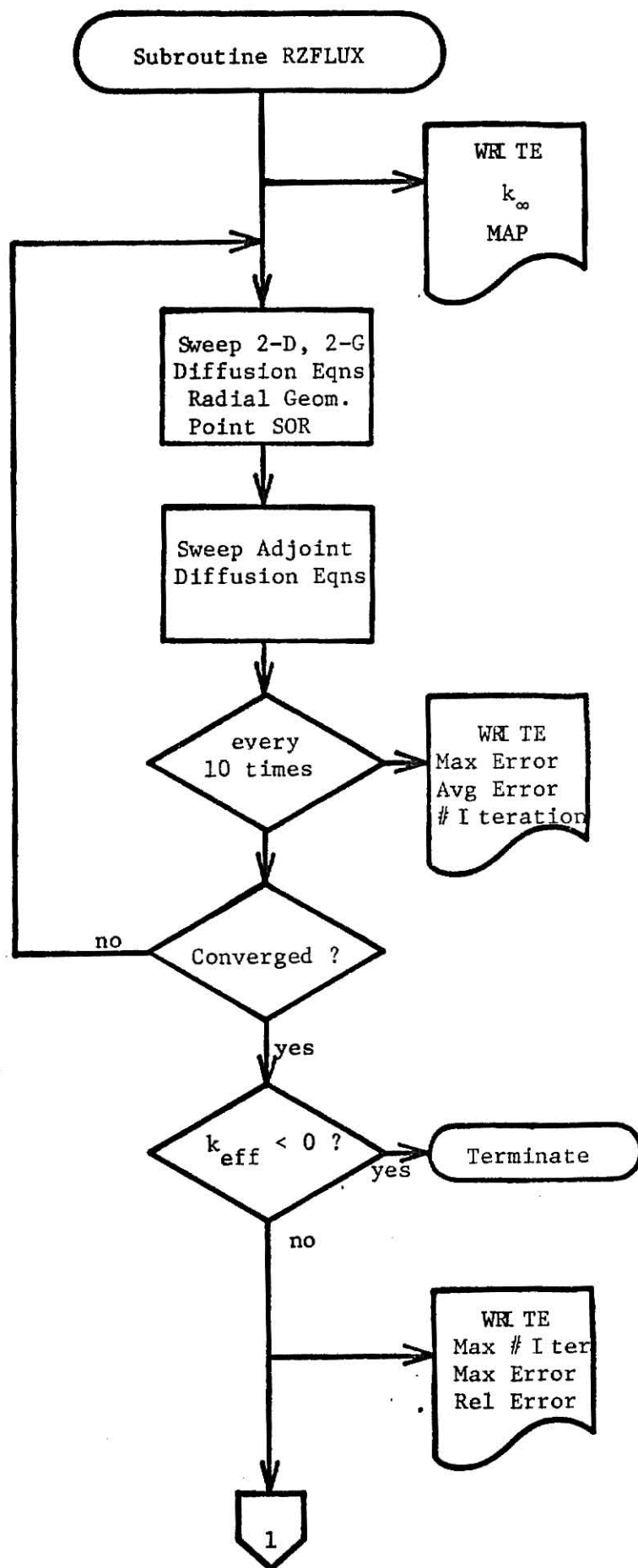


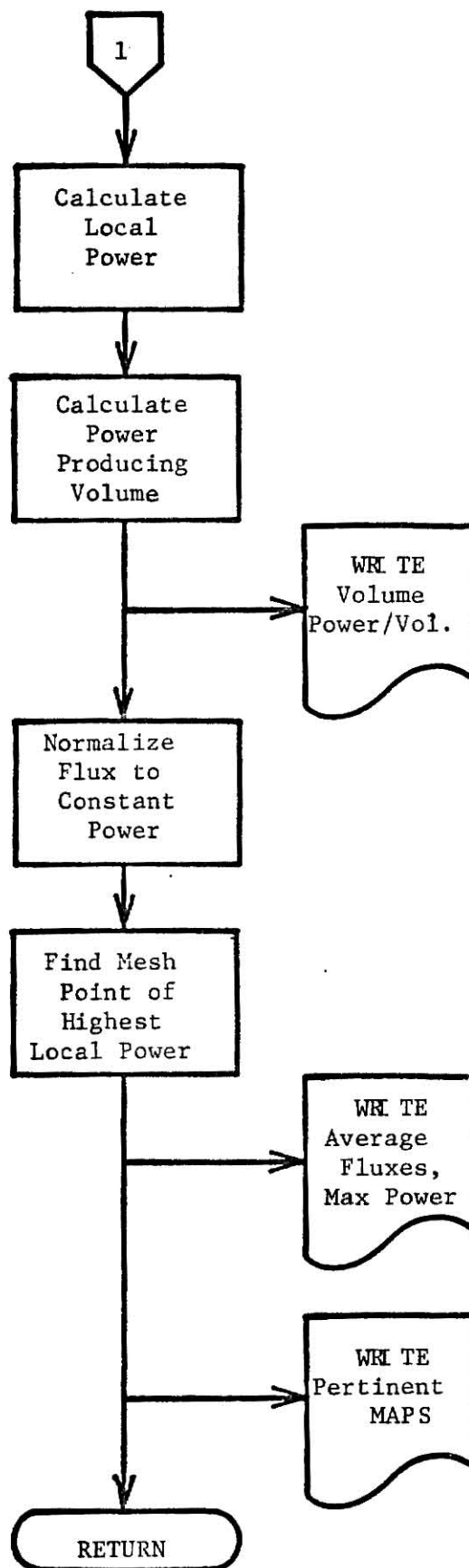


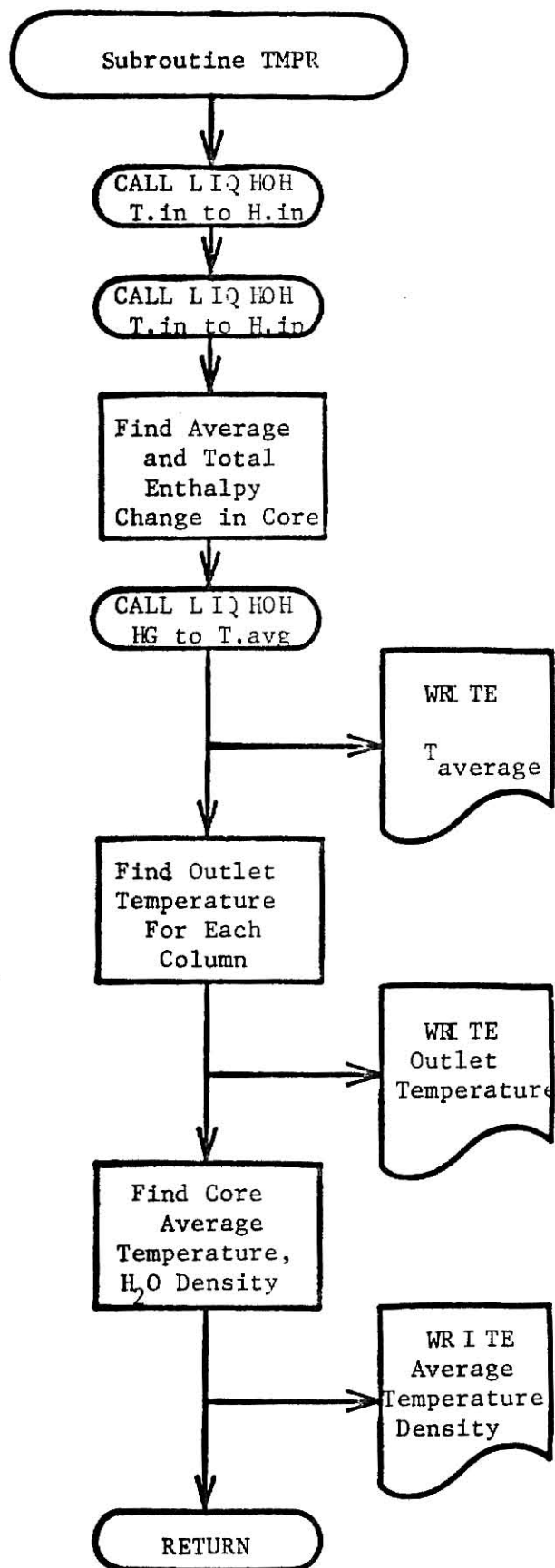












A.3 Computer Listing of the Code BURNRZ

**THIS BOOK IS OF
POOR LEGIBILITY
DUE TO LIGHT
PRINTING
THROUGH OUT IT'S
ENTIRETY.**

**THIS IS AS
RECEIVED FROM
THE CUSTOMER.**

```

C *** INDIVIDUAL ASSEMBLY BURNUP PROGRAM
C *** READS SEPARATE INPUT DECK OF ATOM DENSITIES
C *** A-Z FOR
DATA P1/D.141502657
DATA L4/M/276.77, A9/6.02225/
DATA KPNV7/7
DIMENSION SP(10,10), FEN(10)
DIMENSION WSH(100), RECH(100), ACNT(20), ENRICH(10,20),
1F25(10,20), F23(10,20), EPPM(10), SAE(10,190), IPSUM(190),
2SCS(10,150), SOD(10,10,20)
COMMON/ATM/ATM5,PL,P,CT5,TPAF(190),VNU5,VNU9,VNU1,VNU3,TAU
COMMON/COE/PCTA(190),PCTF(190),WDCOR(190),CK(190),RE(190),
1TCSQ(190),TAUF(190),SISF(1190),SIF12(1190),P2AV(190),NCOG,
2AFEX, XACF, P2L2,SAC,XNUF,XNUJ(1190),EPP
COMMON/COE2/P2RC,NCL,NSV,JMAX,H,DI,DZ,EIGHT,ALPHA,CCRT,
1CCMCO,FWATT,VOLMF,VOLLM,V(10),FCFRAC
COMMON/TP/SACL,SAC2,CENHDA,XDCO,TINERATOR(10),KAT,TOUT(30,10),
1IRUN,LEAR
COMMON/SECT/XST(2,9),VNU(9),GAM15,GAMX5,GAM19,ALAMX,SAX,GAMP5,
1GAMP9,VCELL
EQUIVALENCE (XST(1,1),S45),(XST(1,2),S46),(XST(1,3),SC3),
1(XST(1,4),S40),(XST(1,5),S41),(XST(1,6),S41),(XST(1,9),SABP),
2(XST(2,1),S55),(XST(2,4),S55),(XST(2,6),S51)
3,(XST(1,7),S5P),(XST(1,8),S4S)
2 FORMAT(24PRE10.4)
33 FORMAT(10,5X,10.0/(10,10.0))
9753 FORMAT(20A4)
1 FORMAT(16I5)
READ(L,5753)ACMT
READ(L,11)CYCLE,MODE,IXELSE
C *** MODE = 0: NORMAL
C > 0: PRINT ATOM DENSITIES
C < 0: PRINT ATOM DENSITIES IN ADDITION
DO 2588 I=1,2
DO 2589 J=1,9
2588 XST(I,J)=0.
READ(L,2)S45,S46,S49,S40,S41,SC5,SC8,SC9,SC0,SF5,SF9,SF1,SPP,SAS,
1SAX,SAB2
READ(L,3)S4SS,SAM
DO 7712 I=1,9
7712 VNU(I)=0.
READ(L,3)EPS,P1,P,ET5,ET9,GAM05,GAMP9,GAMX5,GAM15,GAM19,ALAMX,VNU5,ABP00130
1,VNU9,VNU1,VNU8,B2,TAU,D12,VOLMF,VOLSF,DICMF,DICSF,ALAMP
VNU(1)=VNU5
VNU(4)=VNU9
VNU(5)=VNU1
VNU(6)=VNU1
READ(L,3)PCT,PERIOD,FUELD,BELOY,FCFRAC
READ(L,3)JNGO,NCL,NSV,MKEY,JMAX
READ(L,1)NRICH
DO 2828 I=1,NSRICH
2828 J=1,NSICH
DO 100 J=1,N-CM
COMMON/DRAG/1,1132,2235,1*(1,10.#ENRICH(I,J)/235+1)
F25(I,J)=ENRICH(1,J)*0.0007235
F28(I,J)=0.07*(1.-ENRICH(1,J))/235.
100 CONTINUE
2828 CONTINUE
C *** SOLUTION FORM ONLY OF B-10 ALLOWED

