Boiling Heat Transfer Phenomena During Rapid Decompression

by

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NOMENCLATURE

A _L	Area of liquid column
$A_{\overline{\mathbf{T}}}$	Total area of heat transfer surface
$\mathbf{A_{v}}$	Area of vapor column
c	Wave velocity
CHF	Critical heat flux
C _p	Heat capacity at constant pressure
D jet	Diameter of vapor jet flow
f	Frequency of bubble release from heater surface
g	Gravitational acceleration
g _o	Universal gravitation constant
ħ,	Convection heat transfer coefficient
h _L	Enthalpy per unit mass of saturated liquid
$^{ m h}{}_{ m extbf{V}}$	Enthalpy per unit mass of saturated vapor
h _{Lv}	$h_{\mathbf{v}}^{-h}$ or the latent heat of liquid
I	Electric current through heater element
k _u	Kutateladze empirical constant
k	Thermal conductivity
М	Molecular weight
N	Number of active nucleation sites on bubbling surface
$^{ ext{P}}_{ ext{L}}$	Liquid pressure
P _r	Prandtl number
ģ	Heat generation rate per unit volume
(Q/A) _b	Heat flux from boiling
(Q/A) _{CHF}	Critical heat flux
(Q/A) _{CHF,k}	Critical heat flux after Kutateladze
(Q/A) _{cond}	Heat flux from conduction

(Q/A)_{nc} Heat flux from natural convection (Q/A)_{op} Operating heat flux (Q/A)_{Total} Total heat flux Universal gas constant R Bubble departure radius R Electrical resistance of heating wire at 0°C R Resistance of precisive shunt resistor R sys Resistance of heater element and connections Theoretical resistance of the test wire Rth Resistance of the test wire Radius of active site on heating surface r Bubble radius r The radius for which N would be one per unit area of heating surface Liquid temperature T, Saturation temperature Tg Vapor temperature Heater element temperature V_{T.} Liquid Velocity Average volume of microlayer evaporated V_{ME} Voltage of the battery V_P Vapor velocity Voltage across heater element and connections Transient voltage signal from the pressure transducer V pressure Thermal diffusivity Œ Three phase contact angle in Table 1. β δ Bubble thermal boundary layer thickness

	ε	Accommodation coefficient, dimensionless
•	η	Wave amplitude
	λ	Wave length
	ν	Viscosity
	$^{ m ho}_{ m L}$	Liquid density
	$^{p}\mathbf{v}$	Vapor density
	σ	Surface tension between liquid and its vapor

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INTRODUCTION

New developments in engineering applications often require that heat be transferred under not only steady-state, but also transient conditions. For example, during the operation of light water nuclear reactors, heat generated in the fuel rods is increased exponentially as the reactor is started up or is required to change power levels within a short time. If the period is significantly shorter than that at which the coolant can transfer all the heat generated in the fuel rods, the fuel rod cladding may be exposed, for short periods, to nucleate or to film boiling and rapid temperature fluctuations. Rapid decompressions are potentially more serious. If a reactor primary coolant pipe ruptures, the high pressure inside the reactor vessel will be released to the atmosphere in seconds. The coolant becomes superheated rapidly which in turn induces flashing throughout the reactor core and alters the local heat transfer mechanisms. An accurate knowledge of the boiling behavior during pressure transients of varying periods is, therefore, necessary for the designer to estimate the worst possible heat transfer during normal operation or abnormal occurrences.

Boiling heat transfer has been studied comprehensively for several decades. The majority of previous analyses and experiments have been concerned with steady-state, pool boiling. Most conventional applications of boiling processes involve steady-state phenomena, in which the liquids are at a certain pressure and only slow, incremental changes in heat flux are encountered. As a result of this extensive research, there are numerous equations, either experimentally or theoretically derived, available for the prediction of steady-state nucleate boiling heat fluxes and critical heat fluxes under various thermodynamic conditions and with different solid/liquid combinations.

Boiling heat transfer during both power and pressure transients has been studied more recently. Most effort has been devoted to power transients because of their more frequent occurrence in light water reactors. Empirical correlations have been developed which predict the transition from nucleate to film boiling as a function of transient power period, initial heat flux, and thermodynamic conditions. Pressure transients, on the other hand, have been studied only recently because of the increasing concern about the Loss of Coolant Accident (LOCA). However, no systematic study has been performed.

Power transient studies to date have reported that if the transient time constant is longer than about 100 milliseconds, the peak heat flux during the transient can be predicted adequately by steady-state conditions. If, however, the time constant is shorter than 20 milliseconds, the heat flux during the transient may reach a higher value than that predicted by steady-state correlations before the boiling transits from nucleate to film boiling. It has been observed that even under a step increase in power, nucleate boiling always precedes the transition to film boiling. Pressure transients have been much less thoroughly studied. Investigations in rod bundles have been performed more to provide design information rather than fundamental knowledge. In view of the few studies of relatively slow pressure transients, it has been suggested that the critical heat flux can be predicted with no significant errors by steady-state correlations. But the results of recent and more basic research into the heat flux and transient bubble growth indicate that the heat transfer mechanisms may be quite different. Because of the paucity of data on this topic, no solid conclusions have been made about the actual effects of pressure transients on boiling heat transfer.

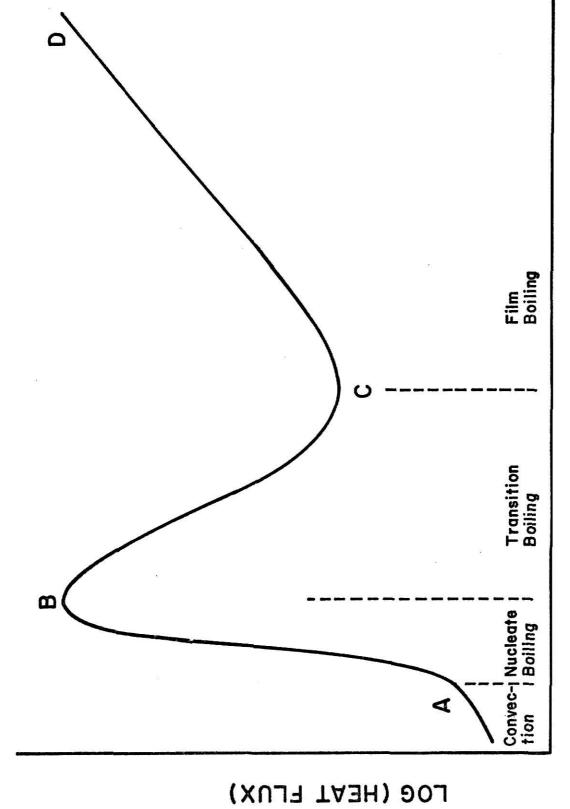
Accordingly, as a first step in a comprehensive research program the effect of very rapid pressure transients on boiling heat transfer has been studied. The effects of two parameters, temperature of bulk liquid and initial heat flux level, on the temporal variation of heat flux and wire temperature are investigated during the depressurization of water from 760 to 420 torr in approximately 10 milliseconds. The review of pertinent literature, design of experiments, and experimental results are discussed in turn in the following sections.

II. LITERATURE REVIEW

In this section, steady-state boiling will be discussed briefly and applicable Critical Heat Flux correlations will be reviewed. Subsequently, critical heat flux during power transients and decompressions will be covered.

A. The Mechanism and Current Theories of Steady-State Pool Boiling

Steady-state boiling heat transfer studies have shown conclusively that different levels of wall superheat, the wall temperature excess above the saturation temperature, result in very different heat transfer phenomena, which are fundamental to the understanding of the boiling mechanism. The existence of several regimes of boiling was first clearly discussed by Nukiyama (41) in 1934 and followed by many papers. A typical steady-state boiling curve is shown in Fig. 1. The entire curve is generally divided into four regions: free convection, nucleate boiling, transition, and stable film boiling regions followed by film boiling with radiation augmentation. Nucleate boiling is characterized by the formation, growth, and detachment of vapor bubbles in the liquid near the heater surface. It results in a sharp increase in heat flux, but only a moderate increase in surface temperature. At low values of wall superheat, bubbles are far apart and are nearly independent of one another. As the superheat increases, more nucleation sites appear and bubbles form more rapidly. At high superheats, no discrete bubbles can be seen as they interfere with one another and unite to form bigger vapor patches. As point B is approached, a further increase in wall superheat will cause the heat flux to decrease since all heat must be transferred through a vapor layer. This is called transition boiling.



LOG (TEMPERATURE)

Fig. 1. A Typical Pool Boiling Curve.

Heat transfer during this region is unstable. Part of the surface is covered by vapor film, the rest is still undergoing stable nucleate boiling until the superheat is increased high enough to cause the coalesced bubbles to form a layer of vapor film covering the entire surface. This is called film boiling. The critical point B is known widely as DNB - depature from nucleate boiling, or the burnout point. Hereafter, this point will be referred to as the boiling transition and the heat flux at this point as the CHF - critical heat flux - as used most commonly in recent literature.

There are several plausible physical explanations for the high rate of heat transfer in nucleate boiling. Earlier researchers tended to focus on the agitation induced by the bubble formations on the heater surface. For example, Rohsenow and Clark (45) proposed that the bubbles behave as agitation agents which increase the heat transfer. This proposal has been studied by Robinson and Katz (43), by Grose et al. (19), and by Mixon et al. (38). These experiments showed that the bubble agitation does increase the heat transfer, but the heat fluxes in these studies were in the lower range of nucleate boiling. Other mechanisms which are not substantially different from the above include that of the "bubble action," whereby the bubbles are assumed to push the superheated liquid layer away from the heater surface (17), and the mechanisms which explain the high heat flux by the increased turbulence imparted to the boundary layer by the growing bubbles (21, 25).

More recently, latent heat transport mechanisms have been put forth by several investigators. Moore and Mesler (39), in a pioneering experiment, provided new data explaining the mechanism by which latent heat is transported away from the heated surface. A fast response thermocouple was used during the nucleate boiling of water. Surface temperature drops of 11 °C to 16.7 °C within 2 milliseconds were observed indicating a rapid extraction of heat during a very short time. The authors suggested that there is a thin liquid layer underneath the growing bubbles and that the evaporation of this layer into the bubbles brings the surface temperature down during bubble growth. Further studies of this model have shown that it is insufficient to explain the very high heat flux (54). Most recently, Kirby and Westwater (32) and Katto and Yokoya (31) experimentally found that a thin liquid film exists beneath irregular vapor masses with the principal heat transfer mechanism at high heat flux being the evaporation from this liquid film. It appears that these different conceptual models reflect different stages of the boiling mechanism; consequently, no overall analytical solution or empirical correlation is available to accurately predict the CHF.

In a recent review of current nucleate boiling theory, Labuntsov (34) pointed out that future analyses should concentrate on the features and structure of a very thin, liquid-rich film directly on the heater surface which sustains the principal temperature drop. The existence of such a film over the entire range of nucleate boiling has been confirmed by probing local voids with micro thermocouples (7,26). Therefore, the bubble agitation model appears to be more satisfactory for lower heat fluxes in nucleate boiling, while the latent heat transfer model is more applicable at high heat fluxes. None of the models previously mentioned explains the results of the latter experiments. Hence present nucleate boiling models will have to be modified or a new model developed in order to describe the entire nucleate boiling regime adequately.

B. Steady-State Pool Boiling Critical Heat Flux (CHF)

Tong (52) and Balzhiser (4) have given comprehensive reviews of pool boiling heat transfer. They concluded that although the problem has been approached from many different points of view, no overall analytical solution has been obtained, and the design equations for CHF predictions are strictly empirical. For instance, Gambill (18) has noted that approximately fifty equations have been developed for predicting CHF in pool and flow boiling situations. Only those generalized CHF correlations that apply to saturated pool boiling will be discussed here.

In boiling, bubbles appear to form repeatedly at nucleation sites (cavities) on the heating surface, thus forming a "column" of bubbles.

As the surface heat flux increases, visual observations indicate that the number of active nucleation sites per unit area increases (44). In addition, observations suggest that, as the heat flux increases to just below the critical value, the successive bubbles coming from a nucleation site converge into each other to form an undulating column of vapor. Rohsenow and Griffith (46) formulated a semi-empirical relation for the critical heat flux by combining observations from experiments in organic liquid and water with the conceptual model noted above. Their relation,

$$(Q/A)_{CHF} = 143 \rho_v h_{Lv} g^{1/4} \left(\frac{\rho_L - \rho_v}{\rho_v}\right)^{0.6}$$
 (2.1)

was derived by tacitly assuming that the heater surface is covered by vapor bubbles at CHF.

Addoms (2) proposed a dimensionless correlation:

$$(Q/A)_{CHF} = 2.5 \lambda \rho_{v} (g\alpha)^{1/3} \left(\frac{\rho_{L} - \rho_{v}}{\rho_{v}}\right)^{0.5}$$
 (2.2)

which also agrees well with data from experiments in water and organic fluids.

Other researchers have developed semi-empirical correlations using a more universal approach. For instance, Zuber (55,56), basing his arguments on wave motion, postulated that the boiling crisis is due to a combination of Taylor's and Helmholtz's instabilities which deal with the instability of a plane interface and the relative velocity of the liquid and its vapor respectively. The resulting equation is:

$$(Q/A)_{CHF} = 0.18 \rho_{v}^{h}_{Lv} \left[\frac{gg_{o}(\rho_{L} - \rho_{v})}{\rho_{v}^{2}}\right]^{1/4} \left[\frac{\rho_{L}}{\rho_{L} + \rho_{v}}\right]^{1/2}$$
 (2.3)

Here the constant, 0.18, was determined empirically from the data shown in Fig. 2.

Kutateladze (33), Chang and Snyder (12), and Cobb and Park (14), have also developed equations which are nearly the same as those of Zuber and Addoms. Their equations, Berenson's modification of Zuber's analysis, in addition to some other pertinent correlations are listed in Table 1.

The examination of the above correlations does not make the effects of different parameters on the critical heat flux entirely obvious. However, from a thermodynamic point of view, it is apparent that the system pressure will be a dominant factor since it determines the latent heat of vaporization, the saturated vapor density and affects the surface tension. In fact, its effect has been well documented by many authors (13,15,18,24). Yet the CHF is far more complicated than this.

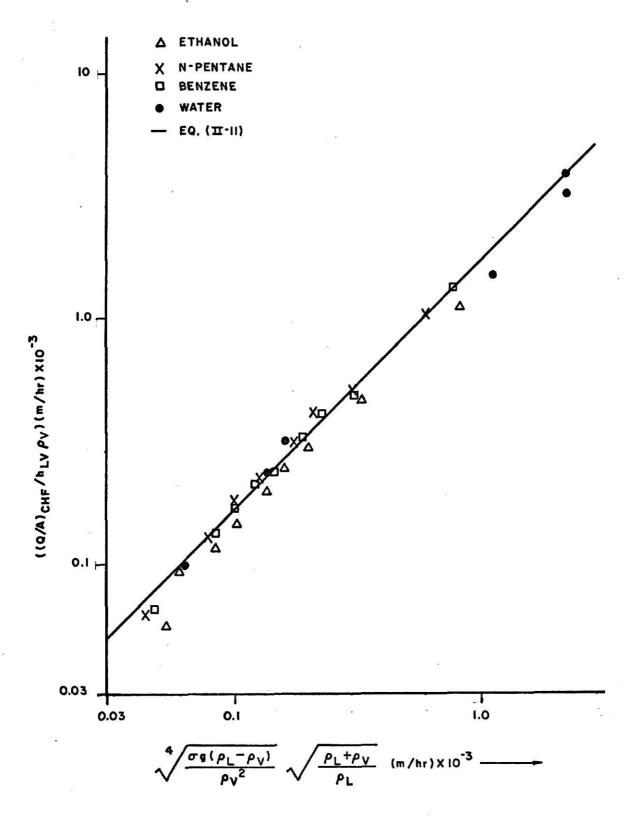


Fig. 2. Correlation of Critical Heat Flux in Pool Boiling. [After Zuber (56)].

