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CONCEPTUAL DESIGN OF A COMMERCIAL-
TOKAMAK-HYBRID-REACTOR FUELING SYSTEM

by

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SUMMARY

Possible fueling concepts were examined to determine which would provide the highest probability of success for application in the CTHR (Commercial Tokamak Hybrid Reactor). Cold fueling was chosen with this idea in mind. Thereafter, the particular scheme chosen was frozen-pellet fueling, since it provides the advantage of maximum fuel-particle density. Several methods of frozen-pellet injection were studied to determine their capabilities, their experimental verification, and projections for the future. Any characteristics which might adversely affect the normal operation of the tokamak were taken into account. To this effect, the light-gas gun was tentatively chosen to provide pellet acceleration.

The ORNL Neutral Gas Shielding Model is the basic theory used to determine the required pellet velocity. It has been modified, however, to account for operation in the commercial temperature regime (as opposed to experimental devices with temperatures around 1 keV). The required pellet velocity is a function of the depth at which the pellet has disappeared as a solid entity.

From this, the pressure level and other essential estimates to be made on the fuel injector design have been made. This leaves only the design of a fuel handling system to implement operation. The fuel handling system has been designed so that a sufficient fuel-pellet supply is produced and quality control systems may be integrated into the system at a later date.

Various ways were studied to locate the fuel injection system so as to minimize the total-system perturbation. The suggested design incorporating the fuel injector into CTHR accounts for both this aspect and that of providing the shortest possible path to the plasma center.

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NOMENCLATURE AND NOTATION

- $a \equiv$ Distance from the outer edge of the plasma to the pressure center (in m). Note that this quantity is taken as 0.9 m for CTHR.
- $c_p \equiv$ Specific heat (constant pressure) of ablatant.
- $c_s \equiv$ Sound speed (in m/s).
- $e \equiv$ Charge of the electron (in C) $\doteq 1.6021892 \times 10^{-19}$.
- $h \equiv$ Ablatant enthalpy variable.
- $(\ell/a) \equiv$ Fraction of pressure center of tokamak to which pellet is to achieve penetration; i.e., if $(\ell/a) = 1$, pellet penetrates to the pressure center of the reactor.
- $m \equiv$ Mass of fuel molecule (in kg).
- $m_e \equiv$ Mass of electron (in kg) $\doteq 9.109537 \times 10^{-31}$.
- $m_{s+p} \equiv$ Sum of the mass of a pellet injector shell and the mass of the fuel pellet (in kg). In the present design, this is approximately 1.947×10^{-3} kg.
- $m_D \equiv$ Mass of deuteron (in kg) $\doteq 3.344549 \times 10^{-27}$.
- $m_H \equiv$ Mass of proton (in kg) $\doteq 1.673560 \times 10^{-27}$.
- $n_\infty \equiv$ Number density of electrons in the tokamak (in m^{-3}). Note: This quantity is position dependent, as in Eqn. 52b.
- $\langle n_\infty \rangle \equiv$ Average electron number density in fusion plasma (in m^{-3}) = 9.8×10^{19} .
- $p \equiv$ Ablatant enthalpy variable.
- $\dot{q} \equiv$ Net heat transfer rate to the ablatant, by electron collisions.
- $r \equiv$ Position variable in fusion plasma, i.e., (r/a) is dimensionless.
- $\hat{r} \equiv$ Ratio of distance from pellet center to pellet radius.
- $r_p \equiv$ Initial pellet radius (in m) = 3.073×10^{-3} .
- $r_o \equiv$ Radius of pellet as a function of time (in m).
- $t \equiv$ Time variable. Used in context of dr_o/dt .
- $u \equiv$ Required pellet velocity to achieve penetration as a solid entity to (ℓ/a) , (in m/sec).

u_{\max} \equiv Maximum attainable velocity of light-gas-gun-projectile (in m/sec).

v \equiv Ablatant velocity variable.

v_o \equiv Speed of the ablatant (in m/sec). For the purposes of the determination of the value of ξ (see Eqn. 20), this quantity is taken to be 400 m/sec.

A_s \equiv Surface area of the pellet injector shell exposed to the propellant gas (in m^2), (numerically, $\pi/4 (0.03)^2 m^2$).

E_o \equiv Thermal energy (in eV) of electrons at pellet surface. Note: This quantity has been taken as 10 eV in Eqn. 51.

E_{∞} \equiv Thermal energy (in eV) of electrons in fusion plasma. Note: This quantity is position dependent as illustrated in Eqn. 52a.

J \equiv Electron current, further defined in Eqn. 16.

L \equiv Distance over which the fuel injector shell and pellet are accelerated (in m), taken here to be 0.30 m.

$L(E)$ \equiv Electron energy loss function, approximated by Milora and Foster¹ as shown in Eqn. 51, (in $eV m^2$), where E is the electron energy (in eV).

$\langle L(E) \rangle$ \equiv Average value of $L(E)$ in the ablatant.

M \equiv Mach number of the ablatant.

R \equiv Ideal gas constant (in J/(kg K)).

T \equiv Ablatant temperature.

$\langle T_{\infty} \rangle$ \equiv Average electron temperature of fusion plasma (in eV) = 1.3×10^4 .

γ \equiv Ratio of specific heats.

μ_s \equiv Ratio of mass of fuel molecule to the mass of a proton.

ρ \equiv Ablatant density variable, function of r .

$\hat{\rho}$ \equiv Ratio of the mass density of the ablatant to the mass density of the ablatant at the pellet surface.

ρ_s \equiv Mass density of fuel pellet (in kg/m^3).

ρ_o \equiv Mass density of ablatant as a function of time (in kg/m^3).

ρ_D \equiv Mass density of solid deuterium (in kg/m^3) $\doteq 2.059 \times 10^2$.

NOTE: All units are assumed to be MKS unless otherwise stated.