```

ABP00020

ABP00030

ABP00100

ABP00120

ABP00130

ABP00140

ABP00150

ABP00160

ABP00170

ABP00190

ABP00190

ABP00200

ABP00210

ABP00220

ABP00230

ABP00250

ABP00260

ABP00270

ABP00280

ABP00290

ABP00300

ABP00310

ABP00320

ABP00330

ABP00340

ABP00350

ABP00360

ABP00370

ABP00380

ABP00510

ABP00460

```

READ(L,33)NDLE,BENHCH,(BPPM(I),I=1,NDLE)
CENS=VCLMF*DIS*FAC=0.1978E-6*BENHCH/10.
READ(L,3)AMODT,TIN
DO 2205 I=1,NBIF
2205 BPPM(I)=BPPM(I)/CEN3
READ(L,3)PHI,D1,D2,HIGHT,ALPHA,CORIT,CORPON,FINATT
READ(L,1)MRAF
IF(MRAF.NE.0)READ(L,3)RATOR
  ABOVE IS EIGHT PER CARD
  NC=IABS(NCYCLE)
  WRITE(K,1071)NC
  WRITE(K,2216)DENHCH,(BPPM(I),I=1,NDLE)
  WRITE(K,2217)XMODT,TIN
2217 FORMAT(//5X,'CORE FLOW RATE =',E11.4,' #/HR/5X,
  17-INLET =',F6.1,' DEG. F.)
  DO 1212 I=1,NRICH
  WRITE(K,1200) I, (ENRICH(I,J), J=1,NROW)
1200 FORMAT(//5X,' ELEMENT TYPE ',14,' ENRICHMENTS',/20F6.3)
1212 CONTINUE
  WRITE(K,1047)SC5,SAS,SAG,SAD,SAL,SC5,SC8
  WRITE(K,1048)SC9,SCU,SFS,SF9,SF1,SPP,SAS
  WRITE(K,1049)SAX,SADP,SASS,SAM
  WRITE(K,1050)EPS,P1,P7,ET5,ET9
  WRITE(K,1051)GAMPS,GAMP9,GAMX5,GAM19
  WRITE(K,1052)ALANX
  WRITE(K,1152)ALAMP
  WRITE(K,1053)VNUS,VNUR,VNHL,VNUS
  WRITE(K,1054)TAU,D12,B2
  WRITE(K,1055)DISSE
  WRITE(K,1056)VOLMF,VOLSF
  WRITE(K,942)JNPRW,NCOL,VSYM,MKEY,JMAX,H,D1,D2,HIGHT,ALPHA,CORIT,
  1CCORPON,PIKATI
9420 FORMAT(//5X,'ADDITIONAL PARAMETERS NEEDED FOR FLUX CORE CALCULATION)
  1N//5X,'NO. OF ROWS =',13,5X,'NO. OF COLUMNS =',13,5X,'SYMMETRY FLAG=00710
  2C10P =',13//5X,'FLUX PRINTOUT KEY =',13,5X,'MAXIMUM NO. OF ITERATIONS=0720
  3CONS =',15//5X,'MESH SIZE =',F7.2,5X,'D1 =',F7.2,5X,'D2 =',F7.2,5X,'ABP00730
  4CORE HEIGHT =',F10.2,5X,'OVER-RELAXATION PARAMETER =',F7.2//5X, ABP00740
  5CONVERGENCE CRITERIA =',F8.5,5X,'CORE POWER =',E12.4,5X,'FISSIONS=ABP00750
  6PER WATT =',F12.4)
  1071 FORMAT(//5X,'***** BURNUP PERIOD CYCLE NUMBER ',12,3X,'*****')ABP00770
  2216 FORMAT(//5X,'BURNABLE POISON TABLE DURING CYCLE (PMR CASE)'/10X, ABP00780
  1'WATER DENSITY =',F9.4/10X,'ATOM DENSITIES BASED UPON FUEL VOLUME',ABP00790
  2//10X,8F14.4))
  1047 FORMAT(//5X,'CROSS SECTIONS IN CM**2/I1,' ,115,'SA25',125,'SA2ABP00910
  16,'TBS,'SA49',145,'SA+D',155,'SAH',165,'S025',175,'SC25//',1,10X,ABP00920
  2710.2/)
  1049 FORMAT(11,' ,115,'XE-135',125,'SABP//',10X,2E10.3/10'.8X,
  1CROSS SECTIONS IN CM**1 STRUCTURE - CLAD =',F5.3,5X,'MODERATOR',ABP00930
  2CR =',F5.3)
  1050 FORMAT(11,' ,17,'REACTOR PHYSICS PARAMETERS',/10'.8X,'EAST FISSION',ABP00970
  1N FACTOR =',F7.4,5X,'FISSION-TO-RESONANCE NONLEAKAGE =',F7.4//',3ABP00980
  2X,'RESONANCE ESCAPE =',F7.4,5X,'ETA U-235 =',F7.4,5X,'ETA PU-239 =ABP00990
  3',F7.4)
  1051 FORMAT(//5X,'FRACTIONAL FISSION PRODUCT YIELDS (GAMMA)'/10X,ABP00910
  1'PM-149 (U-235) =',F7.4,5X,'PM-149 (PU-239) =',F7.4//',10X,'XE-134',ABP00920
  25 (U-235) =',F7.4//',10X,'I-135 (U-235) =',F7.4,5X,'I-135 (PU-239)',ABP00930
  339) =',F7.4)
  1052 FORMAT(//5X,'DECAY CONST. FOR XE-135 =',E10.2,'/SEC.',)
  1152 FORMAT(//5X,'DECAY CONST. FOR PM-149 =',E10.2,'/SEC.')

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ABP00470
ABP00480ABP00490
ABP00500
ABP00520ABP00530
ABP00540
ABP00550
ABP00560ABP00570
ABP00580
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ABP00600
ABP00610
ABP00620
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ABP00640
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ABP00680
ABP00690ABP00700
ABP00710
ABP00720
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ABP00880
ABP00890
ABP00900
ABP00910
ABP00920
ABP00930
ABP00940
ABP00950
ABP00960

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1053 FORMAT('0',8X,'AVERAGE NUMBER OF NEUTRONS RELEASED PER FISSION IN',ABP00970
10235='F7.4/11.',T57,IN U-239='F6.7/4/11.',T57,IN PU-2390980
2241='F6.7/4/11.',T57,IN U-238='F6.7/4/11.',T57,IN PU-2380990
1054 FORMAT('0',8X,'FUEL AGE',F5.1,'CM',F2.5X,'THERMAL DIFFUSION LEAD',F3.00
INSH IN MODERATOR (F2)='F5.1,'CM',F2.5X,'THERMAL DIFFUSION LEAD',F3.00
23,'/CM',F2.1)
1055 FORMAT('0',8X,'DISADVANTAGE FACTOR IN MODERATOR',F5.2,5X,'DISADVANTAGE',F3.00
LANTAGE FACTOR IN CLAD',F5.2)
1056 FORMAT('0',8X,'PHYSICAL DESCRIPTION OF CORE',/10',8X,'VOLUME RATIOS',F3.00
1 - MODERATOR/FUEL',F5.2,5X,'CLAD/FUEL',F5.2)
WRITE(K,2057)FUEL,PCF,PERIOD
8057 FORMAT('0',8X,'UEZ DENSITY',F6.2,'GRAMS/CUBIC CM OF FUEL',/
'0',8X,'THE PLANT CAPACITY FACTOR',F5.3,'FOR A PERIOD OF',F3.00
2,'F5.2,'YEARS')
1058 FORMAT('11.',T15,'SC49',T25,'SC40',T35,'SF25',T45,'SF49',T55,'SF4ABP01110
11',T64,'SAPP',T75,'SM-149',/,'10X',F10.3/)
ACCR0=ACCL*NROW
DO 4000 I=1,NCGR0
CK (I)=0.
SE (I)=0.
ELSO (I)=0.
TAUF (I)=0.
SIGF11(I)=0.
SIGF12(I)=0.
WGPCW(I)=0.
PRAV (I)=0.
XOUT (I)=0.
TPEF (I)=0.
DO 4000 J=1,10
SC4(J,I)=0.
4000 SAB(J,I)=0.
AVE=0.
READ(L,5001)(MESH(I),I=1,NCOR3)
5001 FORMAT(8C11)
DO 6050 I=1,NCOR0
IF (MESH(I).EQ.0) GO TO 6050
IF (MESH(I).EQ.2) GO TO 6050
READ(L,5002)(SAB(J,I),J=1,10)
6050 CONTINUE
CALL XYTCR(SAB,F25,F28,AROW,NCCL,SR,SDD,NCYCLE)
DO 6065 I=1,NCORC
IF (SAB(I,I).NE.0.160) GO TO 6036
6035 CONTINUE
READ(L,5001)(MESH(I),I=1,NCOR3)
KI=0
WRITE(K,6065)
FORMAT('1')
DO 6062 I=1,NROW
DO 6060 J=1,NCOL
FEM(J)=0.0
KI=KI+1
IF (MESH(KI).EQ.0) GO TO 6060
IF (MESH(KI).EQ.2) GO TO 6060
LLU=I
IF (NCYCLE.NE.1) LLU=1
FEM(J)=100.0*SDD(1,J,LLU)/(SDD(1,J,LLU)+SDD(2,J,LLU)+SDD(3,J,LLU))
DO 6061 JP=1,10
6061 SAB(JR,KI)=SDD(JR,J,LLU)
6060 CONTINUE
WRITE(K,6063) (FEM(J), J=1,NCCL)

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6063 FORMAT(/,10X,10F7.3)
6062 CONTINUE
      WRITE(K,6064)
6064 FORMAT( 10X, ' RACIAL AVG FUEL ENRICHMENT MAP, IN 3,
1 COLUMN 9 HAS BEEN MODIFIED')
6065 DO 770 I=1,NCOL
      TPSUM(I)=MESH(I)
      IF (MESH(I).EQ.1) GO TO 77C
      IF (MESH(I).EQ.3) GO TO 77D
      DO 771 J=1,13
        SAR(J,I)=0.
771   SAR(J,I)=0.
770   AVE=AVE+SAR(I,I)
      CALL MAP(NRQW,NCOL,TPSUM,I)
      VOLT=F*H*HIGHT
6000   AFLX=CCNFCW*FLWAT/(VOLT*AVE*SF5)
      V(I)=PI*H*H*H*0.25
      VCLF=2.*PI*H*H
      VCLUM=V(I)
      DO 6110 I=2,NCOL
        V(I)=VCLF*I*H
        IF (I.EQ.9) V(I)=V(I)*FCFRAC
6110   VCLUM=VCLUM+V(I)
      VCLUM=VCLUM*(NRQW-4)
      SAC1=SASS*VCLSF*DISSF
      SAC2=SAV*VCLMF*DISMF
      SAC=SAC1+SAC2
      VCELL=1.*VCLSF+VCLMF
      NBN=PERIOD/REFLUX+1.0001
      DF1=1./NBN-1)
      DF2=1./(NDLF-1)
      READ(L,I)N4
      WRITE(K,6065)N4
699   FORMAT(/,5X,'BORDON ITERATIONS =',I2)
      WRITE(K,6067)N4
87   FORMAT(/,5X,'MRAT =',I2)
      IF (MRAT.NE.0) WRITE(K,86)RATOR
86   FORMAT(/,5X,'RATER VALUES',/10X,10F12.5)
      CALL LICHCHIDEN,HW,IT,Q)
      CPRE=2.*RPPM(I)
      DO 7020 I=1,NBN
        IBDN=1
        MOR=0.0
        IF (MOD(I,2).EQ.0) MOR=2
        IF (I.EQ.NBN) MOR=0
        WRITE(K,8254)
8254  FORMAT('1',5X,'AXIAL BURNUP AVERAGES BY COLUMN POSITION',/T21,'J',
1T30,'BURNUP',/
        DO 8255 J=1,NCOL
          BI=0.
          BU=0.
          DO 8256 KJ=1,NFCW
            KK=NCOL*(KJ-1)+J
            IF (SAB(I0,KK).EQ.0) GO TO 8256
            BN=BN+1.
            BU=BU+SAB(I0,KK)
8256   CONTINUE