TABLE 1-a. Equations for the Prediction of Nucleate Boiling Heat Flux

Equation	Theoretical Background	Validity	Author	Year
$(Q/A)_{Total} = \frac{\binom{C_{pL}(T_{w}-T)^{3}}{k_{Lv}^{1.7}} \binom{L_{h}L_{v}}{\binom{C_{3}}{c_{Sg}}} [\frac{g_{o}^{\gamma}}{g(\rho_{L}-\rho_{g})}]$]-1/2 Dimensional Analysis	It predicts the shape of the nucleate boiling regime with good accuracy.	W.M. Rohsenow (44)	1952
$ (Q/A)_{Total} = \frac{A_{T}}{A_{T}} (Q/A)_{nc} + (Q/A)_{b} $ where $(Q/A)_{b} = C_{1} \frac{C_{2}C_{3}^{1/2}r_{m}^{m}}{\sqrt{\pi} 2^{m-1}} (kp_{C}_{p})_{L}^{1/2} \cdot \frac{h_{Lv}^{p}v_{v}^{m}}{T_{S}^{\sigma}} \cdot [\frac{Q_{S}C_{L}^{-p}v_{v}}{\rho_{L}^{2}}]^{3/4} (Ja^{*})^{1/8} $ $ \cdot [\frac{Q_{S}C_{L}^{p}v_{v}}{T_{S}^{\sigma}}]^{3/4} (Ja^{*})^{15/8} $ $ \cdot (T_{w}^{-}T_{S})^{m+1} $	Assumes the main mechanism of heat transfer in nucleate boiling is transient heat conduction to, and subsequent replacement of, the superheated layer around boiling sites associated with bubble departure and considers effects of heat transfer surface characteristics.	For checking the equation, the required information about the boiling surface was found from q/A versus (T -T) data at one pressure and applied to all. The results are satisfactory.	Mikic and Rohsenow (37)	1969
	=	æ		

TABLE 1-a. (Continued)

Year	*	1976
Author		R. L. Judd (29)
Validity		Microlayer evaporation heat transfer is the prominent feature of this model. More data needed to support this equation.
Theoretical Background	onstants. $\frac{rg(T_w-T_g)^5\alpha^3}{\sqrt{A_T}\alpha}^{1/4}$ 0^7 $\frac{rg(T_w-T_g)^4\alpha^2}{[w^T_g)^4\alpha}^{1/2}$ x 10 ¹⁰	Incorporation of microlayer evapo- ration, natural convection, and nucleate boiling mechanisms
Equation	$c_{1},c_{2},c_{3}, \text{ and m are constants.}$ $J_{a*} = \frac{\rho_{L} p_{L} T_{S}}{\rho_{v} h_{Lv}}$ $(Q/A)_{nc} = 0.54 \ \rho_{L} c_{pL} \left[\frac{rg(T_{w} - T_{S})}{\sqrt{A_{T}}} \right]$ $for \ 10^{5} < R_{a} < 2 \times 10^{7}$ $for \ 2 \times 10^{5} < R_{a} < 3 \times 10^{10}$	$(Q/A)_{Total} = 4.35 \times 10^8 (N/A_t)^{\sharp} \bar{V}_{ME}$ + $184 (T_W - T_L)^{4/3}$. $[1 - K^{\sharp} R_b^2 (N/A_T)]$ + $1543 kR_b^2 \sqrt{\sharp}$. $(N/A_T)^{\sharp} T_W - T_L)$

TABLE 1-b. Equations for the Prediction of Critical Heat Flux

Equation	Theoretical Background	Validity	Author	Year
$(Q/A)_{CHF} = K_{u} p_{v}^{1/2} [g\sigma(\rho_{L} - \rho_{v})]^{1/4}$ ku: 0.14 v a/b	Dimensional Analysis	Agrees well with large amount of data for water and organic liquids	Kutateladze (33)	1953
$(Q/A)_{CHF} = 2.5 \rho_{V}(g\alpha)^{1/4} (\frac{\rho_{L} - \rho_{V}}{\rho_{V}})^{1/2}$	Hydrodynamic Analysis	Agrees well with data of water and organic liquids	Addoms and improved by Ivey in 1961 (2)	1954
$(Q/A)_{CHF} = 0.18 \ \rho_{\rm v} \left[\frac{\sigma_{\rm g} (\rho_{\rm L} - \rho_{\rm v})}{\rho_{\rm v}^2} \right]^{1/4}.$ $\left[\frac{\rho_{\rm L}}{\rho_{\rm L} + \rho_{\rm v}} \right]^{1/2}$	Taylor and Helmholtz Instabilities	Agrees well with data of water and organic liquids.	Zuber (55,56)	1958
$(Q/A)_{CHF} = 0.145 \rho_{V}^{1/2} \left[\frac{\rho_{L} + \rho_{V}}{\rho_{L}} \right]^{1/2} .$ $[g\sigma(\rho_{L} - \rho_{V})]^{1/4}$	2-Dimensional interfacial stability analysis	Agrees with Zuber's and Kutateladze's equations	Chang and Snyder (12)	1960

TABLE 1-b. (Continued)

Equation	Theoretical Background	Validity	Author	Year
$(Q/A)_{CHF} = 0.16 \ \rho_{\rm v}^{\rm h}_{\rm Lv} \cdot \left(\frac{\rho_{\rm L}^{+}\rho_{\rm v}}{\rho_{\rm L}^{\rho}\rho_{\rm v}}\right)^{1/2}.$ $[gg_{o}q(\rho_{\rm L}^{-}\rho_{\rm v})]^{1/4}/[1+\left(\frac{\rho_{\rm v}}{\rho_{\rm L}}\right)^{1/2}+\left(\frac{\rho_{\rm v}}{\rho_{\rm L}}\right)^{1/2}+\left(\frac{\rho_{\rm v}}{\rho_{\rm L}}\right)]$	A modification of Zuber's equation	Same as that of Zuber's but with more satisfactory theoretical consideration	Berenson (6)	1960
$(Q/A)_{CHF} = 143 \ \rho_{\rm v} h_{\rm Lv} g^{1/4} (\frac{\rho_{\rm L} - \rho_{\rm v}}{\rho_{\rm v}})^{0.6}$	Hydrodynamic Analysis	Good agree- ment with data of water and organic liquids	Rohsenow and Griffith (46)	1961
$ (Q/A)_{CHF} = \frac{2570 \rho_{\rm w}}{[1 + (\frac{\rho_{\rm w}}{\rho_{\rm L}})^2]^2} (1+0.318 \frac{\rho_{\rm L}}{\rho_{\rm w}})^{1/2} . $ $ \frac{gg_{\rm o}\sigma(\rho_{\rm L}-\rho_{\rm w})}{[\frac{g^2\rho^2}{\rho^2}]^2} 1/4 $. 3-Dimensional analysis of a cylindrical vortex sheet	Limited because of the limited availability of data on the three phase contact angle for various combinations of fluids and solids	Adams (1)	1963

TABLE 1-b. (Continued)

Year	1963	1965	1971
Author	Noyes 1 (40)	Casell 1 and Balzhiser (10)	Balzhiser, let al.
Validity	0.003 < Pr < 11	Agrees well with liquid metals, water, and organic liquids	Needs more supporting experiments
Theoretical Background	A correction of Zuber's equation takes viscosity and conductivity of sodium into consideration	Dimensional Analysis	Hydrodynamic Instability plus temperature pro- file model
Equation	$(Q/A)_{CHF} = 0.144 \rho_{v} \left[\frac{^{0}L^{-\rho}v}{^{0}v} \right].$ $\frac{gg_{o}\sigma}{[^{\rho}L]} \frac{1/4 P_{r}^{-0.245}}{P_{r}}$	$(Q/A)_{CHF} = 1.02 \times 10^6 \frac{\rho_v k}{C_p \sigma}$. $\left[\frac{\rho_L - \rho_v}{\rho_v}\right]^{0.65} P_r^{0.71}$	$ (Q/A)_{\text{CHF}} = \frac{\pi}{24} \lambda \rho_{\text{v}} \left[\frac{\sigma g (\rho_{\text{L}} - \rho_{\text{v}})}{g_{\text{o}} \rho_{\text{v}}^{2}} \right] ^{1/4} . $ $ (\frac{\rho_{\text{L}}}{\rho_{\text{L}} + \rho_{\text{v}}}) ^{1/2} $ $ + \int_{-\lambda/4}^{\lambda/4} \frac{(T_{\text{w}} - T_{\text{g}})}{\cos(\frac{2\pi x}{\lambda})} dx $ $ \lambda = 2\pi \left[\frac{g_{\text{o}} \sigma}{g(\rho_{\text{L}} - \rho_{\text{v}})} \right] ^{1/2} $

TABLE 1-b. (Continued)

Year	1974	
Author	Bankoff, et al. (5)	
Validity	Needs to be proven experi- mentally	
Theoretical Background	Kutateladze equation condensing heat flux within bubbles	
Equation	$(Q/A)_{CHF} = (Q/A)_{CHF,k} + 2(Q/A)_{cond}$ = $k_u h_{Lv} (\rho_v g)^{1/2} [\sigma(\rho_L - \rho_v)]^{1/4}$	$+\frac{2(T_{\mathbf{v}}-T_{\mathbf{g}})}{\frac{\delta}{k}}+\frac{\frac{2\pi RT}{M}\sqrt{\frac{T_{\mathbf{g}}}{\epsilon h_{\mathbf{L}\mathbf{v}}}\rho_{\mathbf{v}}}}$

Many other variables can influence the CHF experimentally and confuse data interpretation. For instance, Tong (53), commenting on the effects of wetting agents and surface conditions on the CHF in pool boiling, noted that the maximum heat flux is practically independent of surface material, cleanliness, and roughness. Berenson (6) found that a smooth surface has a higher superheat at boiling crisis than a rough surface, although both have approximately the same CHF. On the other hand, other investigators have obtained different results. Costello and Frea (53) found that deposits on stainless steel heater surfaces result in at least a 50% increase in the CHF. Ivey and Morris (53) stated that oxidized surfaces appear to yield a higher CHF than that associated with a clean metallic surface. Other parameters such as the diameter, the surface orientation geometry and system acceleration all have small effects on the CHF (53).

C. Boiling Heat Transfer in Transients

A clear interpretation of boiling behavior and CHF in both power and pressure transients remains elusive even though these phenomena have become of increasing importance. Because knowledge of the boiling behavior during power transients is required for nuclear reactor design and operation, basic studies of this phenomena are numerous while investigations of pressure transients have been more recent, relatively few, and unsystematic.

1. Boiling Heat Transfer in Power Transients

Transient thermal behavior originally became of interest because of the possible effects of a prompt supercritical period on fuel cladding integrity in nuclear reactors. Many studies have been performed on

this topic (27,28,30,48,49,51,52,53). Notable early work was performed by Rosenthal and Miller (47) and Johnson and Schrock (28), who studied exponential increases in power to simulate transient conditions in light water reactors. The experimental results have, in general, been consistent and are reviewed by Tong (52). Transient CHF increases as the power impulse time decreases. In fact, when the impulse time is larger than 200 milliseconds, the transient CHF does not deviate significantly from the steady-state CHF for similar subcooling. According to the results of Rosenthal and Miller (47), the power excursion effect on CHF decreases as the initial exponential period increases and gradually becomes insignificant at periods greater than 14 to 30 milliseconds. When time constants are extremely short the transient heat flux can surpass the steady-state value before the boiling transitions to film boiling. These effects are shown clearly by the data of Tachibana, et al. (50) in Fig. 3. These authors have speculated that the CHF during rapid power transients may increase because of a corresponding increase in the number of surface nucleation sites. This supposition agrees with the observation of Hall and Harrison (20) who observed that in extremely rapid exponential power impulses with periods as short as 0.7 millisecond, film boiling was invariably preceded by a short burst of nucleate boiling. Furthermore, peak heat fluxes 5 to 10 times the steady-state CHF values under the same thermodynamic conditions were noted.

On the other hand, experiments of power transients with relatively long time constants have been conducted also. Fontana (16), for instance, studied transients with periods of up to 2 minutes, and Sakurai et al. (48) studied transients with periods up to 1 second. In both cases, the transient CHF was observed to be the same as the steady-state CHF for the same liquid temperature and pressure. Although Fontana found

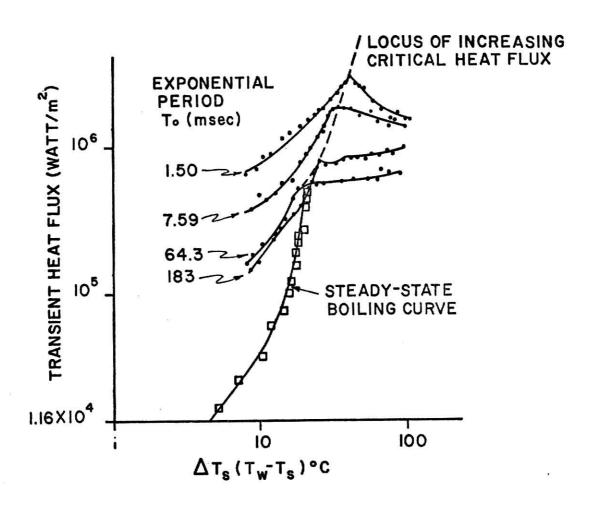


Fig. 3 Effect of Power Transient on CHF. [After Tachibana (50)].

this to be true only for periods greater than 10 seconds, Sakurai observed that transient values of CHF corresponded with steady state levels for periods down to 20 milliseconds. This agrees with earlier power transient experiments conducted by Rosenthal and Miller (47).

It is well known that nucleation and subsequent bubble growth, which are highly dependent on the thermodynamic conditions of the heater and the liquid in the vicinity of the heater surface, directly influence the heat flux. The effects of both power and pressure transients or the CHF may be discussed on these bases. Equation (2.4), which is derived from the surface tension of a spherical bubble and the Clausius - Clapeyron equation, gives an estimation of the surface superheat required to initiate nucleation,

$$T_{V} - T_{S} \approx 2(\frac{R T_{S}^{2} \sigma}{h_{Lv}^{P} L^{r} b})$$
 (2.4)

After nucleation and a few milliseconds after growth has begun, i.e., the later part of the bubble's life, according to Tong's analysis (53), inertia of the surrounding liquid and the surface tension forces can probably be neglected. The pressure can be considered uniform throughout the bubble and liquid. At this stage, bubble growth is governed by the rate at which heat can be supplied from the superheated liquid to the bubble interface to facilitate the vapor formation associated with growth.

If the liquid in the vicinity of the heating surface is very subcooled initially, only a small amount of natural convection heat transfer
occurs. During a subsequent power transient, the heater surface temperature will surge because the heat input is rapidly increasing, but the
surrounding liquid, which remains cool, is not favorable to bubble growth.

However, the fluid temperature lags behind at a value below that necessary for incipient bubble growth. Nucleation may occur, however, on the heater surface due to rapid increase in surface temperature. Transition to film boiling results, as shown by the data of Rosenthal and Miller (49), Johnson and Schrock (28), in Fig. 4. The smaller the impulse power period, the more the temperature overshoot and the shorter the delay time.

If the heater is initially hot enough that the liquid layer surrounding it is close to the saturation temperature, the increasing power during the transient is more effectively transferred because of bubble nucleation and growth. In this case, the transient CHF is significantly different. The effect of the initial heat flux level on CHF is to lower the transient CHF considerably at low initial heat flux and to approach steady-state CHF at high initial heat flux (51).

2. Boiling Heat Transfer in Pressure Transients

The situation during a pressure transient is substantially more complex due to the changing thermodynamic conditions in the bulk fluid. As the pressure decay is initiated, the surface superheat increases due to the decrease of $T_{\rm sat}$. The bulk coolant, instead of only the liquid layer in the vicinity of the heater, becomes warmer relative to the decreasing $T_{\rm sat}$. The magnitude of the pressure decrease, of course, determines the corresponding decrease in $T_{\rm sat}$. If the transient is sufficiently fast that it completes the depressurization before bubbles can detach from the heater surface, the heat transfer cannot increase until after the transient is completed. If the transient is slow, for example, 7.07×10^4 pascal/sec, there is ample time for many generations of bubbles to grow on the surface and equilibrium may be reached. Therefore the heat flux can follow the pressure change, and in this limiting case, the CHF can be predicted by steady-state correlations.

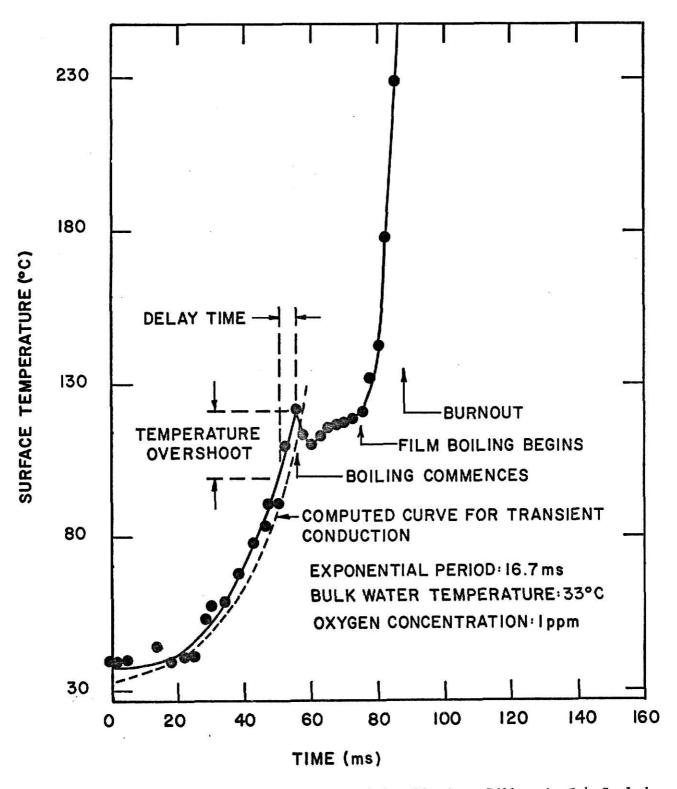


Fig. 4. Typical Temperature-Time Record for Platinum Ribbon in Sub-Cooled Water during Power Transient. [After Rosenthal and Miller (47)].

In practice, the pressure transient is more complicated because bulk flashing is often involved. The heater boundary will be disturbed by bubble agitation and heat transfer mechanisms observed in steady state circumstances may no longer apply shortly after the transient is triggered. No theoretical work has been done on this subject, and only a few experiments have been conducted toward increasing our understanding of this phenomena.

Howell and Bell (23) investigated the pressure transient effect in pool boiling experimentally. The pressurized saturated water was allowed to decay from moderate pressures to atmospheric via a quick-opening valve. Decompression was relatively slow. For example, it required 12 seconds to reach atmospheric pressure from 4.2 \times 10^5 pascal. Because of this slow pressure transient, a link between steady-state and transient cases was attempted. A typical heater temperature behavior during transient boiling is shown in Fig. 5. The relation between transient events and the transient pressure is shown in Fig. 6. The ribbon temperature started to decrease once the decompression began and continued until point c, the inception of film boiling, was reached. The heater temperature then increased until the burnout point was detected. The magnitude of transient CHF appears higher than steady-state CHF by about 20 to 30 percent but independent of the decompression rate. During this experiment, only the initial heat flux was measured; hence, the magnitude of the fluctuation of the heater temperature and the resulting variation in heat flux were not indicated. The time from the start of decompression to the burnout point was discussed more fully. Thus, knowing the initial conditions and the decay rate, the time from point A to point B on Fig. 10b can be calculated. The period between point B until film boiling begins (from

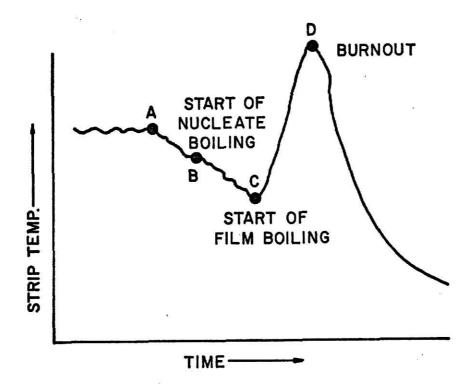


Fig. 5 Ribbon Temperature Behavior During Transient Boiling in Decompression ($P_0 \simeq 4$ atm). [After Howell and Bell (23)].

