

BEAM-FOIL LIFETIMES OF CrII

by

2115-6574R

JAMES R. METTLING

B.A., Southwestern College, 1971

A THESIS

submitted in partial fulfillment of the

requirements for the degree

MASTER OF SCIENCE

Department of Physics

KANSAS STATE UNIVERSITY
Manhattan, Kansas

1974

Approved by:



Major Professor

LD
2668
T4
1974
M48
C.2
Document

1

Table of Contents

List of Tables	ii
List of Figures	iii
Introduction	1
Spectroscopic Background	2
The Energy Loss: Experimental Procedure	7
The Energy Loss: Results	11
The Spectra: Experimental Procedure	14
The Spectra: Analysis	26
The Lifetimes: Experimental Procedure	28
The Lifetimes: Results	29
The Solar Photospheric Abundance of Chromium	34
Acknowledgements	41
References	42
Appendix I: Computer Analyses of Lifetimes	44

List of Tables

Table I.	Theoretical energy loss calculations for 89 keV Ar, 85 keV Cr, and 100 keV Cr	8
Table II.	Lines of CrI and CrII identified in this spectrum . . .	27
Table III.	Possible cascades into levels of interest	31
Table IV.	The lifetimes measured at 85 and 100 keV for five lines	32
Table V.	A comparison of the measured lifetimes to the results of CB, Shackleford, and Byard	34
Table VI.	Transitions from the $z^4\text{H}^0 5\frac{1}{2}$ ($3d^4(a^3\text{H})4p$) level	35
Table VII.	Transitions from the $z^4\text{H}^0 6\frac{1}{2}$ ($3d^4(a^3\text{H})4p$) level	36
Table VIII.	Transitions from the $z^4\text{F}^0 3\frac{1}{2}$ ($3d^4(a^3\text{D})4p$) level	37
Table IX.	Transitions from the $z^4\text{F}^0 4\frac{1}{2}$ ($3d^4(a^3\text{D})4p$) level	38
Table X.	A comparison of the measured lifetimes to the gA values of others	40

List of Figures

Fig. I.	The apparatus for the energy loss experiment . .	9
Fig. II.	The energy spectrum of Ar after passage through a carbon foil	12
Fig. III.	The apparatus for the spectral scans and for the lifetime experiment	15
Fig. IV.	The spectrum: 1900 to 2580 A	17
Fig. V.	The spectrum: 2580 to 3230 A	18
Fig. VI.	The spectrum: 3220 to 3880 A	19
Fig. VII.	The spectrum: 3870 to 4100 A	20
Fig. VIII.	The high resolution spectra of CrI and CrII: 2910 to 2950 A	21
Fig. IX.	The high resolution spectra of CrI and CrII: 2970 to 3010 A	22
Fig. X.	The high resolution spectra of CrI and CrII: 3110 to 3150 A	23
Fig. XI.	The high resolution spectra of CrI and CrII: 3170 to 3220 A	24
Fig. XII.	The high resolution spectra of CrI and CrII: 3950 to 4000 A	25
Fig. XIII.	A partial Grotrian Diagram of CrII	30
Fig. XIV.-XVIII.	The computer output	44

Introduction

In 1962, the first formal proposal for a beam-foil experiment was made by Kay¹. Beam-foil Spectroscopy (BFS) has since come into wide use. Along with other applications BFS provides a new method for measuring the lifetimes of atomic states from which radiative transition rates may be derived. Previously arc, spark, and absorption spectra were the primary sources for this information. BFS presents problems of its own, of course, but it does not have the problem associated with it of the assumption of thermal equilibrium in the light source. In addition, one need not worry about plasma stratification, which is a difficulty encountered in shock tube experiments.²

A BFS lifetime measurement is performed by passing a beam of accelerated ions through a thin foil. As atoms pass through the foil the valence electrons are excited. Downstream of the foil one may observe the number of photons emitted as a function of distance from the foil. The velocity of the beam after the foil is then used to convert the distance scale to a time scale. Thus the final result is a plot of the number of photons emitted versus time. The slope of this plot is the lifetime of the emitting level.

Lifetimes of atomic energy levels may be used with branching ratios to find atomic transition probabilities, which are important to the study of elemental abundances in the solar and stellar atmosphere. Recent calculations of the chromium abundance, based on BFS lifetimes of CrI,³ make it desirable to calculate a similar abundance based on the lifetimes

of CrII. The reason for this is that at solar photospheric temperatures the principle ionic species is CrII. Therefore, a small error in the temperature-pressure stratification might lead to a large error in the fractional abundance of CrI and would produce a significant error in determining the total chromium abundance from CrI. The ratio of CrII to total chromium abundance should be much less sensitive to small errors in the temperature-pressure stratification.

The work presented in this thesis may be conveniently divided into four phases:

1) Some question as to the accuracy of BFS measurements has been raised in connection with the energy loss incurred by the beam.⁴ The problem has been examined to some extent in this work. The energy loss for 89 keV Ar⁺ has been measured and agrees surprisingly well with stopping power calculations.

2) In preparation for the lifetime experiments, spectral scans of the BFS spectrum have been taken between 1900 and 4100 Å. Lines originating in CrI and CrII exist in this region and have been identified.

3) Lifetimes of five lines representing four levels in CrII have been measured. These measurements were made at two energies to ensure adequate knowledge of the energy loss.

4) A preliminary calculation of the solar abundance of chromium has been made based on our CrII results.

Spectroscopic Background

The following discussion will define common spectroscopic terminology and relate the Einstein coefficient (A) to the oscillator strength (f), and to the atomic radiative lifetime (τ). The following definitions are based on the shell model of the atom and the LS coupling

scheme, and will consider only radiative transitions of the valence electrons.⁵ In this coupling scheme, a spectroscopic "term" is characterized by the total orbital angular momentum (L) and the total spin (S) of the valence electrons. The conventional notation for a term is ^{2S+1}L , where $2S+1$ is called the "multiplicity" of the term. The multiplicity is the number of possible orientations of S relative to L . One may think of the multiplicity as the number of ways L and S can couple to form the total angular momentum, J , but this is only true for $L \geq S$. For example, if $S = \frac{1}{2}$ and $L = 0$ as in the hydrogen atom, the term is 2S and there are two possible orientations of S , but only one possible coupling to form J ($J = \frac{1}{2}$). A "level" designates one possible value for J , and is written $^{2S+1}L_J$. One specifies the "sublevel" that an atom is in by specifying the orientation of the atom, M . Sublevels are degenerate in energy except in the presence of a magnetic or electric field. A transition between two sublevels produces a "line component". The blend of line components produced by all transitions between two particular levels form a "line". The set of all lines produced by transitions between two terms is called a "multiplet".

The lifetimes of atomic levels are the principal quantities with which this work is concerned. However, the literature often refers to measurements of "gA-values or gf-values" as a measure of electromagnetic transition probabilities. Therefore, it is appropriate to discuss these two quantities as an alternative means of expressing transition probabilities.⁶

Let us express the initial sublevel as $i(\alpha'J'M')$, where α' represents all internal quantum numbers, and the final sublevel as $f(\alpha JM)$. Also let us define $A(\alpha'J'M' \rightarrow \alpha JM)$ to be the transition rate for the transition

$i \rightarrow f$, that is the transition probability per unit time for $i \rightarrow f$.

Since, in the absence of an electric or magnetic field the components blend to produce a single line, it is important to know the transition rate for the line. For each level there will be $2J+1$ sublevels, corresponding to the $2J+1$ possible values of M . To calculate the transition rate for the level ($A(\alpha'J' - \alpha J)$) it is necessary to sum the sublevel transition rate over all M and M' . Thus, $A(\alpha'J' - \alpha J)$ is defined by the equation

$$N(J') A(\alpha'J' - \alpha J) = \sum_{M'M} N(J'M') A(\alpha'J'M' - \alpha JM),$$

where $N(J')$ is the population of the level and $N(J'M')$ is the population of each of the sublevels $i(\alpha'J'M')$. Since the transition rate of the sublevel is not dependent on the orientation of the atom (M'), it may be factored from the sum over M' . Thus, suppressing M' in the transition rate,

$$N(J') A(\alpha'J' - \alpha J) = \sum_M A(\alpha'J' - \alpha JM) \sum_{M'} N(J'M').$$

Since it is clear that

$$\sum_{M'} N(J'M')/N(J') = 1,$$

we may write

$$A(\alpha'J') = \sum_M A(\alpha'J' - \alpha JM).$$

In a thermally excited source the sublevels are assumed to be equally populated. Thus, we may write

$$N(J') = (2J'+1) N(J'M') = gN(J'M'),$$

which is the defining equation for the statistical weight g . The observed intensity I is equal to the number of atoms multiplied by the transition rate. The intensities and populations are the quantities normally measured in thermal sources. One may express the intensity as

$$I = N(J') A(\alpha' J' - \alpha J) = gA(\alpha' J' - \alpha J) N(J'M').$$

The gA -values may be found experimentally by measuring the populations of the sublevel and the intensity of the line.

In BFS the lifetimes of atomic levels are measured directly. Thus a relation must be sought that will provide the total transition rate ($A(\alpha' J')$) from one particular level. This may be done by summing all the transition rates out of the level of interest. Therefore one has

$$A(\alpha' J') = \sum_{J, \alpha} A(\alpha' J' - \alpha J).$$

The lifetime (τ) of a level is defined to be the inverse of the transition rate, and it is the time needed for the population of the level to decrease by $1/e$ of its original value.

By observing the lifetime of a single line originating in the level of interest, when all transitions from that level are taking place, one is provided with the total lifetime of the level. This is true because

the number of transitions per unit time from the level of interest to a particular final level is proportional to the number of atoms ($N(\alpha'J')$) in the level at that time. But N obeys the time dependant equation

$$dN/dt = -N/\tau ,$$

and will be given by $NA(\alpha'J' - \alpha J)$. Therefore, when all allowed transitions occur each transition out of level ($\alpha'J'$) will be characterized by a decay constant which is the total lifetime of the level. In order to transform the measured lifetime into the corresponding gA -values for a specific transition. More experimental information is needed. This is the branching ratio which is defined as the ratio of the transition rate of $i(J') \geq f(J)$ to the total transition rate.

As has been shown, atomic lifetimes and transition rates are different quantitative measures of the same phenomena. The oscillator strength or f -value is yet another way of describing the rate of an atomic transition.⁷ The oscillator strength may be defined by comparing the strength of absorption by a quantum mechanical atom with that which would be incurred by an electron moving in a potential characterized by the resonant absorption frequency (ν). If, for a particular frequency, there are N such classical absorbers per unit volume, and N' atoms per unit volume, then f is defined by $N = N'f$, and is the number of classical absorbers per atom. In terms of the transition rate the f -value is given by

$$gf = gA \text{ mc}^3 / 8 \pi^2 e^2 \nu^2 = 1.4493 \times 10^{-5} \text{ m}^{-2} gA / \nu^2 ,$$

where m is the mass of the electron and e is the charge on the electron in esu. Since the natural lifetime (T) of a classically absorbing electron is

$$T = 3mc^3/8e^2\pi^2\nu^2,$$

the preceding equation may be written as

$$gf = gA T/3.$$

This relationship indicates the intrinsically classical nature of f -values. It is important to realize that the two values of g are not the same and may not be factored from the equation.

The Energy Loss: Procedure

The energy loss suffered by the chromium ions traversing the foil is of great importance in the calculation of the level lifetimes, as it is necessary to express the distance from the foil as a function of time. This distance is related to the time of flight of the ion by

$$t = d (m/2(E-\Delta E))^{\frac{1}{2}},$$

where E and m are the energy and mass respectively of the incident ion, d is the perpendicular distance between the foil and the axis of the lens, t is the time it takes for the ion to traverse d , and ΔE is the energy loss.

The problem of calculating the stopping power of a low energy ion traversing matter has been explored by Lindhardt.^{8,9} Calculations

of energy losses were made using the method of Lindhardt and a program written by Brand, Fox, and Keller.¹⁰ The results for several relevant choices of ion, energy, and thickness of carbon foil are listed in Table I.

Table I
Theoretical energy loss calculations for 89 keV Ar,
85 keV Cr, and 100 keV Cr

Ion	x (gm/cm ²)	E(keV)	E(keV)
Ar	2.4	89	11
Cr	2.0	85	13
Cr	2.0	100	13

In order to investigate the validity of these calculations, the energy loss for an 89 keV Ar ion was measured. Argon was used rather than chromium because the argon beam is easily produced. Since it has an atomic number ($Z=17$) close to that of chromium ($Z=24$), the results of the experiment provide a useful comparison to the calculation.

The beam for the energy loss experiment was supplied by a 150 kV linear accelerator, equipped with a Physicon universal ion source. The Physicon source is capable of producing ionized gas from either gaseous or solid materials. The argon beam was produced by allowing argon gas to enter the source at a rate which produced a partial pressure in the accelerator of approximately 6×10^{-6} mm of Hg. The gas was ionized in the source and extracted from it by a -10 kV potential on an extractor electrode. The beam was then accelerated to an energy of 89 keV, mass analyzed at the magnet, and directed into the target chamber, as shown

**THIS BOOK
CONTAINS
NUMEROUS PAGES
WITH DIAGRAMS
THAT ARE CROOKED
COMPARED TO THE
REST OF THE
INFORMATION ON
THE PAGE.**

**THIS IS AS
RECEIVED FROM
CUSTOMER.**

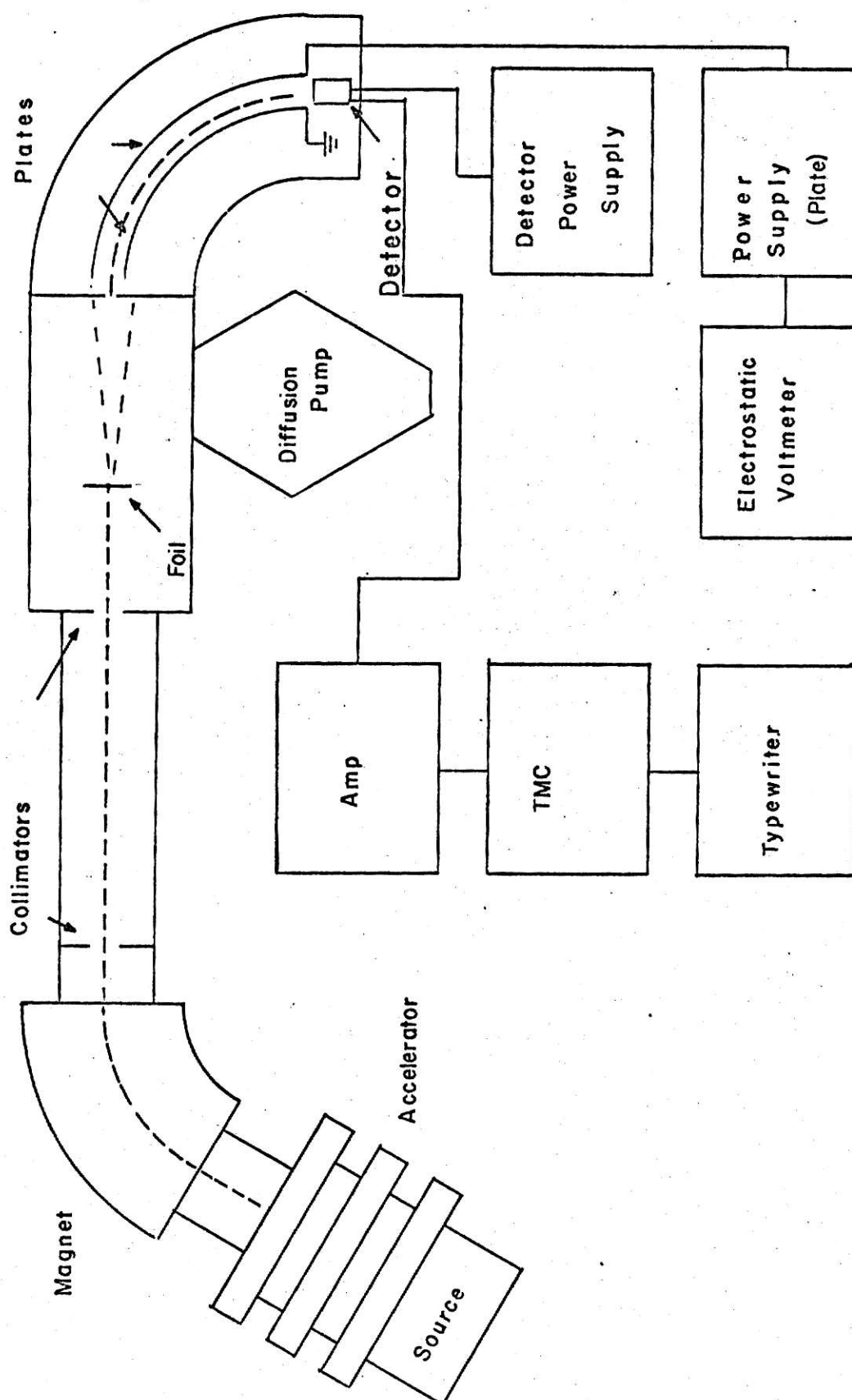


Fig. 1.--The apparatus for the energy loss experiment

in Figure I. The argon beam was identified by maximizing the beam current of the most intense beam at a Faraday cup just downstream of the magnet. The argon leak was then closed; if the beam current dropped dramatically, it was assumed that the argon beam had been identified.

The size of the beam spot was restricted by the use of two $\frac{1}{4}$ in. collimators, and vertical control of the beam was accomplished with the use of electrostatic deflection plates powered by a 400 V variable voltage supply. (The plates and power supply are not shown in Figure I.) A vacuum was maintained in the target chamber by a 4 in. diffusion pump located several inches downstream of the foil.

An electrostatic analyzer, built by Curnutte, was used to measure the energy of the beam before and after passage through the foil, the energy loss being the difference between the two values. The analyzer consisted of two curved parallel plates preceded by a 0.5 mm aperture. Particles were detected by a Bendix Spiraltron electronmultiplier located at the exit aperture of the analyzer.

The signal from the Spiraltron was amplified by a spectroscopic amplifier and recorded on a TMC (Technical Measurements Corporation) multi-channel analyzer, model 404-6. The voltage for the analyzer plates was supplied by a Spellman high voltage power supply, and was calibrated with an electrostatic voltmeter to ± 0.1 kV (corresponding to approximately ± 1.06 keV in particle energy). In order to record the energy spectrum on the TMC, the plate voltage was varied manually at a rate of .03 kV/sec, while the TMC address was advanced internally by 1 channel/sec.

The beam was electrostatically analyzed without a foil from which the ratio of the beam energy to the plate voltage of the electrostatic analyzer was determined to be 0.0945 kV/keV of particle energy. Several energy loss measurements were then made on a single foil. It was noted that as time elapsed the energy loss of the beam increased. Therefore energy loss measurements were made using a new foil noting carefully the time elapsed during each run. The foil used for these runs had been damaged in mounting and thus small holes were present in the foil. Under these circumstances, the energy spectrum contained a peak for the beam of particles that passed through the foil and a second peak for the unperturbed particles, which passed through the holes. Three additional runs were made with unused foils to make possible a calculation of the average energy loss.

The foils used throughout the experiments were provided by the Arizona Foil Company. The thicknesses quoted here, ranging from $10 \mu\text{gm}/\text{cm}^2$ to $2.0 \mu\text{gm}/\text{cm}^2$, are those measured by the supplier. Thin carbon foils are not durable if self-supporting, and therefore were mounted on 80% transmission electroplated nickel mesh. Several foils could be placed in the chamber simultaneously, and could be moved in and out of the beam while the system was under vacuum.

The Energy Loss: Results

The results of the energy loss experiment may be divided into two parts: first, the thickening of the foils; second, the actual energy loss incurred by the beam.