1.0 Introduction

Westinghouse Fusion Power Systems Department is developing a design for a Commercial Tokamak Hybrid Reactor, referred to henceforth as CTHR. In the course of the design of the reactor, it became apparent that fueling system design must be given considerable attention. This is due to the fact that the fuel must cross magnetic-field lines in order to effect penetration of the reacting plasma body. In spite of this, the subject of refueling mechanisms is new to the tokamak-reactor conceptual design. The idea of crossing magnetic field lines leads, basically to two concepts of fueling. These are hot fueling, wherein the fuel consists of either ionized particles or neutral particles with a high thermal energy, and cold fueling, wherein the fuel consists of particles with very small thermal energy.

In the hot fueling concept, the particle temperatures are generally above 100 eV. Some examples of this concept are cluster beams, neutral beams, and plasma guns. The primary disadvantage to using hot ionized-particle fueling lies in the requirement that the refueling "mixture" must have at least as high a thermal energy as the reacting body; clearly, in order to allow electrically charged fuel particles to enter and penetrate the confining magnetic field, their average energy must exceed the energy which the magnetic field was designed to contain. The primary disadvantage to employing hot neutral-particle fueling is due to the rather small probability for interaction with the plasma; obviously, a wastefully excessive amount of fuel would have to be injected to maintain the necessary plasma density. These requirements also cause a secondary disadvantage. The inherent inefficiency of heating the fuel by an electromechanical device includes both the inefficiency of the device and the

Carnot/Rankine inefficiencies accompanying the electrical-power generation upon which the device is operated. Thus, it is probable that cold fueling, if feasible, will be a more efficient and a safer process.

There are, basically, three different schemes for implementation of a cold-fueling concept. These are fueling by frozen D+T pellets, squirting a jet of liquid D+T, and puffing gaseous D+T from the surface of the reactor. All three have in common the fact that the fuel particles are unaffected by the magnetic field since they are electrically neutral. The requirement for refueling here is that the fuel must penetrate the pressure "wave" of the reacting body. Obviously, the greater the density of the fuel, the greater is the probability of penetration, provided that the fuel-injection velocity is held constant. (Remember that, upon injection into the reacting body, the fuel will tend to become very quickly ionized by the surrounding plasma.) This, effectively, justifies the concentration of the study on frozen-pellet fueling while, at least temporarily, eliminating further consideration of the liquid-jet and surface-gas schemes.

There have been many methods suggested by which pellet injection may take place and, presently, some of these will be discussed briefly. However, a far more pressing question arises to the surface. How fast, i.e., with what velocity, must a pellet be injected in order to achieve penetration of the reacting body and thereby refuel the tokamak? The model used in this paper to describe the behavior of the pellet is the ORNL Neutral Gas Shielding Model of Foster and Milora.¹ The model has been modified, here, to fit more closely the realities of the CTHR temperature regime of operation. Explanation of the modification will follow in a later section.

The first method for consideration is that of linear-resonance acceleration. In this method,² a pellet of frozen D+T receives a metallic coating

and is, subsequently, electronically charged. It is then accelerated through an electric field to achieve terminal velocity. There are several disadvantages to this system, pertaining to use in a fusion-power facility, as the following illustrates:

1. The metallic coating, received by the pellet initially, enhances bremsstrahlung losses, which, in turn, increases the value of the critical n_T required for fusion breakeven quite dramatically; for example, a 10% increase in bremsstrahlung would be realized from a pellet containing only 0.01% of iron.
2. Very high accelerating potentials are projected for use in such a reactor (~ 100 MV). To assure that there is no electric-field breakdown, such a device is expected to require a length of around 100 m.² The drift tube could be expected to be enormous.

In spite of these disadvantages, however, velocity achievable by this type of system is predicted to be as high as 10^4 m/sec, although no experimental data are available to confirm the theoretical projection.

Another method for consideration is that of freezing a jet of liquid into a rod which is subsequently shattered into pellets.³ The pellets are then allowed to drift towards the target plasma without further acceleration. This method has been tested experimentally to speeds of 100 m/sec and is expected to deliver an ultimate pellet velocity of 2000 m/sec.³

A third possibility for accelerating the pellet lies in striking it with a laser beam on one side.⁴ The resulting ablatant mass causes the pellet to recoil and to be accelerated into the tokamak with an estimated pellet velocity of 10^4 m/sec. Unfortunately, no experiments have been made to test this idea.⁴

A fourth possibility is to accelerate the pellet mechanically through the use of rotating arbor.^{5,6} The primary disadvantage of this method lies in the hazard to the integrity of the vacuum caused by the rapidly moving

arbor. This would necessitate a protective shield for the vacuum vessel. Terminal velocity is anticipated to exceed 1000 m/sec.⁵

The final method to be discussed here is that of using a light-gas gun to accelerate the pellet.^{6,7,8} In this method, a high-pressure propellant gas strikes a disk and accelerates disk and pellet towards the reacting body. The experimental velocity achieved thus far by this method is ~330 m/sec,⁷ which is higher than that for any other method studied. This method has been projected to achieve velocities of 6000 m/sec.⁸

As this is a conceptual-design paper for a possible fueling system of CTHR, it is requisite that a choice be made of methods by which the scheme of pellet-injection may be achieved. On the basis of inherent **advantages and disadvantages** (including lack of experimental evidence) and projected capabilities, a choice has been made for study in this paper. The system selected and recommended for use in the CTHR design is that of the light-gas gun for pellet injection.

2.0 Fuel Scheme Modeling

Several models were considered for the purpose of modeling the pellet ablation process as it pertains to CTHR.^{1, 11-19, 21} These models were studied to determine the most appropriate one for determination of the pellet-injection velocity requirement. Since the ORNL Neutral Gas Shield Model¹ has a strong theoretical base of physical principles, e.g., conservation of mass, momentum, and energy laws; and since it has a record of adequacy in the lower temperature regimes (1 keV) for which it was formulated, this model was chosen to form the basis of the theory by which the ablation process of a pellet in CTHR is described. Due to the nature of the model, considerable reconstruction was required before it could be applied to the problem at hand. The ORNL Neutral Gas Shield Model was derived to form order of magnitude estimates based on discrete points at which the pellet would disappear as a solid entity. The requirement of this study is to examine the continuum of points at which the pellet will disappear for a particular reactor, namely, the Commercial Tokamak Hybrid Reactor of Westinghouse. By this, a conceptual design of a fuel injection system for CTHR may be formulated.