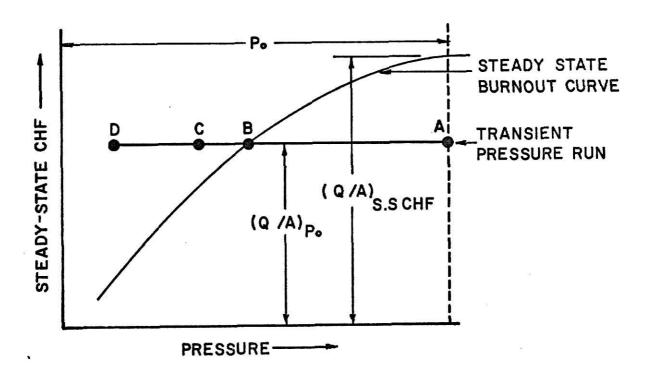


Fig. 6 Relation Between Transient Events and the Pressure. [After Howell and Bell (23)].

point B to point C) is then found from a corrected experimental curve.

The final time segment, between the inception of film boiling and actual burnout (point C to point D) can then be found analytically by assuming complete insulation of the heating surface and calculating the time required to reach the temperature of material failure. The time to burnout was observed to increase with decreasing rate of pressure release. Flashing, encountered in these experiments, was deemed responsible for the delay of the boiling transition corresponding to the steady-state prediction.

Cermak, et al. (11) examined the effect of pressure blowdown on flow CHF in rod bundles. It was observed that for pressures of 6.3×10^5 to 10.5×10^5 pascal and relatively short time constants of 1 second, steady-state CHF agreed with the observed transient CHF values to about 5 percent.

Studies of fast pressure transients are very few. Aoki, et al. (3) investigated the boiling and burnout phenomena during very fast pressure transients of time periods of approximately 25 milliseconds from atmospheric pressure down to at most 7.3 x 10⁴ pascal or higher. Heater surface temperature behaviors for different initial heat flux levels were studied. According to their results, the initial heat flux level determines the overall transient behavior, as shown in Fig. 7. The most interesting findings are for type B behavior (see Fig. 7). Specifically, the investigators noted a spike characteristic of heat fluxes slightly lower than the steady-state CHF of the final equilibrium state, a secondary boiling during which boiling transition occurred, (see Fig. 7), and an initial temperature drop for all levels of heat flux. The third observation was also reported by Howell and Bell, but the first two findings are

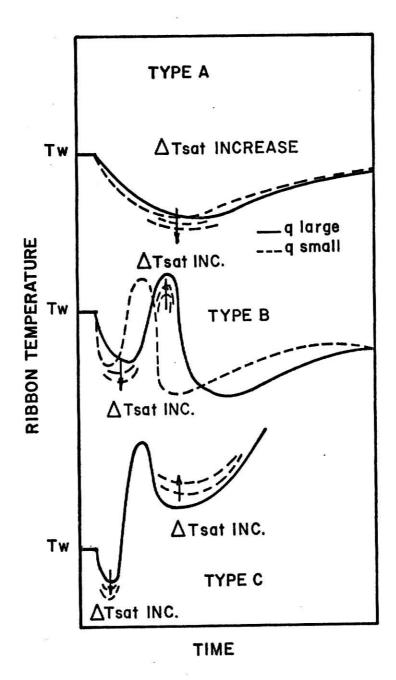


Fig. 7. Ribbon Temperature Behavior during Transient Boiling in Rapid Decompression from Atmospheric Pressure. [After Aoki, et al. (3)].

unique. By the use of high speed photography, Aoki, et al. were able to explain the bubble growth behavior in more depth than Howell and Bell. Thus, at medium heat flux level, the temperature falls during the initial sudden decompression but its drop ceases when the generated bubbles near the wall begin to coalesce. After the heater wall is covered with a large single coalesced bubble, the wall dries out and the temperature increases steeply. However, as the large bubble is removed from the surface by buoyancy, the temperature begins to drop and stable nucleate boiling insues.

When the initial heat flux was higher than the steady-state CHF, boiling transition and subsequent stable film boiling were observed. The transition, however, did not occur during the pressure drop but during the secondary boiling instead. By incorporating the thin liquid layer model for the bubble generation and detachment at high heat flux in nucleate boiling with their observations, the authors contend that boiling transition occurs when a second generation of bubbles is formed inside the thin liquid layer underlying the bubble layer from the primary boiling.

Summarizing previous decompression work, the following points should be noted. Howell and Bell tried to predict transient boiling transition on the basis of the initial conditions, i.e., operating heat flux, pressure decay rate, and range. And perhaps more importantly, they attempted to answer, if boiling transient occurs, at what stage in the decompression will it happen? Aoki, et al. (3), on the other hand, were interested in determining whether the boiling transition occurs during or after the pressure transient. The small time constant and comprehensive recording of heater temperatures accompanied by photographic studies of bubble growth made their investigation more fundamental.

In steady-state boiling, the CHF is defined at the critical point as indicated in the beginning of this chapter. In power transients, the CHF is similarly defined as the point at which the heat flux decreases with further increase in heater temperature. Following the same line of thought and considering that the variation of the heat flux during the transient event is not controllable, perhaps the best definition of the CHF in pressure transients is that minimum initial heat flux which is sufficient to cause boiling transition during or shortly after the pressure transient event. Howell and Bell obtained some data for this purpose, but more data are necessary to establish a comprehensive relation. Aoki, et al. were unable to provide a comprehensive data set since the degree of superheat in their apparatus also affected the rate of decompression and the transient pressure hehavior. Their observation of different temperature behaviors is of potential importance in the better understanding of pressure transient boiling phenomena.

The research reported herein is the initial stage in a comprehensive study of boiling transition under decompression. As an extension to the work described above, subcooled as well as saturated decompressions were studied to better elucidate the effect of flashing on the wire temperature behavior and boiling transition. In subsequent chapters, the experimental apparatus, technique, results, and conclusions are covered.

III. EXPERIMENTAL EQUIPMENT AND PROCEDURES

A. Experimental Equipment

Because only low liquid temperatures (less than 100 °C) and low pressures (less than atmospheric) were involved in this experiment, pyrex glass and aluminum were used to construct the entire apparatus. In order to produce very fast transients, a diaphragm bursting technique was used instead of a quick opening valve. When burst by a spring-loaded plunger, this method has yielded consistent results. A fast response temperature sensor, in the form of a platinum wire, was used for following the transient temperature history for a period of up to 500 milliseconds. The platinum wire was also used as the heater. Details of the test assembly and the measuring circuit are described in the following sections.

1. Test Assembly

A schematic diagram of the test assembly is shown in Fig. 8. The test section was composed of an 800 ml Pyrex glass beaker that was fitted with an aluminum flange. The reduced pressure, or blow-down tank, was made of an eleven - liter glass vessel which has an aluminum flange and valving to provide access to a vacuum pump. Aluminum pipe, 5.08 cm I.D., was used to connect the two sections through a diaphram section. Prior to each transient run, an aluminum foil was clamped between the two flanges of the diaphram section, and the pressure in the dump tank was reduced to sub-atmospheric, in most experiments to 400 torr. An aluminum tank was built originally to provide a constant temperature environment for the test section during steady-state boiling experiments. This was found to be unnecessary, and later experiments were performed

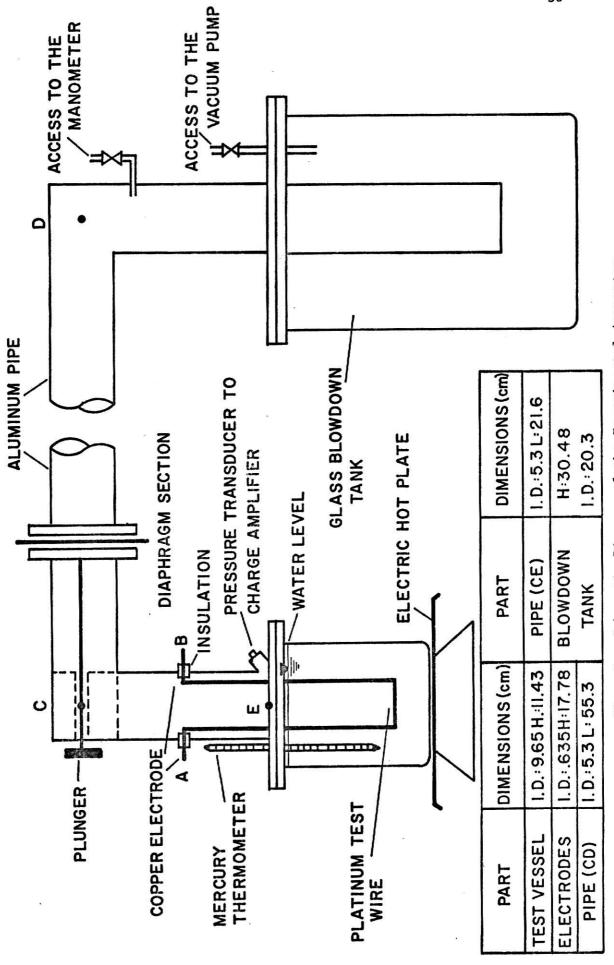


Fig. 8. Schematic Diagram of the Experimental Apparatus.

using a hot plate with a magnetic stirrer as the bath temperature control. Distilled and deionized water was used in the experimentation.

No further filtration or purification was performed.

2. Electrical Circuit

The circuitry used in making steady-state and transient measurements is shown in Fig. 9. Resistors R_1 , R_2 , and R_3 are one 500 ohm and two-1 ohm high wattage resistors respectively for the control of current through the test wire. Resistor R_s is a 0.226 ohm precision shunt resistor. Current through the test wire is proportional to the voltage across R_s . Resistor R_s is the test wire resistance. The test specimen consisted of a 0.0127 cm-dia platinum wire approximately 2.0 cm in length. It was soldered on the ends of two copper electrodes and immersed in the water bath in the test vessel.

In steady-state measurements, as shown in 9a, the test wire was kept 8 cm below the water level to avoid surface effects as suggested by Aoki (3). The steady-state boiling curves for various pressures have been obtained by the use of a digital voltmeter to read wire voltage across A-B and current during operation up to the burnout point.

Two 12-volt lead batteries provided stable wire current during each test run. Bulk water temperature was monitored by a mercury thermometer inserted inside the test vessel. Static pressure was read directly on a u-tube mercury manometer. Steady state resistances were measured by a Leeds and Northrup Type S bridge circuit.

For the transient experiments, a Tektronix type 551 dual-beam oscilloscope with a four-trace preamplifier was used to record the voltage variation across and the current through the test wire, in addition to the output signal from a Kistler type 603 piezoelectric

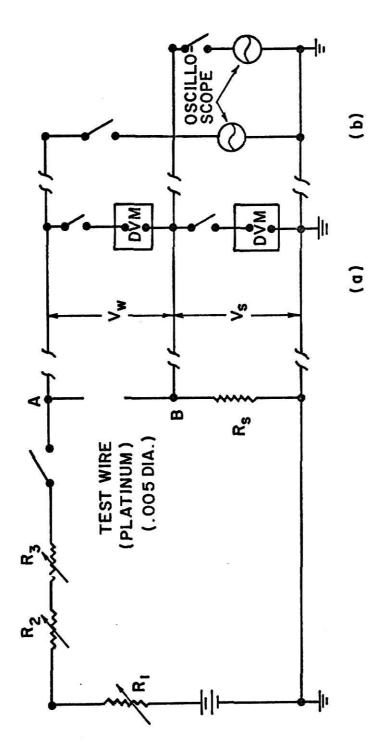


Fig. 9. Circuitry for Steady-State and Transfent Measurements.

transducer with Kistler type 503 charge amplifier. The transient sweep on the oscilloscope was recorded by a Polaroid scope camera.

B. Experimental Procedures

1. Steady-State Experiments

The following procedures were followed to produce steady-state boiling curves at different pressures:

- (1) The test vessel was cleaned and filled with deionized water which had been boiled for at least one hour. The test vessel was then placed in the constant temperature tank which had been previously brought up to the desired temperature. Later experiments were performed using a hot plate as the bath temperature control.
- (2) The test wire was inspected by a 50X optical microscope and cleaned with benzene. The wire was heated in air with a two amp. current for twenty minutes after it was soldered to the ends of the electrodes. The length of the wire was measured to ± 0.05 cm.
- (3) The electrical test assembly was inserted into the beaker. The system was closed and brought to the desired temperature at which time the resistance between points A and B on Fig. 8 was recorded. A vacuum pump was used to lower the pressure to the desired level.
- (4) Power was applied to the test wire, and the wire voltage, V_w , the voltage across the precision resistor, V_s , and the voltage across the power supply, V_p , were recorded.
- (5) The power level was raised incrementally in order that at least 10 data points were taken prior to boiling transition.
- (6) Information in the steady film boiling region was obtained by increasing the power until the wire burned out physically.

2. Transient Experiments

The following procedures were followed in recording the wire temperature behavior during and shortly after the pressure transient.

- (1) The test assembly was prepared prior to each run in accordance with steps 1 and 2 from section 1 above.
- (2) The test assembly was inserted into the test vessel and the whole system was brought to the desired temperature. The resistance between A and B in Fig. 8 was then recorded.
- (3) The aluminum was inserted in the diaphram section and clamped.

 The pressure in the blow-down tank was reduced to 400 torr by the use of a mechanical vacuum pump.
- (4) The DC power supply was turned on and the current through the test wire was brought up gradually to a pre-planned value by adjusting the variable resistors.
- (5) The oscilloscope was DC balanced and proper sweep speed and sensitivity were selected.
- (6) The triggering level on the oscilloscope was pre-set. The triggering level ajustment was the determing factor in the system sensitivity. The beginning of the transient was not recorded, since the scope was triggered by the pressure response. The voltage traces actually represent the transient behavior after several milliseconds into the transient.
- (7) The voltage across the wire, V_w , the voltage across the precision resistor, V_s , and the voltage across the battery, V_p were recorded. Final adjustment of the system pressure was made.
- (8) The shutter control line of the scope camera was pressed. The plunger was released, breaking the diaphragm and then the shutter control was released.

- (9) DC power supply was turned off and the access valve was opened.
- (10) The Polaroid picture was developed.

IV. RESULTS AND DISCUSSION

A relatively straightforward but carefully controlled series of experiments have been performed. It was confirmed that the boiling mechanism is affected significantly during and shortly after fast transients. For instance, if flashing occurs in the bulk coolant, the mechanism of bubble formation and growth may be quite different than that expected when only local boiling around the wire is present. It is surprising, therefore, to observe that even when there was no bulk flashing, i.e., in the subcooled blowdown case, wire temperature behavior was similar to that in the presence of flashing. The understanding and the explanation of these boiling phenomena must be based in part, on a clear understanding of the steady state nucleate boiling mechanism, especially at high heat fluxes near the CHF. Unfortunately, as indicated in Chapter 2, this region is not yet understood fully. In this chapter, the results of both steady-state and transient boiling studies will be covered. Discussion is based on currently accepted boiling models, due to the lack of high speed photographic data in these first experiments.

A. Steady-State Experiments

In order to gain operating experience and to assure the consistency of the apparatus, steady-state boiling experiments were performed as a preliminary test. Data were collected under pressures from atmospheric to 200 torr in a nearly saturated pool of water to construct steady-state boiling curves. The circuitry used for steady-state measurements has been shown previously in Fig. 9a.

1. Interpretation of Steady-State Data

To calculate the temperature of the test wire from its resistance, the resistance-temperature relation of pure platinum wire was used:

$$R_W(T_W) = R_W(0^{\circ}C)(1 + \alpha T_W + \beta T_W^2)$$
, (4.1)

where

$$\alpha = 3.95 \times 10^{-3} \text{ deg}^{-1} \text{ °K},$$

$$\beta = 5.85 \times 10^{-7} \text{ deg}^{-1} \text{ °K} \quad (\text{Ref. 9}),$$

 $T_{\mathbf{w}}$ = the temperature of test wire in degrees of Celsius,

 $R_{\widetilde{W}}(T_{\widetilde{W}})$ = the resistance of test wire at $T_{\widetilde{W}}$ in ohms.

 $R_W^{}(T_W^{})$ can be calculated from the experimentally obtained values of $V_{1}^{}$ and $V_{W}^{}$ was shown below. Due to the existence of a connection resistance and the small resistance of the copper electrodes, $V_{W}^{}$ must be corrected by an amount $V_{Stray}^{}$

where

$$V_{\text{stray}} = I R_{\text{stray}},$$
 (4.2)

and

$$I = V_I/R_s . (4.3)$$

The stray resistance, $R_{\rm stray}$, is assumed to be equal to the difference between the system resistance, $R_{\rm sys}$, and the theoretical reistance, $R_{\rm th}$,

$$R_{\text{stray}} = R_{\text{sys}} - R_{\text{th}} , \qquad (4.4)$$

The resistance of the wire under operating conditions is then calculated as

$$R_{w} = V_{w}'/I = (V_{w} - V_{stray})/I$$
 (4.5)

By solving Eq. IV.1, it is obtained:

$$T_{W} = \frac{-\alpha \pm \sqrt{\alpha^2 - 4\beta\gamma}}{2\beta} , \qquad (4.6)$$

where

$$\gamma = 1 - \frac{R_{w}(T_{w})}{R_{w}(0^{\circ}C)}.$$
 (4.7)

The operating heat flux $\left(Q/A \right)_{\mathrm{op}}$, can be calculated by the formula

$$(Q/A)_{OD} = IV_{W}^{\dagger}/A , \qquad (4.8)$$

where A is the surface area of the test wire.