Figure II shows the gradual thickening with time of a carbon foil, which was initially $2.4 \mu\text{gm}/\text{cm}^2$ in thickness. This time dependence is probably better described as a dependence on the amount of charge

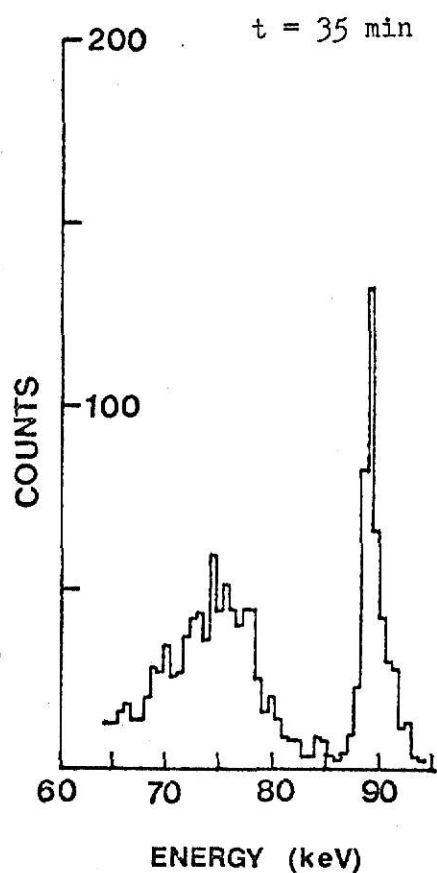
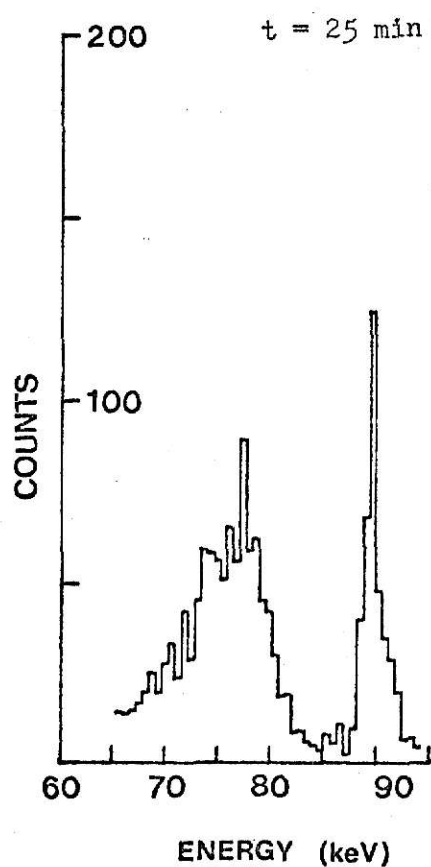
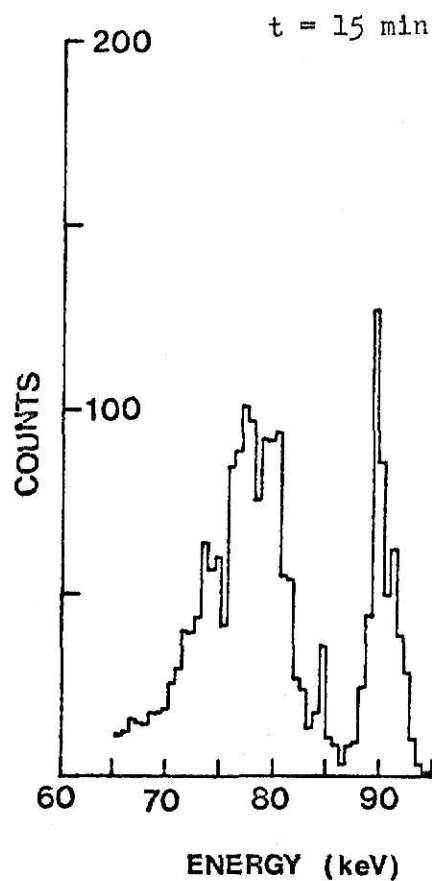
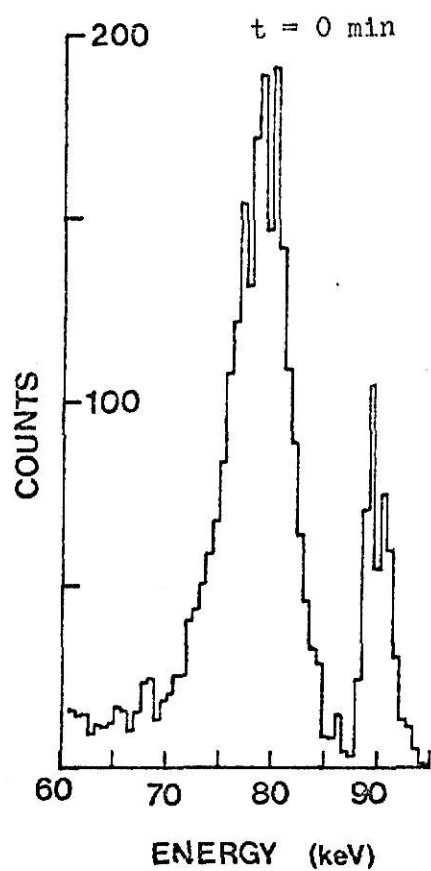


Fig. II.--The energy spectrum of Ar after passage through a carbon foil

delivered to the foil. The incident beam was 89 keV Ar^+ , and corresponds to the peak of higher energy in each graph. The pressure in the system was 10^{-5} mm of Hg. The time indicated is the time at which each run began, a typical run duration being 3 minutes. The beam current could not be measured directly while the electrostatic analyzer was in use, but it was measured at a forward Faraday cup before the experiment. The percentage of current transmitted to the foil from this cup is normally about 30%, but transmission has been as high as 50%. Therefore from the measurement of 13 μA at the Faraday cup, one may conclude that the foil actually received a current on the order of several microamperes. It was found that the energy loss increased from 10 keV to 14 keV in a period of 35 minutes. This implies that the foil thickened at a rate of $0.015 \mu\text{gm}/\text{cm}^2\text{-min}$. This effect is attributed to the presence of hydrocarbons (pump, oil, etc.) in the system. The build up probably occurs when a molecule very close to the foil is hit by an incoming ion. Dissociated, the molecule may form new bonds with the foil or other molecules. It is hypothesized that better vacuum would reduce the size of the effect. Precautions were taken both in subsequent energy loss experiments and in the lifetime experiments to avoid this change in foil thickness. For the energy loss measurements each foil was used only once, while in the lifetime experiments the time each foil was exposed to the beam was kept below 15 minutes.

An average energy loss of 10 keV was measured using $2.4 \mu\text{gm}/\text{cm}^2$ foils. The foils were bombarded for approximately 2 min. The agreement between this measurement and the previous calculation of 11 keV (see Table I) strongly supports the validity of our use of the Lindhardt calculations for the chromium lifetime analyses.

The Spectral Scans: Apparatus

The same acceleration, collimation, and pumping systems were utilized for the spectral scans as for the energy loss experiment (see Figure III). Solid chromous chloride was used in the source to produce the chromium beam. The Cr^+ beam was identified by first producing an Ar^+ beam as in the energy loss experiment. A Hall probe was then used to measure the magnetic field of the deflection magnet. The voltage measured across the Hall probe is proportional to the magnetic field. The magnetic field needed to turn an ion of given energy through a specified radius of curvature is proportional to the square root of the mass of the ion. Thus, for Ar^+ and Cr^+ beams the following formula gives the projected Hall probe voltage (V_2) for the chromium beam,

$$V_2 = 1.14 V_1,$$

where V_1 is the Hall probe voltage measured when argon was used.

An adjustable LiF lens, with a focal length of 4.5 in. at 4000 Å, was placed perpendicular to the beam line and downstream of the foil. The purpose of this lens was to focus the light emitted by the decaying ions on the slits of the monochromator. Survey scans were made using $6 \mu\text{gm}/\text{cm}^2$ foils. The position of the foil was such that a maximum amount of light was focused on the slits of the monochromator. The monochromator used was a McPherson model 218 vacuum ultraviolet monochromator. Two gratings were available for use with the monochromator; a 1200 line per mm grating blazed at 3000 Å which produced a dispersion in the system of 26.5 Å/mm of slit width, and a 2400 line per mm grating blazed at 1500 Å which produced a dispersion of 13.3 Å/mm.

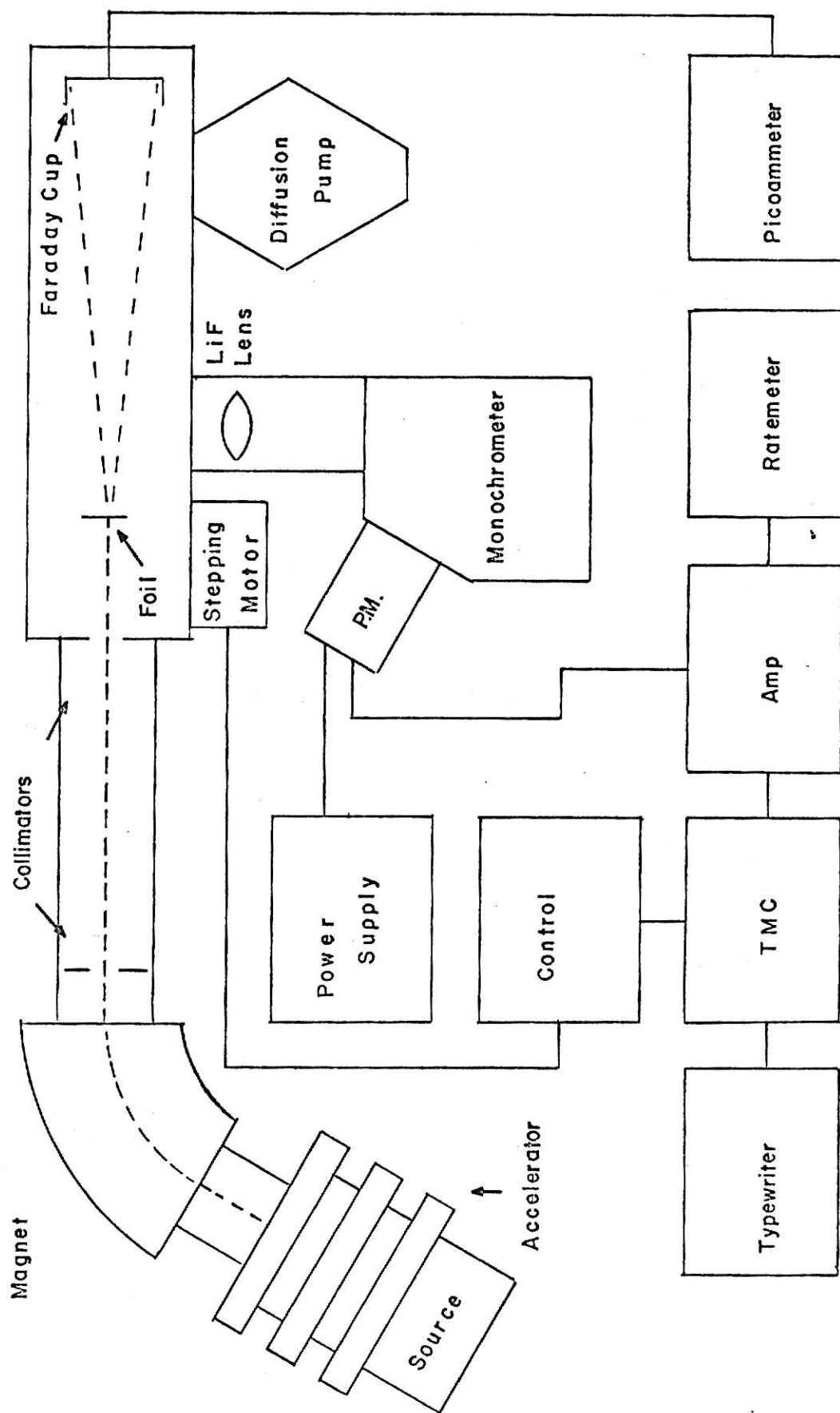


Fig. III.--The apparatus for the spectral scans and for the lifetime experiment

Photons of the proper wavelength were detected with an EMI 6256S photomultiplier tube. The tube was enclosed in a copper case and was cooled by means of a thermoelectric cooling device. A voltage of 1.5 kV was supplied to the tube by a Fluke model 404M power supply. A Nuclear Data Corporation dual channel amplifier and single channel analyzer model ND500, was used to amplify the photomultiplier signal and discriminate against low level noise. This amplifier supplies two output signals, positive pulses which were counted by the TMC, and negative pulses which were monitored on a count rate meter. The electrostatic analyzer had to be removed, and was replaced by a Faraday cup which was used to monitor the beam current.

Two sets of scans were made:

1) Low resolution scans were made using the 2400 line/mm grating. The monochromator slits were set at 250 microns to produce an instrumental resolution of 3 Å. The scanning rate was 50 Å/min. and the TMC was operated in the internal channel advance mode at one channel per second. Thus the TMC recorded all counts produced within a wavelength range of $5/6$ Å in a single channel. Seven overlapping scans of 400 channels (330 Å) each were taken to cover the wavelengths between 1900 and 4100 Å. The results of these scans are presented in Figures IV - IX.

2) High resolution scans of selected regions were taken in a similar manner using the 1200 line/mm grating. Slit widths for these scans were reduced to 150 or 200 microns depending on the resolution needed. The TMC was paced at 1 channel/sec., and the monochromator scan rate was 5 Å/min. The length of these scans was approximately 50 Å. The resulting spectra are shown in Figures IX - XIII.