            IF (BN.EQ.0.150 TO 8258
            BU=BU/BN
8258   WRITE(K,8257)J,BU
8257   FORMAT(I20,I2,I30,F7.0)

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ABP01540
ABP01550

ABP01570
ABP01580
ABP01590
ABP01600
ABP01610
ABP01620
ABP01630

ABP01640

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8255 CONTINUE
      CPPM=3PPM(1)
      IF(1.EQ.1)GO TO 7037
      CPPM=ROPM(NOLF)
      IF(1.EQ.NRAN)GO TO 7037
      DO 7036 J=1,NOLF
        N=J
        IF((J-1)*DF2+1.E-6.GE.(1-1)*DF1)GO TO 7039
7036 CONTINUE
7039 CPPM=3PPM(N)-(3PPM(N)-3PPM(N-1))*((N-1)*DF2-(1-1)*DF1)/DF2
7037 N4J=N4
      IF (N4P .EQ. 2) N4J=1
      DO 700 NJ=1,N4J
        IF(NJ.EQ.1)GO TO 701
        IF(ABS(XKEFF-1.)*LT.3.0005)GO TO 702
        IF(NJ.GT.3)GO TO 703
        IF((1.GT.1).AND.(NJ.GT.2))GO TO 703
        COLD=CPPM
        SUM=0.
        DO 800 J=1,NCGRO
          SUM=SUM+CTA(J)
          CPPM=CPPM+(XKEFF-1.)/(SUM*SABP*2.)
          IF((1.EQ.1).AND.(NJ.EQ.1))CPPM=CPPM*0.95
          SVK=XKEFF
          GO TO 701
703 DC=CPPM-COLD
          COLD=CPPM
          CPPM=CPPM-DC*(XKEFF-1.)/(XKEFF-SVK)
          SVK=XKEFF
          IF(CPPM.GT.CPRE)CPPM=CPRE*0.95
701 DO 7038 J=1,NCCPO
7038 SAR(9,J)=CPPM
          CPN =CPPM/CCN3
          WRITE(K,7040)J,CPM
7040 FORMAT(//5X,'AT BURNUP INCREMENT NUMBER ',I2,', THE NATURAL BORON
          100 CONCENTRATION IS ',F8.0,' PPM.')

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ABP01650
ABP01660
ABP01670
ABP01680
ABP01690
ABP01700
ABP01710
ABP01720
ABP01730

ABP01750

ABP01780

ABP01820
ABP01830
ABP01840
ABP01850
ABP01860
ABP01870

ABP01880
ABP01890
ABP01900
ABP01910
ABP01920
ABP01930

ABP01990

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702 CPPF=CPPM
   IF(1.NE.1)GO TO 710
   IF(1.EQ.1)GO TO 710
   N01=NDLF-1
   DO 709 J=2,N01
     BPPM(J)=BPPM(NDLF)+(BPPM(J)-BPPM(NDLF))*(CPPM-BPPM(NDLF))/
     1 (BPPM(1)-BPPM(NDLF))
     BPPM(1)=CPPM
709 CONTINUE
710 CALL MAP(AROW,NCOL,WCCPDW,4)
   IF(1.EQ.NBN)GO TO 7022
   DO 7021 J=1,NCGR0
     FLU=0.
     IF(CK(J).EQ.0)GO TO 7021
     FLU=P2AV(J)*PFLUX*3.1536E7*PCF*AFLX
7021 TPAT(J)=FLU
7022 CALL RLFAE(SAB ,NCORG)
7022 CONTINUE
   ASITE(K,7599)
7909 FORMAT(1,'5X',AVERAGE CYCLE ORFICE RATIOS')
   NC2=NCCL-2
   DO 7910 I=1,NC2
     SUM=0.
     DO 7920 J=1,NBN
       7920 SUM=SUM+TOUT(J,I)
       7910 RATE(I)=(TOUT-TIN)/(SUM/NBN-TIN)
       WRITE(K,7599)(I,RATOR(I),I=1,NC2)
7908 FORMAT(//15X,13,7X,F10.5)
     IF(MODE.EQ.0)STOP
     WRITE(K,7599)
7599 FORMAT(//5X,RATEM DENSITIES BY LOCATION')
     DO 7601 I=1,NCGR0
       IF (MESH(I) .EQ. 0) GO TO 7601
       IF (MESH(I) .EQ. 2) GO TO 7601
       WRITE(K,7603)I
7603 FORMAT(//5X,'LOCATION' =,13/)
       WRITE(K,7602)(SAB(I),I,J=1,10)
       IF(MODE .LT.0)WRITE(M,5002)(SAB(J),I,J=1,10)
7601 CONTINUE
7602 FORMAT(5E20.8)
5002 FORMAT(15X,5E15.8)
     STOP
   END
C
SUBROUTINE XYTOR(SAZ,F25,F23,ARCWZ,NCOLZ,SB,SDU,NCYCLE)
  DIMENSION SB(10,10),SC(10,10),SAB(10,100),SAB(10,100),MESH(100),
  1MESH(100), F25(10,20), F23(10,20),SDU(10,10,20),XYMAP(100),
  2SAZ(10,100)
  C X-Y AVERAGING INTO DELTA-R RINGS.
  DIMENSION F(100,2),JC(10,20)
  DATA LI,LC,LP/5,6,7/
  DATA N,M/13,17,MAX/0/
  EQUIVALENCE (ACORC,NA)
  C N X N AS QUARTER SYMMETRY
  C N=1, CELL CENTER; N=2, CELL BOUNDARIES
  WRITE(LC,11)N,M
  11 FORMAT(// 5X,'ORDER' =,13,5X,'TYPE' =,12)
  NBN=N
  DO 50 I=1,NBN
    DO 50 J=1,2
      F(I,J)=0.
50

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ABP02010
ABP02020
ABP02030
ABP02040
ABP02050
ABP02060
ABP02070

ABP02080
ABP02090
ABP02100
ABP02110
ABP02120
ABP02130
ABP02140
ABP02150
ABP02160

ABP02170
ABP02180
ABP02190

ABP02210
ABP02220
ABP02230

ABP02250
ABP02260

ABP02270
ABP02280

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NC=NN-N+1
F(KC,1)=1.
F(KC,2)=3.1+159265/4.
Y=0.
PY=1.
IF(M.EQ.1)HY=0.5
DO 200 I=1,N
IR=I-1+1
X=0.
HX=1.
IF(M.EQ.1)HX=0.5
DO 100 J=1,N
KN=(IR-1)*N+J
IF(F(KN,1).NE.0.16C TO 99
PLN=SQRT(X*X+Y*Y)
R1=1.
IF(M.EQ.1)R1=0.5
DO 150 L=1,N
R2=R1+1.
IF(R2.GT.PLN)GO TO 151
R1=21+1.
GO TO 99
151 IF=L+1
PUF=SQRT((X+HX)**2+(Y+HY)**2)
IF(PUF.LE.R1)GO TO 95
IF((PLN.GE.R1).AND.(PUF.GT.R2))GO TO 101
FC=1.
GO TO 98
101 PUN=SQRT(X*X+(Y+HY)**2)
PLP=SQRT((X+HX)**2+Y*Y)
IF(PUN.GE.R2)GO TO 110
IF(PLP.LE.R2)GO TO 120
CASE B FOLLOWS:
ALPHA=SQRT(R2*R2-Y*Y)-X
GAMMA=SQRT(R2*R2-(Y+HY)**2)-X
FC=0.5*(ALPHA+GAMMA)/HX
GO TO 98
110 IF(PLP.GE.R2)GO TO 111
CASE A FOLLOWS:
BETA=SQRT(R2*R2-X*X)-Y
DELTA=SQRT(R2*R2-(X+HX)**2)-Y
FC=0.5*(BETA+DELTA)/HY
GO TO 98
C CASE A FOLLOWS:
111 ALPHA=SQRT(R2*R2-Y*Y)-X
BETA=SQRT(R2*R2-X*X)-Y
FC=0.5*ALPHA*BETA/(HX+HY)
GO TO 98
C CASE C FOLLOWS:
120 GAMMA=SQRT(R2*R2-(Y+HY)**2)-X
DELTA=SQRT(R2*R2-(X+HX)**2)-Y
FC=1.-0.5*(HX-GAMMA)/(HY-DELTA)/(HX+HY)
98 F(KN,1)=IR
F(KN,2)=FC
99 HX=1.
X=X+1.
IF((V.EQ.1).AND.(J.EQ.1))X=0.5
100 CONTINUE
HY=1.
Y=Y+1.

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IF((X.EQ.1).AND.(I.EQ.1))Y=0.5
200 CONTINUE
DO 750 I=1,N
DO 249 J=1,20
249 JS(I,J)=C
KN=1
DO 251 J=1,NN
IP=F(J,1)+.001
IF(IP.RE.1)GO TO 251
JS(I,KN)=J
KN=KN+1
251 CONTINUE
250 CONTINUE
IF(MX.EQ.0)GO TO 409
WRITE(C,19)
19 FORMAT(/5X,'ASSEMBLY FRACTIONS IN RADIAL ZONES',//Y18,'POSITION',
1T40,'75NE',T62,'FRACTION'/)
DO 410 I=1,NN
IF(F(I,1).EQ.0.)GO TO 410
J=F(I,1)+.001
WRITE(C,20)I,J,F(I,2)
410 CONTINUE
20 FORMAT(T20,I3,T40,I3,T60,F10.6)
409 CONTINUE
DO 695 I=1,NCORO
DO 699 J=1,10
SCB(J,I)=C.
699 SAB(J,I)=C.
READ(LI,5001)(MESH(I),I=1,NCORO)
C *** FOR CYCLE ONE USE ZEROS FOR FUEL POSITIONS OF NEW FUEL
5001 FORMAT(80I1)
DO 5005 I=1,NCORC
IF(MESH(I).NE.1)GO TO 5009
READ(LI,5002)(SCR(J,I),J=1,10)
5002 FORMAT(5X,F5E15.8)
5009 CONTINUE
8776 READ(LI,8776)(MESH(I),I=1,NCORO)
MCHK=0
NZZ=NCMZ
IF(MCYCLE.NE.1) NZZ=1
DO 8760 L=1,NZZ
DO 8750 I=1,NCORC
XYMAP(I)=IARS(MESH(I))
SAB(3,I)=SCR(8,I)
KI=MESH(I)
IF(KI)8749,8750,8748
8749 SAB(1,I)=F25(-KI,L)
SAB(3,I)=F28(-KI,L)
MESH(I)=1
DO TO 8750
DO 8743 J=1,10
8743 IF(J.EQ.0)GO TO 8744
SAB(J,I)=SCB(J,KI)
MCHK=1
8744 CONTINUE
8750 CALL ARAD(N, MESH,F,SAB,SB,J6)
DO 2727 LMN=1,10
DO 2727 LMN=1,10

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ABP01340
ABP01350

ABP01370
ABP01380

ABP01400

ABP01410

ABP01420
ABP01430
ABP01440

ABP01480
ABP01490
ABP01500

ABP01510
ABP01530

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2727 SUB(LYN,LMA,LJ)=SB(LMP,LMN)
CONTINUE
8760 CALL MAP(N,N,XYMAP,8)
      K1=(LG,710)
710  FORMAT(/5X,'ABOVE MAP IS X-Y CORE POSITIONING FOR FORMING RADIAL
      11CVES')
      IF(MCHK.EQ.0)RETURN
      NG=N
      DO 700 I=1,NROWZ
      DO 701 J=1,N
      KI=(I-1)*N+J
      DO 701 KK=1,10
      SC(KK,J)=SAZ(KK,KI)
701  CALL ARAC(NG,MESH,F,SCB,SC,JG)
      DO 709 JR=1,NCORC
      DO 709 JR=1,10
      SAB(JR,IR)=0.
709  DO 711 IR=1,NCORO
      SAB(8,IR)=SCB(8,IR)
      KI=MESHN(IR)
      IF(KI)712,711,713
712  SAB(1,IR)=F25(-KI,1)
      SAB(3,IR)=F28(-KI,1)
      GO TO 711
713  DO 714 JR=1,10
      IF(JR.EQ.8)GO TO 714
      SAB(JR,IR)=SCB(JR,KI)
714  CONTINUE
711  CALL ARAD(N,MESH,F,SAB,SC,JG)
      DO 715 J=1,N
      KI=(I-1)*N+J