The wire temperature was assumed uniform both axially and radially.

The radial temperature distribution in a long cylindrical rod with internal heat generation is

$$\frac{T_{w(r)}}{T_{max}} = 1 - \frac{\dot{Q} r_o^2}{4kT_{max}} (\frac{r}{r_o})^2.$$
 (4.9)

where r is the radial distance from the centerline.

Because r_0 is only 0.0635 cm, and the thermal conductivity of platinum, k, is 0.722 (Watt/cm °K) at 27 °C or 0.775 (Watt/cm °K) at 927 °C, the right hand side of Eq. IV.9 is not far from unity, i.e., $T_w \approx T_{max}$. Thus, the temperature distribution from the centerline to the surface is nearly uniform. However, neglecting any end effects may be an invalid assumption in the case of short wires, e.g. length of 1 cm or less. In this experiment, because the length of test wire is close to 2.5 cm and the temperature behavior of the entire wire was measured instead of the local temperature, the end effects were not considered. However, if the temperature distribution along the wire is known, the constant which relates the maximum temperature and the mean temperature of the wire can be calculated.

2. Discussion of Steady-State Results

The data were processed by a short computer program listed in Appendix A. All the results are presented in Appendix B. The boiling curve under atmospheric pressure was plotted along with Petersons' (42) and McAdom's (36) data and is shown in Fig. 10. Rohsenow's equation in Table I was also used and the calculated results are presented in the same figure for comparison.

The free convection region of the boiling curve in the present study deviated slightly from the previous results. This may indicate that an error was introduced because of uncertainty in wire length. For low heat fluxes, the heater temperature is very close to the bulk fluid temperature; therefore, the errors made in the measurement of the length and end effects, will be substantially more important than in the boiling region. One other factor that may have contributed to the uncertainty is the possible influence of augmented free convection on the heat transfer. Since the pool was held at constant temperature by a bottom hot plate, the resulting free convection may increase the heat transfer over that normally present from that induced only by the thermal boundary layer in the near-wire region.

The boiling mechanism becomes very unstable as the boiling transition is approached. Because the controlled variable in this experiment was the heat flux (as in most other studies except Peterson's (42)), it was impossible to precisely detect the critical point. As the point of transition was passed, the physical appearance of the wire changed noticeably. At one location, the wire began to grow red. This local behavior subsequently spread over the wire's entire surface except near

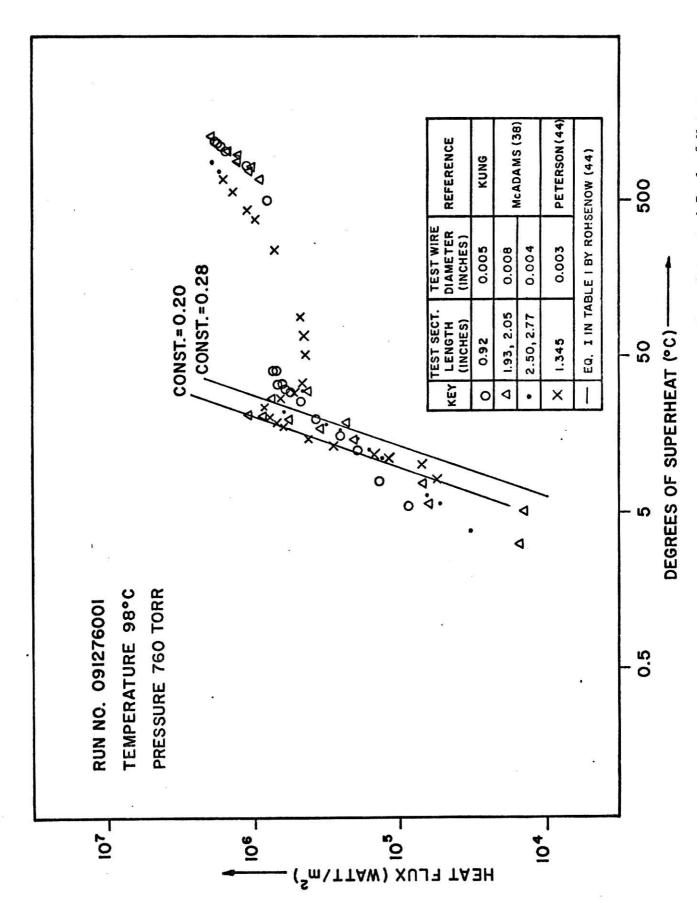


Fig. 10. Steady-State Boiling Curve at Atmospheric Pressure in a Saturated Pool of Water.

the ends where the surface temperature was lower. This phenomena corresponds to the case where the wire becomes engulfed in vapor and the boiling shifts to position D on the boiling curve (see Fig. 1), film boiling with radiation augmentation. The heat flux, as this transition was observed, was taken as the CHF. The results, shown in Fig. 11, are in good agreement with Aoki's data for a similar apparatus.

Frequently, the boiling transition was passed through without a clear indication of nucleate boiling. This behavior has been observed also by Aoki, et al. (3).

The steady-state experimentation described above gave confidence that the measurement system was reliable. Following this, transient data were gathered.

B. Boiling Heat Transfer During and Shortly after the Pressure Transient

1. The Interpretation of Transient Experimental Data

During decompression testing, the pressure decayed to 420 torr from atmospheric. If the water reached its saturation pressure during the run (as determined by the initial bulk fluid temperature), the final pressure was at most 30 torr higher than 420 torr. All the transient data were recorded by Polaroid pictures of oscilloscope traces. The test conditions examined during the experimentation are shown in Table 2. Four levels of heat fluxes were tested. At low heat flux, no bubbles appeared before or after the pressure transient. The ratio of the initial heat flux, $(Q/A)_{op}$, to the CHF corresponding to the final state, $(Q/A)_{CHF}$, was about .6. Medium heat flux level was when $(Q/A)_{op}/(Q/A)_{CHF}$ was about .8. High initial heat flux levels were close to the Q_{CHF} , and very high initial heat flux was higher than Q_{CHF} . For each level of heat flux, runs of different superheats and subcoolings were tested as recorded in Table 2.

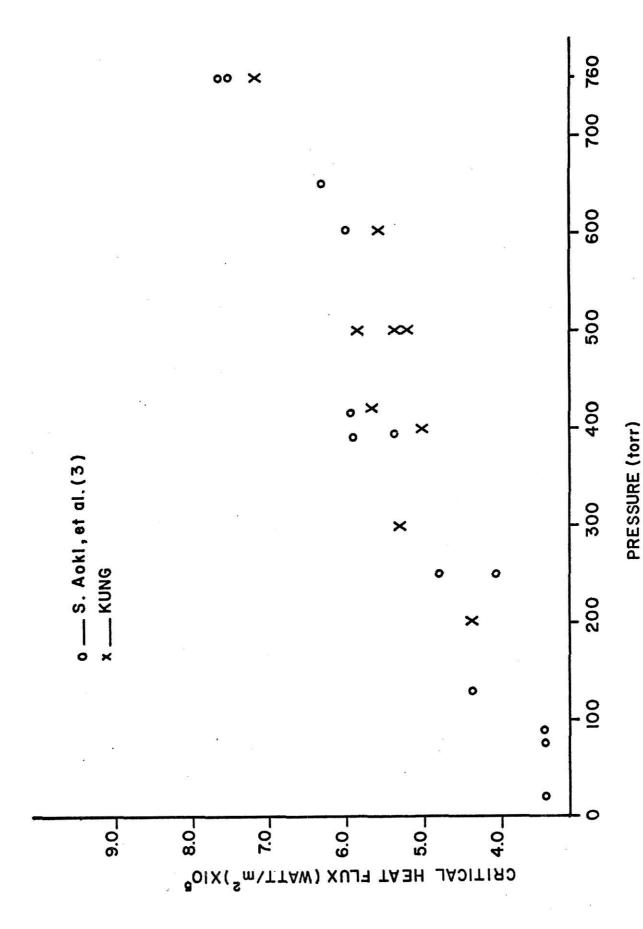


Fig. 11. Steady-State CHF at Sub-Atmospheric Pressures.

Summary of the Test Conditions of the Transfent Experiments TABLE 2.

Note	Water became superheated after the transfent	from ~ 0 to 7°C. $[(Q/A)_{op}/(Q/A)_{CHF}]$ was	approximately 0.5, 0.8, 1.0, and greater than 1.0. Corresponding levels are defined	as low, medium, high, and very high heat fluxes respectively.	Water remained subcooled after transient was initiated. Degrees of subcooling ranged from 1 to 34°C.				
Initial Water Temp. ${ m T_L}({}^{\circ}{ m C})$	× 84				× 84				
Number of Runs	10	7	21	12	7	5	8	6	
Initial Heat Flux Level $[(Q/A)_{Op}/(Q/A)_{CHF}]$	۰ 0.5	8.0 ∿	~ 1.0	> 1.0	۰ 0.5	8.0 ~	۰ 1.0	> 1.0	

(Q/A) $_{
m CHA}$ in saturated water under 420 torr is approximately 5.7 x 10^6 W/m 2 according to Figure 11. NOTE:

It was assumed that the power supply provided a constant current source throughout the run; therefore, the voltage across the test wire, V_w , is proportional to the wire resistance, R_w ,

$$V_{w} \simeq I R_{w}$$
 (4.10)

The time dependence of the current, I, introduces significant errors only at the end of the run, i.e. at the incidence of film boiling and later, and its neglect is therefore justified. From Eq. 4.10 it is apparent that

$$\Delta V_{W} \simeq I \Delta R_{W}$$
 (4.11)

Differentiating Eq. IV.1, it is obtained,

$$\frac{\Delta R_{W}}{R_{Q}} = \alpha \Delta T_{W} + 2\beta T_{W} \Delta T_{W} \simeq \alpha \Delta T_{W}$$
 (4.12)

Therefore,

$$\Delta R \approx k \Delta T_{W}$$
 since $\alpha >> \beta$, (4.13)

where

$$k = \alpha R_{O} = \alpha R_{W}(0 \text{ °C}) . \qquad (4.14)$$

Hence the voltage trace of the wire during the pressure transient gives a direct indication of the heater temperature variation. A typical oscilloscope trace is shown in Fig. 12. The top trace is the pressure transient behavior and indicates that the transient from 760 torr to 420 torr was complete within 10 milliseconds of initiation. The middle curve is the voltage across the test wire with respect to time. The lowest is the current trace which demonstrates that the current was constant during the time frame of interest. The small fluctuations are typical of medium and high heat flux levels because of nucleation and

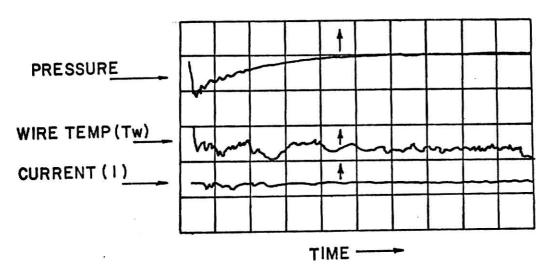


Fig. 12. Typical Pressure Transient Data.

Note: Arrows in the picture indicates positive increase direction.

subsequent bubble detachments altering the surface temperature at high frequency (39). In superheated experiments, the response of the pressure transducer becomes unstable. The oscilloscope trace indicated that after the first dicompression, the pressure would decay further and went beyond the zero pressure level which is impossible. Hence, the pressure in superheated cases is assumed similar to that in subcooled cases. A summary of the experimental conditions and oscilloscope parameters for all the polaroid pictures cited here and after are listed in Table 3.

During the transient part of the experiment, 119 pictures were taken. The first 25 pictures were not useful for analysis because of the immature experimental technique. Out of the 94 following pictures, only 64 proved useful. Examples of these runs are shown and discussed in the following sections. Reduction of transient data is covered in Appendix C.

2. Discussion of the Results of Pressure Transient Experiments

i. Subcooled Case

When the bulk water remained subcooled during and after the pressure transient, no bulk flashing was induced. Typical results at different heat flux level are shown in Fig. 13, 14, and 15.

For low heat fluxes, no bubbles were generated on the heater surface before or after the transient. The heater temperature, $T_{\rm w}$, was observed to drop almost 10 °C immediately after the transient. In this case, because natural convection was not affected significantly by the pressure, $T_{\rm w}$ always drifted back toward its original value at the end of the transient period.

Summary of the Experimental Conditions and Oscilloscope Parameters for Fig. 12-20. TABLE 3.

Estimated Initial Temp.	Drop of Tw	7	8.5	8.2	10.13	11.9	15.7	12.9	6.5	11.8
Observed Initial Voltage	Drop or V _w (volts)	0.035	0.03	0.04	0.058	090.0	0.065	0.064	0.03	0.05
ltage Sensi (Volts/Di	Vpressure	0.5 (AC)	0.5 (DC)	0.5 (AC)	0.5 (AC)	0°5 (DC)	0.05 (DC)	0.05 (AC)	0.5 (AC)	0.5 (DC)
	$^{ m V}_{ m I}$	0.05 (DC)	0.02 (DC)	0.05 (DC)	0.05 (DC)	0.05 (DC)	0.05 (DC)	0°02 (DC)	0.05 (DC)	0.05 (DC)
	V	0.05 (DC)	0.05 (DC)	0.05 (DC)	0.05 (DC)	0.1 (DC)	0.05 (DC)	0.05 (DC)	0.1 (DC)	0.1 (DC)
Sweep Speed	(msec/Div.)	50	50	50	50	20	50	50	20	20
I.	(± 0.2-c)	11	02	62	82	69	06	06	92	06
(Q/A) (Q/A) (Q/A) (Q/A)	CHE	8°0 ~	9.0 ∿	8°0 ~	8°0 ~	1.1 س	∿ 0.8	۰ 0.8	~ 1.0	~ 1.0
No.	Figure	12	13	14	15	16	17	18	19	20

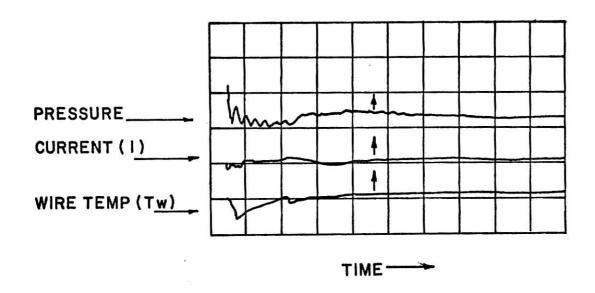


Figure 17.

Typical Sub-Cooled Data at Low Heat Flux.

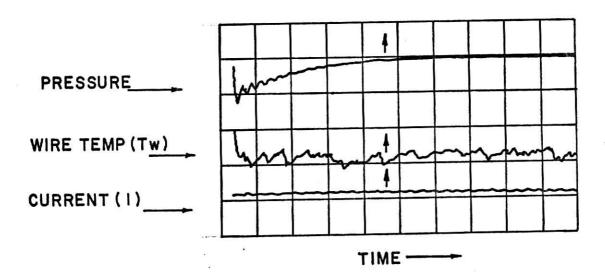


Figure 18.

Typical Sub-Cooled Data at Medium Heat Flux.

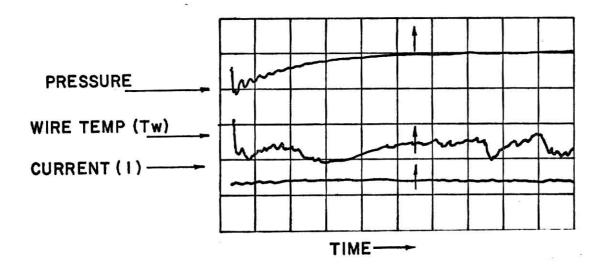


Fig. 15. Typical Sub-Cooled Data at Medium Heat Flux.

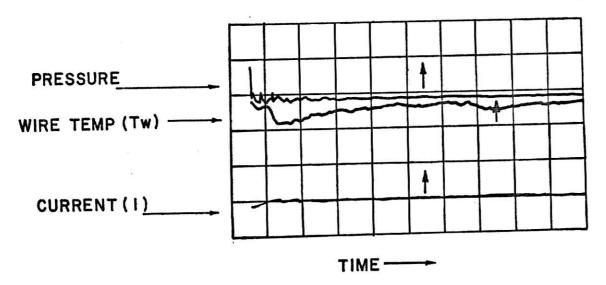


Fig. 16. Typical Sub-Cooled Data at High Heat Flux.

At medium levels of initial heat flux, nucleation was observed on the heater surface before the pressure transient. This bubble formation was enhanced by the reduced pressure during and after the decompression. Therefore, the heater temperature, which dropped initially during the decompression, remained at the reduced level due to the more vigorous nucleate boiling. The same result was obtained when another series of experiments with large initial subcoolings was conducted. In this case, a higher heat flux was applied. Consequently, even though the bulk fluid was well below the saturation temperature initially, nucleation on the wire was still present. Two then stayed at the new reduced level of temperature after the transient. However, if the degree of initial subcooling was greatly increased, only natural convection contributed to the heat transfer. Therefore even if the applied heat flux was relatively high, Tw, in this case, drifted back to its original level. These results were shown in Fig. 15.