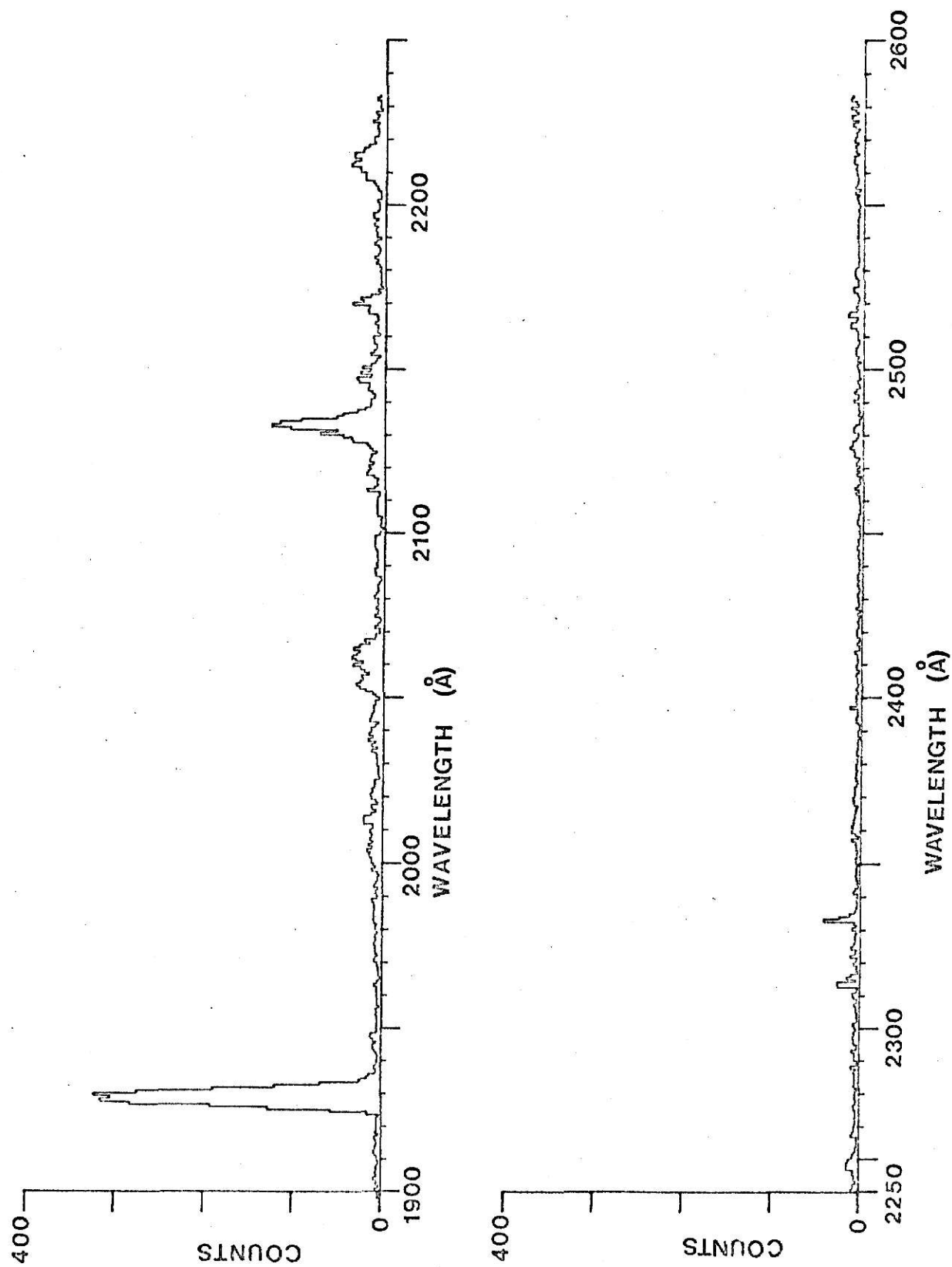


Fig. IV.--The spectrum: 1900 to 2580 Å

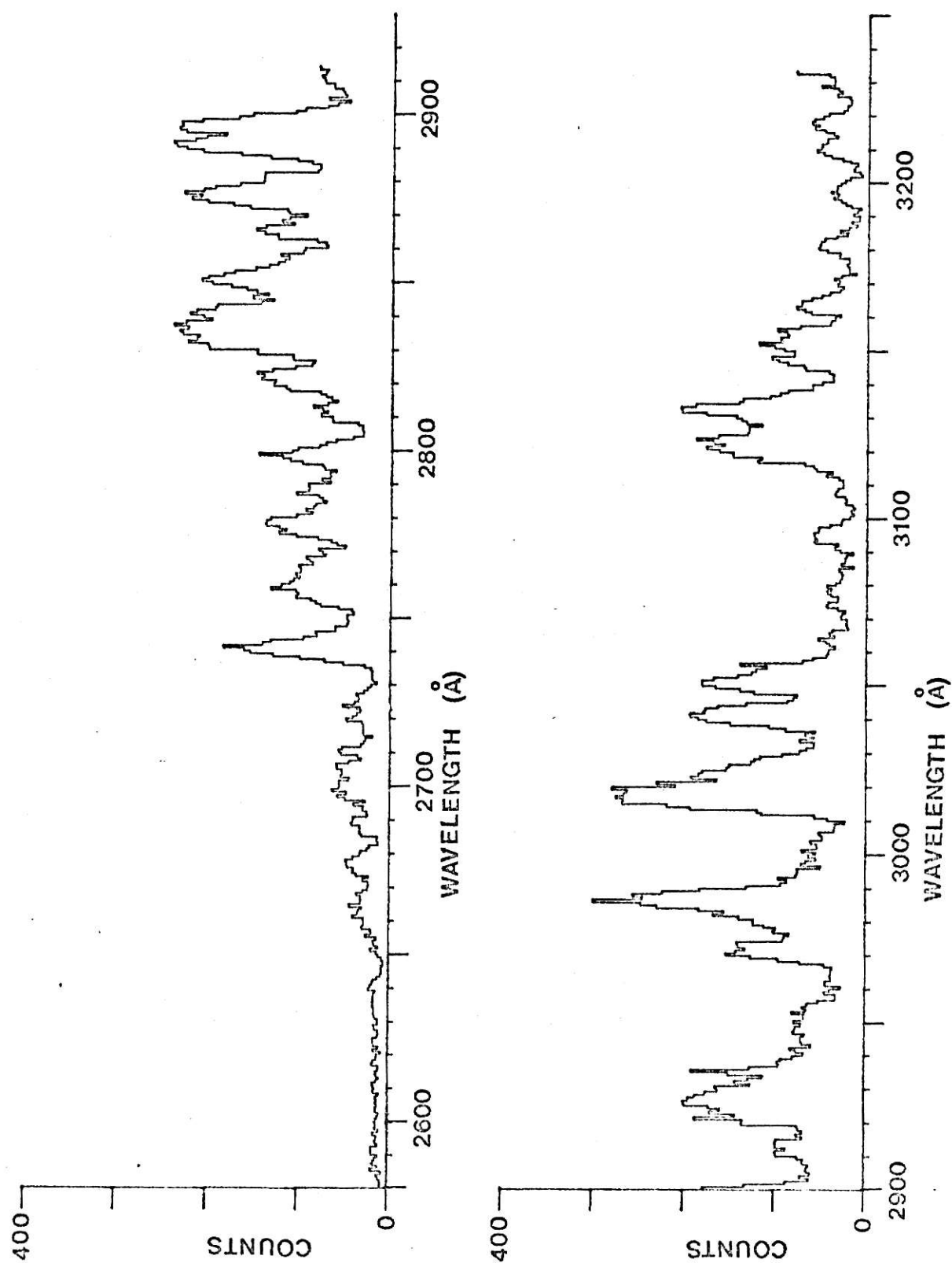


Fig. V.--The spectrum: 2580 to 3230 Å

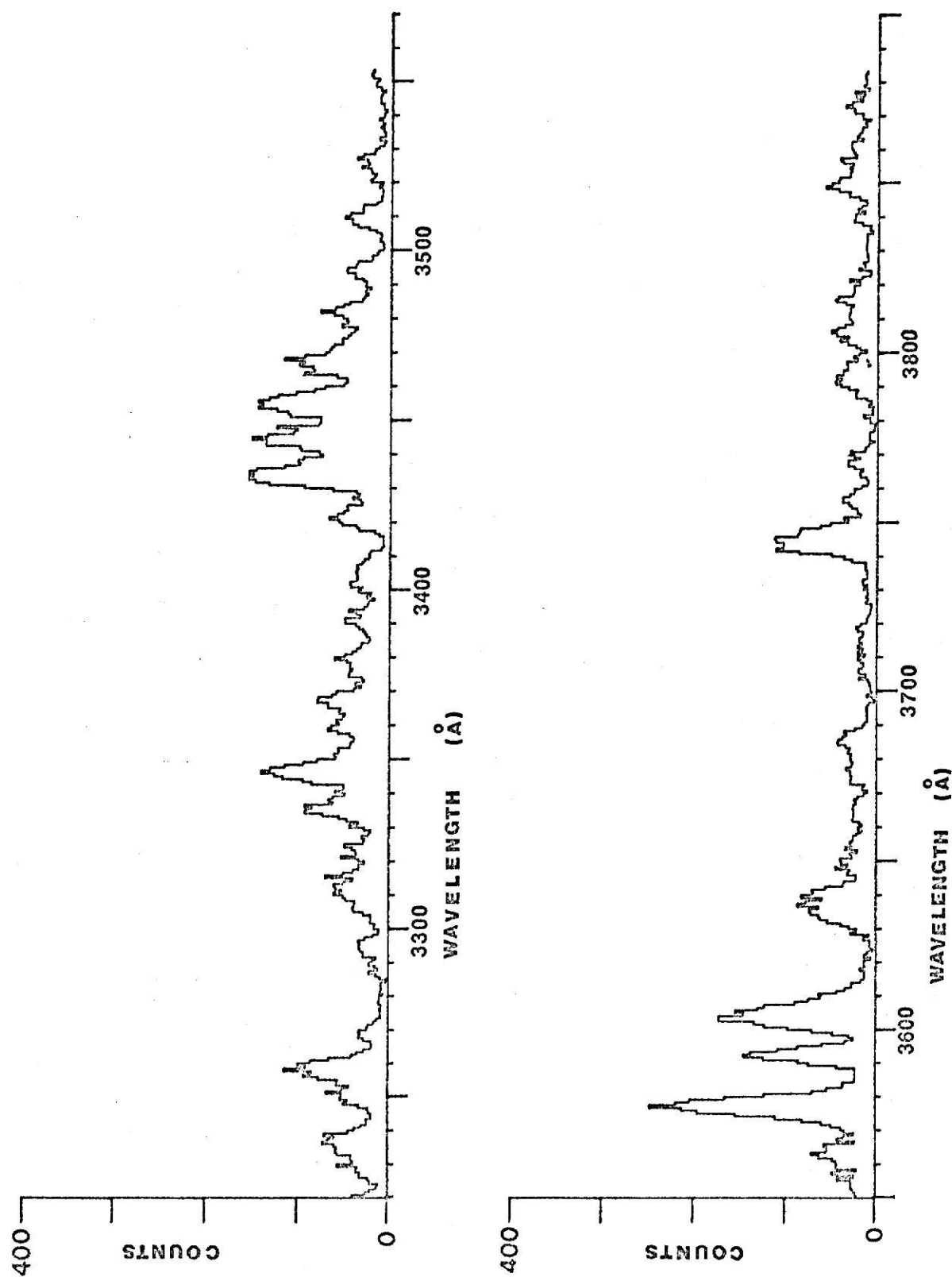


Fig. VI.--The spectrum: 3220 to 3880 Å

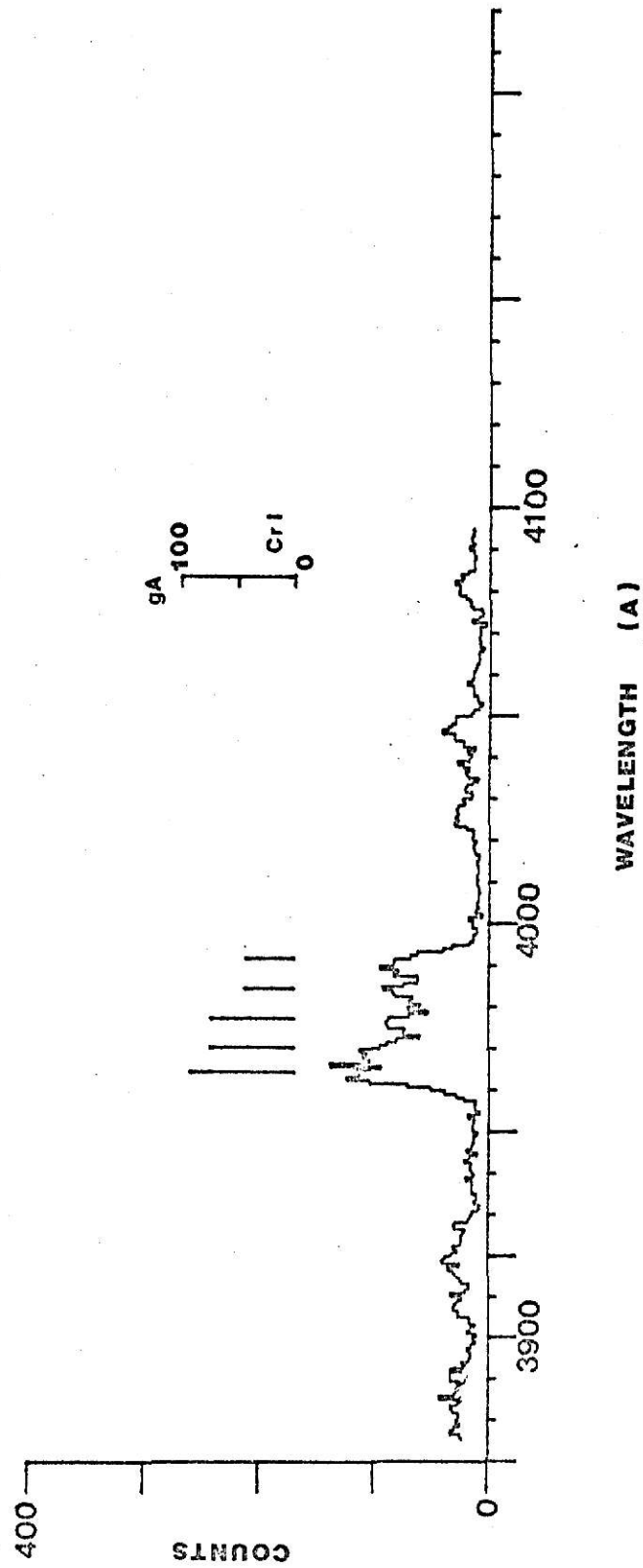


Fig. VII--The spectrum: 3870 to 4100 Å

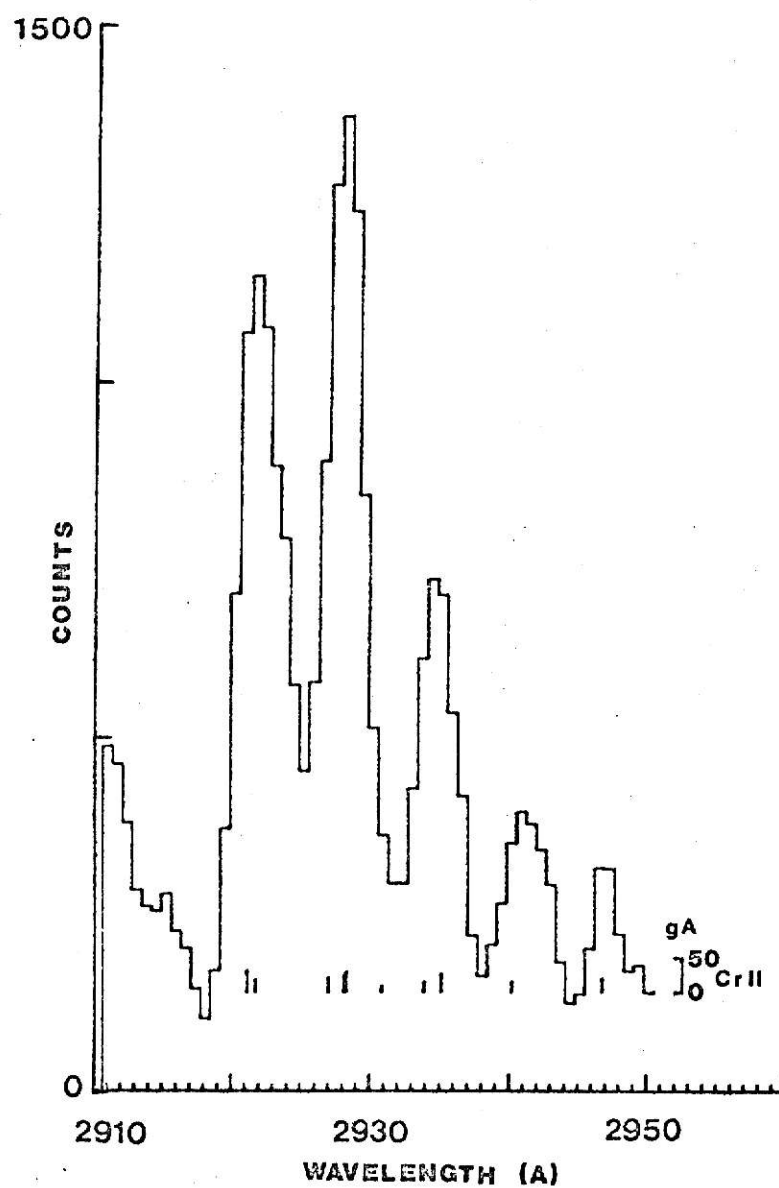


Fig. VIII.--The high resolution spectra of CrI and CrII:
2910 to 2950 Å

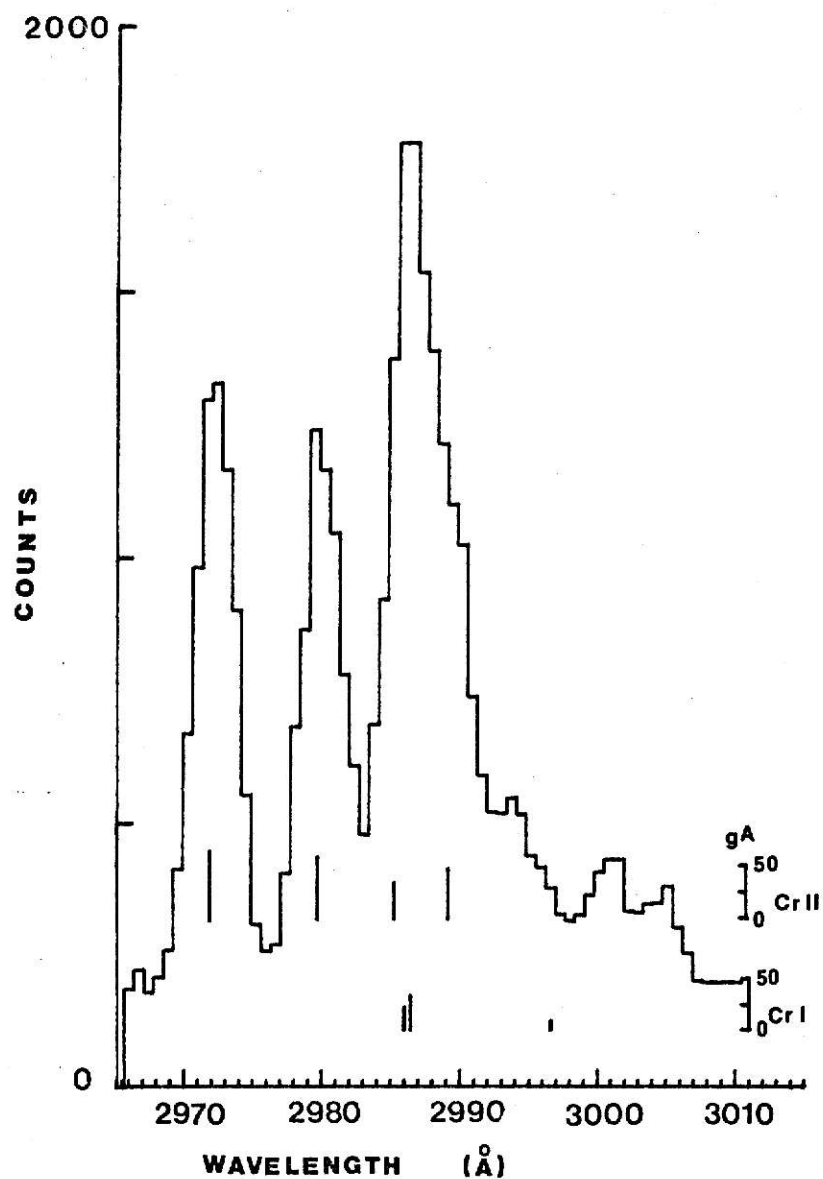


Fig. IX.--The high resolution spectra of CrI and CrII:
2970 to 3010 Å

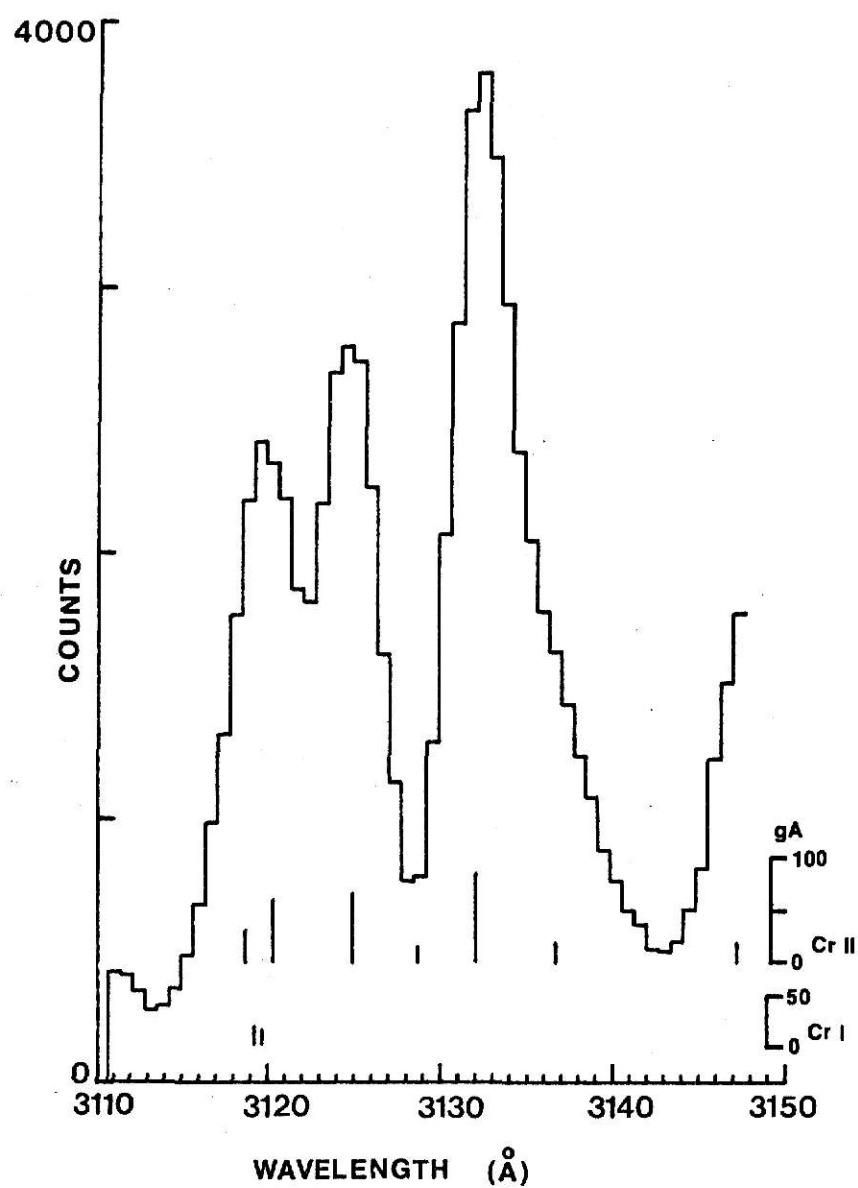


Fig. X.--The high resolution spectra of CrI and CrII;
3110 to 3150 Å

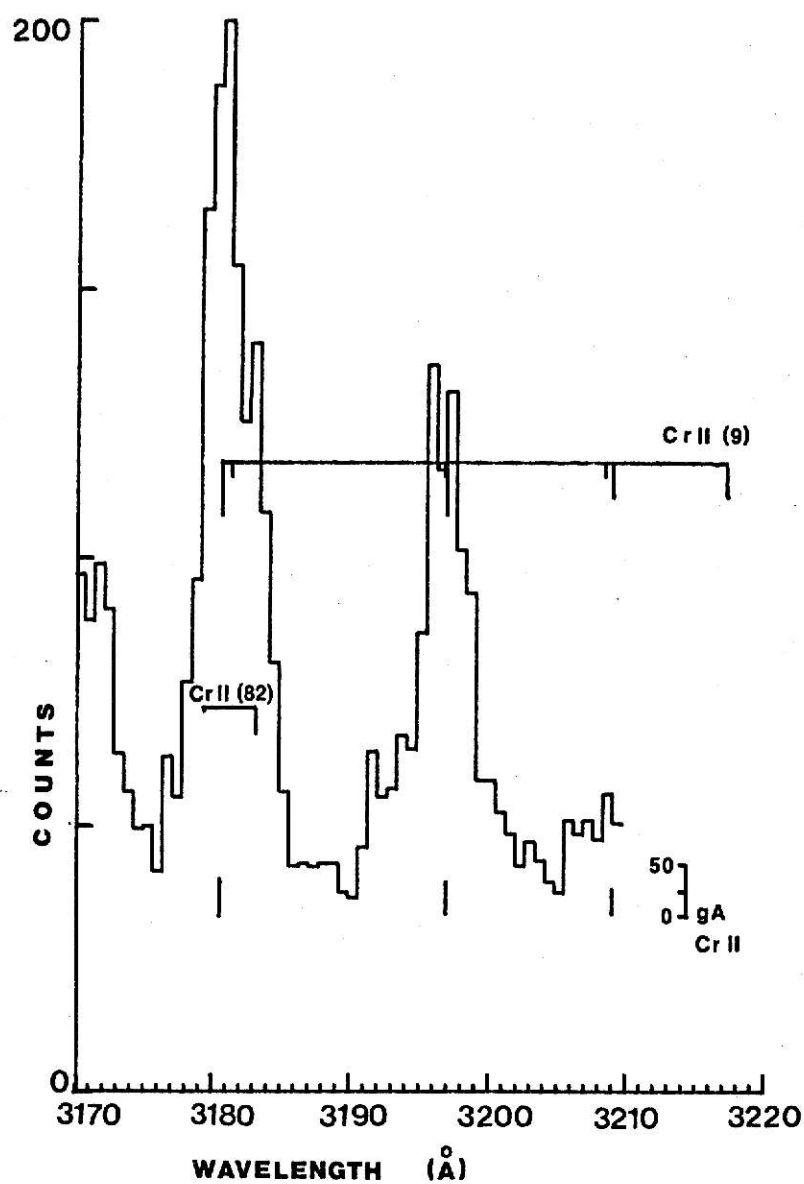


Fig. XI.--The high resolution spectra of CrI and CrII:
3170 to 3220 Å

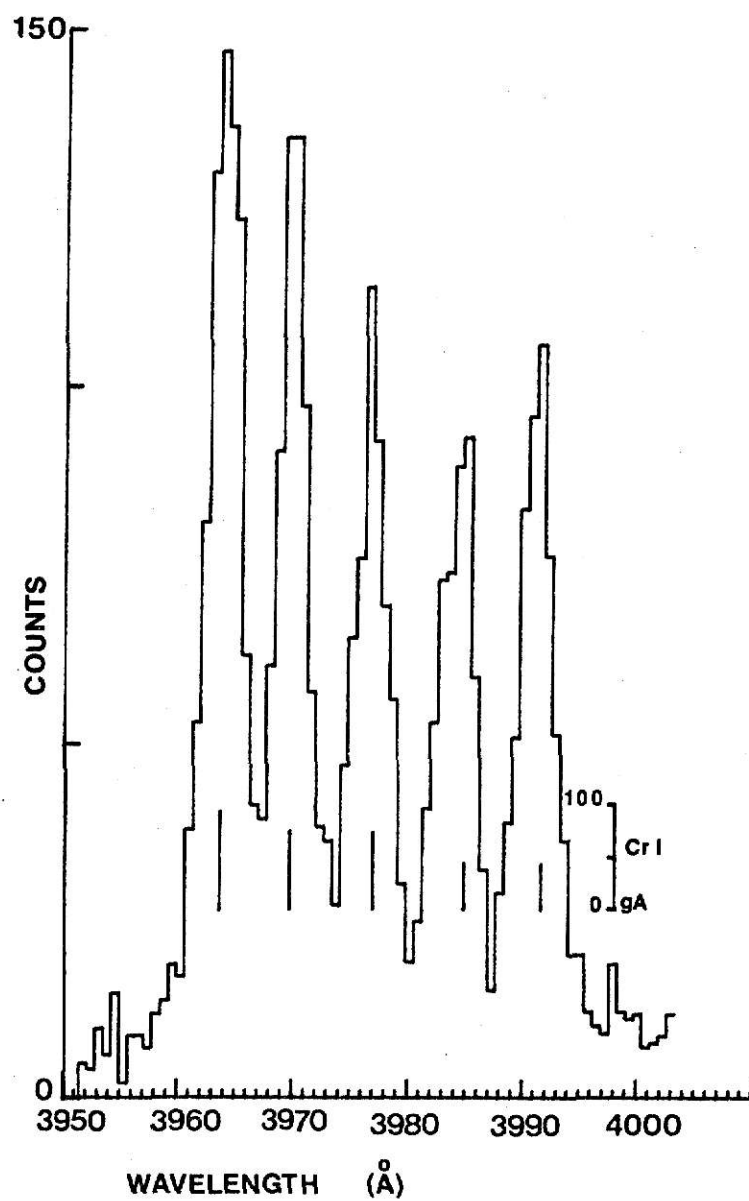


Fig. XII.--The high resolution spectra of CrI and CrII:
3950 to 4000 Å

The Spectra: Analysis

The spectral analysis provides the experimenter with the basic information needed to decide which lines are appropriate for lifetime analysis. The procedure used in this analysis consisted of three steps:

1) A graph of the TMC output was made (counts versus channel, and the appropriate wavelength scale was superimposed on the graph. This produced a graph of counts versus wavelength. Selected gA values ($gA > 10^9 \text{ sec}^{-1}$) of CrI and CrII, taken from the tables of Corliss and Bozman CB),¹⁰ were then plotted on the same graph.

2) Preliminary line identifications were made on the basis of the correspondence between intensity and gA -value. The multiplets to which these lines belong were then found by consulting Moore's A Multiplets Table of Astrophysical Interest¹² and An Ultraviolet Multiplet Table.¹³ If other members of identified multiplets were present with the required intensity, the identification was assumed to be correct. Those CrII multiplets identified in our spectra which seemed to be well resolved and of sufficient intensity for lifetime analysis were noted.

3) High resolution scans were taken in the vicinity of those multiplets chosen to detect any unresolved structure that could influence the lifetime analysis. Table II gives a list of CrI and CrII lines identified.

**THIS BOOK
CONTAINS
NUMEROUS PAGES
WITH THE ORIGINAL
PRINTING BEING
SKEWED
DIFFERENTLY FROM
THE TOP OF THE
PAGE TO THE
BOTTOM.**

**THIS IS AS RECEIVED
FROM THE
CUSTOMER.**

Lines of CrI and CrII identified in this spectrum					
Wavelength (Å)	g_A (10^8 nsec^{-1})	Wavelength (Å)	g_A (10^8 nsec^{-1})	Wavelength (Å)	g_A (10^8 nsec^{-1})
CrI					
2739.38	26.	3639.80	77.	3050.14	44.
2741.07	36.	3743.58	42.	3118.65	32.
2742.17	37.	3743.88	50.	3120.37	60.
2986.00	23.	3757.66	20.	3124.94	66.
2986.47	34.	3768.24	18.	3132.06	85.
3014.92	21.	3804.80	74.	3136.68	20.
3017.57	44.	3963.69	93.	3147.23	20.
3021.56	47.	3969.75	75.	3180.70	35.
3039.78	29.	3976.66	74.	3197.08	32.
3053.88	11.	3983.91	45.	3209.18	25.
3109.34	10.	3991.12	44.	3217.40	20.
3110.86	10.			3295.43	12.
3119.25	18.			3360.30	35.
3119.71	15.			3368.05	46.
3148.44	29.			3378.34	12.
		CrII			
3155.15	35.	2055.52	9.1	3379.83	21.
3163.76	35.	2061.49	7.3	3421.21	18.
3237.73	20.	2065.42	4.8	3422.74	25.
3251.84	41.	2677.16	132.		
3257.82	41.	2800.77	24.		
3259.98	31.	2812.01	27.		
3346.02	29.	2818.36	21.		
3346.74	28.	2822.01	15.		
3433.60	44.	2822.37	56.		
3436.19	26.	2830.47	55.		
3445.62	28.	2840.02	36.		
3453.33	30.	2851.36	39.		
3455.60	20.	2897.67	19.		
3467.02	11.	2911.68	16.		
3467.72	19.	2921.24	33.		
3469.59	13.	2927.08	23.		
3481.30	14.	2928.15	25.		
3481.54	15.	2930.85	11.		
3494.97	15.	2935.14	28.		
3510.54	21.	2940.22	17.		
3566.16	38.	2946.84	22.		
3578.69	8.3	2971.91	66.		
3593.49	7.0	2979.74	61.		
3605.33	5.2	2985.32	36.		
3636.59	43.	2989.19	49.		