2.1 Laws of Conservation of Mass, Momentum, and Energy

The ORNL Neutral Gas Shield Model is based upon the laws of conservation, as mentioned previously. These may be expressed as:

$$\frac{d}{dr} \left(\rho v r^2 \right) = 0; \quad (1a)$$

$$\frac{d}{dr} \left(\rho v^2 r^2 \right) + r^2 \frac{dp}{dr} = 0; \quad (1b)$$

$$\left(\frac{1}{r} \right)^2 \frac{d}{dr} \left(\rho v r^2 \left[h + \frac{v^2}{2} \right] \right) = \dot{q}. \quad (1c)$$

These represent conservation of mass, conservation of momentum, and conservation of energy, respectively. Here, spherical geometry has been assumed and thus, the mass flow rate, w may be represented by $4\pi\rho vr^2$. First, let us manipulate equation (1a).

2.1.1 Equation of Conservation of Mass

The solution to the mass conservation equation is almost trivial, but will prove useful in the development of the model. Therefore, we give it full consideration here. The solution may be expressed in one of two ways, due to the solid-gas interface. Thus, the solution becomes, at the boundary (see Figure 1):

$$\rho vr^2 = \rho_o v_o r_o^2.$$

But, due to the solid-gaseous interface, this may also be written as:

$$\rho vr^2 = \rho_s r_o^2 \left| \frac{dr_o}{dt} \right|.$$

Hence, one has the following relationship:

$$\rho f^2 v = \left(\frac{\rho_s}{\rho_o} \right) \left| \frac{dr_o}{dt} \right| = v_o. \quad (2)$$

Now, we are ready to transform the energy equation into something more easily solvable (although numerical techniques are still required).

2.1.2 Conservation of Energy

The energy equation was previously described by:

$$\left(\frac{1}{r} \right)^2 \frac{d}{dr} \left[\rho vr^2 \left[h + \frac{v^2}{2} \right] \right] = \dot{q}.$$

Making use of the chain rule, we find:

$$\rho v \frac{d}{dr} \left[h + \frac{v^2}{2} \right] + \frac{\left[h + \frac{v^2}{2} \right]}{r^2} \frac{d}{dr} [\rho v r^2] = \dot{q}.$$

However, the last term on the left-hand side of the equation is zero since mass is conserved (see equation (1a)) and thus, the energy equation is given by the following:

$$\rho v \frac{d}{dr} \left[h + \frac{v^2}{2} \right] = \dot{q}.$$

If we divide both sides of the above equation by the quantity, $\frac{\rho v}{dr}$, the result is that:

$$d \left[h + \frac{v^2}{2} \right] = \left(\frac{\dot{q}}{\rho v} \right) dr.$$

Or expressing the above equality in a slightly different manner:

$$d \left[h + \frac{v^2}{2} \right] = \left(\frac{\dot{q} r_o}{\rho v} \right) d\hat{r}, \quad (3)$$

Shapiro²⁰ gives as a general solution to the energy equation, the following, adopting Shapiro's notation temporarily (except for the ratio of specific heats):

$$\begin{aligned} \frac{dM^2}{M^2} = & - \frac{2 \left[1 + \frac{\gamma-1}{2} M^2 \right]}{(1-M^2)} \left(\frac{dA}{A} \right) + \frac{1+\gamma M^2}{(1-M^2)} \frac{\{ dQ - dW_x + dH \}}{c_p T} \\ & + \frac{\gamma M^2 \left[1 + \frac{\gamma-1}{2} M^2 \right]}{(1-M^2)} \left\{ 4f \frac{dx}{D} + \frac{2dx}{\gamma p A M^2} - 2y \frac{dw}{w} \right\} \\ & + \frac{2(1+\gamma M^2) \left(1 + \frac{\gamma-1}{2} M^2 \right)}{(1-M^2)} \left(\frac{dw}{w} \right) - \frac{1+\gamma M^2}{(1-M^2)} \left(\frac{dW}{W} \right) - \left(\frac{d\gamma}{\gamma} \right). \quad (4) \end{aligned}$$

The sixth bracketed term is due to a possible change in the ratio of specific heats, γ . The fifth term in brackets accounts for a change in the molecular weight, W . Both of these are zero since we neglect both disassociation of the molecule and ionization of the gas in the shield. Since we are dealing with a diatomic ideal gas, in the shield, $\gamma = 7/5$. The fourth term in brackets is zero by virtue of mass conservation. The third bracketed term accounts for frictional losses. These are negligible in the shield itself. Hence, only the first two terms are left to consider. The second term in brackets may be interpreted as²⁰:

$$dQ - dW_x + dH = c_p dT + d(v^2/2). \quad (5)$$

But, the differential specific enthalpy of an ideal gas is simply:

$$dh = c_p dT.$$

Thus equation (5) may be rewritten as:

$$dQ - dW_x + dH = d\left(h + \frac{v^2}{2}\right). \quad (6)$$

By substituting the above relationship into equation (4), we obtain:

$$\frac{dM^2}{M^2} = - \frac{[2 + (\gamma - 1)M^2]}{(1 - M^2)} \left(\frac{dA}{A}\right) + \frac{1 + \gamma M^2}{1 - M^2} \left[\frac{d\left(h + \frac{v^2}{2}\right)}{c_p T}\right]. \quad (7)$$

Now let us evaluate the term, dA/A . The surface area, A , of a sphere is given by:

$$A = 4\pi r^2.$$

Upon differentiating the above equation, it is discovered:

$$dA = 8\pi r dr.$$

By combining the above two results, one finds that:

$$\frac{dA}{A} = \frac{8\pi r dr}{4\pi r^2} = 2 \frac{dr}{r} = 2 \frac{df}{f}. \quad (8)$$