At high and very high heat flux levels, although $T_{\overline{W}}$ tended to drop initially as in the former cases, the boiling transited to film boiling and $T_{\overline{W}}$ increased due to reduced heat transfer. However, the delay from the start of the transient to the time of boiling transition was found to be very random. It ranged from 20 milliseconds to several hundred milliseconds. No boiling transition occurred during the decompression under any condition.

Two important results of subcooled depressurization deserve special emphasis. The first is the initial temperature drop under all conditions, even the most subcooled test in which no bubble mechanism existed during the entire event. This can be explained, perhaps, by turbulence in the surrounding liquid caused by the pressure wave. The rate of heat transfer

by convection between a solid boundary and a fluid is generally evaluated by means of Newton's Law of Cooling:

$$q_{\text{surface to fluid}} = A\overline{h}_{c} (T_{w} - T_{\infty})$$
, (4.15)

where

A = heat transfer surface

 \overline{h}_c = the convection coefficient,

 T_{∞} = the water temperature,

 T_{w} = the temperature of test wire.

Once the pressure transient was initiated, the heat balance between the heat generation, which follows Ohm's law as

$$q_{generated} = I^2 R_w(T_w)$$
, (4.16)

and the convection heat flux was disturbed. The convection coefficient was increased because of the turbulence, and the wire temperature started to drop. This forced convection heat flux continued to be larger than the heat generation; therefore, the wire temperature continued to drop until thermal equilibrium was reestablished. In order to evaluate the magnitude of the change in \overline{h}_c , the condition prior to the pressure transient is referred to as 1, and 2 is the state of the reestablished equilibrium. If the following equations are used:

$$I^{2}R_{W}(T_{1}) = (T_{1} - T_{L})\overline{h}_{c1} A$$
, (4.17)

$$I^{2}R_{W}(T_{2}) = (T_{2} - T_{L})\overline{h}_{c2} A$$
, (4.18)

or

$$\overline{h}_{c1} = \frac{I^2}{A} \frac{R_w^{(T_2)}}{(T_1 - T_1)}$$
, (4.19)

$$\overline{h}_{c2} = \frac{I^2}{A} \frac{R_w(T_2)}{(T_2 - T_L)},$$
 (4.20)

where current, I, is assumed constant, then from Eq. IV.1,

$$R_{w}(T_1) \simeq R_0(1 + \alpha T_1)$$
, (4.21)

$$R_{W}(T_{2}) \simeq R_{O}(1 + \alpha T_{2})$$
 (4.22)

It may be easily found that the ratio of the heat transfer coefficients is:

$$\frac{\overline{h}_{c2}}{\overline{h}_{c_1}} = (\frac{1 + \alpha T_2}{1 + \alpha T_1})(\frac{T_1 - T_L}{T_2 - T_L}) . \qquad (4.23)$$

As an example, for run #64, T_1 was 170.5 °C, T_2 - T_1 = -10 °C, and T_1 = 82 °C; therefore,

$$\frac{\overline{h}_{c2}}{\overline{h}_{c1}} = \frac{1 + 0.00395 \times 165}{1 + 0.00395 \times 175} \frac{88.5}{78.5},$$

$$= 1.10.$$

The convection heat transfer caused by turbulence is estimated, in this case, to have been increased by 10%. Of course, further studies of the velocity field next to the wire are necessary to investigate this phenomena more completely. But, it can be seen that violent agitation by bubbles is not necessary to promote a heater surface temperature drops.

The second point of importance is the wire temperature behavior after the transient. It remained at a depressed level at high heat fluxes, but drifted back to its initial value at low heat fluxes. The explanation given in the first part of this section, which notes the presence of nucleate boiling for medium or high heat fluxes after the transient is applicable. But further study is needed because the heat transfer

immediately after the pressure transient is of special importance to the understanding of the effect of flashing, which will be discussed in the next section.

ii. Superheated Cases

Superheated blowdown is much more complex than the preceeding subcooled cases, since bulk flashing is generated which influences the wire heat transfer. In this series, water in the test vessel was slightly subcooled initially so that vapor bubbles were generated throughout the bulk volume once the decompression was initiated. Hence, the objective was to determine the direct effect of the pressure transient and the indirect effect of the flashing on the heat transfer and boiling transition.

When the initial heat flux was low, T_w dropped during the pressure transient and remained at the lower level. By increasing the heat flux, the effect of flashing could be seen by comparing the transient wall temperature behavior with the subcooled cases. Instead of aiding the heat transfer, the flashing actually caused the heat transfer from the wire to deteriorate which resulted in the temperature peaks shown in Fig. 17.

As the initial heat flux was increased, the results, shown in Fig. 18, became unpredictable, because during some runs nucleate boiling persisted during and after the transient while during others, the boiling transitioned to film boiling with radiation augmentation. The characteristic fluctuation observed by Aoki, et al. as shown in Fig. 7, however, was not duplicated in these runs.

High initial heat fluxes were also studied. In all cases, film boiling was evidenced by the rapid temperature rise, as shown in Fig. 19, 20.

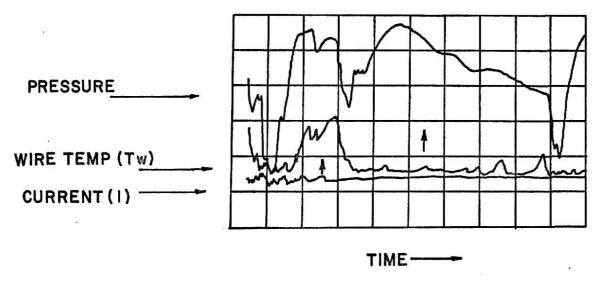


Fig. 17. Typical Superheated Data at Low Heat Flux.

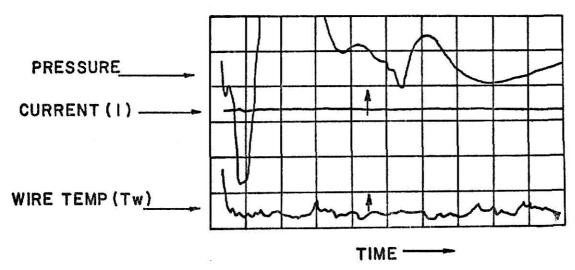


Fig. 18. Typical Superheated Data at Medium Heat Flux.

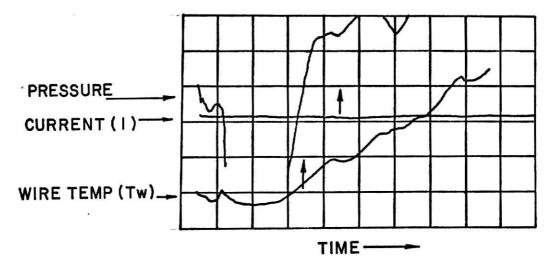


Fig. 19. Typical Superheated Data at High Heat Flux,

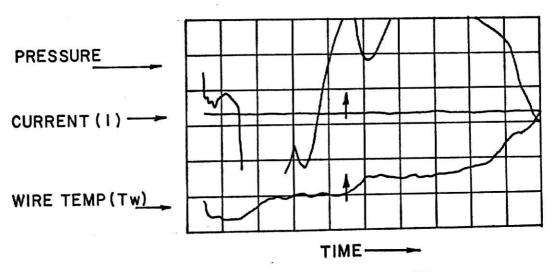


Fig. 20. Typical Superheated Data at High Heat Flux.

In all the runs, T_w dropped during the decompression, and rose rapidly as film boiling began. The secondary boiling reported in Aoki's work

(3) could not be clearly identified. However, there was a delay time to boiling transition under all conditions, a result that should be regarded as one of the most important findings of the present study.

It has been explained previously, see for instance Howell and Bell (23), that the drop in wire temperature during the decompression is due to the violent flashing accompanying the decompression. That this is not the case is shown clearly by comparing the present subcooled blowdown results. The wire temperature drop during the transient was observed to be remarkably insensitive to not only flashing, but also the bulk water temperature. No significant dependence on the bulk water temperature can be found. This is contrary to what Aoki, et al. had reported, wherein the drop of T_w increased with increasing superheat as shown in Fig. 7. These results were generated, however, from a limited amount of data from saturated decompression only. Hence it must be concluded that $\Delta T_{_{\rm tr}}$, the initial temperature drop, is a result of the pressure transient and not unique to flashing. Flashing, which only occurred in the superheated experiments, probably did not become significant during the rapid pressure decay and, therefore, could have had little effect on ΔT_{w} . Hopper and Abdelmessih's investigation (22) supports this supposition. For decompression times of 5 milliseconds, they observed no bubble nucleation with less than 4 milliseconds delay beyond the termination of the decompression period. Moreover, it was found that several milliseconds were required for bubbles to grow to a significant size.

iii. Conclusions of Transient Experiments

There are several interesting findings from this experiment, some of which had not been reported before. Among these are:

- (1) Boiling transition occurred after the completion of the pressure transient. This observation, as shown in Fig. 19 and 20 was also noted by Aoki, et al (3). But the "secondary boiling" postulated by them to explain this phenomenon was not identified in the current series of experiments. The time to the boiling transition after transient initiation and the characteristic behavior of the heater need to be studied further in order to make a definitive conclusion about the importance of "secondary boiling" and the pressure transient on transition.
- (2) The heater temperature dropped immediately after the pressure transient was initiated in both superheated and subcooled cases.

 The temperature depression in the superheated case was expected, but the temperature drop in subcooled transients has not been reported before. Furthermore, the amount of drop in the subcooled case was comparable to that in superheated depressurization. This, perhaps, indicates the importance of the convection current induced by the pressure wave propagation.
- (3) The amount of the initial heater temperature drop was found to be insensitive to the initial heat flux level and the surrounding coolant temperature. This is different from what Aoki, et al (3) had observed (see Fig. 7). However, the decompression rate in this experiment is approximately twice as fast as that used in the previous work. Since neither bubble nucleation in the bulk liquid nor that on the heating

surface is instantaneous, a critical decompression rate may exist which significantly alters the bubble growth history and, therefore, affects the heat transfer. This obviously places considerable importance on the study of bubble growth phenomena under pressure transients of varying time periods in future research.

(4) The flashing may inhibit rather than inhance the heat transfer.

Temperature peaks, observed only in the superheated cases, may imply that the heater surface experienced partial vapor insulation for a short period of time. The large amount of vapor generated during the process of flashing provides favorable conditions for the formation of a temporary vapor film on the heating surface.

C. Suggestions for Future Research

Since part of the purpose of this study was to be a first step in a comprehensive program, suggestions for future apparatus and areas of future interest are:

- (1) Since the input heat flux dominates the temporal behavior of the heater, as discussed in previous sections, a constant heat flux is desirable during and after the pressure transient. Therefore, a heating source isolated from the heat transfer surface is suggested.
- (2) At high heat fluxes in the nucleate boiling region, high frequency detachments from heating surface cause rapid fluctuations in both local and overall surface temperatures. Efficient averaging methods need to be developed to obtain reliable temperature measurements.
- (3) For rapid decompression work, the pressure history at the location of the heat transfer surface needs to be precisely recorded. A

- pressure transducer for transient hydraulic pressure measurement mounted close to the heat transfer surface is recommended.
- (4) The pressure range needs to be extended, and a variable decompression rate is required to determine the actual effect of pressure transients on boiling heat transfer and transition.
- (5) Fast photography is recommended to investigate transient bubble growth. Such data, in turn, provides part of the information necessary for fundamental qualitative and quantitative analysis of the boiling process.

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APPENDICES

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21	REAL IS(100), LENTH DIMENSION XARRAY(1 DIMENSION VS(100), DIVENSION IBUF(200 DIMENSION TITLE(10) DIMENSION TITL(10) DIMENSION NN(3)	MHERE LENTH: MHERE LENTH: RSYS: PRES: VS: VM: DT: T: T	YEIS=10**3 YCYPIM=4.c/6.0 XFIRS=0.0 XFIRS=0.1 XLEN=6.0 XCYPIN=5.0/8.0 XCYPIN=6.0/9.0 XEAD(5.4000) (TYL(1),1=1,7) READ(5.10000) (TYL(1),1=1,5) READ(5.10000) (TYL(1),1=1,5) READ(5.1000) (TYL(1),1=1,5) READ(5.1) RS.20 FCRMAT(2F6.4) READ(5.1) RS.20 FCRMAT(5.3,F5.2,F5.4,F5.1,F8 READ(5.1) WAITE(6.900) TRMP WRITE(6.9000) TRMP WRITE(6.1000) TRMP
G LEVEL		_	- N
FORTRAN IV	0000 0003 0003 0000 0005 0005	*	0009 0010 00111 00112 00123 00124 00129 0023 0023 0023 0023 0023 0023 0023 00

FORTRAN	IV G LEVEL	21 MAIN	DATE = 77175	10/48/18	PAGE 0002
9630		WRITE(6,20000)			
S		WRITE(6,3000)			
3		WRITE(6,4000)			
63		X Z R R AY (M+1) =-1.0			
63		XARRAY (M+2)=5.0/XLEN	7-2:		
3		YARRAY(M+1)=3.0			
40		YAPRAY (M+7)=4.0/YLFN			
2000		NUM Y=28			
2		NUMX=28			
	U				
	U	CALCULATION			
	U				
0044		A=0.00392			
50		B=-0.J000CC588			
5		RTI=1+A+TEMP+B*TEMP*+2.0			
2		RTI=RTI/RT120			
2		CTUEL FRITHADONADII			
5 8		1 1 1 1 2 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1			
5 0		ADDA - CHITTLE TAILS ADDA DOCTO			
2		AKT 4 H L L L L L L L L L L L C C C C C C C C	•		
S		DC 100 1=1.7		ī	
02		IS(1)=VS(1)/RS			
CS	3	PH(I)=VH(I)/IS(I)-SR			
5		H1(1)=PW(1)+1S(1)**2.0/AREA			
50		H2[[]=H][] *3176.2			
2		7017 (717) 177			
3 8		021107	C 11 1 1 C 2 1 1 1 1 1 1 1 1 1 1 1 1 1 1		
S		1(1)=(-A+(A**2*6-4*0*B*(1*0-H11*K1120*KM(11/K1HEU))**0*5)/(Z*0*B)	0**(103H X/(I)MX#02 IX*	.5)/(Z.0*b)	
05		CT(1)=(T(I)-TEMP)			
53		WRITE(6.5000) I. IS(I). H2(I). H3(I). DT(I). RW(I). T(I)) . DT (1) . RW(1) . T(1)		
3	100	BINI INCO			
3 3	2				
5		** I E C + I TOOO			
S		DC 3C0 1=1,W			
8		XARRAY(I)=ALGGIO(DT(I))			
C		YARRAY([) = ALCG10(H2(I))			
3	30.)	HINT LACT			
2		TECM-1110-10-20			
3					
	, (i			
	. (
1000	3	CALL LIMITS (130-+11-0+23+8+-3)			
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5900	07				
0/00			,		
0071			-1-1-		
2730			- NO FX		
0073		CALL OLDGAX(00., VLEN, YCYPIN, 90., +1)	0.,+1)		
9200			(トモコマー		
0075			,0,,+1)		
97,00			,90.,-1)		
0.77			1		
9730		CALL SYMBCL(0.5,5.0,0.125,TITLE,0.0,20)	,0.0,201		
9100			3.0,20)		ře
0000			.0,20)		
00.81					
00.92	400	CONTINUE			
00.63		CALL PLOT (XLEN+5.0.0.,999)			
0084	3000	FORMAT(//4x, 'ND. ',3X, CLRRENT(AMP)',15X, 'HEAT FLUX',16X, "WALL	AMP) . 15x, . HEAT FLUX . 1	6X. WALL SU	
		IPERHEAT (C) . 3X, "RESISTANCE OF WI	RE (OHM) ., 3X, . WIRE TEMP	(//.(3)*	
00.85	4000	FCRMAT (32X, BTU/FT# #2HR , 5X, KC	AL/M*+2HR*//)		

TAGE COOS																					
01/04/01	13.6,12X,E13						,				ZX, V. ACROS					, 1X, F 5.4, 1X,					
2111	, E13.6, 5x, E13.6, 9X, E1										CSS SHUNT RESISTOR		2x, F4.1, C./)	F4.0,1X,'TCRR'/)	,F5.3,1X, INCH'/)	YSTEM TEMPERATURE ="		ICN./)			
	5000 FCQMAT(3x,12,6X,E13,6,7X,E13,6,3X,E13,6,5X,E13,6,9X,E13,6,12X,E13,12X,E13	6000 FCRMAT(7A4)	12)	313)	2x, 'RUN NJ. ', 2X, 313/)	5141	141)	F4.11	1X,6E13.6,1X,214)	3X, 12, 14X, F6.4, 2CX, F6.41	FICHMAT(1X, CATA POINT", 2X, "V. ACROSS SHUNT RESISTOR", 2X, "V. ACROS	(186./)	/2X, "WATER TEMPERATURE =",	ZX, SYSTEY PRESSURE = ,2X,	2X, 'TEST WIRE LENGTH =", 2X	FORWAT (2X, SYSTEM RESISTANCE AT SYSTEM TEMPERATURE = 11X, FS.4.1X,		20000 FCRMAT(///2X, RESULTS OF CALCULATION //)			
17	FCQMAT	FERMAT	FCPMAT	FUEMAT	FOF "AT	FCRMLT	FCPMAT	FCRMAT	FCP MAT	FCRMAT	FCPMAT	S TEST 1	FCRMAT	FURMAT	FORMAT	FORWAT	1.0HM.	FCRMAT	STOP	GNB	
17 72427 9 41	5000	6009	7000	9000	9000	10000	11033	12000	13933	14000	15000	-	16000	17000	19333	19300	-4	20000		E	
TURINA IN C LEV	9830	00.97	CCBB	9C P S	0600	1600	26.00	6530	7500	2600	9500		1500	0:58	9696	0100		0101	0102	0103	

0.949165E 02 0.977937E 02 0.977933E 02 0.103402E 03 0.11636E 03 0.11636E 03 0.123649E 03 0.133649E 03 0.134247E 03 0.137080E 03 0.137080E 03 0.137080E 03

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SYSTEM RESISTANCE AT SYSTEM TEMPERATURE = .2390 OHM C-860 INCH

SYSTEM PRESSURE = 420. TORR WATER TEMPERATURE = 83.0C

TEST WIRE LENGTH =

8 0.9690 10 0.5510 11 0.2590 12 1.0250 13 0.7600 14 0.7470 15 0.8850 16 0.8850 17 0.236283E 00 2 0.445903E 00 3 0.445903E 00 4 0.173049E 00 4 0.17304E 01 6 0.355752E 01 7 0.45504E 01	► I	KCAL/M**2H 0.128102E 0.461935E 0.207842E 0.19495E 0.310430E	the state of the s		WIRE TEMP.(C) 0.949165E 02 0.977937E 02 0.11536E 03 0.111536E 03 0.11536E 03 0.11536E 03
CF CALCUI	งัญ ตั ตั ตั				·
CURRENT (AMP	HEAT FLU: BTL/FT**2HR			٠, ۴	WIRE TEMP.(C)
500046			674		
7 3 3 3 3 3 3 3		107867E 930026E 770686E			0.124345E 00.00.137080E 03.00.137080E 03.00.137080E 03.00.135408E 03.00.135408E 04.296E 04.00.135408E 04.00.135408

			IRE = .2190 CHM	V. ACRESS TEST WIRE	658	0.0668	0.92	293	1CE	123	144	0.166	175	191	218	254	265	290	358	405	436	.465	500	536	566	566	5,540	635	.651	679	735	. 806	. 839	646	613		303	500-	55	223	260	
PERATURE = \$3.5C	ESSURE = 600. TORR	LENGTH = 0.760 INCH	SISTANCE AT SYSTEM TEMPERATURE	V. ACROSS SHUNT RESISTOR	.057	0.0658	081	092	136	121	163	162	171	153	213	237	257	282	344	364	413	442	0440	500	525	.531	545	557	.613	637	6.70	752	2775	. 789	811	222	346	545	. 551	65.	781	
 WATER TEM	SYSTEN PRE	TEST WIRE	SYSTEM RES	ATA PCINT	-	۳ م	1 4	· In	9	_	nc o	10	11	12	13	51.	15	17	81	10	21	22	23	25	56	27	28	30	IE.	32	76	35	36	37	38	3.4	3 7	42	43	* "	7 9	!