The Lifetimes: Experimental Procedure

In addition to the apparatus used for the spectral scans, the lifetime experiment required the use of a stepping motor and a control box (see Figure III). The stepping motor was used to step the foil up and down stream, the length of one step being 0.258 mm. The control box provided stepping voltage to the motor and a channel advance to the TMC. Since all lines chosen for analysis were in the region of 3000 Å, the grating blazed at 3000 Å was used for the lifetime runs. The foils used were measured by the Arizona Foil Company to be 2.0 $\mu\text{gm}/\text{cm}^2$ in thickness.

A lifetime measurement was made at 85 and 100 keV for each of the lines. Spectral lines were located by setting the monochromator for the appropriate wavelength and maximizing the count rate meter in that wavelength region. The slit widths were set at 100 or 150 microns depending upon the resolution required. The foil and the TMC were stepped at the same rate, i.e. 1.04 sec/channel and 1.04 sec/motor step. At this rate one run of 400 channels took approximately six minutes.

One run consisted, in most cases, of stepping the foil upstream a distance of approximately 3 cm and back downstream the same distance. This produced two decay curves for one run. In some cases the counting rate was such that the reversal was unnecessary. During each run the chamber was isolated from the accelerator for several seconds in order to measure the dark count. The dark count was measured to be from four to six counts per channel.

The beam current was monitored continuously in order to note any large fluctuations in the amount of beam delivered to the foil. When

changes of more than 15% occurred the run was discontinued. During most runs the beam current was quite steady and fluctuated no more than 5%. Because of the hydrocarbon buildup mentioned in the energy loss experiment the foils were not used for more than 15 min. The time elapsed on each foil was noted and is stated in the heading of each computer output (see the Appendix).

Lifetimes: The Results

Lifetime data was taken on five lines originating from four atomic levels in CrII. Figure XIII is a partial Grotrian Diagram showing the transitions whose lifetimes were measured and possible cascades into these levels of interest. A measurement was made on each line at two energies (85 and 100 keV). Each decay curve was analyzed separately, and an average over all of the data for each level was used to obtain the stated lifetimes.

The lifetimes were calculated from the decay curves using a four parameter least-squares fit to the double exponential function

$$F(t) = A \exp(-t/\tau_1) + B \exp(-t/\tau_2) + D,$$

where D is the measured dark count, and A and B are the amplitudes of the decays. The computer program used was written by Brand, Fox, and Keller. The Appendix is composed of photographic reductions of the final computer output.

Brand¹⁴ has indicated that the fitting program will resolve a short component and a longer lived "cascade" component, but that the program will not resolve two short lived components. Therefore it is imperative that one know whether or not strong cascading occurs. Table III provides

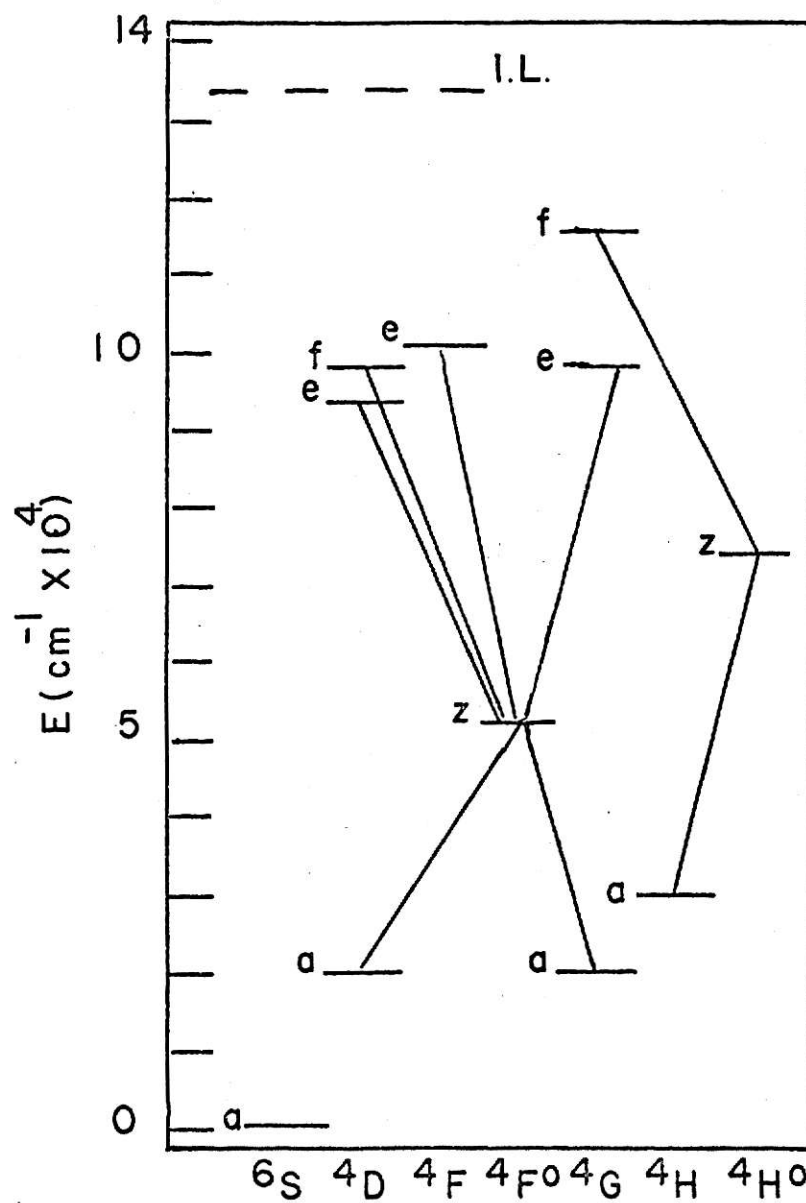


Fig. XIII.--A partial Grotrian Diagram of CrII

Table III

Possible cascades into levels of interest

Wavelength	I	Level	Configuration
Transitions into		$z^4H^o_{6\frac{1}{2}}$	$3d^4(a^3H)4p$
2415.23	5w	$f^4G_{5\frac{1}{2}}$	$3d^4(a^3G)5s$
Transitions into		$z^4H^o_{5\frac{1}{2}}$	$3d^4(a^3H)4p$
2404.72	2w	$f^4G_{5\frac{1}{2}}$	$3d^4(a^3G)5s$
2408.02	3w	$f^4G_{4\frac{1}{2}}$	$3d^4(a^3G)5s$
Transitions into		$z^4F^o_{4\frac{1}{2}}$	$3d^4(a^5D)4p$
3049.49	10w,1	$e^4D_{3\frac{1}{2}}$	$3d^4(a^5D)5s$
2569.40	152,1	$e^4F_{4\frac{1}{2}}$	$3d^4(a^5D)4d$
2577.74	10w	$e^4F_{3\frac{1}{2}}$	$3d^4(a^5D)4d$
2661.22	50w	$e^4G_{5\frac{1}{2}}$	$3d^4(a^5D)4d$
2674.26	7w	$e^4G_{4\frac{1}{2}}$	$3d^4(a^5D)4d$
Transitions into		$z^4F^o_{3\frac{1}{2}}$	$3d^4(a^5D)4p$
3056.66	8w,1	$e^4D_{2\frac{1}{2}}$	$3d^4(a^5D)5s$
2642.60	2w	$f^4D_{3\frac{1}{2}}$	$3d^4(a^5D)4d$
2652.78	3w,1	$f^4D_{2\frac{1}{2}}$	$3d^4(a^5D)4d$
2567.50	5w	$e^4F_{3\frac{1}{2}}$	$3d^4(a^5D)4d$
2576.45	2w	$e^4F_{2\frac{1}{2}}$	$3d^4(a^5D)4d$
2663.28	30w,1	$e^4G_{4\frac{1}{2}}$	$3d^4(a^5D)4d$
2674.07	8w	$e^4G_{3\frac{1}{2}}$	$3d^4(a^5D)4d$

a listing of all transitions into the levels studies given by Keiss.¹⁵ There are no lines listed that occur in the spectral scans with a significant counting rate. The generally low intensities given by Keiss for these lines also suggest that one should not expect short lived cascades. Under these circumstances it was assumed that no strong cascades need be considered.

Table IV provides a summary of the lifetime analyses. The average value for each line are given for 85 and 100 keV, and an average lifetime for each level is stated. Note that there is no systematic relationship between the measured lifetime and the energy at which it was taken. This indicates that if a trend of this kind is present, its magnitude must be within the limits of the experiment. The errors quoted in Table V are based on the standard deviations of the lifetime measurements, and on an estimated possible error in the energy loss of 50%. The standard deviations ran between 0.1 and 0.4 nsec., while the error in energy loss produced an error in the lifetimes of 0.2 nsec.

Table IV.

The lifetimes measured at 85 and 100 keV for five lines

Wavelength (Å)	Level	Lifetimes (nsec.)		Level
		85 keV	100 keV	
2971.9	${}^4\text{H}^{\text{O}}6\frac{1}{2}$	5.3	4.6	5.0 ± 0.6
2979.7	${}^4\text{H}^{\text{O}}5\frac{1}{2}$	4.9	4.7	4.8 ± 0.3
3124.9	${}^4\text{F}^{\text{O}}3\frac{1}{2}$	5.0	4.9	5.0 ± 0.4
3132.1	${}^4\text{F}^{\text{O}}4\frac{1}{2}$	4.9	5.2	4.9 ± 0.4
3180.8	${}^4\text{F}^{\text{O}}4\frac{1}{2}$	4.8	4.7	

Since branching ratio measurements are needed before gf values can be derived for the transitions, the BF data cannot be compared directly to previous measurements of gf-values. To avoid this difficulty, lifetimes were derived from all gA-values given by CB for transitions out of the level of interest. A comparison was then made on the basis of these lifetimes (see Table V). These lifetimes were derived by summing all gA-values for transitions out of the level of interest and dividing by the g of the upper level. This is the total transition rate out of the level and its inverse is the lifetime of the level. The ratio of these lifetimes to BFS lifetimes should be equivalent to a ratio of gf-values. Thus the ratios of gf-values given in Table V compared to the ratio of lifetimes provides a method of inter-comparing our lifetime results with gf-values measured by other authors. Shackleford¹⁴ has measured a gf-value for one line and Byard¹⁵ has measured the gf-values for two lines originating from the $z^4F^o_{4\frac{1}{2}}$ level. It may be seen that these results give gf/gf(CB) in fair agreement with our $\tau(CB)/\tau(BF)$ for that level.

One should note that the ratio of lifetimes is a factor of two higher for the lines at 2971.9 Å and 2979.7 Å. The reason for this discrepancy may be that some transitions out of the z^4H^o term were not seen by Corliss and Bozman. If this were true $\tau(CB)$ would be reduced, producing a more uniform ratio of lifetimes. In particular, lines at 2297.17 and 2307.19 Å were not observed in the present spectrum nor does CB list gA-values for them, but Keiss lists them with substantial intensities (Table VII-VIII). This result may indicate experimental problems in observing this part of the spectrum (2250 to 2650 Å) since there seems to be no structure in this region (see

Figure V). Although, the same arguments could be made for the z^4F^o levels, it is believed not to be the case since there is good agreement with gf-values measured by Schackelford and Byard gf-values for those transitions.

Table V.

A comparison of the measured lifetimes to the results of CB, Shackleford, and Byard

Wavelength (A)	$\tau^{(BF)}$ (nsec)	$\tau^{(CB)}$ (nsec)	$\tau^{(CB)}/\tau^{(BF)}$	gf(OSU)/gf(CB)	gf(S)/gf(CB)
2971.9	5.0	1.8	.36		
2979.7	4.8	1.7	.38		
3124.9	5.0	0.8	.16		
3132.1				.09	.10
	4.9	0.7	.14		
3180.8				.16	

The Solar Photospheric Abundance of Chromium

Two measurements are necessary before the calculation of the chromium abundance is possible. First, the branching ratios for the measured levels must be found, and second, the equivalent widths for the lines out of these levels must be measured in the solar spectrum. The branching ratios can easily be measured with conventional light sources. Problems exist in the measurement of the equivalent width of ultraviolet lines in the solar spectrum. The intensity of the continuum in the UV region is low and the number of lines is large. Thus the exact height of the continuum is not well known. This is not as much of a problem in the visible region and, as can be seen in Tables VI-IX,

Table VI
Transitions from the $z^4\text{H}^{\circ}_{5\frac{1}{2}}$ ($3d^4(a^3\text{H})4p$) level

Wavelength	I	gA	Level	Configuration	Solar Eq. Width
3552.42	2		$b^2\text{H}_{5\frac{1}{2}}$	$3d^5$	+
3540.23	2		$b^2\text{H}_{4\frac{1}{2}}$	$3d^5$	+
3442.98	1		$a^2\text{H}_{5\frac{1}{2}}$	$3d^4(a^3\text{H})4s$	+
3421.62 ^(a)	5		$a^2\text{H}_{4\frac{1}{2}}$	$3d^4(a^3\text{H})4s$	2.3 min.
3315.28	12		$b^4\text{G}_{5\frac{1}{2}}$	$3d^4(a^3\text{G})4s$	blended
3307.02	50	9.2	$b^4\text{G}_{4\frac{1}{2}}$	$3d^4(a^3\text{G})4s$	blended
3225.44	8		$b^4\text{F}_{4\frac{1}{2}}$	$3d^5$	blended
3063.82	7		$a^4\text{F}_{4\frac{1}{2}}$	$3d^4(a^3\text{H})4s$	78 m
2988.04	12		$a^4\text{H}_{6\frac{1}{2}}$	$3d^4(a^3\text{H})4s$	blended
2979.73*	80	61	$a^4\text{H}_{5\frac{1}{2}}$	$3d^4(a^3\text{H})4s$?
2972.67	7w		$a^4\text{H}_{4\frac{1}{2}}$	$3d^4(a^3\text{H})4s$?
2307.19	35		$a^4\text{G}_{4\frac{1}{2}}$	$3d^5$	+
2306.81	10		$a^4\text{G}_{5\frac{1}{2}}$	$3d^5$	+

*Lifetime measured in this experiment.

+Not observed according to MMH.

?Principal contributor to blend.

(a) $a^2\text{H}_{5\frac{1}{2}}-z^4\text{H}^{\circ}_{6\frac{1}{2}}$ may be present.

Table VII
Transitions from the $z^4\text{H}^{\text{O}}_{6\frac{1}{2}}$ ($3d^4(a^3\text{H})4p$) level

Wavelength	I	gA	Level	Configuration	Solar Eq. Width
4195.40	4		$b^2\text{I}_{6\frac{1}{2}}$	$3d^4(a^1\text{I})rs$	11 m
3529.73	2		$b^2\text{H}_{5\frac{1}{2}}$	$3d^5$	+
3421.62 ^(a)	5		$a^2\text{H}_{5\frac{1}{2}}$	$3d^4(a^3\text{H})4s$	23
3295.42	50	12	$b^4\text{G}_{5\frac{1}{2}}$	$3d^4(a^3\text{G})4s$	53
2971.90*	75	66	$a^4\text{H}_{6\frac{1}{2}}$	$3d^4(a^3\text{H})4s$?
2950.69	7		$a^2\text{I}_{6\frac{1}{2}}$	$3d^5$	+
2950.10 ^(b)	10		$a^2\text{I}_{5\frac{1}{2}}$	$3d^5$	+
2297.17	50		$a^4\text{G}_{5\frac{1}{2}}$	$3d^5$	+

*Lifetime measured in this experiment.

+Not observed by M.M.H.

?Principal contributor to blend.

(a) $a^2\text{H}_{4\frac{1}{2}} - z^4\text{H}^{\text{O}}_{5\frac{1}{2}}$ transition may also be present.

(b) $b^4\text{G}_{4\frac{1}{2}} - z^2\text{I}^{\text{O}}_{5\frac{1}{2}}$ transition may also be present.

Table VIII
Transitions from the $z^4F^o_{3\frac{1}{2}}$ ($3d^4(a^3D)4p$) level

Wavelength	I	gA	Level	Configuration	Solar Eq. Width
5502.07	40		$b^4G_{4\frac{1}{2}}$	$3d^4(a^3G)4s$	23
5472.60	12		$b^4G_{3\frac{1}{2}}$	$3d^4(a^3G)4s$	+
5280.08	30		$b^4F_{2\frac{1}{2}}$	$3d^5$	blended
5279.88	50		$b^4F_{4\frac{1}{2}}$	$3d^5$	18
5274.99	75		$b^4F_{3\frac{1}{2}}$	$3d^5$	blended
4860.20	20		$a^4F_{4\frac{1}{2}}$	$3d^4(a^3F)4s$	9
4848.24	75		$a^4F_{3\frac{1}{2}}$	$3d^4(a^3F)4s$	52
4836.22	25		$a^4F_{2\frac{1}{2}}$	$3d^4(a^3F)4s$	blended
3738.38	35		$c^4D_{2\frac{1}{2}}$	$3d^4(a^3D)4s$	blended
3736.56	1		$b^4D_{3\frac{1}{2}}$	$3d^5$	+
3336.16	2		$a^4P_{2\frac{1}{2}}$	$3d^5$	+
3197.08	75	32	$a^4G_{4\frac{1}{2}}$	$3d^5$	blended
3196.93	20		$a^4G_{3\frac{1}{2}}$	$3d^5$	blended
3196.35	3		$a^4G_{2\frac{1}{2}}$	$3d^5$	39
3147.22 ^(a)	50	20	$a^4D_{3\frac{1}{2}}$	$3d^4(a^5D)4s$	blended
3124.94*	40	66	$a^4D_{2\frac{1}{2}}$	$3d^4(a^5D)4s$	149
2544.26	15		$a^6D_{4\frac{1}{2}}$	$3d^4(a^5D)4s$	+
2531.84	25		$a^6D_{3\frac{1}{2}}$	$3d^4(a^5D)4s$	+

*Lifetime measured in this experiment.

+Not observed by MMH.

?Principal contributor to blend.

(a) $b^4G_{4\frac{1}{2}} - z^4G_{4\frac{1}{2}}$ transition may be present.

Table IX
Transitions from the $z^4F^o_{4\frac{1}{2}}$ ($3d^4(a^3D)4p$) level

Wavelength	I	gA	Level	Configuration	Solar Eq. Width
7311.60	2		$c^4D_{3\frac{1}{2}}$	$3d^4(a^3D)4s$	5
6478.37	50		$b^4G_{5\frac{1}{2}}$	$3d^4(a^3G)4s$	+
5455.86	8		$b^4G_{4\frac{1}{2}}$	$3d^4(a^3G)4s$	+
5237.35	100		$b^4F_{4\frac{1}{2}}$	$3d^5$	49
5232.54	20		$b^4F_{3\frac{1}{2}}$	$3d^5$	12
4824.12	100		$a^4F_{4\frac{1}{2}}$	$3d^4(a^3F)4s$	94
4812.34	25		$a^4F_{3\frac{1}{2}}$	$3d^4(a^3F)4s$	41
3715.18	25		$b^4D_{3\frac{1}{2}}$	$3d^5$	58
3181.42	20	4.6	$a^4G_{4\frac{1}{2}}$	$3d^5$	blended
3180.70*	75	35	$a^4G_{5\frac{1}{2}}$	$3d^5$	blended
3132.05*	100	85	$a^4D_{3\frac{1}{2}}$	$3d^4(a^5D)4s$	137
2534.33	40	6.4	$a^6D_{4\frac{1}{2}}$	$3d^4(a^5D)4s$	+
2522.01	4		$a^6D_{3\frac{1}{2}}$	$3d^4(a^5D)4s$	+

*Lifetime measured in this experiment.

+Not observed by M.M.H.

?Principal contributor to blend.

which are a compilation from Keiss¹⁴ and Moore, Minnaert, and Houtgast¹⁸, lines from the levels measured do exist in the visible. Thus, a relatively straight forward experiment to measure the branching ratios will produce many gf values for lines which may be useful in the solar abundance problem.

Although, the measured lifetimes alone do not determine the abundance unambiguously, one may derive an approximate abundance by comparison. Warner^{19,20} has measured gf values in CrII and has derived an abundance of $\log N_{\text{Cr}} = 5.47$, where $\log N_{\text{H}} = 12$, based on those gf-values. Warner's gf-values may not be compared to the BF lifetimes directly, as there are no lines in common, but a comparison of both measurements to CB will produce a correction to the abundance given by Warner.

Table X compares gf-values taken from CB with those of several authors and with the results of this experiment. The value of 0.82 dex for the CB-BF comparison was arrived at on the assumption that strong transitions out of the $z^4\text{H}^0$ levels do exist and were not observed by CB. The correction is thus the log of the average of the two ratios listed in Table VI for the $z^4\text{F}^0$ transitions.

The value obtained by Cocke, Stark, and Evans³ based on BFS lifetimes for neutral chromium is 5.80. The correction of 0.82 dex from the present study would indicate a value of 5.54 for $\log N_{\text{Cr}}$. Although the branching ratios must be measured before any firm conclusions are possible, one might seek an explanation for the large disparity in the gf-values of $\log N_{\text{Cr}}$, either in the ionization equilibrium for chromium or in the abundance of 5.47 given by Warner.