      DO 715 KK=1,10
715  SAZ(KK,KI)=SC(KK,J)
700  CONTINUE
      SUM=0.
      SUMN=0.
      DO 721 J=1,NROWZ
      KI=(J-1)*N+1
      IF(SAZ(1,KI).EQ.0.)GO TO 721
      SUM=SUM+SAZ(1,KI)
      SUMN=SUMN+1.
721  CONTINUE
      IF(SUMN.EQ.0.)SUMN=1.
      SUM=SUM/SUMN
720  WRITE(LG,722)1,SUM,SB(1,1)
722  FORMAT(/5X,'U-235 CHECK FOR COLUMN',I3,5X,'P-Z AVG =',E13.6,
      5X,' X-Y AVG =',E13.6,'/5X,' X-Y AVG NOT VALID FOR AXIAL VARIATIO
      2NS IN FUEL ENRICHMENT')
      WRITE(LG,730)
730  FORMAT(1,'5X,'U-235 CORE MAP BEFORE CORRECTION, COLUMN 9 IS MOD
      #IFIED WITH WATER'//)
      AF1=NSCWX+1
      DO 740 I=1,NRI
      KI=(I-1)*NCOLZ+1
      K2=I*NCOLZ
740  WRITE(LG,731)(SAZ(1,J),J=K1,K2)
731  FORMAT(/5X,10E12.3)

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```

      RETURN
    END
    SUBROUTINE TUPR(DLT,SXC,MESH,DIN,DEA)
      DATA KKK/6/
      DIMENSION DLT(190),SXC(190),T(10),MESH(190)
      COMMON/CORE/ACTA(190),ACTF(190),WCCPOK(190),CK(190),RE(190),
      1ELSCU(190),TAUF(190),SIGFI(190),SIGF2(190),P2AV(190),RCRO,
      2ZFLX, XKEFF,DL2,SAC,XNUP,XNUT(190),EPQ
      COMMON/CORE2/NRCW,NCCL,NSYN,JMAX,H,DI,DZ,HIGHT,ALPHA,CCRIT,
      1CCNPOK,PIWATT,VOLMF,VOLUV,VI(10),FCFRAC
      COMMON/TP/SAC1,SAC2,DENHOH,XMDOOT,TIN,PAIOR(10),MRAT,POUT(30,10),
      1IBURN,TBAR
      TCUT=0.
      VA=0.
      VG=0.
      TAVG=0.
      NCL=NCCL-1
      CALL LICHCH(DIN,HIN,TIN,1)
      CALL LICHCH(DIN,HIN,TIN,4)
      HAVG=CCNPOK*3.413/XMDOOT
      HG=HAVG*HIN
      CALL LICHCH(DG,HG,TBAR,3)
      WRITE(KKK,27)TBAR
      FORMAT(//5X,'OVERALL AVG. EXIT TEMP. =',F7.1)
      DO 100 I=1,NCL
        T(I)=TIN
      100 I=1,NCL
      HZ=HIN
      CCN=HAVG/(NRCW-4)
      IF (MRAT.NE.0) CCN=CCN*RATOR(I)
      DO 200 J=1,NRCW
        JP=NRCK-J+1
        K=NCCL*(JP-1)+1
        DLT(K)=DL2
        SXC(K)=SAC
      200 K=1,NCL
      IF (MESH(K).EQ.0) GO TO 200
      IF (MESH(K).EQ.2) GO TO 200
      HZ=HZ+CCN*WCCPOK(K)
      CALL LICHCH(DZ,HZ,TZ,2)
      DI=DZ/DIN
      DLT(K)=DL2/DI
      SXC(K)=SAC1+SAC2*DI
      CALL LICHCH(DZ,HZ,T(I),3)
      TAVG=TAVG+T(I)*V(I)
      VA=VA+V(I)
      CONTINUE
      POUT(1BURN,I)=T(I)
      IF (I.EQ.NCL) GO TO 100
      TCUT=TCUT+T(I)*V(I)
      VC=VC+V(I)
      CONTINUE
      WRITE(KKK,25)T(I),I=1,NCL
      FORMAT(// 5X,'OUTLET TEMPERATURES',/(10F11.1))
      TCUT=TCUT/VC
      WRITE(KKK,23)TCUT
      25 28 26
      FORMAT(//5X,'AVERAGE EXIT TEMPERATURE =',F7.1)
      TAVG=TAVG/VA
      CALL LICHCH(DEN,HA,TAVG,4)
      DENHOH=DEN/62.4
      WRITE(KKK,26)DENHOH,TAVG
      FORMAT(//5X,'AVERAGE CORE WATER SPECIFIC GRAVITY =',F7.4)

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15X,AVERAGE CORE WATER TEMPERATURE =,F7.1)
WRITE(KKK,50J)
600 FORMAT(//5X,'CALCULATED CRFICE RATIOS (IN TO OUT)://)
DO 601 I=1,NC1
PT=0.
IF(CRUIT(I)-TIN).LT-.0.1 )GC TO 601
PT=(TEAR-TIN)/(T(I)-TIN)
601 WRITE(KKK,502)I,PT
602 FORMAT(I10X,I3,5X,F8.4)
RETURN
END
SUBROUTINE RZFLUX(MODE,MKEY,MPCRP)
***** 2-0, R-Z, 2 GROUP CORE PROGRAM *****
REAL*4 LCCPCN
DATA KKK/6/
DIMENSION P1(190),P2(190),OLD(190),PIA(190),PZA(190)
COMMON/COPEZ/NROW,NCOL,NSYM,JMAX,H,D1,D2,HIGHT,ALM1,CCRIT,
ICCNPCN,FINATT,VOLMF,VOLUM,V(10),FCFRAC
COMMON/CCRE/RECTA(190),RECTF(190),LCCPCN(190),C(190),P(190),
IELSQ(190),TAU(190),SIGFI(190),SIGFI2(190),PZU(190),NCORO,
2RZFAG,XKEFF,BL2,SAC,XNUF,XNUT(190),EPQ
EQUIVALENCE (LCCPCN(1),OLD(1))
C
C CODE IS: C IS K-INF.
C P IS RESCANCE ESCAPE PROBABILITY.
C ELSQU IS DIFFUSION LENGTH SQUARED.
C TAU IS TAU THERMAL.
C SIGFI-J IS SIGMA FISSION FOR FAST, J=1, AND THERMAL,
C J=2, GROUPS.
C R-Z GEOMETRY WITH DELTA-R = DELTA-Z = H
TCD=4.0
IF (MODE .EQ. 2) GC TO 1126
HS=H*H
IF (MPCRP.NE.0)CALL MAP(NROW,NCOL,C,7)
WRITE(KKK,801)
861 FORMAT(//5X,'FLUX CALCULATION'///)
DO 49 I=1,NCORO
CLO(I)=0.
SET=1.0
IF(P(I).EQ.0.)SET=0.
P1(I)=SET
P2(I)=SET
PIA(I)=SET
PZA(I)=SET
49
C ***
C NSYM IS SYMMETRY KEY
C NSYM = 0: NO SYMMETRY
C IF NSYM = 1,2: MAP IS UPPER CORE SO SYMMETRY ABOUT Z=0
C IF NSYM = 1, SYMMETRY IS THROUGH CELL BOUNDARY
C IF NSYM = 2, SYMMETRY IS THROUGH CELL CENTER OR MESH POINT
ZNCORM=1.0
SUM1=0.
SUM2=0.
NCO=NSYM+1
ALPHA=1.0
XKEFF=1.0
EVP=0.01
ALT=1.0
DO 123 J=1,JMAX
ANUM=0.

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ABP02320

ABP02390
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      ERVAV=C.
      RI=J
      AK=0
      ERVAV=0.
      NA=NCCL*1
      NP=CCGRJ-NCCL
      IF(NSYM.NE.O)NR=NCGRQ-1
      SWEEP THE DIFFUSION EQUATIONS
      DO 113 I=NA,NR
        IF(I.EQ.0.1)GO TO 113
        PP=DI*HS*P(I)/(D2*TAU(I))
        S=1.0/(TCD+HS/TAU(I))
        TC=1.0/(TCD+HS/ELSCU(I))
        CC=02*HS*C(I)*ZNRN/(DI*ELSCU(I)*P(I))
        IP=I+NCCL
        IM=I-NCCL
        IWL=I-1
        IPI=I+1
        GO TO (109,609,608),NQC
      609 IF(IP.GT.NCGRQ)IP=I
        GO TO 109
      608 IF(IP.GT.NCGRQ)IP=IM
      108 ICR=NCCL(I-1,NCCL)
      CIP1=1.+1./I2*IDR
      CIM1=1.-1./I2*IDR
      GO TO 1081
    1080 CIP1=2.
      CIM1=0.
    1081 CONTINUE
      P1(I)=(CC*P2(I)+CIP1*P1(IPI)+CIM1*P1(IWL)+P1(IPI)+P1(IWL))*S-P1(I)
      I*ALPHA+P1(I)
      P2(I)=(PP*P1(I)+CIP1*P2(IPI)+CIM1*P2(IM1)+P2(IPI)+P2(IM1))*TC
      1-P2(I))*ALPHA+P2(I)
      ACUJAT CALCULATIONS FOLLOW.
      P1A(I)=(P2*P2A(I)+CIP1*P1A(IPI)+CIM1*P1A(IM1)+P1A(IPI)+P1A(IM1))*S
      1-P1A(I))*ALPHA+P1A(I)
      P2A(I)=(CC*P1A(I)+CIP1*P2A(IPI)+CIM1*P2A(IM1)+P2A(IPI)+P2A(IM1))*TC
      1-P2A(I))*ALPHA+P2A(I)
      IF(SIGF12(I).EQ.0.1)GO TO 113
      ERR=ABS(P2(I)-OLD(I))/P2(I)
      IF(ERR.GT.CCRIT)NR=NR+1
      ERVAV=AMAX1(ERVAV,ERR)
      ANUM=ANUM+1.
      ERVAV=ERVAV+ERR
      CLO(I)=P2(I)
    113 CONTINUE
      IF(-ALPHA.GT.-1.)JERM=0.005
      IF(J.EQ.1)GO TO 115
      IF(AMVIERAV/ESV).GT.1.)GO TO 115
      ALTN=2.0/(1.+SQRT(1.-(ERVAV/ESV)**2))
      IF(ABS(1.-ALTN/ALT).LT.ERM) ALPHA=AMIN1(ALM1,ALTN)
      ALT=ALTN
      ESV=ERVAV
    115 ERVAV=ERVAV/ANUM
      IF(NK.EQ.0)GO TO 125
      IF(MD(J-60).EQ.0)ALPHA=ALPHA-0.1
      ALPHA=AMAX1(ALPHA,1.0)
    119 IF(J-1) 123,120,1120
    120 DO 121 I=1,NCGRQ

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ABP02800
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ABP02990
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 ABP03020

ABP03040

ABP03070

ABP03120
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 ABP03140
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 ABP03160
 ABP03170
 ABP03180
 ABP03190
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 ABP03310
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SUM1=SUM1+P1(I)+P2(I)
121 CONTINUE
GO TO 123
1120 DO 1121 I=1,NCGR
SUM2=SUM2+P1(I)+P2(I)
1121 CONTINUE
ZNCGR=ZNCGR*SUM1/SUM2
XKEFF=XKEFF*SUM2/SUM1
SUM1=SUM2
SUM2=0.
IF(MOD(I,10).EQ.0)WRITE(KKK,2350)J,ERMAX,ERAV,ALPHA
2350 FORMAT(5X,1ITERATION NC. =,I3.5X,MAXIMUM ERROR =,E11.4,5X,
1,AVERAGE ERROR =,E11.4,5X,ALPHA =,F9.6)
IF(XKEFF.GT.0.1GC TO 123
WRITE(KKK,5405)XKEFF,J
5405 FORMAT(//5X,7HK-EFF =,F11.6,5X,3HJ =,I4,5X,8HTERMINATE )
STOP
123 CONTINUE
125 WRITE(KKK,127)JMAX,NIT,ERMAX,CRIT