•	WIRE TEMP.(C)		0.111763E 03 0.111763E 03 0.111305E 03 0.112064E 03 0.112064E 03 0.112519E 03 0.112599E 03 0.112599E 03 0.112599E 03 0.112791E 03 0.115791E 03
	RESISTANCE OF WIRE (OHM)		0.215447E 00 0.215125E 00 0.2154865E 00 0.215795E 00 0.21511FE 00 0.21571FE 00 0.21574E 00 0.216576E 00 0.216774E 00 0.216774E 00 0.21774F 00 0.21774F 00 0.21774F 00 0.21774F 00 0.21774F 00 0.21774F 00 0.21774F 00 0.21774F 00 0.2254F 00
	WALL SUPERHEAT(C)		0.170738E 02 0.182629E 02 0.185641E 02 0.185641E 02 0.18584E 02 0.185849E 02 0.190986E 02 0.190027E 02 0.2029189E 02 0.202913E 02 0.202913E 02 0.222913E 02 0.247208E 02 0.247208E 02 0.247208E 02 0.247208E 02 0.247208E 02 0.247208E 02 0.256619E 02 0.256619E 02 0.256619E 02 0.265619E 02 0.365619E 02 0.365619E 02 0.365619E 02 0.365619E 02 0.365619E 02 0.365619E 02 0.365619E 02 0.366619E 02
	×	KCAL/M**ZHR	0.154453E C4 0.203609E 04 0.2381109E 04 0.312611E 04 0.402346E 04 0.402346E 04 0.402346E 04 0.402346E 04 0.125151E 05 0.125151E 05 0.125151E 05 0.215832E 05 0.215832E 05 0.21476 05 0.215832E 05 0.21476 05 0.215832E 05 0.21674E 05 0.21674E 05 0.21674E 05 0.21674E 05 0.21632E 05 0.21
လိုဆိုဆိုဆိုဆိုလ်လ်လ်လ်လ်လ်ကိုက်ကိုက်ကိ •	HEAT FLUX	81L/FT**2HR	0.56535E 03 0.750605E 03 C.877794E 03 0.115244E 04 0.14835E 04 0.14835E 04 0.25505B 04 0.314247E 04 0.461368E 04 0.514503E 04 0.514503E 04 0.514503E 04 0.514503E 04 0.514503E 04 0.514503E 04 0.116915E 05 0.116915E 05 0.116915E 05 0.116915E 05 0.116916E 05 0.10711E 05 0.267760E 05 0.26776
CF CALCULATION	CUPRENT (AMP)		0.2539826 00 0.301150E CO 0.315046E CO 0.3676E CO 0.407282E 00 0.537164E 00 0.537164E 00 0.537164E 00 0.537164E 00 0.75728E 00 0.17912E 00 0.17912E 00 0.17928E 01 0.10521E 01 0.10521E 01 0.10521E 01 0.10521E 01 0.15212E 01 0.15386E CI
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RUN NO. 270776 5

NATER TEMPEFATURE = 86.5C

SYSTEM PRESSURE = 500. TORR

TEST LIRE LENGTH # 0.800 INCH

SYSTEM RESISTANCE AT SYSTEM TEMPERATURE . 2450 CHM

DATA PCINT V. ACROSS SHUNT RESISTOR V. ACRCSS TEST WIRE

355	:63	17		2	112	136	158	178	951	502	236	256	283	214	343	380	157	187	521	565	610	665	702	745	603	822	841	875	6 6 5	910	937	251	175	60	3	C54	CE1	105	133	15	000	780	480	.2800	16	30
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	RESISTANCE OF WIRE (OHM)		224265E 224338E 223821E	.224442E .224949E .225328E	226153E 225939E 226049E	0,226402F U0 0,227415F Q0 0,227478F Q0	2291136	233568E	.232461E	.232053E .231828E	232805E	.233063E
	WALL, SUPERHEAT(C)		196263E 197495E 189851E	199217E 207673E 213975E	224149E 224149E 225993E	0.2437826 02 0.2437826 02 0.2457396 02	277119E 293068E	351616E	333093E	326260E 322493E	338848E	343166E
1550 17500 2.5000 2.5000 2.7600 2.7600 2.5900 3.7600 3.7600 3.7600 3.7600 3.7600 4.6500 4.6500	צ	KCAL/M**2HR	109101E 144023E 181569E	.299672E .457850E .671593E	113895E 140986E	0.156987E C5 0.156987E O5 0.233866E O5	417487E	730558E 831261E	.952972E	1308106	194640E	236756E 248314E
	 PEAT FLUX	BTU/FT**2HR				0.133566E 04 0.133566E 04						
0.7040 0.6710 0.6710 0.7570 0.7670 0.8630 0.8650 0.8650 0.8650 0.8650 0.9130 0.9130 0.9130	OF CALCULATICN CURPENT(AMP)		214159E 0 246319E 0 276549E 0	13335E 0 13335B 0 53338E C	5115C4E C 5443E15 O 765814E O	0,410175 00 0,509850E 00 0,984513F 00	123030E 0	171726E 0	196553E 0	233531F 0 251661E 0	2637525 C	309469£ 0
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HEAT FLUX

CURRENT (AMP)

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RESULTS OF CALCULATION

WALL SUPERHEATIC) RESISTANCE OF WIRE (OHM) WIRE TEMP.(C)

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2701
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RCN

			CHM
			.2210
			Ħ
	_	L)	SYSTEM RESISTANCE AT SYSTEM TEMPERATURE = .2210 CHM
ပ္က	ORB	Ē	¥
88.		0.800	SYS
18	2	_	A
WATER TEMPERATURE = 88.5C	SYSTEM PRESSURE # 500. TORR	TEST WIRE LENGTH = 0.800 INCH	RESISTANCE
WATER T	SYSTEM	TEST WI	SYSTEM

V. ACRESS TEST WIRE	573	050	134	129	148	163	199	.228	512	328	C-3835	511	.563	.668	146	. £13	.636	£98°	. 887	968.	.916	\$39	595		.C5E	960.	515	130	.560	457	.380	360	.780	.889
V. ACROSS SHUNT RESISTOR	073	CEE	110	123	143	158	152	.223	.268	.3CA	0.3634	414	517	.625	859ª	156	.178	. 8CE	£23	832	. 64B	.876	\$58.	dio.	.582	.019	.813	169	.683	.670	.661	.657	669.	552.
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88	SCO. TCRR	6.72	SYS
p			A
MATER TEMPERATURE = 88.5C	SYSTEM PRESSLRE =	TEST HIRE LENGTH = 0.720 INCH	SYSTEM RESISTANCE AT SYSTEM TEMPERATURE = .2120 CHM

T WIRE																				
TEST	CI	m	635	970	354	940	681	955	654	805	073	645	315	135	1 29	005	340	4	240	0
ACRCS S	o	-	7	-	~	2	4,	4	4	5	-	-	0.8	5.	s.	5.	12	9	9	'n
•																				
RES ISTOR																				
SHUNT	145	048	131	585	370	240	£37	370	7.00	593	626	473	.8135	1 7 6		919	270	575	000	260
ACROSS SHUNT	O	-	-	-		0.29	0.3	0.4	0.4	0.5	9.0	0.7	9.0	0.9	6.0	0.5	0.8	0.7	0.7	9.0
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PCIAT																				
DATA	1	~	1	4	S	•	~	80	0	10		12	13	14	15	16	17	18	19	20

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Š	CURRENT (AMP)	HEAT FLUX	×	WALL SUPERHEAT(C)	RESISTANCE OF WIRE COMM)	WIRE TEMP.(C)
		ETU/FT**2HR	KCAL/M**2HR			
-	00 3399665	0.9473366 03	0.2569748 04	0.174295E 02	0.200551E 00	0.105929E 03
. ^		0 18P387E 04	0.513205E 04	0.186667E 02	0.201320E 00	0.107167E 0
u r		70 577770	0-124086F C5	0. 202070E 02	0.202153E 00	0.108708E 0
n 4		0.6337926 04	0.1853536 05	0.218556E 02	0.203043E 00	0.1103565 0
ru			0.263993E 05	0.230295E 02	0.203676E 00	0.111529E 0
٠.		0-150878F 05	0.40927JE C5	0.258581E 02	0.205202E 00	0.114358E 0
, r			0-634674E C5	0.309317E 02	0.237936E 00	0-119432E 0
۰ «			0.521549E 05	1000	0.209132E 00	0.121656E 0
0		0.405629E 05	0.110031E 06	U.323534F 02	0.208700E 00	0-120853E 0
			0.156786E 06	0.341145E 02	0.209648E 00	0.122514E C
2 =			0.232442E C6	0.347667E 02	0.209998E 00	0.123267E 0
: 2		0.1000125 06	0.271293E 06	0.352321E 02	0.213249E 00	0-123732E 0
2	0-363303E 01	0.118345E C6		0.350912E 02	0.210174E 00	0.123591E 0

0.123466E 03 0.128140E 03 0.127030E 03 0.214421E 04 0.229354E 03 0.747916E 03 0.593409E 03

0.2126186 00 0.2126186 00 0.2120226 00 0.512026 00 0.5132076 00 0.5132076 00

0.349656E 02 0.396399E 02 0.38530E 02 0.105571E 04 C.840854E 03 0.659416E 03

0.287564E 06 0.426611E 06 0.4551816E 06 0.105301E 07 0.775855E 06 0.580261E 06

0.142876E 06 0.157270E 06 0.166562E 06 0.287075E 06 0.213913E 06 0.162639E 06

0.395619E 01 0.412511E 01 0.425221E 01 0.335177E 01 0.309735E 01 0.290265E 01

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TEMPERATURE	
WATER	

SYSTEM PRESSURE = 400, TCRR

SYSTEM RESISTANCE AT SYSTEM TEMPERATURE = .2030 DHM

DATA	PCINT	>	. ACROSS SHUNT RESISTOR	SISTOR	Ă ,	ACROS 5	TEST	F. R.
-			067				m	
7			.053				67 C	
m			22				149	2
4			.178	đ			676	
5						0.2	2195	
9			.258				845	
7			.352				951	
80			448				375	
o			.570		80		672	
2			.617		•		140	
=			10				000	
			118.			•	130	
13			. 8 7 B				930	8
14			.517				290	
15			.527				360	
16			.94B			•	200	
17			3				016	
18			0.7603			•	226	
19			.677				250	

RESULTS OF CALCULATION

0 0 0	CURRENT (AMP)	HEAT FLUX	žie.	WALL SUPERHEATICS	RESISTANCE OF WIRE COHM)	WIRE TEMP.(C)
		BTL/FT**2HR	KCAL/M**2HR			
-	0.2586735 00	0.758832E 03	0.205841E 04	0.142481E 02	0.217544€ 00	0.969481E 02
7	0.412339E CO	C.144243E C4	0.391274E C4	0.131893E 02	0.216936E 00	0.958893E 02
m	0.542035E UO	0.250357E 04	0.679113E 04	0.148715E 02	0.217919E 00	0.975715E 02
4	0.7876118 00	C. 53C562E C4	0.143926E 05	0.1622A9E 02	0.218736E 00	0.9892895 C2
ĸ	0.102655E 01	0.505574E C4	0.2455465 C5	0.1793735 02	0.219764E 00	0.109637E 03
•	0.1319C3E C1	0.150780E 05	0.4090c6E 05	0.2104046 02	0.221630E 00	0-1037406 03
_	0.173628F C1	C.264727E 05	0.7183588 05	0.259303E 02	0.224568E 00	0.108631E 03
æ	0.153+07E 01	0.348569E 05	0.545529E C5	0.2906296 02	0.226447E 00	0.111763E 03
ው	0.252345E C1	0.574471E 05	0.155831E 06	0.361829E 02	0.230712E 00	0.118883E 03
01	0.2730C9E 01	0.6727E4E 05	0.182499E C6	0.364007E 02	0.230842E 00	0.1191015 03
11	· 0.312832E 01	0.879016F 05	0.238442E 06	0.344964E 02	0.229703E 00	0.117196E 03
12	0.359850E 01	0.117072E 06	0.317569E 06	0.391705E 02	0.232498E 00	0.121870E 03
13	0.368496E 01	0.1391646 06	0.377456E 06	0.4470305 02	0,235802E 00	0-127403E 03
*	0.405752E 01	0.151220E 06	0.410200E 06	0.4318925 02	0.2348586 00	0.125889E 03

TEST WIRE LENGTH = 0.800 INCH

0.124610E 03 0.121736E 03 0.105423E 04 0.856736E 03

0.234135E 00 0.232417E 00 0.708901E 00 0.621493E 00

0.419104E 02 0.390356E 02 0.571533E 03 0.774035E 03 0.552930E 03

0.417833E 06 0.433773E 06 0.102153E C7 0.745488E 06 0.490210E 06

C.154034E 06 0.159910E 06 0.376589E C6 0.274824E 06

0.410177E C1 0.419469E C1 0.368584E G1 0.336283E G1 0.299557E G1

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WATER TEMPERATURE = 76.0C SYSTEM PRESSURE = 300. TCRR TEST WIRE LENGTH = 0.780 INCH SYSTEM RESISTANCE AT SYSTEM TEMPERATURE = .2020 C+M	
NATE SYST TEST SYST	

R V. ACROSS TEST WIRE	072	116	185	243	332	402	0.5363	606	692	783	641	505	515	535	05+5	578	030	580		787	142	273
V. ACROSS SHUNT RESISTOR	673	123	199	256	343	412	0.5100	663	.657	3.779	835	999	615°	930	550	.583	840	.721	. 65E	. 225	.875	899
DATA PCINT	-	7	m	4		•	~	80	o	10	=	12	13	71	15	16	17	18	19	20	21	22

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NC.	CUR 2 ENT (AMP)	FEAT FLUX		WALL SUPERHEATICS	RESISTANCE OF WIRE (CHM)	WIRE TEMO.(C)
						3. 4 3.
		BIL/FT**2HR	KCAL/M**2HR			
-	0.118918F 00	0.568399E 03	0.2626885 04	0.176871E 02	0.210189E 00	0.936871E 02
		40 978 AC 25	0.688151F 04	0.187251E 02	0.210799E 00	0.947251E C2
J (*		40 TO 1200000	0.180257F 05	0. 207630E 02	0.211997E 00	0.967630E 02
۱ ۷		0.109598F 05	0.297295E 05	0.218709E 02	0.212646E 00	0.578709E 02
r v		0.200841F C5	0.5448C1E C5	0.282670E 02	0.216398E 00	0-1042675 03
٠.		0.2917505 05	0.7515C9E C5	0.316712E 02	0.218351E 00	0.107671E 03
		0.4536465 05	0.123056E 06	0.3804815 02	0.222121E CO	0.1140486 03
- α		0.648045F CS	0-1758CZE C6	0. 400277E 02	0.223277E 00	0.116028E 03
9 0		0.8489765 05	0-230293E 06	0.383581E 02		0.1143588 03
· 5		0-1071755 06	0.250722E 06	0,428431E 02		0.1188436 03
:=	0.369469E 01	0.123351E 06	0.334711E 06	0. 436400E 02	0.225384E 00	0-11964DE 03