Table X

A comparison of the measured lifetimes
to the g_A values of others

$\log gf(\text{CB-Warner})$	0.75 dex
$\log gf(\text{CB-Byard})$	0.89 dex
$\log gf(\text{CB-Shackleford})$	0.95 dex
$\log gf(\text{CB-Huber and Tobey})$	0.94 dex
$\log gf(\text{CB-Wolnik et al})$	0.60 dex
$\log (\text{BF-CB})$	0.82 dex

Aknowledgements

The author wishes to thank Dr. C. L. Cocke for his valuable assistance in this study. He wishes to thank Dr. B. Curnutte for his advice and assistance. Dr. J. C. Evans helped immeasurably in the preparation of this work. Mr. Fox, Mr. Keller and Dr. Brand were responsible for the computer programs used in the data analysis and in the stopping power calculations. The author wishes to thank the Atomic Energy Commission whose support he enjoyed for the duration of this research.

FOOTNOTES

- ¹Kay, L., Phys. Letters 5, 36 (1962).
- ²Warner B. Monthly Notices R.A.S. 138, 229 (1968).
- ³Cocke, C. L., A. Stark, and J. C. Evans to be published.
- ⁴Andersen, T., K. A. Jesson, G. Sorensen, Nud. Instr., Methods 90, H1 (1970).
- ⁵Shore, B. W. and D. H. Menzel, Principles of Atomic Spectra (John Wiley and Sons, Inc., New York) 1968, p. 110 ff.
- ⁶Ibid., p. 430 ff.
- ⁷Motz, L. Astrophysics and Stellar Structure. (Ginn and Company, Waltham Mass.) 1970 p. 199 ff.
- ⁸Lindhard, J., M. Scharff, and H. E. Schiott, Mgt. Fgo. Medd. Dan. Vid. Selssk. 33 No. 14 (1963).
- ⁹Northcliffe, L., Ann. Rev. Nud. Sci. 13, 67 (1963).
- ¹⁰Brand, J. H., Thesis, Kansas State University (1972) p. 160 ff.
- ¹¹Corliso, C. H. and W. R. Bozman, N.B.S. Monograph, No. 53 (1962).
- ¹²Moore, C. E., A Multiplet Table of Astrophysical Interest N.B.S. Technical Note No. 36 Nov. (1959).
- ¹³Moore, C. E., An Ultraviolet Multiplet Table, Circular of the National Bureau of Standards 488 (1950).
- ¹⁴Brand, p. 55 ff.
- ¹⁵Kiess, C. C., Journal of Research of the N.V.S. 47, 385 (1951).
- ¹⁶Shackleford, W. L., J.Q.S.R.T., 5, 303 (1965).

¹⁷Byard, P., J.Q.S.R.T., 8, 1543 (1968).

¹⁸Moore, C. E., M. G. J. Minnaert, and J. Houtgast
N.B.S. Monograph, No. 61 (1966).

¹⁹Warner, B., Memoirs R.A.S. 70, 165 (1967).

²⁰Warner, B., Obs. 89, 107 (1969).

Figure XIV. A. Part one of the 100 keV lifetime fit for the
line at 2971.9 Å

ILLEGIBLE DOCUMENT

**THE FOLLOWING
DOCUMENT(S) IS OF
POOR LEGIBILITY IN
THE ORIGINAL**

**THIS IS THE BEST
COPY AVAILABLE**

2071.0A RUN #2 PART 1 FULL #1 TIME ON POIL 9 TO 16 MIN 100 KW
 MASS= 52. ENERGY= 87. KEV 1.7666 NSEC/CM DARK COUNTS= 4.
 THERE WERE 0.29800 CM/CHANNEL
 START FIT IN CHANNEL NO 17

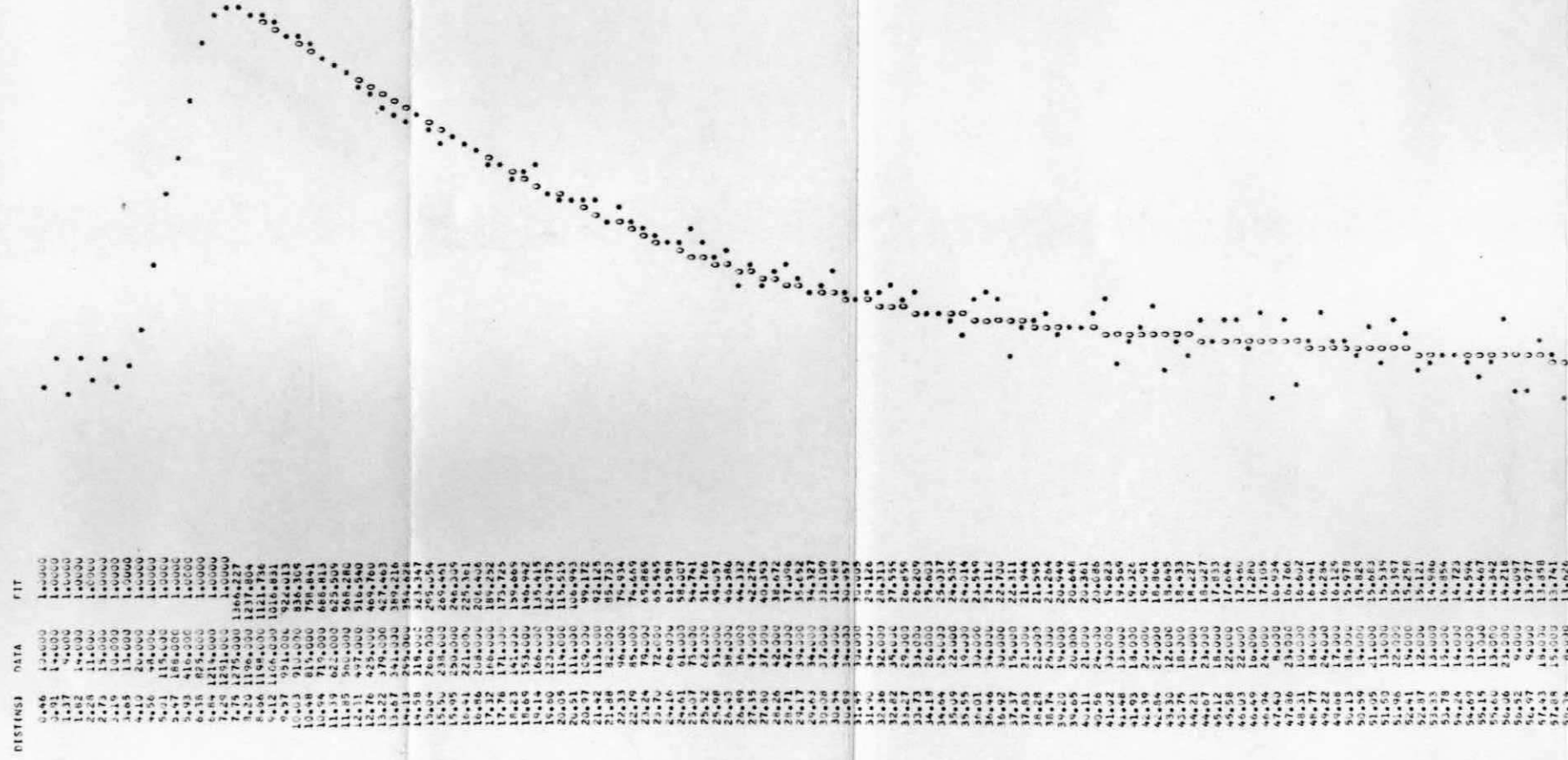


Figure XIV. B. Part two of the 100 keV lifetime fit for the
line at 2971.9 Å

2021-9A RUN #2 PART 2 FOIL #1 TIME ON FOIL 9 TO 10 MIN 100 MV
 52° ENERGY 87. KEV 1/0= 1.7666 NSEC/MV DARK COUNTS= 4.
 THESE VALUES WERE USED TO CORRECT THE DATA
 START FIT IN CHANNEL NO. 13

DIST(INS)	DATA	FIT
0.46	7.000	1.0000
0.51	10.000	1.0000
0.57	12.000	1.0000
1.82	15.000	1.0000
2.28	9.000	1.0000
2.73	9.000	1.0000
3.19	11.000	1.0000
3.65	9.000	1.0000
4.10	23.000	1.0000
4.56	41.000	1.0000
5.02	157.000	1.0000
5.47	187.000	1.0000
5.93	381.000	1.0000
6.38	729.000	1.0000
6.84	995.000	1.0000
7.29	957.000	1.0000
7.75	981.000	1.0000
8.20	958.000	955.105
8.66	877.000	878.451
9.12	787.000	789.117
9.57	717.000	720.557
10.03	644.000	660.049
10.48	590.000	602.289
10.94	554.000	545.789
11.39	499.000	502.068
11.85	428.000	459.687
12.31	386.000	410.251
12.76	351.000	363.397
13.22	351.000	355.372
13.67	330.000	321.163
14.13	285.000	294.212
14.58	253.000	269.703
15.04	231.000	247.412
15.50	223.000	227.137
15.95	176.000	160.694
16.41	176.000	174.645
16.86	156.000	162.756
17.32	152.000	150.113
17.78	120.000	138.603
18.23	160.000	128.125
18.69	105.000	118.583
19.14	112.000	109.893
19.60	104.000	104.747
20.05	108.000	94.767
20.51	108.000	88.190
20.97	100.000	82.156
21.42	87.000	76.731
21.88	75.000	71.746
22.33	106.000	67.196
22.79	78.000	63.043
23.24	68.000	58.231
23.70	48.000	52.619
24.16	40.000	49.723
24.61	40.000	47.073
25.07	42.000	44.648
25.52	62.000	42.426
25.98	38.000	40.389
26.43	61.000	38.521
26.89	49.000	35.230
27.34	35.000	33.782
27.80	35.000	32.445
28.26	35.000	31.221
28.71	36.000	30.090
29.17	35.000	29.045
29.63	40.000	28.079
30.08	32.000	26.355
30.54	31.000	25.592
31.00	31.000	24.880
31.45	30.000	24.217
31.90	27.000	23.601
32.36	27.000	23.025
32.82	27.000	21.998
33.27	22.000	21.514
33.73	26.000	21.072
34.19	31.000	20.656
34.66	30.000	20.265
35.09	30.000	19.895
35.55	28.000	19.546
36.01	27.000	19.215
36.46	28.000	18.904
36.92	21.000	18.321
37.37	26.000	18.051
37.83	28.000	17.792
38.29	16.000	17.546
38.74	27.000	17.309
39.20	16.000	17.082
39.65	16.000	16.864
40.11	17.000	16.654
40.56	22.000	16.451
41.02	23.000	16.256
41.48	16.000	16.067
41.93	16.000	15.884
42.39	27.000	15.706
42.84	17.000	15.534
43.30	14.000	15.367
43.75	14.000	15.204
44.21	26.000	15.046
44.67	14.000	14.891
45.12	20.000	14.741
45.58	16.000	14.593
46.03	13.000	14.450
46.49	13.000	14.309
46.94	13.000	14.172
47.40	11.000	14.039
47.86	18.000	13.905
48.31	16.000	13.776
48.77	24.000	13.649
49.22	12.000	13.524
49.68	12.000	13.402
50.13	11.000	13.281
50.59	18.000	13.163
51.05	16.000	13.047
51.50	10.000	12.932
51.96	16.000	12.820
52.41	12.000	12.710
52.87	7.000	12.601
53.33	20.000	12.493
53.78	6.000	12.388
54.24	4.000	12.283
54.69	21.000	12.181
55.15	8.000	12.080
55.60	6.000	11.980
56.06	15.000	11.880
56.52	20.000	11.780
56.97	25.000	11.680
57.43	10.000	11.580
57.88	14.000	11.480
58.34	8.000	11.380

INTERCEPT-1 IS 9183.6820 SLOPE-1 IS -0.2106 TAU-1 IS 4.7475 NSEC
 INTERCEPT-2 IS 37.7253 SLOPE-2 IS -0.0267 TAU-2 IS 37.4831 NSEC
 THE DATA ARE FROM CHANNEL 17
 THE SLOPE COEFFICIENTS AT THE FOIL IS: 33.03666
 RUN PROCESSED 3 JULY 72.

CHISQR

1 1.5203650 00
 2 1.5061320 00
 3 1.4924210 00
 4 1.4794890 00

Figure XIV. C. Part one of the 85 keV lifetime fit for the
line at 2971.9 Å

2071.0A RUN #11 PART 1, FOIL #6 TIME ON FOIL 0 TO 4 MIN 95.4V
 TIME OFF 0.225901 72.4KV 1V= 1.0e-10 NSEC/CM
 START FIT IN CHANNEL NO 50

DISTINX	DATA	FIT
1.50	3.200	1.0000
1.51	1.000	1.0000
1.52	1.000	1.0000
1.53	2.000	1.0000
2.51	5.000	1.0000
3.51	7.000	1.0000
4.51	9.000	1.0000
5.51	1.0000	1.0000
6.51	2.000	1.0000
7.51	3.000	1.0000
8.51	4.000	1.0000
9.51	5.000	1.0000
10.51	6.000	1.0000
11.51	7.000	1.0000
12.51	8.000	1.0000
13.51	9.000	1.0000
14.51	1.0000	1.0000
15.51	2.000	1.0000
16.51	3.000	1.0000
17.51	4.000	1.0000
18.51	5.000	1.0000
19.51	6.000	1.0000
20.51	7.000	1.0000
21.51	8.000	1.0000
22.51	9.000	1.0000
23.51	1.0000	1.0000
24.51	2.000	1.0000
25.51	3.000	1.0000
26.51	4.000	1.0000
27.51	5.000	1.0000
28.51	6.000	1.0000
29.51	7.000	1.0000
30.51	8.000	1.0000
31.51	9.000	1.0000
32.51	1.0000	1.0000
33.51	2.000	1.0000
34.51	3.000	1.0000
35.51	4.000	1.0000
36.51	5.000	1.0000
37.51	6.000	1.0000
38.51	7.000	1.0000
39.51	8.000	1.0000
40.51	9.000	1.0000
41.51	1.0000	1.0000
42.51	2.000	1.0000
43.51	3.000	1.0000
44.51	4.000	1.0000
45.51	5.000	1.0000
46.51	6.000	1.0000
47.51	7.000	1.0000
48.51	8.000	1.0000
49.51	9.000	1.0000
50.51	1.0000	1.0000
51.51	2.000	1.0000
52.51	3.000	1.0000
53.51	4.000	1.0000
54.51	5.000	1.0000
55.51	6.000	1.0000
56.51	7.000	1.0000
57.51	8.000	1.0000
58.51	9.000	1.0000
59.51	1.0000	1.0000
60.51	2.000	1.0000
61.51	3.000	1.0000
62.51	4.000	1.0000
63.51	5.000	1.0000
64.51	6.000	1.0000
65.51	7.000	1.0000
66.51	8.000	1.0000
67.51	9.000	1.0000
68.51	1.0000	1.0000
69.51	2.000	1.0000
70.51	3.000	1.0000
71.51	4.000	1.0000
72.51	5.000	1.0000
73.51	6.000	1.0000
74.51	7.000	1.0000
75.51	8.000	1.0000
76.51	9.000	1.0000
77.51	1.0000	1.0000
78.51	2.000	1.0000
79.51	3.000	1.0000
80.51	4.000	1.0000
81.51	5.000	1.0000
82.51	6.000	1.0000
83.51	7.000	1.0000
84.51	8.000	1.0000
85.51	9.000	1.0000
86.51	1.0000	1.0000
87.51	2.000	1.0000
88.51	3.000	1.0000
89.51	4.000	1.0000
90.51	5.000	1.0000
91.51	6.000	1.0000
92.51	7.000	1.0000
93.51	8.000	1.0000
94.51	9.000	1.0000
95.51	1.0000	1.0000
96.51	2.000	1.0000
97.51	3.000	1.0000
98.51	4.000	1.0000
99.51	5.000	1.0000
100.51	6.000	1.0000

INTERCEPT IS 7.5550471 TIME=1 IS -0.1986 TIME=15 9.3024 NSEC
 INTERCEPT IS 7.6152 TIME=2 IS -0.0256 TIME=15 30.0156 NSEC
 THE DATA AT OR NEAR CHANNEL 20
 WERE DISREGARDED AT THE FOIL FOR 200.64253
 OWN POSITION 3.8 MV 7.5.

PHOTO

1 2.5602000 00
 2 1.0371670 00
 3 1.0371670 00
 4 1.0371670 00

Figure XIV. D. Part two of the 85 keV lifetime fit for the
line at 2971.9 Å

2071-04 RUN #11 DIST 2 CMTL #6 TIME ON ETL 0 TO 4 MTS 84.4V
 MISS# 877 ENVELOPE 72.4V 100% 1.6410 MSEC/CM MARK COMPARE 4.
 THREE YEAR 0.00000 MSEC/CM
 START PIT IN CHANNEL NO 10

DISTANCE	DATA	PIT
1.50	1.000	1.000
1.51	1.000	1.000
1.52	1.000	1.000
1.53	1.000	1.000
1.54	1.000	1.000
1.55	1.000	1.000
1.56	1.000	1.000
1.57	1.000	1.000
1.58	1.000	1.000
1.59	1.000	1.000
1.60	1.000	1.000
1.61	1.000	1.000
1.62	1.000	1.000
1.63	1.000	1.000
1.64	1.000	1.000
1.65	1.000	1.000
1.66	1.000	1.000
1.67	1.000	1.000
1.68	1.000	1.000
1.69	1.000	1.000
1.70	1.000	1.000
1.71	1.000	1.000
1.72	1.000	1.000
1.73	1.000	1.000
1.74	1.000	1.000
1.75	1.000	1.000
1.76	1.000	1.000
1.77	1.000	1.000
1.78	1.000	1.000
1.79	1.000	1.000
1.80	1.000	1.000
1.81	1.000	1.000
1.82	1.000	1.000
1.83	1.000	1.000
1.84	1.000	1.000
1.85	1.000	1.000
1.86	1.000	1.000
1.87	1.000	1.000
1.88	1.000	1.000
1.89	1.000	1.000
1.90	1.000	1.000
1.91	1.000	1.000
1.92	1.000	1.000
1.93	1.000	1.000
1.94	1.000	1.000
1.95	1.000	1.000
1.96	1.000	1.000
1.97	1.000	1.000
1.98	1.000	1.000
1.99	1.000	1.000
2.00	1.000	1.000
2.01	1.000	1.000
2.02	1.000	1.000
2.03	1.000	1.000
2.04	1.000	1.000
2.05	1.000	1.000
2.06	1.000	1.000
2.07	1.000	1.000
2.08	1.000	1.000
2.09	1.000	1.000
2.10	1.000	1.000
2.11	1.000	1.000
2.12	1.000	1.000
2.13	1.000	1.000
2.14	1.000	1.000
2.15	1.000	1.000
2.16	1.000	1.000
2.17	1.000	1.000
2.18	1.000	1.000
2.19	1.000	1.000
2.20	1.000	1.000
2.21	1.000	1.000
2.22	1.000	1.000
2.23	1.000	1.000
2.24	1.000	1.000
2.25	1.000	1.000
2.26	1.000	1.000
2.27	1.000	1.000
2.28	1.000	1.000
2.29	1.000	1.000
2.30	1.000	1.000
2.31	1.000	1.000
2.32	1.000	1.000
2.33	1.000	1.000
2.34	1.000	1.000
2.35	1.000	1.000
2.36	1.000	1.000
2.37	1.000	1.000
2.38	1.000	1.000
2.39	1.000	1.000
2.40	1.000	1.000
2.41	1.000	1.000
2.42	1.000	1.000
2.43	1.000	1.000
2.44	1.000	1.000
2.45	1.000	1.000
2.46	1.000	1.000
2.47	1.000	1.000
2.48	1.000	1.000
2.49	1.000	1.000
2.50	1.000	1.000
2.51	1.000	1.000
2.52	1.000	1.000
2.53	1.000	1.000
2.54	1.000	1.000
2.55	1.000	1.000
2.56	1.000	1.000
2.57	1.000	1.000
2.58	1.000	1.000
2.59	1.000	1.000
2.60	1.000	1.000
2.61	1.000	1.000
2.62	1.000	1.000
2.63	1.000	1.000
2.64	1.000	1.000
2.65	1.000	1.000
2.66	1.000	1.000
2.67	1.000	1.000
2.68	1.000	1.000
2.69	1.000	1.000
2.70	1.000	1.000
2.71	1.000	1.000
2.72	1.000	1.000
2.73	1.000	1.000
2.74	1.000	1.000
2.75	1.000	1.000
2.76	1.000	1.000
2.77	1.000	1.000
2.78	1.000	1.000
2.79	1.000	1.000
2.80	1.000	1.000
2.81	1.000	1.000
2.82	1.000	1.000
2.83	1.000	1.000
2.84	1.000	1.000
2.85	1.000	1.000
2.86	1.000	1.000
2.87	1.000	1.000
2.88	1.000	1.000
2.89	1.000	1.000
2.90	1.000	1.000
2.91	1.000	1.000
2.92	1.000	1.000
2.93	1.000	1.000
2.94	1.000	1.000
2.95	1.000	1.000
2.96	1.000	1.000
2.97	1.000	1.000
2.98	1.000	1.000
2.99	1.000	1.000
3.00	1.000	1.000
3.01	1.000	1.000
3.02	1.000	1.000
3.03	1.000	1.000
3.04	1.000	1.000
3.05	1.000	1.000
3.06	1.000	1.000
3.07	1.000	1.000
3.08	1.000	1.000
3.09	1.000	1.000
3.10	1.000	1.000
3.11	1.000	1.000
3.12	1.000	1.000
3.13	1.000	1.000
3.14	1.000	1.000
3.15	1.000	1.000
3.16	1.000	1.000
3.17	1.000	1.000
3.18	1.000	1.000
3.19	1.000	1.000
3.20	1.000	1.000
3.21	1.000	1.000
3.22	1.000	1.000
3.23	1.000	1.000
3.24	1.000	1.000
3.25	1.000	1.000
3.26	1.000	1.000
3.27	1.000	1.000
3.28	1.000	1.000
3.29	1.000	1.000
3.30	1.000	1.000
3.31	1.000	1.000
3.32	1.000	1.000
3.33	1.000	1.000
3.34	1.000	1.000
3.35	1.000	1.000
3.36	1.000	1.000
3.37	1.000	1.000
3.38	1.000	1.000
3.39	1.000	1.000
3.40	1.000	1.000
3.41	1.000	1.000
3.42	1.000	1.000
3.43	1.000	1.000
3.44	1.000	1.000
3.45	1.000	1.000
3.46	1.000	1.000
3.47	1.000	1.000
3.48	1.000	1.000
3.49	1.000	1.000
3.50	1.000	1.000
3.51	1.000	1.000
3.52	1.000	1.000
3.53	1.000	1.000
3.54	1.000	1.000
3.55	1.000	1.000
3.56	1.000	1.000
3.57	1.000	1.000
3.58	1.000	1.000
3.59	1.000	1.000
3.60	1.000	1.000
3.61	1.000	1.000
3.62	1.000	1.000
3.63	1.000	1.000
3.64	1.000	1.000
3.65	1.000	1.000
3.66	1.000	1.000
3.67	1.000	1.000
3.68	1.000	1.000
3.69	1.000	1.000
3.70	1.000	1.000
3.71	1.000	1.000
3.72	1.000	1.000
3.73	1.000	1.000
3.74	1.000	1.000
3.75	1.000	1.000
3.76	1.000	1.000
3.77	1.000	1.000
3.78	1.000	1.000
3.79	1.000	1.000
3.80	1.000	1.000
3.81	1.000	1.000
3.82	1.000	1.000
3.83	1.000	1.000
3.84	1.000	1.000
3.85	1.000	1.000
3.86	1.000	1.000
3.87	1.000	1.000
3.88	1.000	1.000
3.89	1.000	1.000
3.90	1.000	1.000
3.91	1.000	1.000
3.92	1.000	1.000
3.93	1.000	1.000
3.94	1.000	1.000
3.95	1.000	1.000
3.96	1.000	1.000
3.97	1.000	1.000
3.98	1.000	1.000
3.99	1.000	1.000
4.00	1.000	1.000