If equations (3) and (8) are substituted into equation (7), the energy equation becomes:

$$\frac{dM^2}{M^2} = - \frac{[2 + (\gamma - 1)M^2]}{(1 - M^2)} \left(\frac{2 df}{f} \right) + \frac{1 + \gamma M^2}{1 - M^2} \left(\frac{\dot{q} r_o}{\rho v c_p T} \right). \quad (9)$$

At this point, consider an ideal gas for which the following holds²⁰:

$$c_s^2 = \left(\frac{\partial p}{\partial \rho} \right)_s = \gamma R T.$$

Then, we have the following relationships:

$$M^2 \equiv v^2 / c_s^2 = v^2 / (\gamma R T); \quad (10a)$$

$$c_p \equiv \frac{\gamma}{\gamma - 1} R. \quad (10b)$$

Dividing (10a) by $v^2/(\gamma-1)$, we find that:

$$(\gamma - 1) \frac{M^2}{v^2} = \frac{1}{\frac{\gamma}{\gamma - 1} R T}. \quad (11)$$

But, the term in the denominator of the right-hand side of equation (11) is given by equation (10b) as simply, $c_p T$. Thus, it is seen that:

$$\frac{1}{c_p T} = (\gamma - 1) \frac{M^2}{v^2}. \quad (12)$$

Substitution of this relationship into equation (9) yields:

$$\frac{dM^2}{M^2} = - \frac{[2 + (\gamma - 1)M^2]}{(1 - M^2)} \left(\frac{2 df}{f} \right) + \frac{1 + \gamma M^2}{1 - M^2} \left\{ \frac{\dot{q}}{\rho} r_o \frac{1}{v^3} (\gamma - 1) M^2 df \right\}. \quad (13)$$

Upon rearranging equation (2), one obtains:

$$\frac{1}{v} = \beta f^2 \left\{ \frac{\rho_o}{\rho_s} \right\} \left| \frac{dr_o}{dt} \right|^{-1}.$$

and thus, equation (13) becomes:

$$\begin{aligned} \frac{dM^2}{M^2} = & - \frac{[2 + (\gamma - 1)M^2]}{(1 - M^2)} \left(\frac{2 df}{f} \right) + \frac{1 + \gamma M^2}{1 - M^2} \left[\left(\frac{\dot{q}}{\rho} \right) r_o (\beta f^2)^3 \left| \frac{dr_o}{dt} \right|^{-3} \right. \\ & \left. \times (\gamma - 1) M^2 df \right]. \quad (14) \end{aligned}$$

Using the chain rule on the right-hand side, and multiplying through by the quantity $M(1-M^2)/(2 df)$, one discovers, upon rearranging terms:

$$\begin{aligned} (1 - M^2) \frac{dM}{df} = & \frac{(\gamma - 1)}{2} \left\{ \frac{\rho_o}{\rho_s} \right\}^3 \left(\frac{\dot{q}}{\rho} \right) r_o \left| \frac{dr_o}{dt} \right|^{-3} (\beta f^2)^3 M^3 (1 + \gamma M^2) \\ & - \frac{2M}{f} \left[1 + \frac{\gamma - 1}{2} M^2 \right]. \quad (15) \end{aligned}$$

At this time, it is necessary to define further the heat deposition rate, \dot{q} . The ORNL Neutral Gas Shield Model assumes that the primary mechanism by which heat is deposited in the shield is electron collision. This is given by¹:

$$\dot{q} = eJ \frac{dE}{dr} = eJ \frac{\rho}{m} L(E), \text{ where } J \equiv n_{\infty} \sqrt{\frac{e E_{\infty}}{3\pi m_e}}. \quad (16)*$$

* $L(E)$ may be approximated by¹:

$$L(E) \doteq (2.35 \times 10^{18} + 4 \times 10^{15} E + 2 \times 10^{21}/E^2)^{-1}.$$

Thus equation (15) becomes:

$$(1 - M^2) \frac{dM}{d\hat{r}} = \frac{\gamma - 1}{2} \left(\frac{\rho_o}{\rho_s} \right)^3 \left(\frac{e L(E)_J}{m} \right) r_o \left| \frac{dr_o}{dt} \right|^{-3} (\beta \hat{r}^2)^3 M^3 (1 + \gamma M^2) - \frac{2M}{\hat{r}} \left(1 + \frac{\gamma - 1}{2} M^2 \right). \quad (17)$$

At this point it becomes useful to define a dimensionless parameter, ξ :

$$\xi \equiv \frac{\gamma - 1}{2} \left(\frac{\rho_o}{\rho_s} \right)^3 \left(\frac{e L(E)_J}{m} \right) r_o \left| \frac{dr_o}{dt} \right|^{-3}. \quad (18)$$

Hence, equation (17) becomes, upon substitution of the definition of ξ :

$$(1 - M^2) \frac{dM}{d\hat{r}} = \xi (\beta \hat{r}^2)^3 M^3 (1 + \gamma M^2) - \frac{2M}{\hat{r}} \left(1 + \frac{\gamma - 1}{2} M^2 \right). \quad (19a)$$

Or, by performing an algebraic manipulation, one finds:

$$\frac{dM}{d\hat{r}} = \frac{M}{1 - M^2} \left(\xi (\beta \hat{r}^2)^3 M^2 (1 + \gamma M^2) - \frac{[2 + (\gamma - 1) M^2]}{\hat{r}} \right). \quad (19b)$$

As it turns out, equation (1b) will yield a differential equation in terms of β and M as a function of \hat{r} ; so, momentarily at least, let us turn our attention to the definition of ξ since we shall have two equations of β and M as a function of \hat{r} (and thus, though the solution must be numerical, we can solve for both variables).

Recalling equation (2), we have the relationship:

$$v_o = \left(\frac{\rho_s}{\rho_o} \right) \left| \frac{dr_o}{dt} \right|.$$

If this result is further applied to equation (18), one discovers:

$$\xi = \frac{\gamma - 1}{2} \frac{e L(E)}{m} J \frac{r_o}{v_o^3} . \quad (20)$$

The practical application of the ORNL Neutral Gas Shield Model demands that ξ be given a numerical value. The resultant set of coupled differential equations is then solved. Thus r_o is taken to be the initial pellet radius, r_p ; v_o is taken to be the sound speed of a low temperature hydrogen gas (about 400 m/sec); and J is taken to be the average electron current of the plasma.