0.119441E 03 0.117276E 03 0.117647E 03 0.11486E 03 0.11486E 04 0.134657E 04 0.74282E 04 0.123042E 04 0.136038E 04

0.225269E 00 0.22400TE 00 0.223524 00 0.22352E 00 0.22261E 00 0.804333E 00 0.618398E 00 0.51229E 00 0.80254E 00

0.434408E 02 0.412762E 02 0.404487E 02 0.404487E 02 0.127057E 02 0.127057E 04 0.115442E 03 0.115442E 04 0.115442E 04

0.3877876 06 0.4020866 06 0.4131086 06 0.4581706 06 0.1234886 07 0.6647196 06 0.4580726 07

0.14295E 06 0.14829E 06 0.15292E 06 0.158400E 06 0.16890E 06 0.16861E 06 0.406832E 06 0.48653E 06

0.397788E 01 0.406155F 01 0.421354E 02 0.434954E 01 0.37564E 01 0.37564E 01 0.36504E 01 0.36504E 01 0.36504E 01

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TEMPERATURE	
MATER	

SYSTEM PRESSURE = 200, TORR TEST WIRE LENGTH = 0.850 INCH

SYSTEM RESISTANCE AT SYSTEM TEMPERATURE = "217C CHM DATA PCINT V. ACROSS SHUNT RESISTOR V. ACROSS TEST WIRE

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NC.	CURRENT (AMP)	HEAT FLUX	×	WALL SUPERHEATICS	RESISTANCE OF WIRE (CHM)	WIRE TEMP.(C)
		BTL/FT**2HR	KCAL/M**2HR	27	e w	
-	0.221239E 03	0.3933455 03	0.1C6617E 04	0.137803E 02	0.223759E 00	0.767803E 0
2	0.363531E CO	0.116278E 04	0.3154175 04	0.137803E 02	0.220759E 00	0.767803E 0;
m		0.377716E 04	0.1024596 05	0.1378035 02	0.220759E 00	0.767303E 02
4	0.128319F 01	0.134JB7E 05	0.343725E CS	0.185690E 02	0.223877E 00	0.815690E 0
S		0.255445E 05	0.692921E C5	0.2250816 02	0.2264385 00	0.855981E 0
•	0.293905E CI	0.742328E CS	0.201364E 06	0.378995E 02	0.2364165 00	0.100899E 0
۲.		0.103970E C6	0.2820305 06	0.561129E 02	0.2481625 00	0.1191136 03
80		0.125939E 06	0.341623E 06	0.463996E 02	0.2419065 00	0.109400E 0
σ		C.139672E C6	0.378874E C6	0.587204E 02	0.249838E 00	0.121720E C
10		0.141338E 06	0.383354E C6	0,563489E 02	0.248314€ 00	0-119349E 0
11		0.317946E 06	0.862459E 06	0.992978E 03	0.762864E 00	0.1055985 04
12		0.228129E 06	0.418322E C6	0.725176E 03	0.63366ZE 00	0.788176E 0:
13		0.376579E 06	0.102151E C7	0.118642E 04	0.847266E 00	0-124942E 0
14		0-433224E 06	0.117516E 07	0.128475E 04	0.887295E 00	0.134775E Q
15		0.488215E 06	0.132433E 07	0.143406E 04	0.944384E 00	0-149706E 0

* €0	* * * * * * * * * * * * * * * * * * *			5•				e.•					\$55		9	WIRE TEMP.(C)				0.805794F C2			D1 I	0.105779E 03
		•	ā				vs s	90		20 E		£		1		RESISTANCE OF WIRE (OHM)				0.22075E 00				0.239571E 00
																WALL SUPERHEAT(C)				0.13/803E 02			0.437266E 02	
			# .2170 CHM	ACRCSS TEST WIRE	0.0450	0.0870	0.2770	0.4000	0.110	0.5240	1.0250	•1100		ij		×	KCAL/M**2HR			0. 102424E 05				0.336744E C6
	ر د د د د د د د د د د د د د د د د د د د	INCH	EMPERATURE	SHUNT PESISTCR V. ACRO	•	00		0	<i>3</i> C		-	-				HEAT FLUX	BTL/FT**2HR	0.00750		0.3///165 04				0.124141E 06
30177 2	WATER TEMPERATURE = 63,0C SYSTEM PPESSURE = 200, TORR	TEST LIRE LENGTH = 0.860 INCH	SYSTEM RESISTANCE AT SYSTEM TO	V. ACROSS	0.0450	0.0870	0.2740	0.3520	3.6700	0.8530	05450	1.0150		RESULTS OF CALCULATION		CURRENT (AMP)				0.645841E 00			339823E	0.377434E 01
RUN NO.	WATER TEMPERATUS SYSTEM PPESSURE	TEST LIR	SYSTEN R	DATA PCINT	-	N M	4	v.	۰ ۰	- 00	٠	91		RESULTS		, U		-	2 '	el A	· 10	•	~	0 0

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WATER TEMPERATURE = 59.0C System pressure = 760. Torr

TEST WIRE LENGTH = 0.800 INCH

SYSTEM RESISTANCE AT SYSTEM TEMPERATURE = .2340 CHM
DATA PCINT V. ACROSS SHUNT RESISTOR V. ACROSS TEST WIRE

• 04	.10	• 20	32	. 42	3	•64	0.7600	• 86	0.9250	1.CCBC	1.0610	1.0900	2.7000	ŏ
.044	.054	.186	257	.354	88	.585	.685	.784	0.8270	.902	.943	558	.820	.857
-	2	m	4	ĸ	9	7	0 0	Φ	01	1	12	13	14	15

RESULTS CF CALCULATION

					E	
0	CURRENT (AMP)	HEAT FLUX	×	WALL SUPERHEAT(C)	RESISTANCE OF WIRE (OHM)	WIRE TEMP.(C)
		BTL/FT**2HR	KCAL/M**2HR		78	
_	0.194690E 00	0.335247E 03	0.909390E 03	0.123297E 02	0.226187E 00	0.111330E 03
2	_		0.417657E C4	0.146580E 02	0.227607E 00	0.113698E 03
m	0.823009E 00	0.603322E 04	0.163657E C5	0.1500196 02	0.227788E 00	0.114002E 03
4	0.131416E 01	0.154160E 05	0.418174E 05	0.1582166 02	0.228279E 00	0.114822E 03
ın		L 0.272957E C5	0.740531E C5	G. 182026E 02	0.229707E 00	0-117203E 03
9	0.215929E 01	L 0.423127E 05	0.114777E 06	0.221722E 02	0.232081E 00	0.121172E 03
_	0.258850E CI	C.6C7909E 05	0.164901E 06	C. 220801E 02	0.232025E 00	0.121080E 03
80	0.303057E 01	1 0.846066E 05	0.229504E 06	0.275351E 02	0.23552E 00	0.126935E 03
0	0.346903E 01	0.110309E (0.299223E 06	0.260786E 02	0.234415E 00	0.125079E 03
10	0.365929E 01	L 0.124387E C6	0.337411E C6	0.213490E 02	0.237559E 00	0.130349E 03
11	0.399115E 01	0.147832E 06	0.401009E 06	0.309764E 02	0.237336E 00	0.129976E 03
12	0.417257E 01	0-162749E	0.441472E 06	0.338620€ 02	0.239057E 00	0.132862E 03
13	0.423894E 01	0.169577E	0.461078E 06	0.3866878 02	0.241917E 00	0-137669E 03
14	0.362832E 01	1 0.375233E 06	0.101786E 07	0.1002936 04	0.728924E 00	0-110193E 04
15	0.379204E 0	1 C.418486E 06	0.113518E 07	0.104018E 04	0.744264E 00	0-113918E 04

Appendix A (continued)

					WIRE TEMP.(C)		0.1035166 0.1035166 0.1055456 0.1105586 0.1136056 0.1280606 0.1291026 0.137376 0.1317376 0.1393696
					RESISTANCE OF WIRE (OHM)		0.251536E 00 0.254719E 00 0.256399E 00 0.255399E 00 0.268834E 00 0.271628E 00 0.272338E 00 0.274184E 00 0.279367E 00
٠					WALL SUPERHEATIC)	n	0.513467E 00 0.551564E 01 0.794492E 01 0.12584E 02 0.156089E 02 0.201299E 02 0.300603E 02 0.317967E 02 0.317967E 02 0.41364E 02 0.409453E 02
		INCH H TEMPERATURE = .2590 CHM	ACRCSS	0.3466 0.5840 0.5840 0.5000 0.5110 1.2080 1.2090 1.3090 1.3000 2.8000 2.8000 3.5000 3.5000 3.5000	HEAT FLUX	BTU/FT**2HR KCAL/M**2HR	0.152421E G5 0.413457E G5 0.269561E 05 0.732295E 05 0.42295E 05 0.15291E 06 0.48437E G5 0.42595E 05 0.163731E G6 0.163731E G6 0.163731E G6 0.163731E G6 0.4134767E 06 0.595536E 06 0.26553E 06 0.595536E 06 0.214703E 06 0.59550E 06 0.514703E 06 0.59550E 06 0.514703E 06 0.59550E 06 0.514703E 06 0.59550E 06 0.59550E 06 0.59550E 06 0.59550E 06 0.59550E 06
1. 91276 1 TEMPERATURE = 98.0C	SSURE = 760, TC	LENGTH = C.920 SISTANCE AT SYSTE	V. ACROSS SFUNT		CT CALCULATION		0.133496E 01 0.176549E 01 0.226756E 01 0.261504E 01 0.256903E 01 0.40170E 01 0.41850E 01 0.41850E 01 0.47564E 01 0.47564E 01
RUN NG.	YSTE	TEST WIRE	-		RESOLTS		10646929333

0.11391835 03 0.1154445 04 0.9470725 03 0.6073325 03 0.1275855 04 0.1323255 04

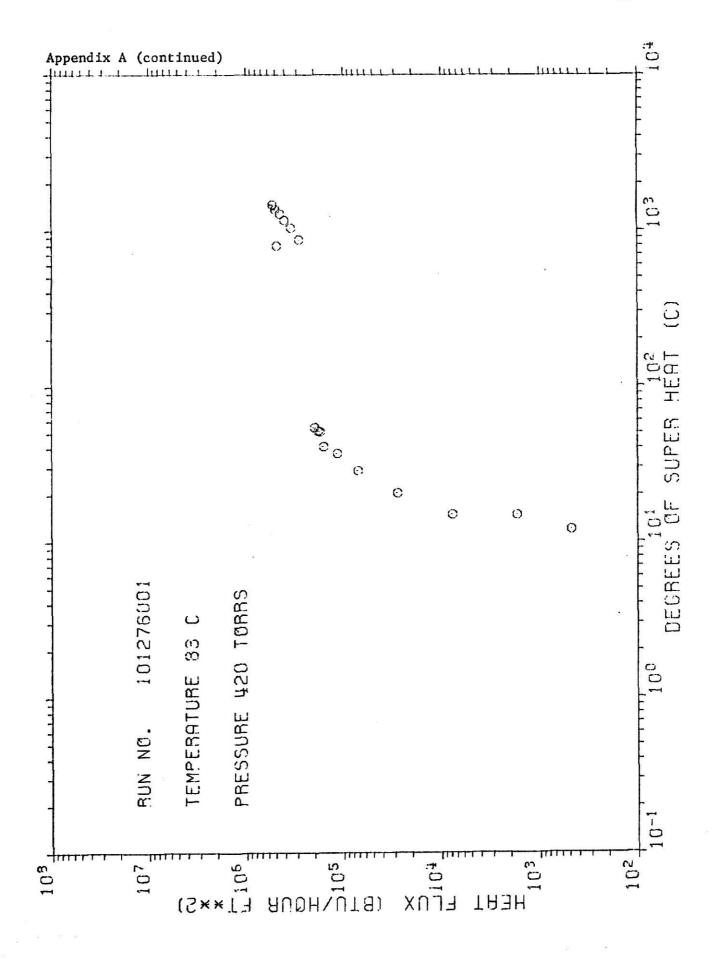
0.2792429 00 0.463047E 00 0.575171E 00 0.57585E 00 0.918091E 00 0.918091E 00

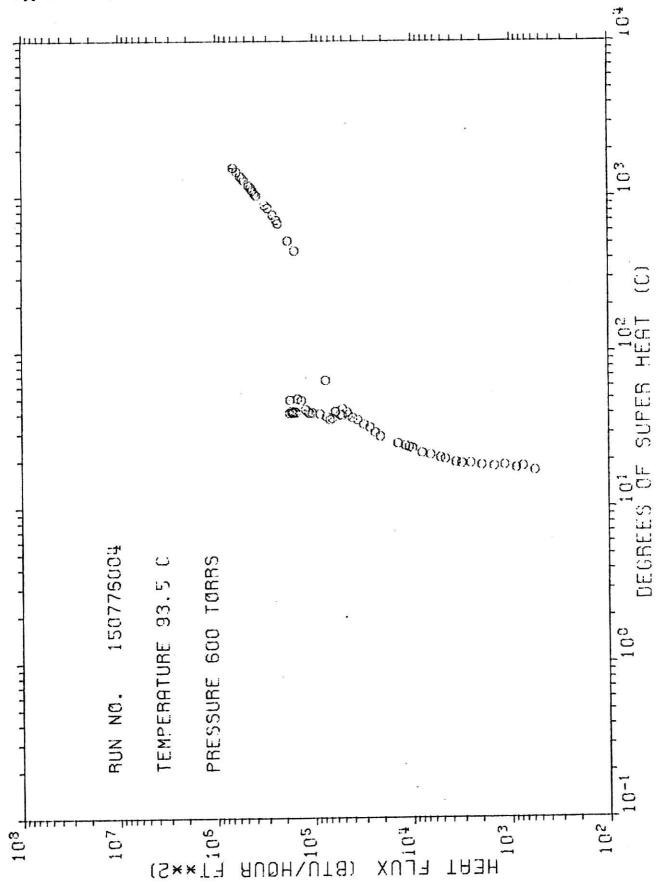
0.411829E 02 0.105644F 04 0.845072E 03 0.509332E 03 0.113786E 04 0.117785E 04

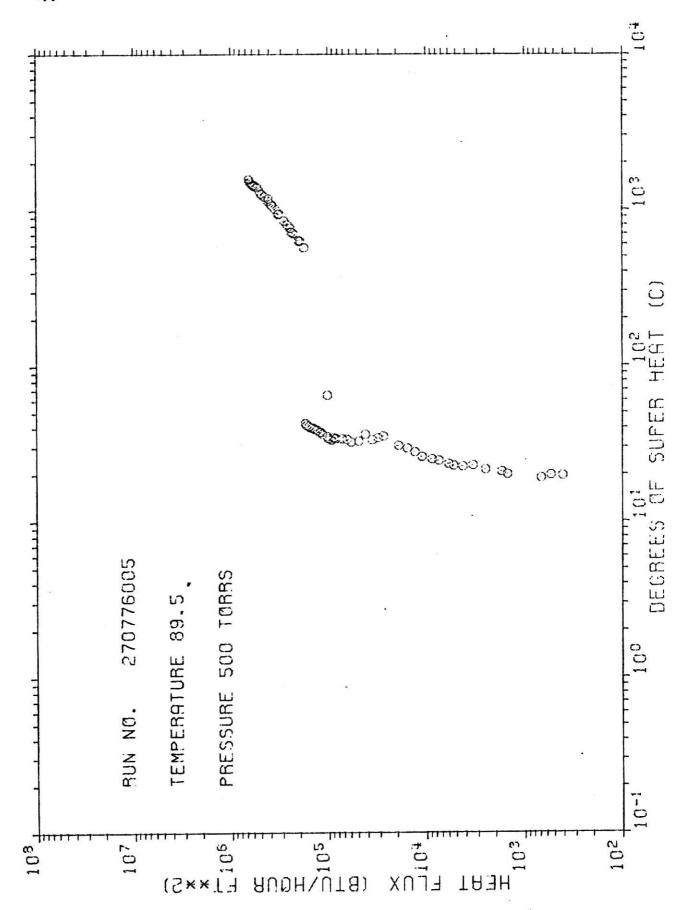
0.6201879 06 0.1234556 C7 0.9294439 06 0.673979 06 0.1385096 C7 0.1529906 C7

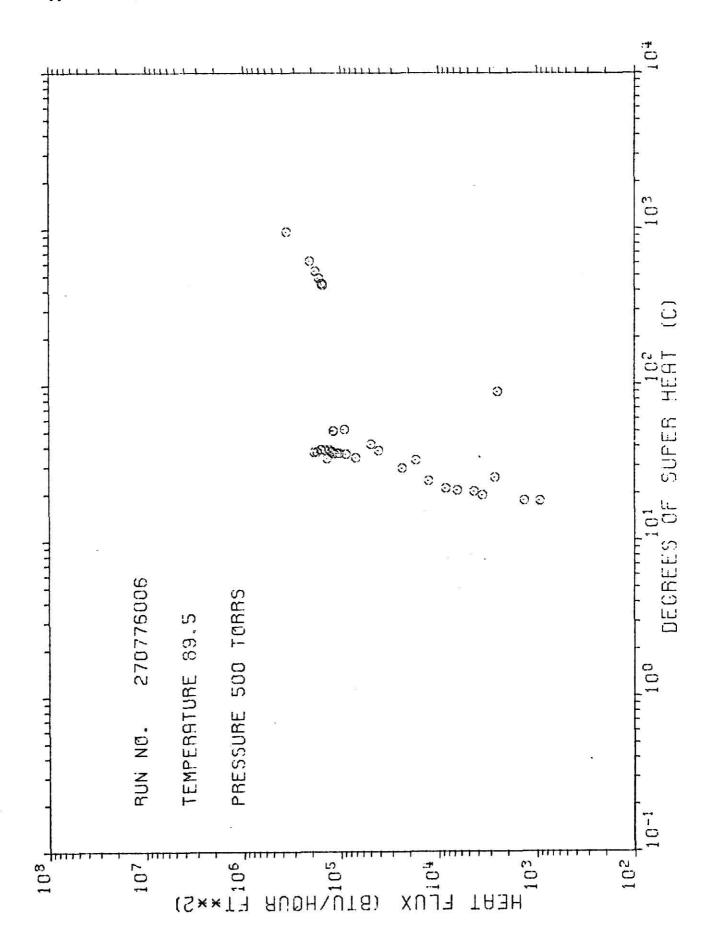
0.226632F 06 0.477696E 06 0.342641F 06 0.746425E 06 0.516425E 06 0.516610F 06 0.556610F 06