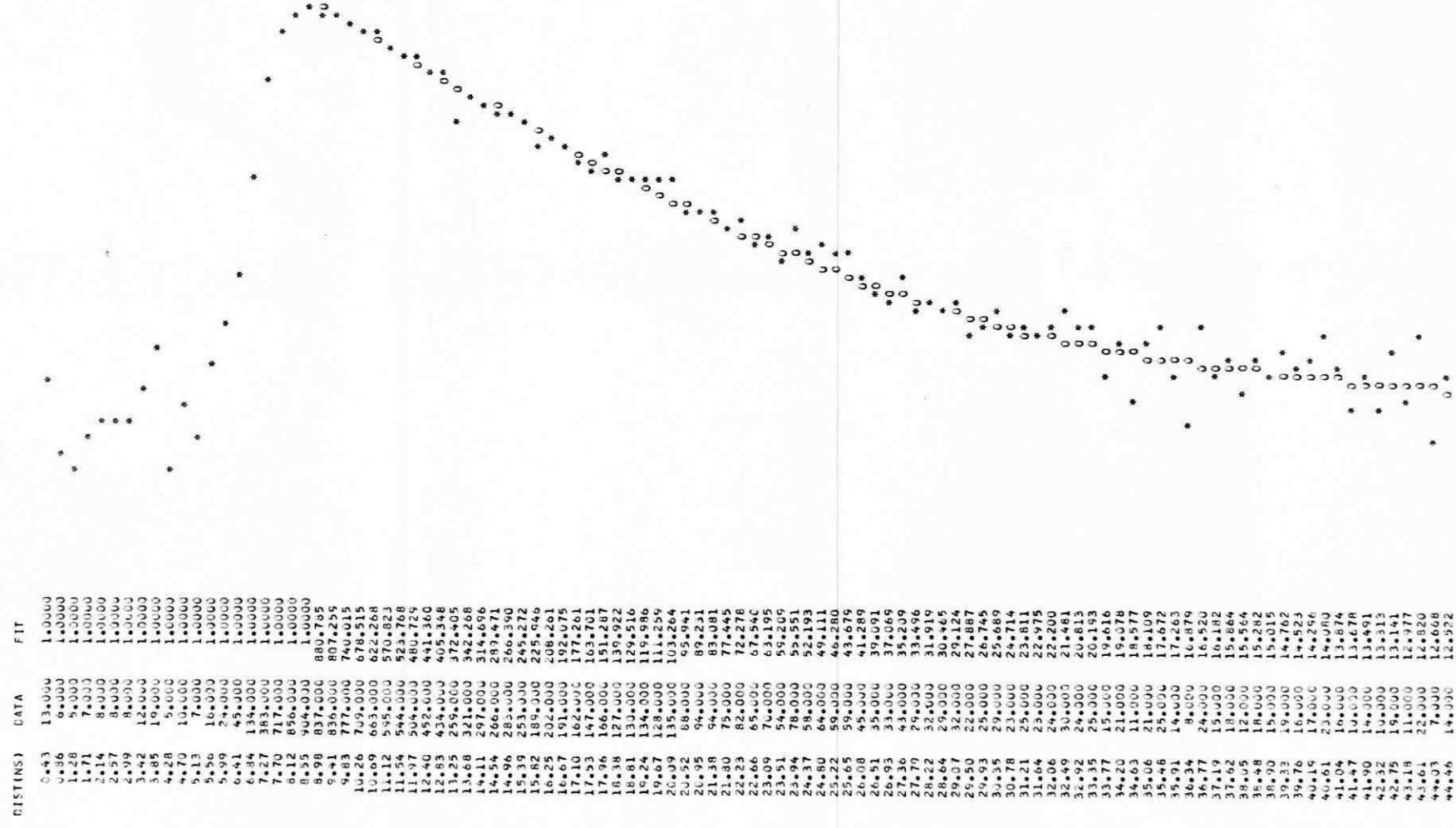
INTERFER-1 IS 26.72/661 SLIGHT-1 IS 40.000000 TIME-1 IS 4.000000
 INTERFER-2 IS 6.6673 SLIGHT-2 IS 40.000000 TIME-2 IS 4.000000
 THE DATA IS BY CH. 2000 CHANNEL 10
 THE RATIO OF COEFFICIENTS AT THE ETL IS 84.50157
 RUN SUCCESSFUL 3 JULY 72.

CHASOP

- 1 2.554000 01
- 2 1.460224 00
- 3 1.340000 00
- 4 1.580750 00

Figure XV. A. Part one of the 100 keV lifetime fit for the
line at 2979.7 Å

2079.7A RUN #3 PART 1 FOIL #2 TIME UN FOIL 0 TO 4 MIN 100 KV
 MASS= 52 ENERGY= 0.24200 MW/CHANNEL DARK COUNTS= 4
 THERE WERE 0.24200 MW/CHANNEL
 START FIT IN CHANNEL NO 21



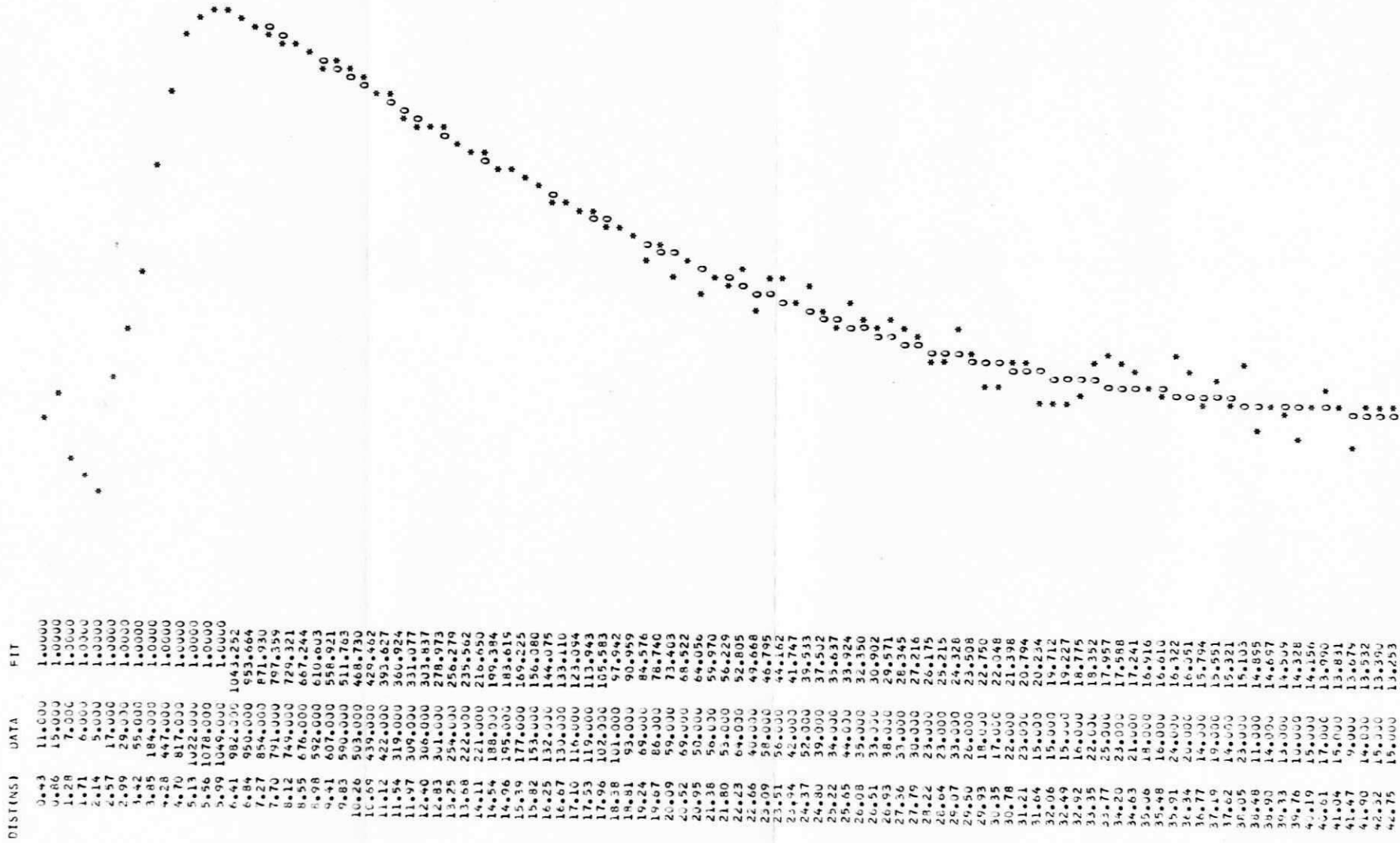
INTERCEPT-1 IS 5611.6818 SLOPE-1 IS -0.2096 TAU-1 IS 4.7704 NSEC
 INTERCEPT-2 IS 28.6466 SLOPE-2 IS -0.0236 TAU-2 IS 34.4206 NSEC
 THE PEAK IS AT OR NEAR CHANNEL 20
 THE RATIO OF COEFFICIENTS AT THE FOIL IS: 41.58316
 RUN PROCESSED 3 JULY 72.

GHISOR

1 2.10516350 JJ
 2 2.1048890 JJ
 3 2.1026920 JJ
 4 2.10151290 JJ

Figure XV. B. Part two of the 100 keV lifetime fit for the
line at 2979.7 Å

2979.7A RUN #3 PART 2 FOIL #2 TIME ON FOIL 0 TO 4 MIN 100 KV
 MASS= 52 ENERGY= 87 KEV 1/V= 1.7666 NSEC/MM DARK COUNTS= 4.
 THERE ARE 15 CHANNELS
 START FIT IN CHANNEL NJ 15



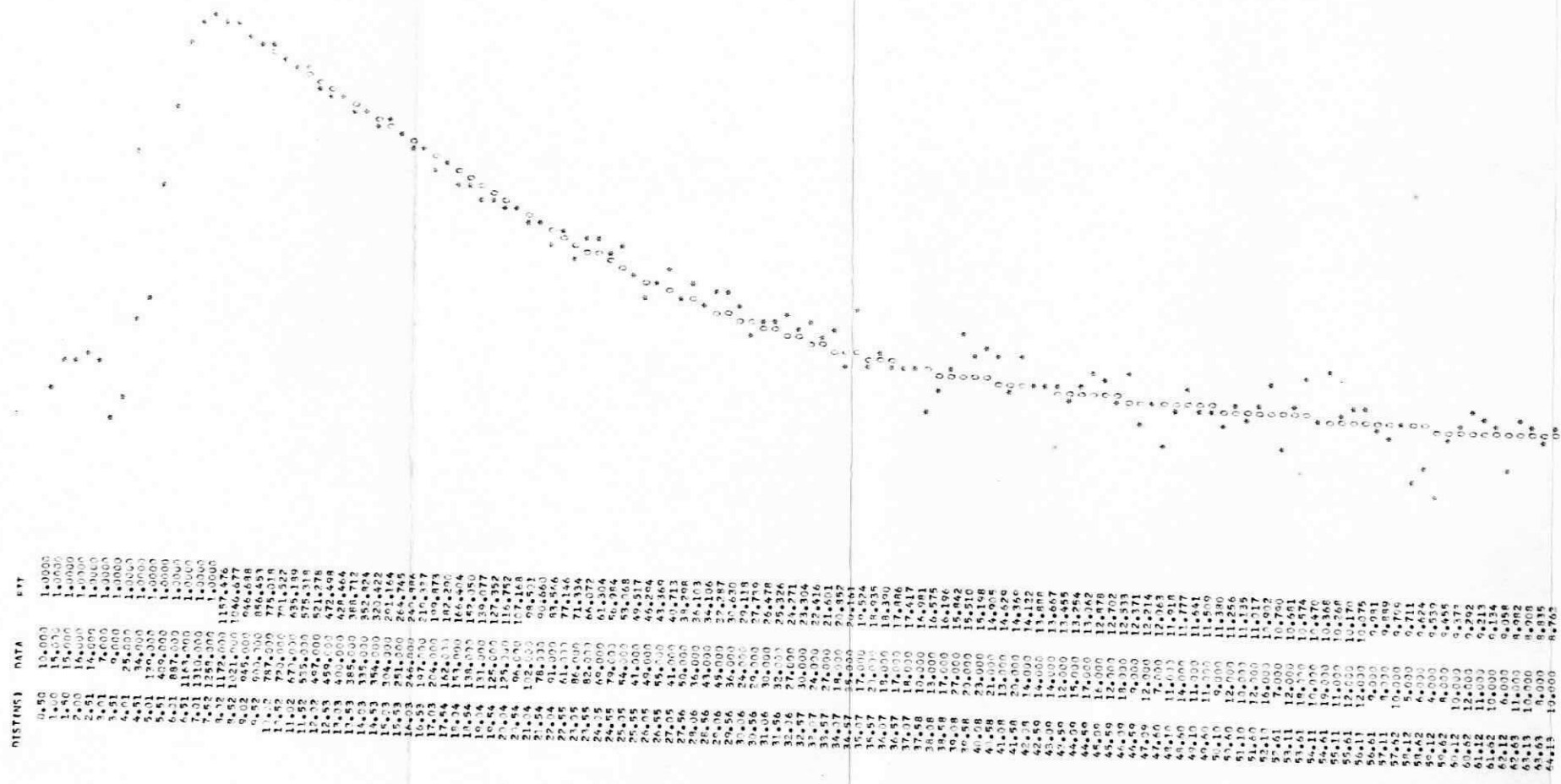
INTERCEPT-1 IS 4039.7828 SLOPE-1 IS -0.2151 TAU-1 IS 4.6484 NSEC
 INTERCEPT-2 IS 4046.71 SLOPE-2 IS -0.0256 TAU-2 IS 38.9981 NSEC
 THE PEAK IS AT OR NEAR CHANNEL 13
 THE RATIO OF COEFFICIENTS AT THE FOIL IS: 53.24791
 RUN PROCEEDED 3 JULY 72.

CHISQ

- 1 1.18099630 00
- 2 1.18099620 00
- 3 1.18099610 00
- 4 1.18099600 00

Figure XV. C. Part one of the 85 keV lifetime fit for the
line at 2979.7 A

2079.7 A RUN #10 PART 1 FOR 45 TIME ON PULS 4 TO 8 MIN
 MASS UP 82. ENERGY 72. KEV 1/4" 1.6410 SEC/CM
 START SET IN CHANNEL NO. 16



INTERCEPT-1 IS 5062.4438 SLOPE-1 IS -0.2095 TAU-1 IS 4.8556 NS/C
 INTERCEPT-2 IS 31.6477 SLOPE-2 IS -0.0294 TAU-2 IS 39.0603 NS/C
 THE PEAK IS AT OR NEAR CHANNEL 15
 THE DETECTOR COEFFICIENTS AT THE PULS 16 40.70569
 RUN PROCESSED 3 JULY 79.

CHISQR
 1 1.22075470 00
 2 1.22136670 00
 3 1.21342720 00
 4 1.20460920 00

Figure XV. D. Part two of the 85 keV lifetime fit for the
line at 2979.7 Å

2070.7 A. CHN 410. DART 2. ENI 45. TIME ON PCL 3 TO 8 MIN. PEAK
 MASS 52. ENERGY 77. KEV
 THREE WERE 0.24000. W/CHANNEL 1
 START FIT IN CHANNEL NO 19

DISTIN DATA FIT

0.50	1.000	1.000*
1.00	1.000	1.000*
1.50	1.000	1.000*
2.00	1.000	1.000*
2.50	1.000	1.000*
3.00	1.000	1.000*
3.50	1.000	1.000*
4.00	1.000	1.000*
4.50	1.000	1.000*
5.00	1.000	1.000*
5.50	1.000	1.000*
6.00	1.000	1.000*
6.50	1.000	1.000*
7.00	1.000	1.000*
7.50	1.000	1.000*
8.00	1.000	1.000*
8.50	1.000	1.000*
9.00	1.000	1.000*
9.50	1.000	1.000*
10.00	1.000	1.000*
10.50	1.000	1.000*
11.00	1.000	1.000*
11.50	1.000	1.000*
12.00	1.000	1.000*
12.50	1.000	1.000*
13.00	1.000	1.000*
13.50	1.000	1.000*
14.00	1.000	1.000*
14.50	1.000	1.000*
15.00	1.000	1.000*
15.50	1.000	1.000*
16.00	1.000	1.000*
16.50	1.000	1.000*
17.00	1.000	1.000*
17.50	1.000	1.000*
18.00	1.000	1.000*
18.50	1.000	1.000*
19.00	1.000	1.000*
19.50	1.000	1.000*
20.00	1.000	1.000*
20.50	1.000	1.000*
21.00	1.000	1.000*
21.50	1.000	1.000*
22.00	1.000	1.000*
22.50	1.000	1.000*
23.00	1.000	1.000*
23.50	1.000	1.000*
24.00	1.000	1.000*
24.50	1.000	1.000*
25.00	1.000	1.000*
25.50	1.000	1.000*
26.00	1.000	1.000*
26.50	1.000	1.000*
27.00	1.000	1.000*
27.50	1.000	1.000*
28.00	1.000	1.000*
28.50	1.000	1.000*
29.00	1.000	1.000*
29.50	1.000	1.000*
30.00	1.000	1.000*
30.50	1.000	1.000*
31.00	1.000	1.000*
31.50	1.000	1.000*
32.00	1.000	1.000*
32.50	1.000	1.000*
33.00	1.000	1.000*
33.50	1.000	1.000*
34.00	1.000	1.000*
34.50	1.000	1.000*
35.00	1.000	1.000*
35.50	1.000	1.000*
36.00	1.000	1.000*
36.50	1.000	1.000*
37.00	1.000	1.000*
37.50	1.000	1.000*
38.00	1.000	1.000*
38.50	1.000	1.000*
39.00	1.000	1.000*
39.50	1.000	1.000*
40.00	1.000	1.000*
40.50	1.000	1.000*
41.00	1.000	1.000*
41.50	1.000	1.000*
42.00	1.000	1.000*
42.50	1.000	1.000*
43.00	1.000	1.000*
43.50	1.000	1.000*
44.00	1.000	1.000*
44.50	1.000	1.000*
45.00	1.000	1.000*
45.50	1.000	1.000*
46.00	1.000	1.000*
46.50	1.000	1.000*
47.00	1.000	1.000*
47.50	1.000	1.000*
48.00	1.000	1.000*
48.50	1.000	1.000*
49.00	1.000	1.000*
49.50	1.000	1.000*
50.00	1.000	1.000*
50.50	1.000	1.000*
51.00	1.000	1.000*
51.50	1.000	1.000*
52.00	1.000	1.000*
52.50	1.000	1.000*
53.00	1.000	1.000*
53.50	1.000	1.000*
54.00	1.000	1.000*
54.50	1.000	1.000*
55.00	1.000	1.000*
55.50	1.000	1.000*
56.00	1.000	1.000*
56.50	1.000	1.000*
57.00	1.000	1.000*
57.50	1.000	1.000*
58.00	1.000	1.000*
58.50	1.000	1.000*
59.00	1.000	1.000*
59.50	1.000	1.000*
60.00	1.000	1.000*
60.50	1.000	1.000*
61.00	1.000	1.000*
61.50	1.000	1.000*
62.00	1.000	1.000*
62.50	1.000	1.000*
63.00	1.000	1.000*
63.50	1.000	1.000*
64.00	1.000	1.000*
64.50	1.000	1.000*
65.00	1.000	1.000*