127 FORMAT(//5X,28HMAXIMUM NUMBER OF ITERATIONS ,I4,18H, ITERATIONS
1USED ,I4,19H TO RELATIVE ERROR ,E10.4,21H. ERROR CRITERIA IS ,
2E10.4)
C
C CALCULATE POWER PRODUCING VOLUME
C
1126 POWER=C.
PIFAG=0.
P2FAG=0.
VOLUME=0.0
DO 1110 I=1,NCGR
LOCPOW(I)=0.
IF(SIGF12(I)) 1110,1110,1221
1221 FOR=LOC(I,NCOL)
VOL=V(1DR)
GO TO 1223,1229,1232,NQO
1232 IF(1.GT.NCGR-NCOL)VOL=VOL*0.5
1220 VOLUME=VOLUME + VOL
LOCPOW(I)=(SIGF11(I)*P1(I)+SIGF12(I)*P2(I))*VOL
POWER=POWER+LOCPOW(I)
PIFAG=PIFAG+P1(I)*VOL
P2FAG=P2FAG+P2(I)*VOL
P2O(I)=VOL
1110 CONTINUE
APD=CFNPGW/VOLUME
WRITE(KKK,111)VOLUME,APD
1111 FORMAT(//5X,8HVOLUME =E11.4,5X,14HPCW/AVOLUME =E11.4//)
POWER=POWER/PIFAG
ENCRM = CCNPGW/POWER
PIFAG=PIFAG*ENCRM/VOLUME
P2FAG=P2FAG*ENCRM/VOLUME
C
C FLUX NORMALIZATION TO CONSTANT POWER
C
DO 131 I=1,NCGR
P1(I) = P1(I)*ENCRM
P2(I) = P2(I) * ENCRM
PIA(I)=PIA(I)*ENCRM
P2A(I)=P2A(I)*ENCRM
IF(LOCPOW(I).EQ.0)GC TO 131
J=MOD(I,NCOL)

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ABP03330
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      LOCPOW(I)=LOCPOW(I)*VOLUM/(PDMR*V(J))
131  CONTINUE
C
C      CALCULATE POWER AT MESH POINTS
C
      SWAX=0.
      DO 2160 I=1,NCORC
      IF(LOCPOW(I)-LT-SMAX)GO TO 2160
      SWAX=LOCPOW(I)
      KPCW=I
2160  CONTINUE
      J=MOD(KPCW,NCOL)
      JR=KPCW/NCOL
      IF(J.NE.0)JR=JR+1
      IF(J.EQ.0)J=NCOL
      WRITE(KKK,2161)JR,J,SWAX
2161  FORMAT(//5X,4HRW,12,9H, COLUMN ,12,46H, IS POSITION OF MAXIMUM
      LOCAL POWER OF VALUE ,F8.5//)
      WRITE(KKK,2133)PIFAG,P2FAG
2133  FORMAT(//5X,1AVERAGE FAST FLUX =,E12.5,10X,1AVERAGE THERMAL FLUX
      I=,E12.5)
      WRITE(KKK,2505)XKEFF
2505  FORMAT(//2X,7HK-EFF =,F11.6)
      IF(MODE.NE.1)GO TO 2590
      MAP KEY PRINTOUT
      MKEY=0, LOCPCW PRINTOUT.
      MKEY=1, P1 PRINTOUT.
      MKEY=2, P2 PRINTOUT.
      MKEY=3, P1 & P2 PRINTOUT.
      MKEY=4, P1 & P2 & LOCPCW PRINTOUT.
      MKEY=0, PUNCH P1 & P2.
      IF(MKEY.NE.0)GO TO 1175
      CALL MAP(NROW,NCOL,LOCPCW,4)
      GO TO 2590
1175  IF (IABS(MKEY).NE.2)CALL MAP(NROW,NCOL,P1,3)
      IF (IABS(MKEY).GE.2)CALL MAP(NROW,NCOL,P2,2)
      IF (IABS(MKEY).EQ.4)CALL MAP(NROW,NCOL,LOCPCW,4)
      IF (IABS(MKEY).NE.2)CALL MAP(NROW,NCOL,P1A,-3)
      IF (IABS(MKEY).GE.2)CALL MAP(NROW,NCOL,P2A,-2)
2590  SUMD=0.
      C *** SYMMETRY CONDITIONS ON MESH DO NOT APPLY TO PERTURBATION
      C *** COEFFICIENTS SINCE THEORY REQUIRES A FULL CORE WITH VACUUM
      C *** BOUNDARY CONDITIONS.
      DO 3001 I=1,NCORC
      P2U(I)=P2(I)/P2FAG
      IF(P1(I).EQ.0.)GO TO 3001
      J=MOD(I,NCOL)
      SUMD=SUMD+(XNUF*SIGF1(I)*PIA(I)*PI(I)*XNUT(I)*SIGF12(I)*PIA(I)*
      1 P2(I)+V(J))
3001  CONTINUE
      DO 3002 I=1,NCORC
      RCTF(I)=0.
      RCTA(I)=0.
      IF(P1(I).EQ.0.)GO TO 3002
      J=MOD(I,NCOL)
      RCTA(I)=(P2A(I)*P2(I)+V(J)/SUMD)*XKEFF*XKEFF
      RCTF(I)=(PEPQ-1.)*P1(I)+P2(I)*PIA(I)*XNUT(I)*XKEFF*V(J)/SUMD
3002  CONTINUE
      IF(MKEY.NE.0)CALL MAP(NROW,NCOL,RCTA,5)
      IF(MKEY.NE.0)CALL MAP(NROW,NCOL,RCTF,6)

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 ABP0447C
 ABP0448C
 ABP0449C
 ABP0450C
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ABP0454C
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 ABP0457C
 ABP0458C

ABP0461C
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85      RETURN
      END
      SUBROUTINE NUDAR(SA9,NCOROI)
      DIMENSION SA3(10,199)
      COMMON/PAF/PA1,P,ET5,ET6,ET7,ET8,ET9,VNU5,VNU6,VNU7,VNU8,TAU
      COMMON/AXSEC/XT(2,4),VNU(9),CAM15,CAMX5,GAN19,ALAMX,SA5,GAMP5,
      IGAMP9,VCELL
      EQUIVALENCE (XT(1,1),SA3),(XT(1,2),SA5),(XT(1,3),SC8),
      IXST(1,4),SA9),(XT(1,5),SAU),(XT(1,6),SAL),(XT(1,9),SABP),
      ZXST(2,1),SF5),(XT(2,4),SF9),(XT(2,6),SFI)
      SC9=SA5-SF5
      SC9=SA5-SF9
      SC9=SAU
      C9S=EPS*PI*(1.-P)*SA5*ET5
      C9P=(EPS*PI*(1.-P)*SA9*ET9)-SA9
      CAS=-SC9/C9P
      CAS=1./((SA9+C9P)
      C10S=SC9/SAU
      C12S=SC9/(SA5-SA0)
      C14S=SC9/SAL
      C15S=SC9/(C9P+SA1)
      C16S=SC9/(SAL-SA5)
      C17S=SC9/(SAL-SA0)
      C22S=(1-2.*SFI)/SAU
      C23S=(1-2.*SFI)/SAL
      DO 200 I=1,NCOROI
      IF (TPAF(I).EQ.0.16C) TO 200
      T=TPAF(I)
      A5I=SA8(1,I)
      A6I=SA8(2,I)
      A3I=SA8(3,I)
      A9I=SA8(4,I)
      A0I=SA8(5,I)
      A1I=SA8(6,I)
      A7I=SA8(7,I)
      A8I=SA8(9,I)
      SAB(9,I)=API*EXP(-SA3P*T)
      SAB(1,I)=A5I*EXP(-SA5*T)
      SAB(2,I)=((A5I*SC9)/(SA5-SA3))*((EXP(-SA5*T))-EXP(-SA6*T))+A6I*EXP(-SA6*T)
      C5P=C9PS*AS1
      C8=C9S*ARI
      C6=C6S*C5P
      SA4(4,I)=C8*(1.-EXP(C9P*T))+C6*(EXP(C9P*T))-EXP(-SA5*T))+A9I*EXP(CAS*P5950)
      I=PST
      C10=C10S*CS
      C11=(SC9*(C6+A9I-C8))/(C9P+SA0)
      C12=C12S*CS
      C13=A5I*(C11+C12)
      SAB(5,I)=C10+C11*EXP(C9P*T)+C12*EXP(-SA5*T)+C13*EXP(-SA0*T)
      C14=C14S*PI
      C15=C15S*C11
      C16=C16S*C12
      C17=C17S*C13
      C18=A1I*(C14+C15+C16+C17)
      SAB(6,I)=C14+C15*EXP(C9P*T)+C16*EXP(-SA5*T)+C17*EXP(-SA0*T)+C18*EXP(-SA0*P5950)
      I=P(-SA1*T)
      C22=C22S*C17
      C23=C23S*C13
      C25=1.0-EXP(-SA5*T)

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F5=2.0*(A51*SF5+C25/SA5
F6=2.0*SF6*(C87-C6*C25/SA5+(C8-C6-A91)*(1.-EXP(C5P*T))/C9P)
F1=2.0*SF1*(C14*I+C16*C25/SA5)-C23*(1.-EXP(-SA1*T)) -2.0*SF1*C15ADP06140
I=(1.-EXP(C9P*T))/C9P-C22*(1.-EXP(-SAQ*T))
C16=(EPS-1.)/(VNU8-1.)
F5=C16*(VNU5*F5-VNU5*F9+VNU1*F1)
SAB(7,1)=F5+F9+1+AF1
SAB(3,1)=ARI*(1.-SGC*T)-0.5*(VNU5*EPS*PI*(1.-P)+F5+VNU9*EPS*PI*(1A3P)6150
1.-P)+F5+P8)
IF(AFI.NE.0.160 TO 160
IF(A51+AE1.EQ.0.160 TO 200
SAB(10,1)=(SAB(7,1)/(2.0*(A51+481)))*9.E5
GO TO 200
CONTINUE
RETURN
END
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IF(NB.EQ.0)GO TO 155
DO 1912 K=J,L
  IF(XAR(K).LT.-J.JXAR(K))=0.
  IF(XAR(K).NE.0.)XAR(K)=XAR(K)+0.0005
1912 CONTINUE
  WRITE(KKK,1913)(XAR(K),K=J,L)
1913 FORMAT(//3X,15F7.3)
  GO TO 192
195 DO 196 K=J,L
196 XVAR(K)=XAR(K)+0.1
  WRITE(KKK,197)(XVAR(K),K=J,L)
197 FORMAT(//3X,15F4)
192 J=J+NUM2
  L=J+NUM1-1
193 CONTINUE
  KCM=KCN+1
  IF(NUM22) 2550,2550,2910
2910 IF (KCN-2) 2920,2550,2950
2920 J=NUM21+1
  L=NUM42
  NUM2=NUM22
  WRITE(KKK,171)
  GO TO 191
2950 CONTINUE
  IF(NB.LT.0)WRITE(KKK,170)
  170 FORVAL(17)FORADJOINT SOLUTION)
  GO TO (1560,1962,1964,1966,1967,1968,1969,1970),NO
1960 WRITE(KKK,174)
  RETURN
1962 WRITE(KKK,175)XNCRM
  RETURN
1964 WRITE(KKK,176)XNCRN
  RETURN
1966 WRITE(KKK,177)XNDRM
  RETURN
1967 XNCRM=-XNCRM
  WRITE(KKK,178)XNCRM
  RETURN
1968 WRITE(KKK,179)XNDRM
  RETURN
1969 WRITE(KKK,1790)XNDRM
  RETURN
1970 WRITE(KKK,1791)
  RETURN
END
SUBROUTINE LTOHRO(D,H,T,MT)
  DATA L,K,H/5,5,117
  DIMENSION P(3,11),G(3)
  IF(MT.GT.0)GO TO 100
  READ(1,10)G,(P(J,I),J=1,3,I=1,N)
  FORMAT(24F10F10.3)
  10 J=1, TEMPERATURE: J=2, SPECIFIC VOLUME: J=3, ENTHALPY.
  C WRITE(K,110),G,(P(J,I),J=1,3,I=1,N)
  11 FORMAT(//5X,'WATER THERMODYNAMIC DATA'/18,3A4,I40,'TEMPERATURE',
  1760,'SPECIFIC VOLUME',180,'ENTHALPY'/(143,F5.0,I63,F8.5,I82,F5.1)
  2)
  RETURN
  C
  C MT=1, T-IN, H-OUT
  C MT=2, H-IN, D-OUT
  C MT=3, P-IN, T-OUT