2.1.3 Equation of Conservation of Momentum

The momentum equation is recalled to be:

$$\frac{d}{dr} \left(\rho v^2 r^2 \right) + r^2 \frac{dp}{dr} = 0 .$$

Using the ideal gas law, one finds:

$$\frac{d}{dr} \left(\rho v^2 r^2 \right) + r^2 \frac{d(\rho RT)}{dr} = 0 .$$

Dividing the above equation by $\rho_o r_o$ yields:

$$\frac{d}{d\hat{r}} \left(\hat{\rho} \hat{v}^2 \hat{r}^2 \right) + \hat{r}^2 \frac{d(\hat{\rho} RT)}{d\hat{r}} = 0 . \quad (21)$$

Now, recall the definition of the Mach Number, equation (10a):

$$M^2 \equiv \frac{v^2}{\gamma RT} .$$

By algebraically rearranging the above equation, one obtains:

$$RT = \frac{v^2}{\gamma M^2} .$$

Upon substitution of the above relationship into equation (21), one finds:

$$\frac{d}{d\hat{t}} \left(\beta \hat{t}^2 v^2 \right) + \hat{t}^2 \frac{d}{d\hat{t}} \left(\frac{\beta v^2}{\gamma M^2} \right) = 0. \quad (22)$$

Equation (2) may be rearranged to provide the following relationship:

$$v = \frac{v_o}{\beta \hat{t}^2}. \quad (23)$$

The substitution of equation (23) into equation (22) produces the following:

$$\frac{d}{d\hat{t}} \left(\frac{v_o^2}{\beta \hat{t}^2} \right) + \hat{t}^2 \frac{d}{d\hat{t}} \left(\frac{v_o^2}{\gamma M^2 \beta \hat{t}^4} \right) = 0.$$

If the above equation is divided through by v_o^2 , one has:

$$\frac{d}{d\hat{t}} \left(\frac{1}{\beta \hat{t}^2} \right) + \frac{\hat{t}^2}{\gamma} \frac{d}{d\hat{t}} \left(\frac{1}{M^2 \beta \hat{t}^4} \right) = 0. \quad (24)$$

The chain rule may be applied to equation (24) to yield:

$$\frac{d}{d\hat{t}} \left(\frac{1}{\beta \hat{t}^2} \right) + \frac{\hat{t}^2}{\gamma} \left[\frac{1}{M^2 \hat{t}^2} \frac{d}{d\hat{t}} \left(\frac{1}{\beta \hat{t}^2} \right) + \frac{1}{\beta \hat{t}^2} \frac{d}{d\hat{t}} \left(\frac{1}{M^2 \hat{t}^2} \right) \right] = 0. \quad (24b)$$

Upon rearranging, one finds:

$$\left(1 + \frac{1}{\gamma M^2} \right) \frac{d}{d\hat{t}} \left(\frac{1}{\beta \hat{t}^2} \right) + \frac{1}{\gamma \beta} \frac{d}{d\hat{t}} \left(\frac{1}{M^2 \hat{t}^2} \right) = 0.$$

Again using the chain rule yields:

$$\frac{d}{d\hat{r}} \left(\frac{1}{\beta \hat{r}^2} \right) \left[\frac{\gamma M^2 + 1}{\gamma M^2} \right] + \frac{1}{\gamma \beta M^2} \frac{d}{d\hat{r}} \left(\frac{1}{\hat{r}^2} \right) + \frac{1}{\gamma \beta \hat{r}^2} \frac{d}{d\hat{r}} \left(\frac{1}{M^2} \right) = 0.$$

This may be reexpressed as follows by differentiating and then multiplying the equation by $-\gamma M^2 \beta \hat{r}^2$:

$$\frac{1}{(\beta \hat{r}^2)} \frac{d}{d\hat{r}} (\beta \hat{r}^2) [1 + \gamma M^2] + \frac{2}{\hat{r}} + \frac{2}{M} \frac{dM}{d\hat{r}} = 0.$$

The following results from using the chain rule of differentiation and combining terms:

$$\frac{d\beta}{d\hat{r}} \left[\frac{1 + \gamma M^2}{\beta} \right] + \frac{2(2 + \gamma M^2)}{\hat{r}} + \frac{2}{M} \frac{dM}{d\hat{r}} = 0.$$

If equation (19b) is substituted into the above relationship, we have:

$$\frac{d\beta}{d\hat{r}} \left[\frac{1 + \gamma M^2}{\beta} \right] = - \frac{2(2 + \gamma M^2)}{\hat{r}} - \frac{2}{M} \left[\frac{M}{1 - M^2} \left\{ \xi (\beta \hat{r}^2)^3 M^2 (\gamma M^2 + 1) - \frac{2 + (\gamma - 1) M^2}{\hat{r}} \right\} \right].$$

Upon rearranging terms further, the following is found:

$$\left(\frac{1 + \gamma M^2}{\beta} \right) \frac{d\beta}{d\hat{r}} = - \frac{2}{1 - M^2} \left[\xi (\beta \hat{r}^2)^3 M^2 (1 + \gamma M^2) - \frac{2}{\hat{r}} + \frac{2}{\hat{r}} (1 - M^2) - \frac{(\gamma - 1) M^2}{\hat{r}} + \frac{\gamma M^2}{\hat{r}} (1 - M^2) \right].$$

This may be reexpressed in the following form by combining terms:

$$\left(\frac{1 + \gamma M^2}{\beta} \right) \frac{d\beta}{d\hat{r}} = - \frac{2M^2}{1 - M^2} \left[\xi (\beta \hat{r}^2)^3 (1 + \gamma M^2) - \frac{2}{\hat{r}} + \frac{1}{\hat{r}} - \frac{\gamma M^2}{\hat{r}} \right].$$