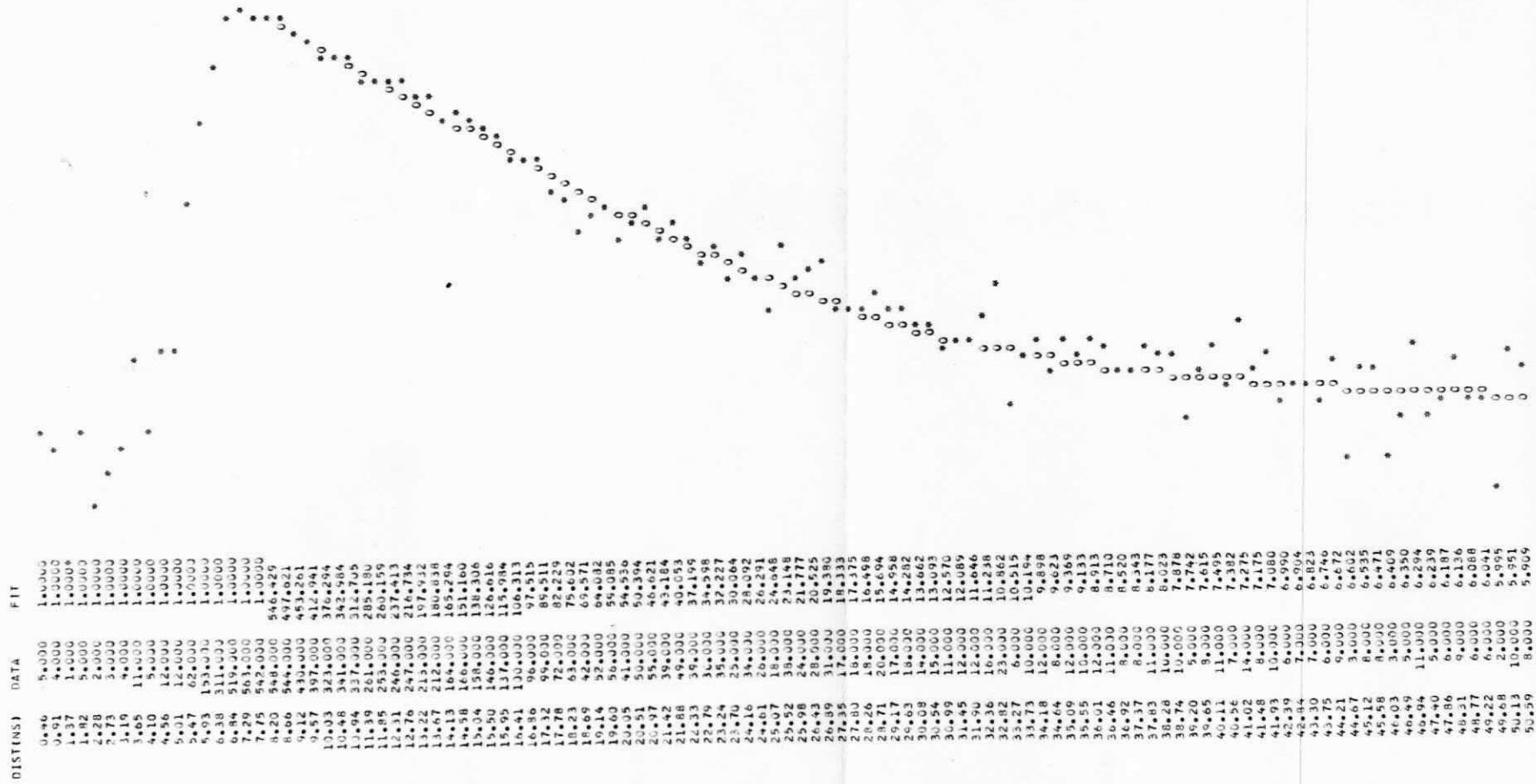
INTERCEPT-1 IS 7364.347 4100-1 IS -0.2058 TAU-1 IS 4.48502 NS/EV
 INTERCEPT-2 IS 39.3647 SLOPE-2 IS -0.0305 TAU-2 IS 32.7947 NS/EV
 THE PEAK IS AT 04.948 CHANNEL 18
 THE RATIO OF COEFFICIENTS AT THE PEAK IS 38.80667
 RUN PROCESSED 3 JULY 72.

CH1508

- 1 1.20461720 00
- 2 1.26416170 00
- 3 1.27373310 00
- 4 1.27276780 00

Figure XVI. A. Part one of the 100 keV lifetime fit for the
line at 3124.9 A

3145A RUN #4 PART 1 FILE #2 TIME ON FOIL 4 TO 10 MIN 100 KV
 MASS= 52. ENERGY= 87. KEV 1762 1.7666 NSIC/MM DARK COUNTS= 4.
 THERE WERE 325600 MM/CHANNEL
 START FIT IN CHANNEL NO 18



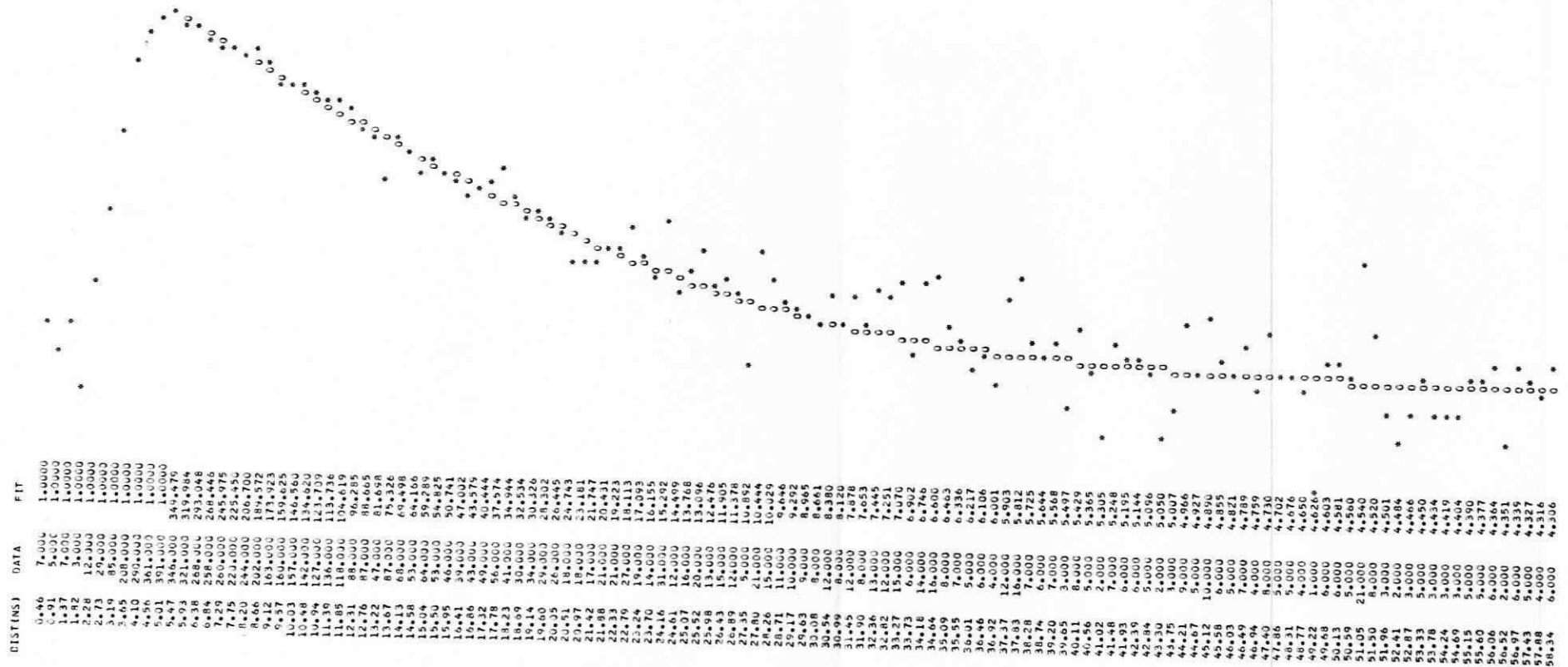
INTERCEPT-1 IS 2936.1560 SLOPE-1 IS -0.2104 TAU-1 IS 4.7537 NSIC
 INTERCEPT-2 IS 15.2183 SLOPE-2 IS -0.0418 TAU-2 IS 23.9313 NSIC
 THE PEAK IS AT OR NEAR CHANNEL 16
 THE RATIO OF COEFFICIENTS AT THE FOIL IS: 57.39299
 RUN PROCESSED 3 JULY 72.

CHISQR

1 1.54311980 00
 2 1.54181400 00
 3 1.54062530 00
 4 1.53949820 00

Figure XVI. B. Part two of the 100 keV lifetime fit for the
line at 3124.9 Å

3124.04 RUN #4 PART 2 FOIL #2 TIME ON FOIL 4 TO 10 MIN
MISS= 52 ENERGY= 87.0 KEV LUM= 1.7066 NSEC/AM DARK COUNTS= 4.
THERE WERE 0.23900 MM/CHANNEL
START FIT IN CHANNEL NO 12



INTERCEPT-1 IS 996.2389 SLOPE-1 IS -0.2011 TAU-1 IS 4.9737 NSEC
INTERCEPT-2 IS 20.4297 SLOPE-2 IS -0.0725 TAU-2 IS 13.7939 NSEC
THE PEAK IS AT OR NEAR CHANNEL 11
THE RATIO OF COEFFICIENTS AT THE FOIL IS: 25.59626
RUN PROCESSED 3 JULY 72.

GHTSQN

- 1 1.67396180 00
- 2 1.66804550 00
- 3 1.66473070 00
- 4 1.66136900 00

Figure XVI. C. Part one of the 85 keV lifetime fit for the
line at 3124.9 Å

Figure XVI. D. Part two of the 85 keV lifetime fit for the
line at 3124.9 A

3125.0A RUN #6 PART 2 CH1 45
 MASS= 52. ENERGY= 72. KEV
 THERE WERE 0.59000 NOISE CHANNELS
 START FIT IN CHANNEL NO 7

88KV
 DARK COUNTS = 6.

TIME ON POLE 0 TO 4 MIN
 1/0= 1.0410 MSEC/CH

DISTANCE DATA FIT

0.50	45.000	1.0000
1.00	215.000	1.0000
1.50	547.000	1.0000
2.00	847.000	1.0000
2.50	1100.000	1.0000
3.00	1225.000	1.0000
3.50	1248.000	1.0000
4.00	1238.000	1.0000
4.50	1176.000	1.0000
5.00	1028.000	1.0000
5.50	898.000	1.0000
6.00	798.000	1.0000
6.50	715.000	1.0000
7.00	624.000	1.0000
7.50	578.000	1.0000
8.00	542.000	1.0000
8.50	485.000	1.0000
9.00	451.000	1.0000
9.50	417.000	1.0000
10.00	377.000	1.0000
10.50	336.000	1.0000
11.00	296.000	1.0000
11.50	264.000	1.0000
12.00	258.000	1.0000
12.50	218.000	1.0000
13.00	228.000	1.0000
13.50	178.000	1.0000
14.00	174.000	1.0000
14.50	171.000	1.0000
15.00	132.000	1.0000
15.50	128.000	1.0000
16.00	101.000	1.0000
16.50	96.000	1.0000
17.00	75.000	1.0000
17.50	71.000	1.0000
18.00	65.000	1.0000
18.50	63.000	1.0000
19.00	51.000	1.0000
19.50	48.000	1.0000
20.00	48.000	1.0000
20.50	32.000	1.0000
21.00	31.000	1.0000
21.50	26.000	1.0000
22.00	24.000	1.0000
22.50	21.000	1.0000
23.00	21.000	1.0000
23.50	15.000	1.0000
24.00	15.000	1.0000
24.50	15.000	1.0000
25.00	15.000	1.0000
25.50	15.000	1.0000
26.00	15.000	1.0000
26.50	15.000	1.0000
27.00	15.000	1.0000
27.50	15.000	1.0000
28.00	15.000	1.0000
28.50	15.000	1.0000
29.00	15.000	1.0000
29.50	15.000	1.0000
30.00	15.000	1.0000
30.50	15.000	1.0000
31.00	15.000	1.0000
31.50	15.000	1.0000
32.00	15.000	1.0000
32.50	15.000	1.0000
33.00	15.000	1.0000
33.50	15.000	1.0000
34.00	15.000	1.0000
34.50	15.000	1.0000
35.00	15.000	1.0000
35.50	15.000	1.0000
36.00	15.000	1.0000
36.50	15.000	1.0000
37.00	15.000	1.0000
37.50	15.000	1.0000
38.00	15.000	1.0000
38.50	15.000	1.0000
39.00	15.000	1.0000
39.50	15.000	1.0000
40.00	15.000	1.0000
40.50	15.000	1.0000
41.00	15.000	1.0000
41.50	15.000	1.0000
42.00	15.000	1.0000
42.50	15.000	1.0000
43.00	15.000	1.0000
43.50	15.000	1.0000
44.00	15.000	1.0000
44.50	15.000	1.0000
45.00	15.000	1.0000
45.50	15.000	1.0000
46.00	15.000	1.0000
46.50	15.000	1.0000
47.00	15.000	1.0000
47.50	15.000	1.0000
48.00	15.000	1.0000
48.50	15.000	1.0000
49.00	15.000	1.0000
49.50	15.000	1.0000
50.00	15.000	1.0000
50.50	15.000	1.0000
51.00	15.000	1.0000
51.50	15.000	1.0000
52.00	15.000	1.0000
52.50	15.000	1.0000
53.00	15.000	1.0000
53.50	15.000	1.0000
54.00	15.000	1.0000
54.50	15.000	1.0000
55.00	15.000	1.0000
55.50	15.000	1.0000
56.00	15.000	1.0000
56.50	15.000	1.0000
57.00	15.000	1.0000
57.50	15.000	1.0000
58.00	15.000	1.0000
58.50	15.000	1.0000
59.00	15.000	1.0000
59.50	15.000	1.0000
60.00	15.000	1.0000
60.50	15.000	1.0000
61.00	15.000	1.0000
61.50	15.000	1.0000
62.00	15.000	1.0000
62.50	15.000	1.0000
63.00	15.000	1.0000
63.50	15.000	1.0000
64.00	15.000	1.0000
64.50	15.000	1.0000

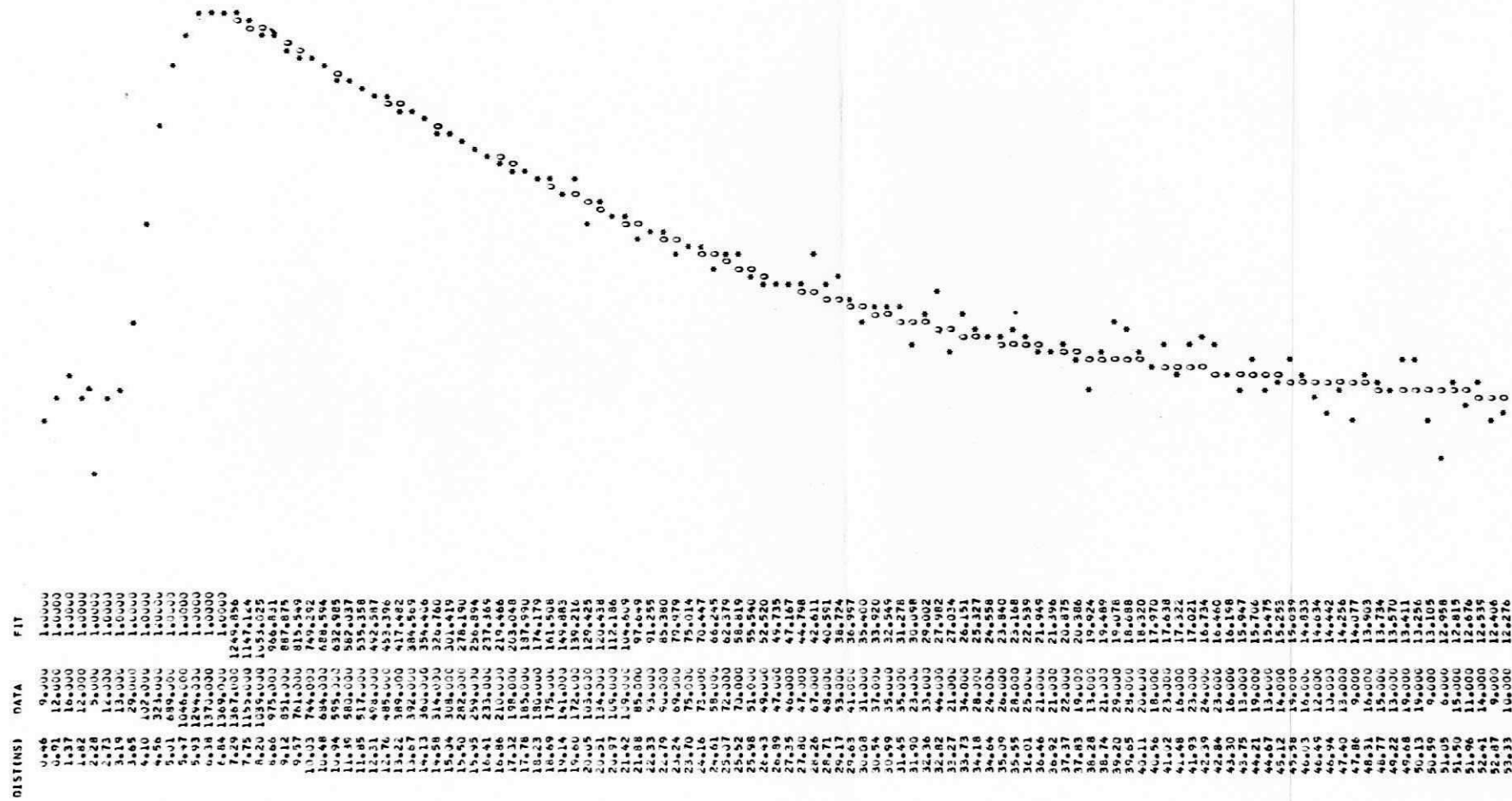
INTERCEPT-1 IS 2415.858
 INTERCEPT-2 IS 26.4263
 THE SLOPE IS 0.000000
 THE RATIO OF COEFFICIENTS AT THE POLE IS 58.80378
 RUN DEPRESSOR 3.000000

CH100

1 1.0013700 00
 2 1.0013700 00
 3 1.0013700 00
 4 1.0013700 00

Figure XVII. A. The 100 keV lifetime fit for the
line at 3132.1 Å

3132.14 RUN #5 ENERGY 17% TIME ON PULS 0 TO 4 MIN 100 KV
 MASS #82 0-25800 NM/CHANNEL DARK COUNTS= 4.
 THERE WERE 17% 1.7666 NSEC/MIN
 START FIT IN CHANNEL NO 10



INTERCEPT-1 IS 4965.0476 SLOPE-1 IS -0.1634 TAU-1 IS 5.1719 NSEC
 INTERCEPT-2 IS 41.9883 SLOPE-2 IS -0.0308 TAU-2 IS 32.4330 NSEC
 THE PEAK IS AT 41.9883 NM/CHANNEL
 THE RATIO OF COEFFICIENTS AT THE FCIL IS 41.92534
 RUN PROCESSED 3 JULY 72.