```

ARP05090
 ARP05100
 ARP05110
 ARP05120
 ARP05130
 ARP05140
 ARP05150
 ARP05160
 ARP05170
 ARP05180
 ARP05190
 ARP05200
 A=005210
 ARP05220
 ARP05230
 ARP05240
 ARP05250
 ARP05260
 ARP05270
 ARP05280
 ARP05290
 ARP05300
 ARP05310
 ARP05320
 ARP05330
 ARP05340
 ARP05350
 ARP05360
 ARP05370
 ARP05380
 ARP05390
 ARP05400
 ARP05410
 ARP05420
 ARP05430
 ARP05440
 ARP05450
 ARP05460
 ARP05470
 ARP05480
 ARP05490
 ARP05500
 ARP05510
 ARP05520
 ARP05530

```

C      MT=4, T-IN, D-CUT
100 GO TO (101,102,103,104),MT
101 J1=1
    J2=3
    PX=T
    GO TO 110
102 J1=3
    J2=2
    PX=H
    GO TO 110
103 J1=3
    J2=1
    PX=H
    GO TO 110
104 J1=1
    J2=2
    PX=T
    GO 111 I=1,N
110 IF(PX-LE-P(J1,I))GO TO 112
111 CONTINUE
    WRITE(*,113)PX,P(J1,I),MT
113 FORMAT(5X,'EXTRAPOLATION, VALUE =',E11.4,2X,'LAST TABLE VALUE =',
112 E11.4,2X,'TYPE =',I2)
    I=I-1
    IF(I-GE-N-I+1)N=2
    IF(I-EG-0)I=1
    F1=(P(J2,I+1)-P(J2,I))/(P(J1,I+1)-P(J1,I))
    F12=(P(J2,I+2)-P(J2,I+1))/(P(J1,I+2)-P(J1,I+1))
    F2=(F12-F1)/(P(J1,I+2)-P(J1,I))
    A=P(J2,I)+PX-P(J1,I)*(F1+(PX-P(J1,I+1))*F2)
    GO TO (201,202,203,202),MT
201 H=1
    GO TO 210
202 D=1./A
    GO TO 210
203 T=A
210 RETURN
    END
SUBROUTINE ARADIN, WESH,F,SAB,SR,JG)
    DIMENSION WESH(100),F(100,2),SAB(10,100),SR(10,10),SN(10,10)
    DIMENSION JG(10,20)
    NN=NA
    IF(N.LT.0)GO TO 800
    DO 598 I=1,N
    DO 598 J=1,10
    SN(J,I)=0.
    SR(J,I)=0.
    DO 700 I=1,NN
    IF(MESH(I).NE.1)GO TO 700
    IF(F(I,1)-EQ-0)GO TO 700
    F=F(I,1)+001
    DO 701 J=1,10
    IF(SAB(J,I)-EQ-0)GO TO 701
    SS(J,IR)=SR(J,IR)+F(I,2)*SAB(J,I)
    SN(J,IR)=SN(J,IR)+F(I,2)
    IF(IR-EG-NN)GO TO 701
    SS(J,IR+1)=SR(J,IR+1)+(1.-F(I,2))*SAB(J,I)
    SN(J,IR+1)=SN(J,IR+1)+1.-F(I,2)
    CONTINUE
701
703 CONTINUE

```

```

DO 702 J=1,10
DO 702 I=1,N
IF(SN(J,I).EQ.0.)SN(J,I)=1.
702 SR(J,I)=SK(J,I)/SN(J,I)
RETURN
M=-N
DO 801 J=1,NN
DO 801 I=1,10
801 SAB(J,I)=0.
DO 810 L=1,M
DO 811 K=1,20
KJ=JG(L,K)
IF(KJ.EQ.0)GO TO 810
DO 812 J=1,10
SAR(J,KJ)=SAB(J,KJ)*SR(J,L)*F(KJ,2)
IF(L.NE.M)SAB(J,KJ)=SAB(J,KJ)+SR(J,L+1)*(1.-F(KJ,2))
812 CONTINUE
811 CONTINUE
810 CONTINUE
RETURN
END
SUBROUTINE CPROP(SAB,MESH,IXEUSE)
DATA KKK/6/
DIMENSION SAB(10,100),MESH(100),DLT(100),SXC(100)
COMMON/CPRE2/PCFA(100),PCFE(100),KOCPOW(100),CK(100),RE(100),
IELSQ(100),TAUF(100),SIGFI(100),SIGFIZ(100),P2AV(100),NCORO,
ZAFIX, XKEFF,DL2,SAC,XNUF,XNUT(100),EPQ
COMMON/XSECT/XST(2,0),VNUT(9),GAMIS,GAMX5,GAM19,ALAMX,SAX,GAMP5,
IGAMP9,VCELL
COMMON/PAPA/EP5,P1,P,ET5,ET9,TPAF(100),VNUS,VNUT,VNUS,TAU
COMMON/CPRE2/NROK,NCL,NSYM,JMAX,H,D1,D2,HIGHT,ALPHA,CCRIT,
ICCPDOW,FINATI,VOLMF,VOLUP,V(10),FCFRAC
CALL
TMPR(DLT,SXC,MESH,DIN,DAVG)
EPQ=EPS
XNUF=VNUT8
CAVG=0.
CNUM=0.
DO 6001 I=1,NCORO
IF(VOLUP(I).EQ.0) DLT(I)=0.0
TAUM=TAU*(DLT(I)/DL2)*2
XNUT(I)=0.
L=MESH(I)
C *** L=1, REGULAR; L>1, WATER; L=0, VACUUM.
IF(L-1)61,60,59
59 IF(L.GT.2)GO TO 60
CK(I)=0.0
SIGFI2(I)=0.
SIGFI1(I)=0.
RE(I)=1.
TAUF(I)=TAUM
IELSQ(I)=DLT(I)
IF(SAB(9,I).NE.0.)IELSQ(I)=D2/(XST(1,9)*SAB(9,I)/VOLMF+D2/DLT(I))
GO TO 6001
61 CK(I)=0.
SIGFI2(I)=0.
SIGFI1(I)=0.
RE(I)=0.
TAUF(I)=0.
IELSQ(I)=0.
GO TO 6001