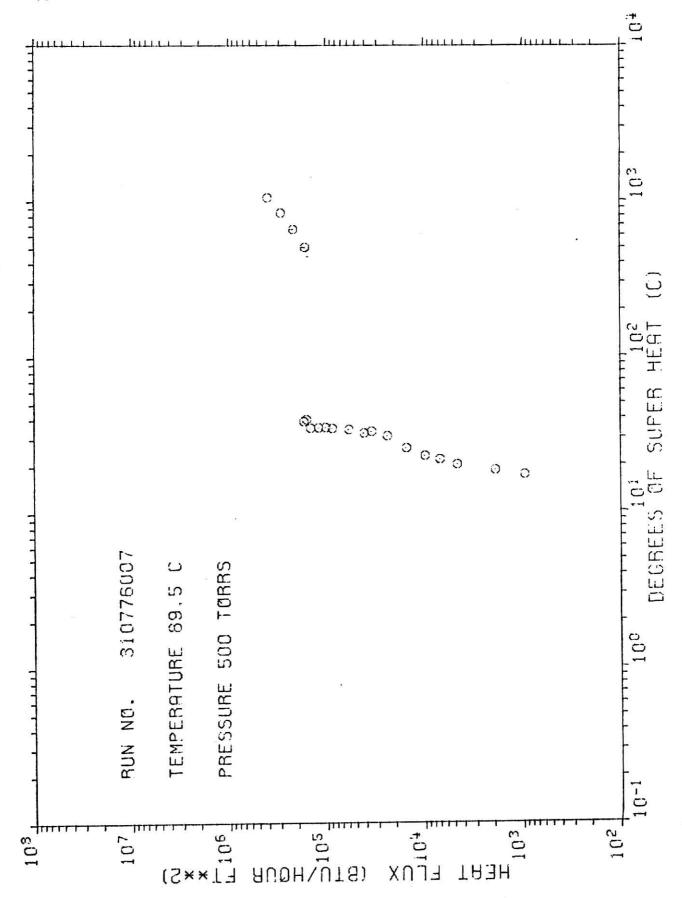
0.\$507638 01 0.\$017708 01 0.\$047175 01 0.\$04075 01 0.\$04075 01 0.\$1\$\$07 01 0.\$4273\$4 01

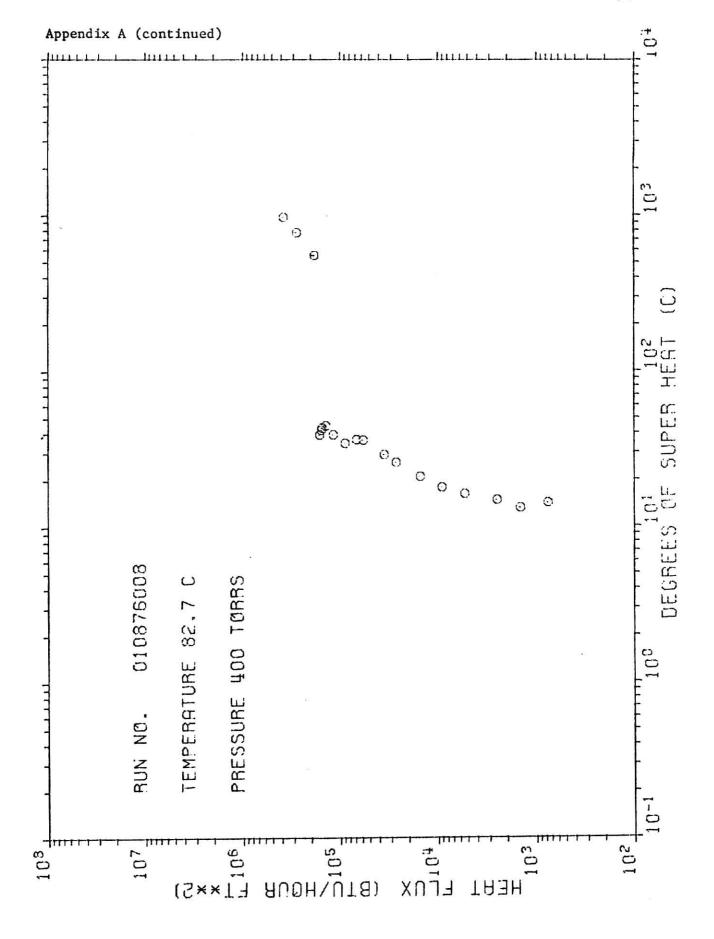


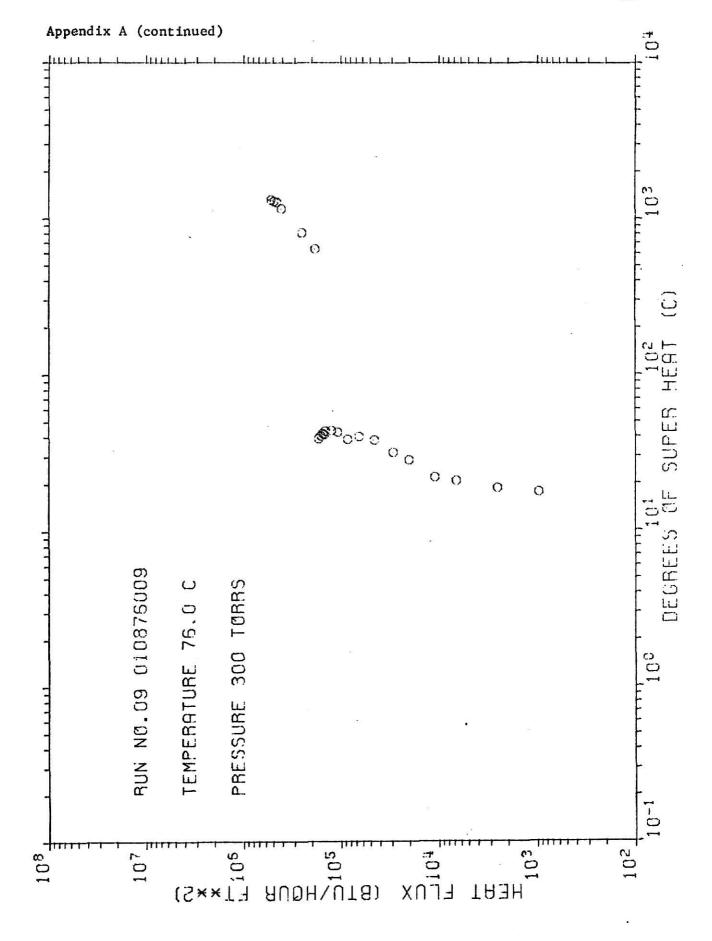


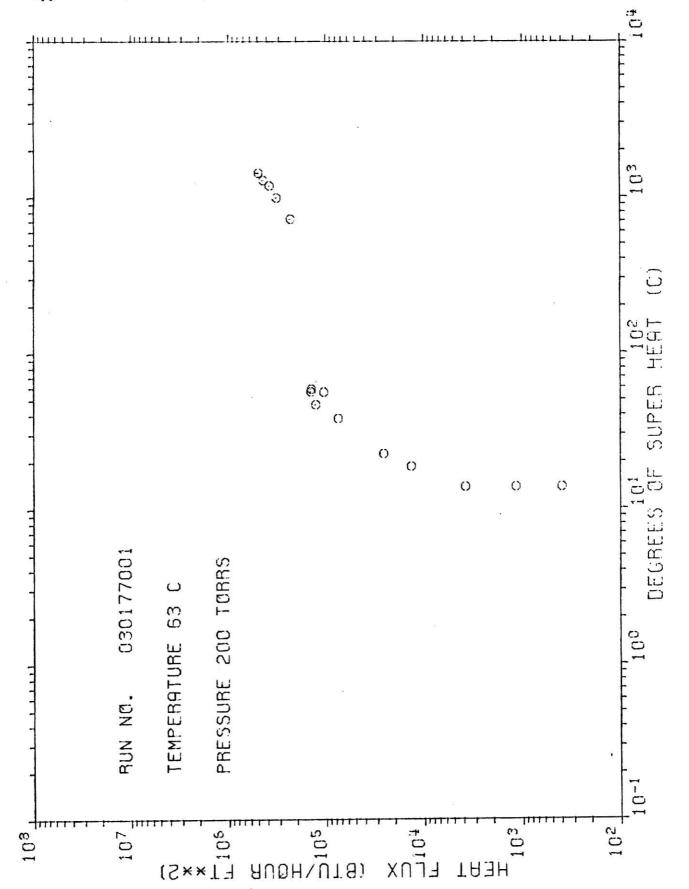


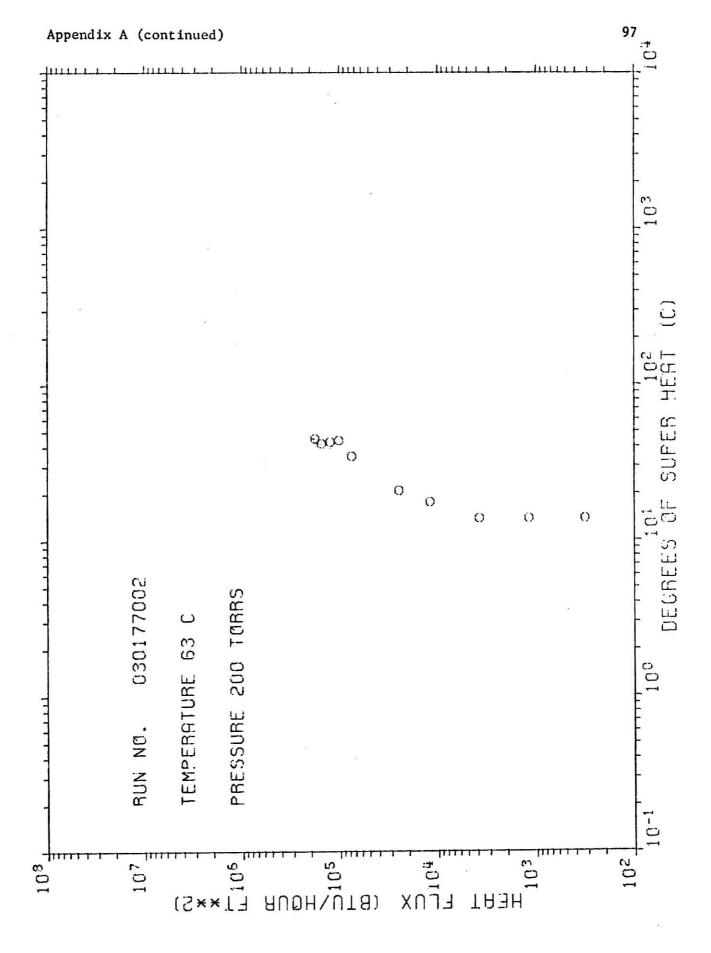


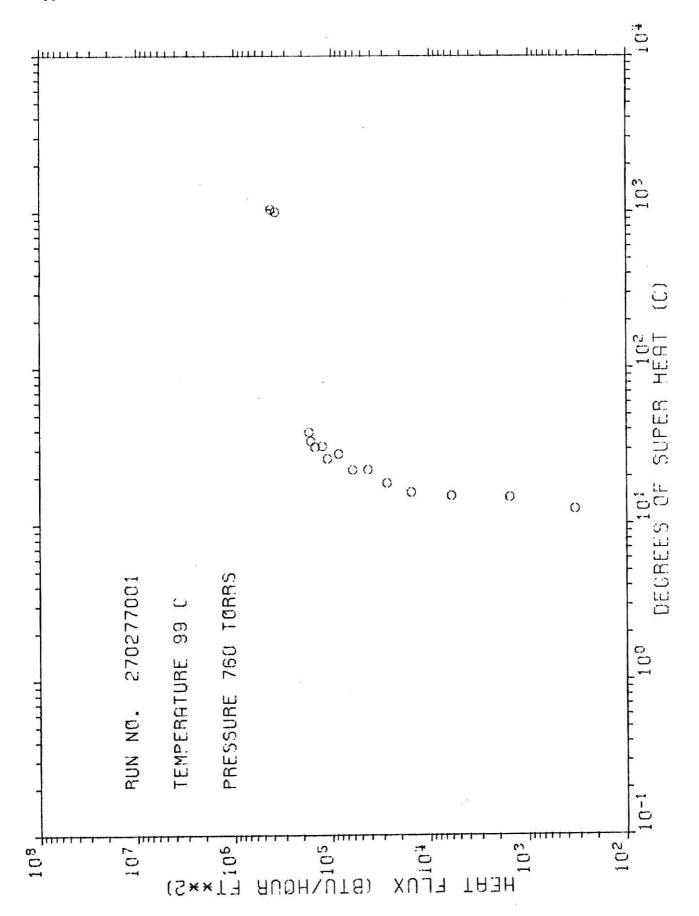












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	-		IS FCR PRO	CESSING THE PRES	MIS PREGRAM IS FCR PROCESSING THE PRESSURE TRANSIENT DATA		8 14	
	ں ن ن ن	WHERE	LENTH TX CT	THE LENGTH OF THE INITIAL WIRE TEMP.	THE LENGTH OF THE TEST WIRE INITIAL WIRE TEMP.			
8	u u u		OV TEMPA RTH	INITIAL WIRE VULIAGE DE TEMPERATURE CF WATER THECRETICAL RESISTANCE	CF WATER RESISTANCE OF TEST WIRE AT	RE AT		
	o o		RSYS	CPERATING TEMPERATURE WIME RESISTANCE INCLU	CPERATING TEMPERATURE MAIR RESISTANCE INCLUDING ELECTRODES AND CANAGE TESTERS	DES AND		
	., ., .,		< < <	VCLTAGE ACROSS THE VOLTAGE ACROSS THE	S THE SHUNT RESISTOR S THE TEST WIRE			
	u u	INTEGER RUND	רעאם					
2000		REAL IC, LENTH	FILTE		•			
4000		WAITE(6,2)	_			*		
0005	21	g	END=101PUN	MKITELOSS) REAC(5,1,END=10)PUNNC,TEMPA,LENTH,RSYS,VI,VW,CV	RSYS, VI, VW, CV			*
9007	e		, F3. 1, F3.	FEFWAT(16,F3.1,F3.2,F4.3,F5.4,F5.4,F4.3)	,F4.3)			
3000		A=0.00352	0					
0010		B=-C.C000CC5E8	3600	12				
0011		IC=VI/RS				¥)		
0313		RTH=0.197	187*LENTH*	RTH=0.19787*LENTH*(1.0+4*TEMPA+8*TEMPA**2)	EHPA**2)	*		
0014		VEPEVM-IC *KSTRAY	*S-RTH					
9100		RC=VWP/1C						
0017		DT=DV/(A*	D1=DV/(A*(Vk-V1*K5)FAT/K5)) HFLUX]=3170*2*IC*VWP/(3*141	DT=DV/(^*(Vk-VI*K>! kAY/*>)) HFLUX]=3170.2*[C*VWP/(3.1416*0.005*LENTH*2.54**2)	#LENTH*2.54**2)			
6013		HFLLX2=HF	HFLLX2=HFLLX1/0,3687	87				
0000		1X=(40*()	4)RUNKG.H	1X=(RO*(1.0+A*1EMFA)/K!H=1.01/A LRITF(6.4)RUNNG.HFLUX2.FFLUX1.TEMPA.TX.DT	PA, TX, DT			
3022		2 FCF MATC//	/3X FUN N	D 13X, INITIAL	FCFMAT(//3x, RUN NO. 13x, INITIAL FEAT FLUX 10x, WATER TEMP.",4	FER TENP. " , 4		
0023		3 FCRMAT(16X,	6X, CKCAL/	HR M**21.,4X,"(3	FORMATILEX, * (COL/HR M**2)*, 4X, * (STU/FT **2 HR)*, 6X, * (C) *, 15X, * (C) *, 15X, * (C)	:) •, 15×, •(C)		
9200			X, 16, 6X, E1	3.6,6X,E13.6,7X,	F[PMAT[3X,16,6X,E13.6,6X,E13.6,7X,F5.2,11X,F7.2,15X,F5.2/]	5.2/)		
0025		GO TO 100	(L					
5522		10 CCNTINUE						
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BOILING HEAT TRANSFER PHENOMENA DURING RAPID DECOMPRESSION

by

Shin-Ping Kung

B.S., National Tsing Hua University, 1972

AN ABSTRACT OF A MASTER'S THESIS

submitted in partial fulfillment of the requirement for the degree

MASTER OF SCIENCE

Department of Nuclear Engineering

KANSAS STATE UNIVERSITY

Manhattan, Kansas

ABSTRACT

Results of an experimental study of the temporal behavior of the transient behavior of local boiling under rapid decompression are presented. Specifically, the heat transfer from a thin wire heating element to water during and shortly after the decompression is discussed. The fast pressure transient is observed to cause an initial heater temperature drop and delay the boiling transition in both totally subcooled depressurization and depressurization with flashing. Flashing, which has been considered responsible for the enhanced heat transfer during and after a pressure transient, is found to be of questionable importance in this experiment, since at this rapid decompression rate, flashing does not become significant until the decompression is complete.