Combining terms again in the above equation yields:

$$\left(\frac{1 + \gamma M^2}{\beta} \right) \frac{d\beta}{d\hat{r}} = - \frac{2M^2}{1 - M^2} \left[\xi (\beta \hat{r}^2)^3 (1 + \gamma M^2) - \frac{(\gamma M^2 + 1)}{\hat{r}} \right].$$

Finally, we have:

$$\frac{d\beta}{d\hat{r}} = - \frac{2M^2\beta}{1 - M^2} \left[\xi (\beta\hat{r}^2)^3 - \frac{1}{\hat{r}} \right]. \quad (25)$$

2.2 Boundary Conditions of the Model

First, let us recall the coupled differential equations which must be solved in order to evaluate the pellet velocity requirement. These are equations (19b) and (25).

$$\frac{dM}{d\hat{r}} = \frac{M}{1 - M^2} \left[\xi (\beta\hat{r}^2)^3 M^2 (\gamma M^2 + 1) - \frac{[2 + (\gamma - 1)M^2]}{\hat{r}} \right]; \quad (19b)$$

$$\frac{d\beta}{d\hat{r}} = - \frac{2M^2\beta}{1 - M^2} \left[\xi (\beta\hat{r}^2)^3 - \frac{1}{\hat{r}} \right]. \quad (25)$$

The difficulty in solving the above set of differential equations stems from the fact that the known boundary conditions are insufficient. Specifically, these boundary conditions are:

- (i) $\beta(\hat{r} = 1) = 1,$
- (ii) $\lim_{\hat{r} \rightarrow \infty} \frac{dM}{d\hat{r}} = 0.$

It can be shown that through the use of equations (19b) and (24b) that the second boundary condition is equivalent to the following:

$$\lim_{\hat{r} \rightarrow \infty} M = \sqrt{(5/\gamma)}.$$

Equation (24b) yields a power law solution for β at large values of \hat{r} but, due to the nature of the equations, integrating from large values of \hat{r} to the

pellet surface, numerically, is unsuccessful in attaining the first boundary condition. Thus, there are only two possibilities left. The first is to guess the boundary value. While this possibility works in principle, it can become very expensive in terms of computer time usage and thus, the second possibility was explored and found to be fruitful. This possibility was to attempt to provide a solution for $M(\hat{r} = 1) = M_0$ and will now be explained here, in detail.

First, however, let us examine the defining equations in more detail to determine further boundary conditions. We note that for all finite \hat{r} :

$$\frac{d\beta}{d\hat{r}} < 0.$$

Any other condition might bring about recoalescence, which is physically unrealistic for a scheme to effect fuel injection into a fusion plasma. It may also be noted that, based on physical grounds, for all finite \hat{r} , the following is true:

$$\frac{dM}{d\hat{r}} > 0.$$

Examination of equation (19b) will reveal that for cases where the Mach Number is subsonic, the quantity in brackets must be positive and that when the Mach Number exceeds unity, this same quantity must be negative. In order to avoid a singularity at $M = M_1 = 1$, the bracketed quantity must be zero. Thus, we have the following:

$$\left\{ \xi (\beta \hat{r}^2)^3 (1 + \gamma) - \frac{(1 + \gamma)}{\hat{r}} \right\}_{M=1} = 0.$$

The above equality reduces to the following relationship upon dividing through the equation written above by the quantity, $(1 + \gamma)$:

$$\left\{ \xi (\beta \hat{r}^2)^3 - \frac{1}{\hat{r}} \right\}_{M=1} = 0. \quad (26)$$

This provided a motivation to introduce the following definition:

$$\beta \equiv \eta \xi^{-1/3} \hat{r}^{-7/3}. \quad (27)$$

Upon substitution of equation (27) into equation (26), one finds:

$$\eta_1 = \eta|_{M=1} = 1.$$

Next, we substitute the definition of η into equation (19b) to obtain:

$$\frac{dM}{d\hat{r}} = - \frac{M}{1 - M^2} \left\{ \frac{\eta^3 M^2 (1 + \gamma M^2)}{\hat{r}} - \frac{2 + (\gamma - 1)M^2}{\hat{r}} \right\}.$$

Upon combining terms in the above equation, one discovers:

$$\frac{dM}{d\hat{r}} = - \frac{M}{\hat{r}} \left\{ \frac{2 + \gamma M^2 - M^2 - \eta^3 M^2 - \gamma \eta^3 M^4}{1 - M^2} \right\}.$$

By rearranging the terms in the above equality, the energy conservation equation becomes:

$$\frac{dM}{d\hat{r}} = - \frac{M}{\hat{r}} \left\{ \frac{2 + M^2[(1 + \eta^3) - \gamma(1 - M^2\eta^3)]}{1 - M^2} \right\} \equiv - \frac{U}{\hat{r}(1 - M^2)}. \quad (28)$$

Again, the definition of η is substituted into one of the conservation equations. Thus, equation (25) becomes, upon carrying out the differentiation of the definition of β :

$$\left\{ \frac{d\eta}{d\hat{r}} - \frac{7\eta}{3\hat{r}} \right\} \hat{r}^{-7/3} \xi^{-1/3} = - \frac{2M^2\eta \xi^{-1/3} \hat{r}^{-7/3}}{1 - M^2} \left\{ \frac{\eta^3 - 1}{\hat{r}} \right\}.$$

The quantity $\xi^{-1/3} \hat{r}^{-7/3}$ may be eliminated from the above equation immediately to yield:

$$\frac{d\eta}{d\hat{r}} - \frac{7\eta}{3\hat{r}} = - \frac{2M^2\eta}{1 - M^2} \left(\frac{\eta^3 - 1}{\hat{r}} \right).$$

If the above equation is solved for $d\eta/d\hat{r}$, directly (placing terms on a common denominator):

$$\frac{d\eta}{d\hat{r}} = \left(\frac{-6M^2(\eta^3 - 1) + 7(1 - M^2)}{1 - M^2} \right) \frac{\eta}{\hat{r}}.$$