```

ABP06310

ABP06350
ABP06360

ABP06380

ABP06400
ABP06410
ABP06420
ABP06430
ABP06440ABP06450
ABP06460
ABP06470
ABP06480ABP06500
ABP06510
ABP06520ABP06560
ABP06570
ABP06580
ABP06590
ABP06600
ABP06610
ABP06620
ABP06630

```

60 ASIGF=C.
  ANUM=C.
  ADEME=C.
  ADEMB=C.
  ADEMA=C.
  FLUX=FLXW22AV(I)
  IF (FLUX.EQ.0.) FLUX=SAB(8,I)
  DO 6915 K=1,6
    ASIGF=ASIGF+XST(2,K)*SAB(K,I)
    ANUM=ANUM+VNU(K)*SAB(K,I)+XST(2,K)
    ADEME=ADEME+SAB(K,I)*XST(1,K)
    IF (FLUX .LE. 0.0) SIGAX = 0.0
  IF (FLUX .GT. 0.0)
    1 SIGAX=((GAM15+GAMVX5)*XST(2,1)*SAB(1,I)+GAM19*XST(2,4)*SAB(4,I))*
    2 FLUX/(ALAMX/SAX+FLUX)
    SIGAS=GAMPS*XST(2,1)*SAB(1,I)+GAMP9*XST(2,4)*SAB(4,I)
C *** SIGAS IS EQUILIBRIUM SM-149 MACRO CROSS SECTION FOR ABS.
C *** SINCE SAB(3,I) NOT NEEDED FOR SM-149, SET IT TO FLUX
  IF (IXDUSE .NE. 0) GO TO 22
  SIGAX=C.0
  SIGAS=C.0
  22 CONTINUE
  SAB(8,I)=FLUX
  DO 6916 K=7,9,2
    ADEME=ADEME+XST(1,K)*SAB(K,I)
    ADEME=ADEME+SIGAX+SIGAS
    TU=ADEMB/(ADEMB+ADEME+SXC(I))
    ELSQU(I)=CLT(I)*(1.-TU)
    CX(I)=TU*EPS*P*ANUM/ADEMB
    SIGF12(I)=ASIGF/VCELL
    SIGF11(I)=((ANUM/VNU8)*{EPS-1.}/VCELL)
    FC(I)=P
    TAUF(I)=TAUM
    XROT(I)=ANUM/ASIGF
    IF (L.LE.2) GO TO 6002
    SIGF11(I)=SIGF11(I)*FCFRAC
    SIGF12(I)=SIGF12(I)*FCFRAC
    RET(I)=1.-FCFRAC+FCFRAC*RET(I)
    CX(I)=CX(I)*FCFRAC
    ELSQU(I)=CLT(I)*(1.-FCFRAC)+FCFRAC*ELSQU(I)
    CAVG=CAVG+CX(I)
    CNUM=CNUM+1.
  6001 CONTINUE
  CAVG=CAVG/CNUM
  IF (CAVG.LI.0.75) GO TO 70
  IF (CAVG.GT.VNU5100 TO 7C
  WRITE(KKK,71)CAVG
  71 FORMAT(/5X,'CROSS CALCULATION: K-AVG =',F8.5)
  RETURN
  70 WRITE(KKK,72) CAVG
  72 FORMAT(/5X,'TERMINATION ON OUT OF BOUNDS FOR K-AVG =',E15.6)
  STOP
  END

```

ARP06640
ARP06650
ARP06660
ARP06670
ARP06680
ARP06690
ARP06700
ARP06710
ARP06720
ARP06730

ARP06740
ARP06760
ARP06770
ARP06780

ARP06790
ARP06800
ARP06810
ARP06820

ARP06850

ARP06880
ARP06900

ARP06920
ARP06930
ARP06940

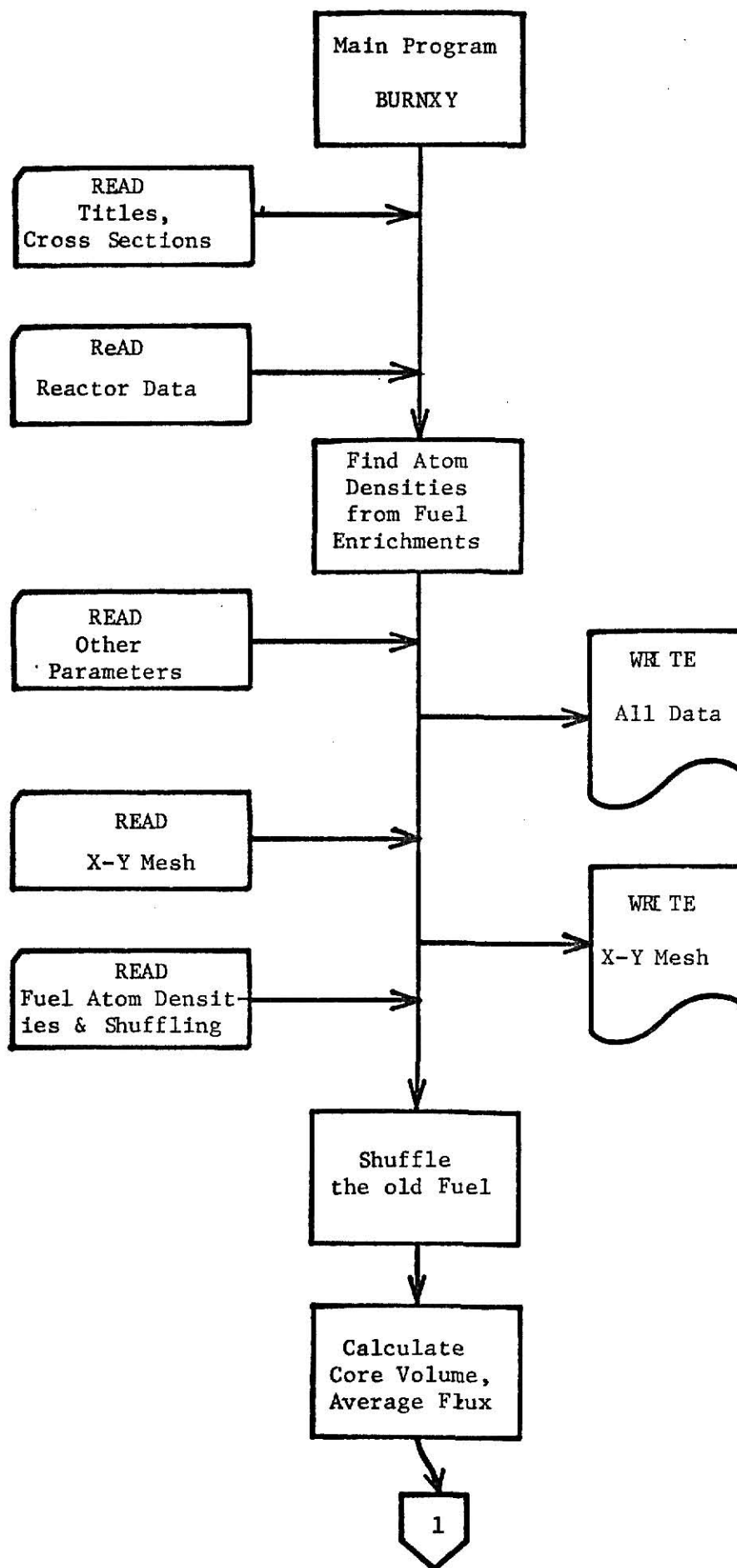
ARP06960
ARP06970
ARP06980
ARP06990
ARP07000
ARP07010
ARP07020
ARP07030