CHISON

1 1.22679380 00
 2 1.22245670 00
 3 1.22819700 00
 4 1.22441120 00

Figure XVII. B. Part one of the 85 keV lifetime fit for the
line at 3132.1 A

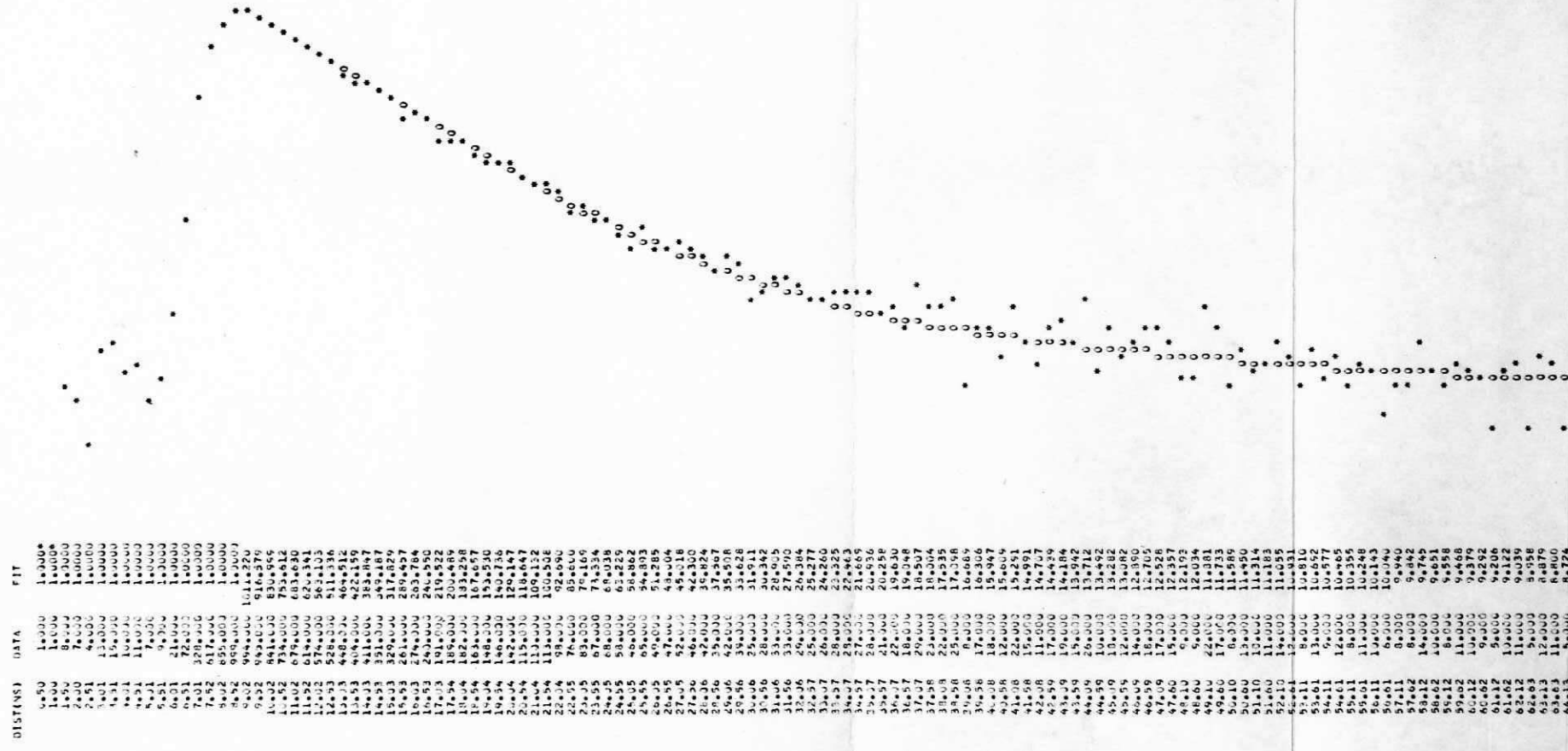
3132A1A RUN 48 PART 1 FOIL #4 TIME ON FOIL 0 TO 12MIN 65 KV 4.
 MASS 42. 0.24000 NSEC/CM NEW COUNTS
 THERE WERE 72. KEV 1/CM 1.0619 NSEC/CM
 START FIT IN CHANNEL NO 14

DISTANCE DATA FIT

0.50	6.000	1.0000
1.00	9.000	1.0000
1.50	12.000	1.0000
2.00	15.000	1.0000
2.50	18.000	1.0000
3.00	21.000	1.0000
3.50	24.000	1.0000
4.00	27.000	1.0000
4.50	30.000	1.0000
5.00	33.000	1.0000
5.50	36.000	1.0000
6.00	39.000	1.0000
6.50	42.000	1.0000
7.00	45.000	1.0000
7.50	48.000	1.0000
8.00	51.000	1.0000
8.50	54.000	1.0000
9.00	57.000	1.0000
9.50	60.000	1.0000
10.00	63.000	1.0000
10.50	66.000	1.0000
11.00	69.000	1.0000
11.50	72.000	1.0000
12.00	75.000	1.0000
12.50	78.000	1.0000
13.00	81.000	1.0000
13.50	84.000	1.0000
14.00	87.000	1.0000
14.50	90.000	1.0000
15.00	93.000	1.0000
15.50	96.000	1.0000
16.00	99.000	1.0000
16.50	102.000	1.0000
17.00	105.000	1.0000
17.50	108.000	1.0000
18.00	111.000	1.0000
18.50	114.000	1.0000
19.00	117.000	1.0000
19.50	120.000	1.0000
20.00	123.000	1.0000
20.50	126.000	1.0000
21.00	129.000	1.0000
21.50	132.000	1.0000
22.00	135.000	1.0000
22.50	138.000	1.0000
23.00	141.000	1.0000
23.50	144.000	1.0000
24.00	147.000	1.0000
24.50	150.000	1.0000
25.00	153.000	1.0000
25.50	156.000	1.0000
26.00	159.000	1.0000
26.50	162.000	1.0000
27.00	165.000	1.0000
27.50	168.000	1.0000
28.00	171.000	1.0000
28.50	174.000	1.0000
29.00	177.000	1.0000
29.50	180.000	1.0000
30.00	183.000	1.0000
30.50	186.000	1.0000
31.00	189.000	1.0000
31.50	192.000	1.0000
32.00	195.000	1.0000
32.50	198.000	1.0000
33.00	201.000	1.0000
33.50	204.000	1.0000
34.00	207.000	1.0000
34.50	210.000	1.0000
35.00	213.000	1.0000
35.50	216.000	1.0000
36.00	219.000	1.0000
36.50	222.000	1.0000
37.00	225.000	1.0000
37.50	228.000	1.0000
38.00	231.000	1.0000
38.50	234.000	1.0000
39.00	237.000	1.0000
39.50	240.000	1.0000
40.00	243.000	1.0000
40.50	246.000	1.0000
41.00	249.000	1.0000
41.50	252.000	1.0000
42.00	255.000	1.0000
42.50	258.000	1.0000
43.00	261.000	1.0000
43.50	264.000	1.0000
44.00	267.000	1.0000
44.50	270.000	1.0000
45.00	273.000	1.0000
45.50	276.000	1.0000
46.00	279.000	1.0000
46.50	282.000	1.0000
47.00	285.000	1.0000
47.50	288.000	1.0000
48.00	291.000	1.0000
48.50	294.000	1.0000
49.00	297.000	1.0000
49.50	300.000	1.0000
50.00	303.000	1.0000
50.50	306.000	1.0000
51.00	309.000	1.0000
51.50	312.000	1.0000
52.00	315.000	1.0000
52.50	318.000	1.0000
53.00	321.000	1.0000
53.50	324.000	1.0000
54.00	327.000	1.0000
54.50	330.000	1.0000
55.00	333.000	1.0000
55.50	336.000	1.0000
56.00	339.000	1.0000
56.50	342.000	1.0000
57.00	345.000	1.0000
57.50	348.000	1.0000
58.00	351.000	1.0000
58.50	354.000	1.0000
59.00	357.000	1.0000
59.50	360.000	1.0000
60.00	363.000	1.0000
60.50	366.000	1.0000
61.00	369.000	1.0000
61.50	372.000	1.0000
62.00	375.000	1.0000
62.50	378.000	1.0000
63.00	381.000	1.0000
63.50	384.000	1.0000
64.00	387.000	1.0000
64.50	390.000	1.0000
65.00	393.000	1.0000
65.50	396.000	1.0000
66.00	399.000	1.0000
66.50	402.000	1.0000
67.00	405.000	1.0000
67.50	408.000	1.0000
68.00	411.000	1.0000
68.50	414.000	1.0000
69.00	417.000	1.0000
69.50	420.000	1.0000
70.00	423.000	1.0000
70.50	426.000	1.0000
71.00	429.000	1.0000
71.50	432.000	1.0000
72.00	435.000	1.0000
72.50	438.000	1.0000
73.00	441.000	1.0000
73.50	444.000	1.0000
74.00	447.000	1.0000
74.50	450.000	1.0000
75.00	453.000	1.0000
75.50	456.000	1.0000
76.00	459.000	1.0000
76.50	462.000	1.0000
77.00	465.000	1.0000
77.50	468.000	1.0000
78.00	471.000	1.0000
78.50	474.000	1.0000
79.00	477.000	1.0000
79.50	480.000	1.0000
80.00	483.000	1.0000
80.50	486.000	1.0000
81.00	489.000	1.0000
81.50	492.000	1.0000
82.00	495.000	1.0000
82.50	498.000	1.0000
83.00	501.000	1.0000
83.50	504.000	1.0000
84.00	507.000	1.0000
84.50	510.000	1.0000
85.00	513.000	1.0000
85.50	516.000	1.0000
86.00	519.000	1.0000
86.50	522.000	1.0000
87.00	525.000	1.0000
87.50	528.000	1.0000
88.00	531.000	1.0000
88.50	534.000	1.0000
89.00	537.000	1.0000
89.50	540.000	1.0000
90.00	543.000	1.0000
90.50	546.000	1.0000
91.00	549.000	1.0000
91.50	552.000	1.0000
92.00	555.000	1.0000
92.50	558.000	1.0000
93.00	561.000	1.0000
93.50	564.000	1.0000
94.00	567.000	1.0000
94.50	570.000	1.0000
95.00	573.000	1.0000
95.50	576.000	1.0000
96.00	579.000	1.0000
96.50	582.000	1.0000
97.00	585.000	1.0000
97.50	588.000	1.0000
98.00	591.000	1.0000
98.50	594.000	1.0000
99.00	597.000	1.0000
99.50	600.000	1.0000
100.00	603.000	1.0000
100.50	606.000	1.0000
101.00	609.000	1.0000
101.50	612.000	1.0000
102.00	615.000	1.0000
102.50	618.000	1.0000
103.00	621.000	1.0000
103.50	624.000	1.0000
104.00	627.000	1.0000
104.50	630.000	1.0000
105.00	633.000	1.0000
105.50	636.000	1.0000
106.00	639.000	1.0000
106.50	642.000	1.0000
107.00	645.000	1.0000
107.50	648.000	1.0000
108.00	651.000	1.0000
108.50	654.000	1.0000
109.00	657.000	1.0000
109.50	660.000	1.0000
110.00	663.000	1.0000
110.50	666.000	1.0000
111.00	669.000	1.0000
111.50	672.000	1.0000
112.00	675.000	1.0000
112.50	678.000	1.0000
113.00	681.000	1.0000
113.50	684.000	1.0000
114.00	687.000	1.0000
114.50	690.000	1.0000
115.00	693.000	1.0000
115.50	696.000	1.0000
116.00	699.000	1.0000
116.50	702.000	1.0000
117.00	705.000	1.0000
117.50	708.000	1.0000
118.00	711.000	1.0000
118.50	714.000	1.0000
119.00	717.000	1.0000
119.50	720.000	1.0000
120.00	723.000	1.0000
120.50	726.000	1.0000
121.00	729.000	1.0000
121.50	732.000	1.0000
122.00	735.000	1.0000
122.50	738.000	1.0000
123.00	741.000	1.0000
123.50	744.000	1.0000
124.00	747.000	1.0000
124.50	750.000	1.0000
125.00	753.000	1.0000
125.50	756.000	1.0000
126.00	759.000	1.0000
126.50	762.000	1.0000
127.00	765.000	1.0000
127.50	768.000	1.0000
128.00	771.000	1.0000
128.50	774.000	1.0000
129.00	777.000	1.0000
129.50	780.000	1.0000
130.00	783.000	1.0000
130.50	786.000	1.0000
131.00	789.000	1.0000
131.50	792.000	1.0000
132.00	795.000	1.0000
132.50	798.000	1.0000
133.00	801.000	1.0000
133.50	804.000	1.0000
134.00	807.000	1.0000
134.50	810.000	1.0000
135.00	813.000	1.0000
135.50	816.000	1.0000
136.00	819.000	1.0000
136.50	822.000	1.0000
137.00	825.000	1.0000
137.50	828.000	1.0000
138.00	831.000	1.0000
138.50	834.000	1.0000
139.00	837.000	1.0000
139.50	840.000	1.0000
140.00	843.000	1.0000
140.50	846.000	1.0000
141.00	849.000	1.0000
141.50	852.000	1.0000
142.00	855.000	1.0000
142.50	858.000	1.0000
143.00	861.000	1.0000
143.50	864.000	1.0000
144.00	867.000	1.0000
144.50	870.000	1.0000
145.00	873.000	1.0000
145.50	876.000	1.0000
146.00	879.000	1.0000
146.50	882.000	1.0000
147.00	885.000	1.0000
147.50	888.000	1.0000
148.00	891.000	1.0000
148.50	894.000	1.0000
149.00	897.000	1.0000
149.50	900.000	1.0000
150.00	903.000	1.0000
150.50	906.000	1.0000
151.00	909.000	1.0000
151.50	912.000	1.0000
152.00	915.000	1.0000
152.50	918.000	1.0000
153.00	921.000	1.0000
153.50	924.000	1.0000
154.00	927.000	1.0000
154.50	930.000	1.0000
155.00	933.000	1.0000
155.50	936.000	1.0000
156.00	939.000	1.0000
156.50	942.000	1.0000
157.00	945.000	1.0000
157.50	948.000	1.0000
158.00	951.000	1.0000
158.50	954.000	1.0000
159.00	957.000	1.

Figure XVII. C. Part two of the 85 keV lifetime fit for the line at 3132.1 Å

3132-1A RUN IN PART 2 FOIL #4 TIME ON FOIL 9 TO 12 MIN 85 KV DARK COUNTS= 4.
 MASS= 52. ENERGY= 72. REV L/V= 1.9419 NSEC/MM
 THERE WERE 0.25900 MM/CHANNEL
 START FIT IN CHANNEL NO 13



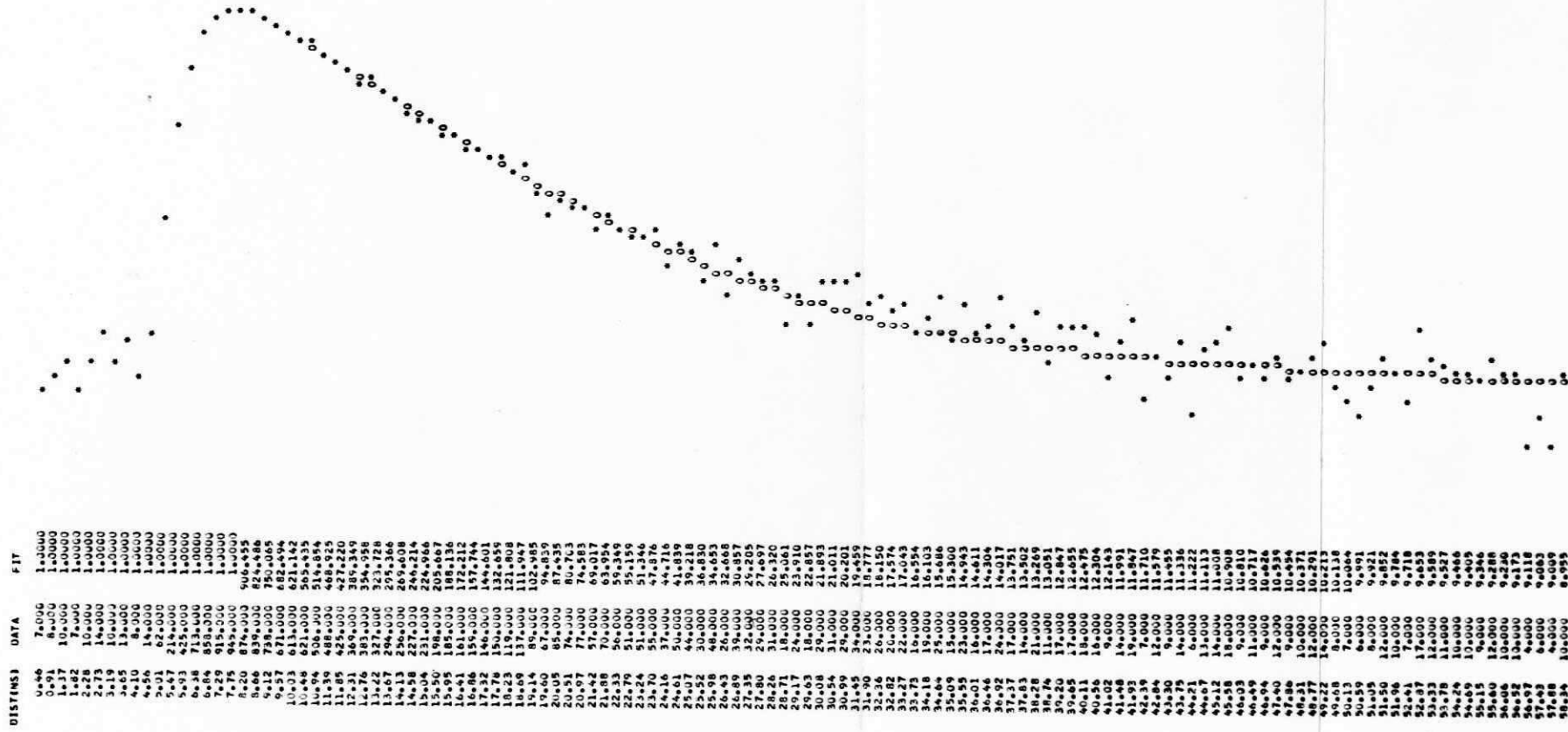
INTERCEPT-1 IS 60.43-6626 SLOPE-1 IS -0.2017 TAU-1 IS 4.0579 NSEC
 INTERCEPT-2 IS 35.8665 SLOPE-2 IS -0.0317 TAU-2 IS 31.9850 NSEC
 THE PEAK IS AT ON NEAR CHANNEL 17
 THE RATIO OF COEFFICIENTS AT THE FOIL IS: 39.95943
 RUN PROCESSED 3 JULY 72.

CHISON

- 1 9.4752240-01
- 2 9.21764540-01
- 3 9.77519690-01
- 4 9.7347740-01

Figure XVIII. A. The 100 keV lifetime fit for the
line at 3180.8 Å

3180-RA RUN #6 ENERGY= 87.1 KEV TIME ON FOIL 4 TO 8 MIN 100 MV
 MASS= 52 THERE WERE 0.25800 MW/CHANNEL 17% 1.7666 NS/C/AM DARK COUNTS= 4.



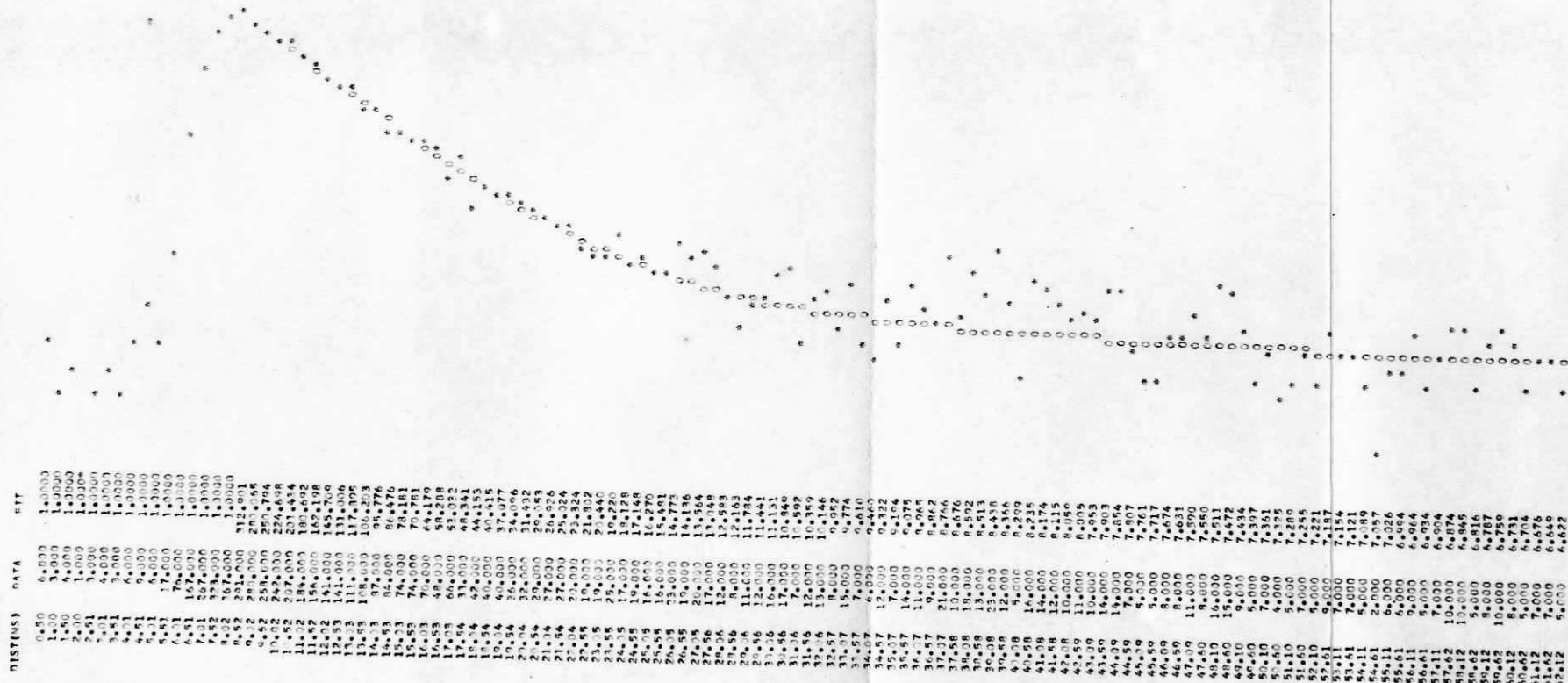
INTERCEPT-1 IS 9.0422964 SLOPE-1 IS -0.2123 TAU-1 IS 4.7104 NS/C
 INTERCEPT-2 IS 18.5500 SLOPE-2 IS -0.0227 TAU-2 IS 44.1066 NS/C
 THE PEAK IS AT 0.46 KEV
 THE RATIO OF COEFFICIENTS AT THE FOIL IS 62.80762
 RUN PROCESSED 3 JULY 72.

CWISON

1 1.3447332D 00
 2 1.3233049D 00
 3 1.3244224D 00
 4 1.3152803D 00

Figure XVIII. B. The 85 keV lifetime fit for the
line at 3180.8 Å

3180-RA RUN 47 PART 1 FOIL #5 TIME ON FOIL 8 TO 104TH AS MV
 MASS= 52 ENERGY= 72.4KV 1/00 1.9410 NECC/AM DARK COUNTS= 4.
 THERE WERE 0.25000 W/CH/AN/CL
 START FIT IN CHANNEL NO 17



INTERCEPT-1 IS 2132.7667 SLOPE-1 IS -0.2298 TAN-1 IS 4.3511 NSFC
 INTERCEPT-2 IS 9.1500 SLOPE-2 IS -0.0201 TAN-2 IS 49.4406 NSFC
 BEAK IS AT 9.1500 NEAR CHANNEL 16
 THE COEFFICIENTS AT THE FOIL IS 43.35765
 RUN PROCESSED 3 JULY 72.

CMISOR

- 1 2.54104000 00
- 2 1.23333440 00
- 3 1.22206480 00
- 4 1.21607420 00

BEAM-FOIL LIFETIMES IN CrII

by

JAMES RICHARD METTLING

B.A., Southwestern College, 1971

AN ABSTRACT OF A MASTER'S THESIS

submitted in partial fulfillment of the

requirements for the degree

MASTER OF SCIENCE

Department of Physics

KANSAS STATE UNIVERSITY
Manhattan, Kansas

1974

The lifetimes of four atomic levels of CrII have been measured at 85 and 100 keV using beam-foil methods. In preparation for the lifetime experiment the beam-foil spectrum of 85 keV chromium from 1900 to 4100 Å was taken. The energy loss incurred by 89 keV Ar⁺ passing through a 2.4 μgm/cm² carbon foil has been measured and found to be in adequate agreement with stopping power calculations. From our lifetime measurements a preliminary solar abundance for chromium of $\log N_{\text{Cr}} = 5.54$ (where $\log N_{\text{H}} = 12.0$) has been calculated.