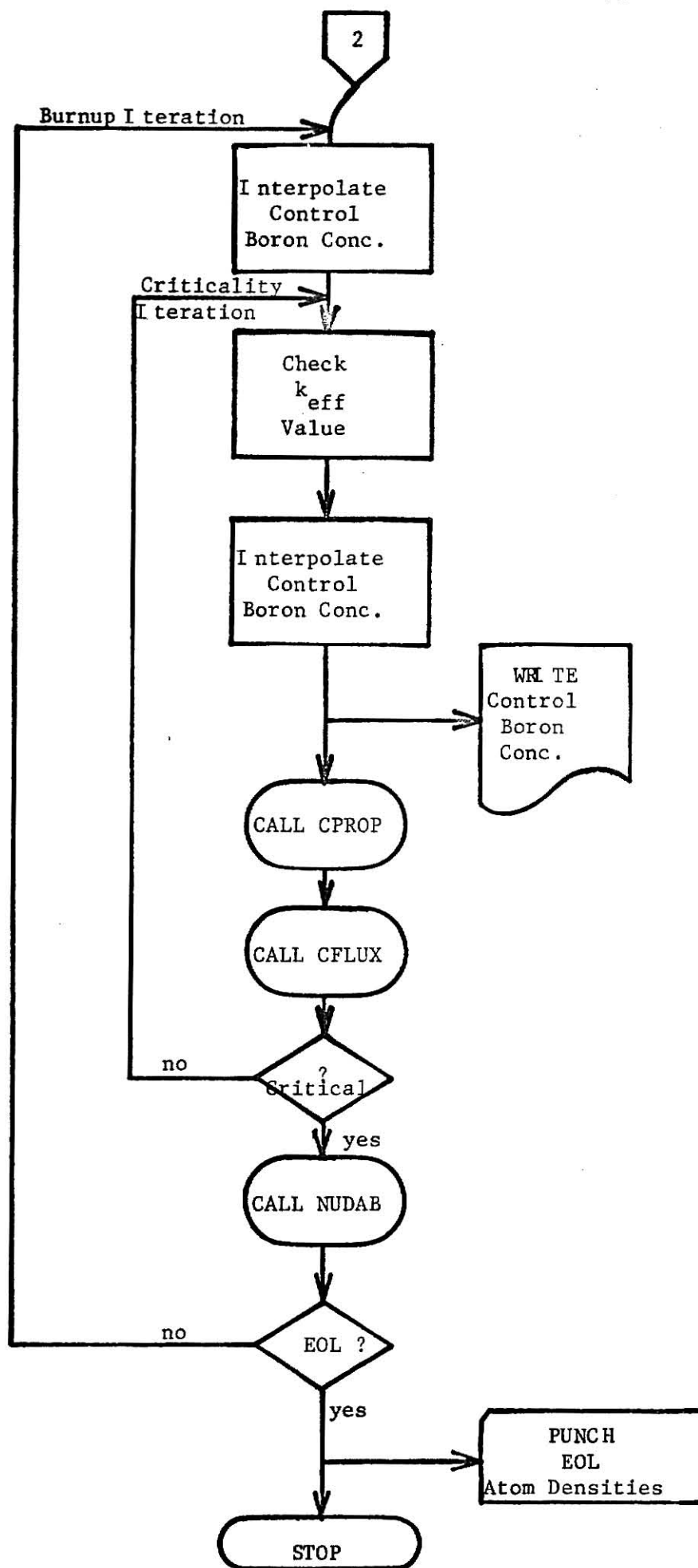
APPENDIX B: THE CODE BURNXY

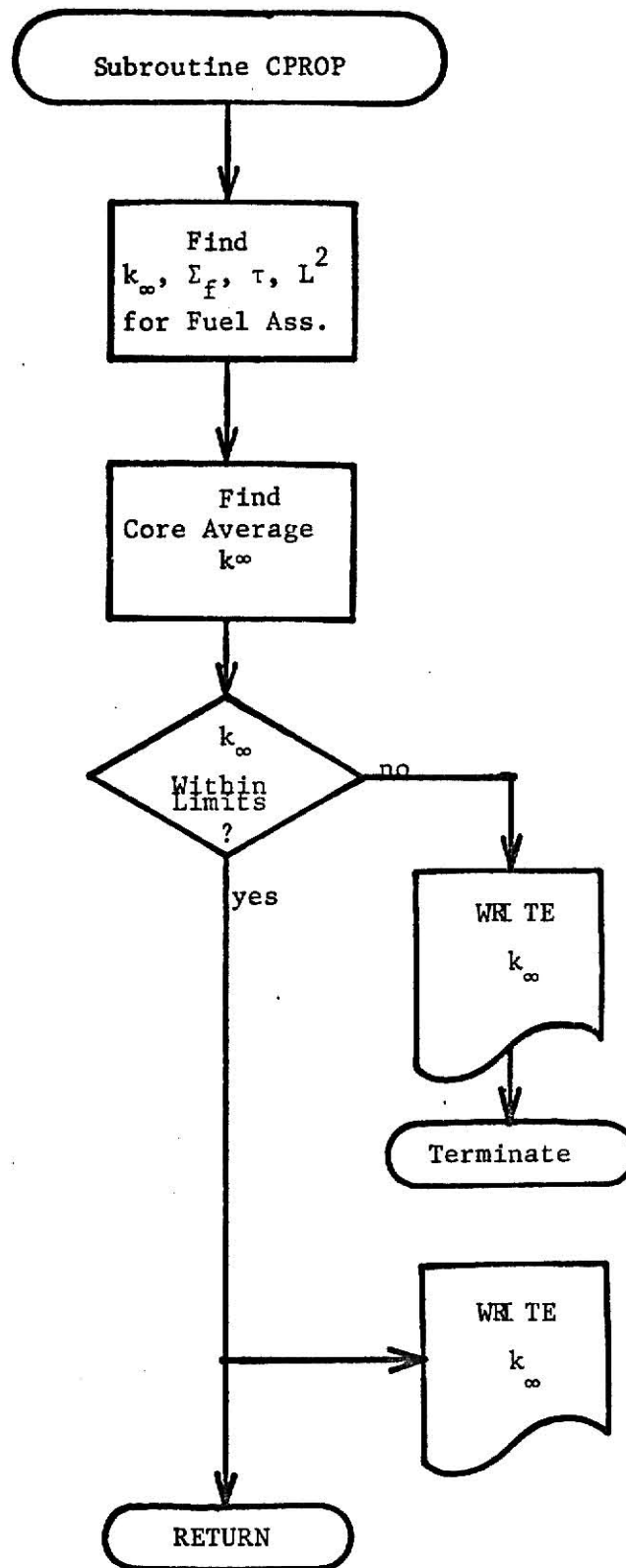
B.1 BURNXY Description

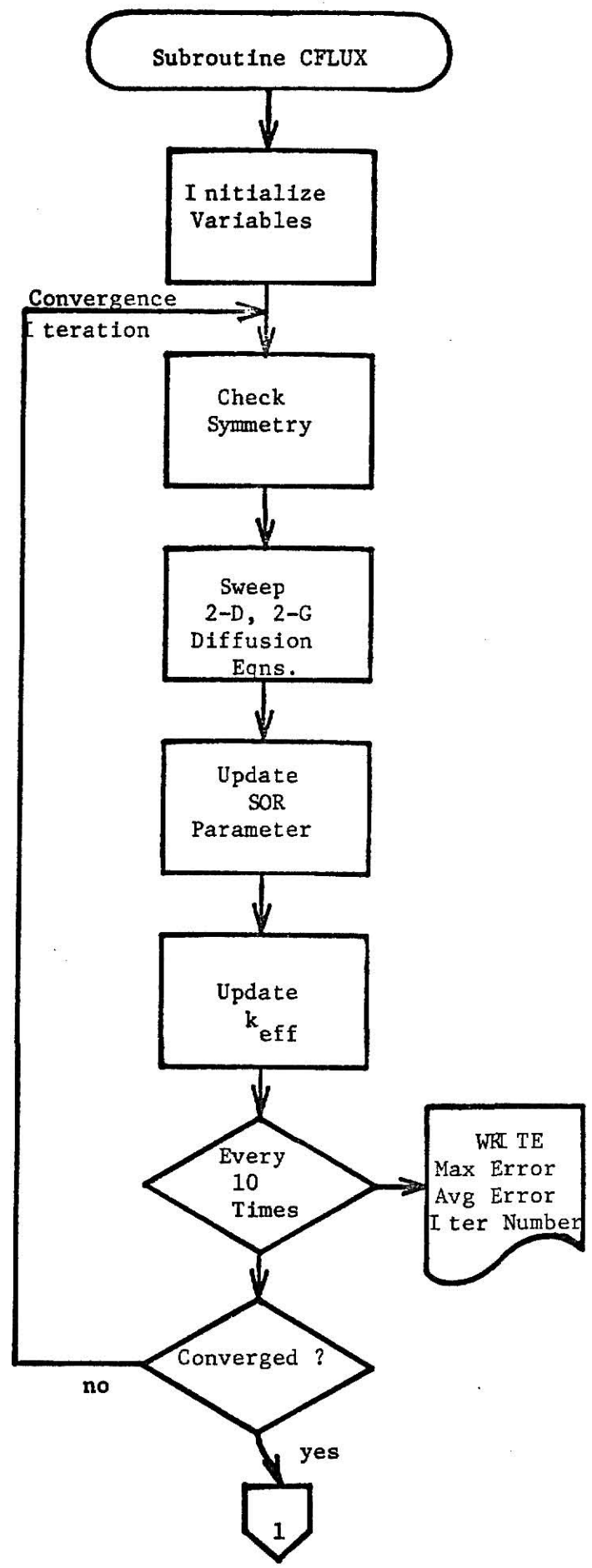
The code BURNXY is similar in function to the previously described code BURNRZ. The codes differ significantly only in the flux calculation routines and in the lack of a subroutine in BURNXY analogous to LIQHOH in BURNRZ. Naturally, the radial averaging routines XYTOR and ARAD are also not present. The BURNXY flux calculation is very similar to that in BURNRZ, but the geometry of the Laplacian operator is X-Y, not R-Z.

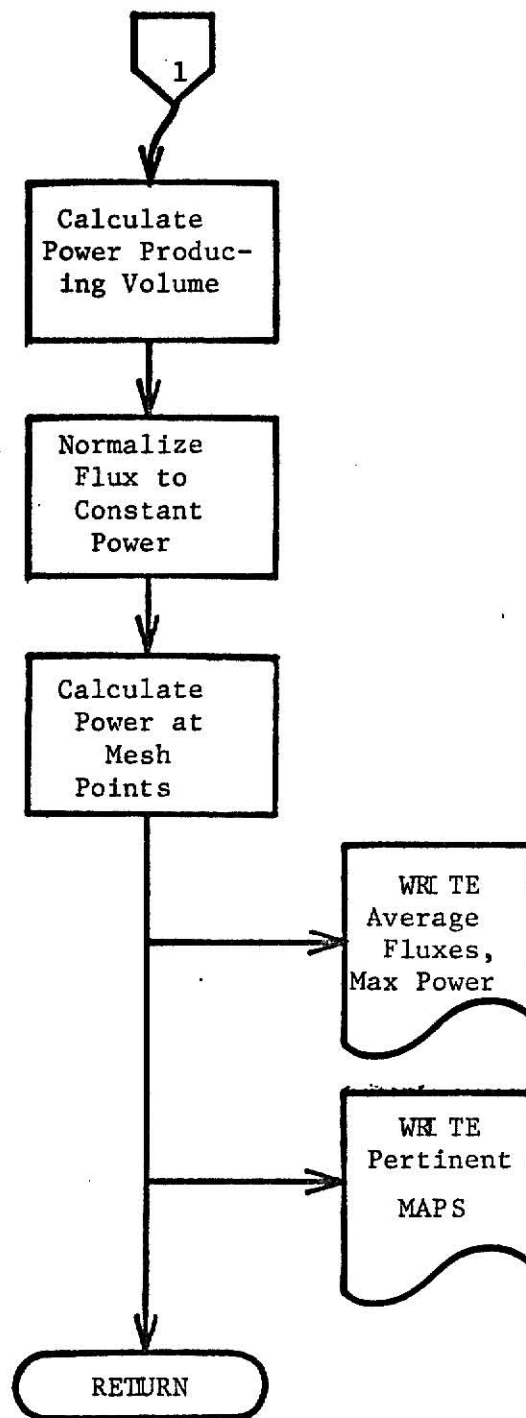
B.2 Flow Diagram of the Code BURNXY











APPENDIX C: LINEAR PROGRAMMING EQUATIONS

The linear programming equations used in the fuel shuffle model are fully described in Chen⁽¹⁹⁾, and are here reproduced for convenience.

$$\text{Objective Function: Deviation of Fuel Distribution} = \sum_{j=1}^{59} d_{3j}$$

(Minimize)

Constraints:

Radial Uniformity:

$$1 = \frac{1}{A_k} \sum_{i=1}^N g_{ik} X_i - D_k \quad k=1,2,3,\dots,14$$

Azimuthal Uniformity:

$$1 = \frac{1}{A_j} \sum_{i=1}^N f_{ij} X_i - D_j \quad j=1,2,3,\dots,12$$

Local Uniformity:

$$1 = \frac{1}{5} (X_i + G_i) - D_i \quad i=1,2,3,\dots,33$$

Normalization:

$$M_\ell = \sum_{i=1}^N a_{i\ell} X_i \quad \ell=2,3$$

Deviation:

$$d_{3j} = d_{1j} + d_{2j} \quad j=1,2,3,\dots,59$$

where the variables are defined as

N = the total number of positions in the core,

D_j = deviation of fuel distribution from ideal, sector j ,

d_{1j} , d_{2j} , d_{3j} = modified deviations for sector j ,

A_j = total area in sector j ,

g_{ik} = fraction of area if position i that lies in the boundary of ring k ,

X_i = state variable,

f_{ig} = fraction of area of position i that lies inside azimuthal area g ,

G_i = summation of four state variables immediately adjacent to i ,

M_ℓ = total number of positions in zone ℓ , and

$a_{i\ell}$ = 1 if i is in ℓ , 0 otherwise.

AXIAL DEPENDENCE OF NUCLEAR FUEL MANAGEMENT

by

BRUCE ALAN NAPIER

B.S., Kansas State University, 1975

AN ABSTRACT OF A MASTER'S THESIS

submitted in partial fulfillment of the
requirements for the degree

MASTER OF SCIENCE

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1977

ABSTRACT

This study represents an investigation of the effects of variations of burnup in the axial direction of nuclear reactor fuel on the operation of an example pressurized water reactor (PWR). The effects of axially changing the fuel enrichments to compensate for burnup and leakage are examined.

Flux and power distributions are neither uniform nor constant through the life of a nuclear reactor. Theoretical considerations can suggest a fuel distribution for which the flux or power distribution is instantaneously uniform, but burnup effects quickly distort this uniformity. Therefore the best that can be obtained is a fuel distribution which results in an acceptable power distribution for an adequate period of time.

A typical PWR core is modeled with a two-group, two-dimensional diffusion code and a linear programming fuel shuffling code. The behavior of the fuel burnup, local power peaking factor, and a related parameter, the core "Figure-of-Merit," are examined for a new clean core and two subsequent burnup cycles. In addition the same parameters are examined for three experimental cases for which the virgin fuel enrichments are varied axially. It is found that a small increase in fuel enrichment, for example from 3.1% to 3.2%, at the reactor top, bottom and outside edge results in lower power peaking factors, higher Figures-of-Merit, and thus in more even burnup of fuel. More even burnup indicates better fuel consumption and eventually lower fuel costs.