Rearrangement of terms in the above equation yields the following simplification:

$$\frac{d\eta}{d\hat{r}} = \left(\frac{7 - M^2(1 + 6\eta^3)}{3(1 - M^2)} \right) \frac{\eta}{\hat{r}} \equiv \frac{V}{\hat{r}(1 - M^2)}. \quad (29)$$

Consider, now, the following transformation:

$$s = \ln \hat{r}.$$

Therefore, by taking the derivative with respect to \hat{r} of both sides of the above equation, one finds:

$$\frac{ds}{d\hat{r}} = \frac{1}{\hat{r}}.$$

Upon dividing both sides of equations (28) and (29) by the above, one obtains:

$$\frac{dM}{ds} = - \frac{M\{2 - M^2[(1 + \eta^3) - \gamma(1 - M^2\eta^3)]\}}{1 - M^2} \equiv - \frac{U}{1 - M^2}; \quad (30a)$$

$$\frac{d\eta}{ds} = \frac{\eta[7 - M^2(1 + 6\eta^3)]}{3(1 - M^2)} \equiv \frac{V}{1 - M^2}. \quad (30b)$$

Thus an autonomous system of differential equations has been developed. Upon the elimination of the parametric variable, s , from equations (30a) and (30b), we discover:

$$\frac{dM}{d\eta} = \frac{-3M\{2 - M^2[(1 + \eta^3) - \gamma(1 - M^2\eta^3)]\}}{\eta[7 - M^2(1 + 6\eta^3)]} = -\frac{U}{V} \quad (31)$$

From equation (27), we can see that when $\hat{r} = 1$, i.e., when $M = M_0$, then $\eta(\hat{r} = 1) = \sqrt[3]{\xi}$. This means that if equation (31) can be integrated from any given point on the physically realistic solution, M_0 may be determined. There are two points where the boundary values are known. These are at $M = M_1 = 1$ with $\eta = \eta_1 = 1$; and at $M = \sqrt{(5/\gamma)}$ with $\eta = \sqrt[3]{(7\gamma - 5)/30}$ (see equation (28)). However, upon inspection, it will be discovered that both of these points result in an indeterminate derivative and thus equation (31) cannot yet be integrated, since it is obvious that equation (31) very probably has no simple analytical solution and for this reason will require a numerical integration.

From numerical considerations, either point would allow for solution of equation (31), but the point $\eta(M = M_1) = \eta_1$ is chosen as the most appropriate since this will preclude numerical discrepancies due to integration through a critical point. Therefore, in a sufficiently small region of the critical point, linearization will be valid and we find, for $M \doteq 1 + \mu$ and $\eta \doteq 1 + \epsilon$, by substitution into equation (31):

$$\frac{d\mu}{d\epsilon} = \frac{-3(1 + \mu)\{2 - (1 + \mu)^2[(1 + (1 + \epsilon)^3) - \gamma(1 - (1 + \mu)^2(1 + \epsilon)^3)]\}}{(1 + \epsilon)[7 - (1 + \mu)^2(1 + 6(1 + \epsilon)^3)]} \quad .$$

Upon linearizing M^y and η^x (where x and y are arbitrary integer powers) it is found that:

$$\frac{d\mu}{d\epsilon} \approx \frac{-3(1 + \mu)\{2 - (1 + 2\mu)[(2 + 3\epsilon) - \gamma(1 - (1 + 2\mu)(1 + 3\epsilon))]\}}{(1 + \epsilon)[7 - (1 + 2\mu)(7 + 18\epsilon)]} \quad .$$

Linearizing the above equation further, one obtains:

$$\frac{d\mu}{d\varepsilon} \approx \frac{-3(1+\mu)\{2 - (1+2\mu)[(2+3\varepsilon) + \gamma(2\mu+3\varepsilon)]\}}{(1+\varepsilon)[7 - (7+14\mu+18\varepsilon)]}.$$

Upon rearranging terms, one discovers:

$$\frac{d\mu}{d\varepsilon} \approx \frac{-3(1+\mu)\{2 - (1+2\mu)[2 + 3(\gamma+1)\varepsilon + 2\gamma\mu]\}}{(1+\varepsilon)[-14\mu - 18\varepsilon]}.$$

If we linearize equation (31) once again, the following relationship will be provided:

$$\frac{d\mu}{d\varepsilon} \approx \frac{-3(1+\mu)\{2 - [2 + 3(\gamma+1)\varepsilon + 2\gamma\mu + 4\mu]\}}{-14\mu - 18\varepsilon}.$$

By rearranging terms in the above equation, we see that:

$$\frac{d\mu}{d\varepsilon} \approx \frac{-3(1+\mu)\{-3(\gamma+1)\varepsilon - 2(\gamma+2)\mu\}}{-14\mu - 18\varepsilon}.$$

A final round of linearization yields:

$$\frac{d\mu}{d\varepsilon} \approx -\frac{9(\gamma+1)\varepsilon + 6(\gamma+2)\mu}{18\varepsilon + 14\mu}. \quad (32)$$

Thus, if η and M are very close to the critical point $\eta(M=1)$, the above equation will prove to be mathematically equivalent to equation (31). If we now consider that for an independent variable τ :

$$\frac{d\mu}{d\tau} \approx [9(\gamma+1)\varepsilon + 6(\gamma+2)\mu]; \text{ and} \quad (33a)$$

$$\frac{d\varepsilon}{d\tau} \approx -(18\varepsilon + 14\mu). \quad (33b)$$

We again have an autonomous system of differential equations which are not only the mathematical equivalent of equation (31) but are also linear. They may be

rewritten as a matrix equation, as shown below:

$$\begin{bmatrix} \frac{d\mu}{d\tau} \\ \frac{d\varepsilon}{d\tau} \end{bmatrix} \approx \begin{bmatrix} 6(\gamma + 2) & 9(\gamma + 1) \\ -14 & -18 \end{bmatrix} \times \begin{bmatrix} \mu \\ \varepsilon \end{bmatrix}. \quad (34)$$

This is a system of equations of the form:

$$\vec{X}' = \vec{A} \times \vec{X}. \quad (35)$$

If the following is hypothesized:

$\vec{A} \times \vec{X} = \vec{\lambda} \times \vec{X}$, where $\vec{\lambda}$ is a diagonal matrix, equation (35) may be solved in terms of the elements of $\vec{\lambda}$. Thus, we have:

$$\begin{bmatrix} 6(\gamma + 2) - \lambda & 9(\gamma + 1) \\ -14 & -18 - \lambda \end{bmatrix} = 0. \quad (36)$$

The characteristic equation, obtained by taking the determinant of both sides of equation (36), becomes, after algebraic simplification:

$$\lambda^2 - 6(\gamma - 1)\lambda + 18(\gamma - 5) = 0; \lambda = \lambda_1, \lambda_2. \quad (37)$$

It is important to note here that since $1 < \gamma \leq 5/3$, by virtue of the fact that $\gamma = (\nu + 2)/\nu$ where the integer $\nu \geq 3$ is the number of degrees of freedom, and since $\lambda_1 \lambda_2 = 18(\gamma - 5)$, the product $\lambda_1 \lambda_2$ must always be negative and therefore, λ_1 and λ_2 are of opposite sign. The net effect of this is that the critical point $M = 1$, $n = 1$ (or $\mu = 0$, $\varepsilon = 0$) is a saddle point.

Now, if the expression $\vec{A} \times \vec{X} = \vec{\lambda} \times \vec{X}$ is written in matrix form we have (for $\lambda_1 = \lambda_1, \lambda_2$):

$$\begin{bmatrix} 6(\gamma + 2) & 9(\gamma + 1) \\ -14 & -18 \end{bmatrix} \times \begin{bmatrix} \mu \\ \varepsilon \end{bmatrix} = \begin{bmatrix} \lambda_i & \\ & \lambda_i \end{bmatrix} \times \begin{bmatrix} \mu \\ \varepsilon \end{bmatrix}. \quad (38)$$

Rewriting the above expression as a set of coupled algebraic equations, one obtains:

$$6(\gamma + 2)\mu + 9(\gamma + 1)\varepsilon = \lambda_i \mu; \quad (39a)$$

$$-14\mu - 18\varepsilon = \lambda_i \varepsilon. \quad (39b)$$

If equation (39b) is solved for the ratio (μ/ε) , one finds:

$$\frac{\mu}{\varepsilon} = -\frac{\lambda_i + 18}{14}; \quad i = 1, 2. \quad (40)$$

Equation (40) is true for both of the roots of equation (37). At the point $\mu = 0, \varepsilon = 0$ (or $M = 1, \eta = 1$), the above equation becomes, through the application of L'Hospital's Rule:

$$\frac{d\mu}{d\varepsilon} = -\frac{\lambda_i + 18}{14}; \quad i = 1, 2. \quad (41)$$

But, $d\mu/d\varepsilon = dM/d\eta$ and ergo:

$$M'_{si} \equiv \left. \frac{dM}{d\eta} \right|_{M=1, \eta=1} = -\frac{\lambda_i + 18}{14}; \quad i = 1, 2. \quad (42)$$

Thus, upon simultaneous solution of equations (42) and (37), we discover:

$$M'_{s1} = -3[(\gamma + 5) - \sqrt{(\gamma - 2)^2 + 7}]/14; \quad (43a)$$

$$M'_{s2} = -3[(\gamma + 5) + \sqrt{(\gamma - 2)^2 + 7}]/14. \quad (43b)$$

At this point, we have both the location of the saddle point and two

possibilities for the slope at the saddle. One of these slopes must be physically correct and the other must be physically impossible.

Consider now Figure 2, serving as an illustrative example for the specific case of our interest where $\gamma = 7/5$. However, the general conclusions remain valid for any physically permissible value of the ratio of specific heats, i.e., $1 < \gamma \leq 5/3$. The isocline of $U = 0$ (see equation (28)) is intersected by the line $M = 1$. Since values of M that are greater than $\sqrt{5/\gamma}$ are not physically possible, another line is drawn at $M = \sqrt{5/\gamma}$. This forms five regions. The uppermost region is Region V. The four remaining areas are defined in the following manner. Starting at the upper left-hand corner of the as yet undefined areas and numbering clockwise, Region I is where $\sqrt{5/\gamma} > M > 1$ and $U > 0$; Region II is where $\sqrt{5/\gamma} > M > 1$ and $U < 0$; Region III is where $M < 1$ and $U < 0$; and Region IV is where $M < 1$ and $U > 0$.

It may be seen from equation (28) (and in light of the fact that $dM/d\hat{r}$ must be positive) that for $M < 1$, U must be negative and that for $M > 1$, U must be positive. This clearly excludes Regions II and IV from the physically realistic solution set. Region V is excluded as well since M must be less than $\sqrt{5/\gamma}$. Upon further examination of Figure 2, it will be discovered that the solution to equation (31) based on M'_{s2} is not valid on either side of the sonic point since it occupies Region II above the sonic point and Region IV below the sonic point. However, the other solution (based on M'_{s1}) never enters either Region II or IV. This results in the conclusion that the solution to equation (31) based on M'_{s1} must be the only physically possible solution. This solution is presented graphically in Figure 3. The computer program that generated these solutions (Figures 2 and 3) is in Appendix I.

Let us now consider equation (27). We know that at the point $\hat{r} = 1$, $M = M_0$ and $\beta = 1$. If equation (27) is evaluated at this point, we have the

seemingly trivial solution of $\eta(\xi = 1) = \xi^{1.3}$. However, if equation (31) is integrated to this value of η , the corresponding Mach Number will be M_0 . Hence, M_0 may be readily determined by performing the aforementioned integration. Finally, having determined M_0 , the equations of conservation may be solved and the pellet velocity requirement may, in principle, be determined.

The equations of conservation may be solved in one of two ways. The first is to integrate directly equations (19b) and (25), utilizing the value of M_0 that has just been determined. The second method is to apply a simple quadrature to the solution set of $M(\eta)$. The method chosen for application here is the former. The equations of conservation were solved for two separate cases. The reference case ($\xi = 1000$) is shown in Figure 4. The second case solved was for CTHR and is shown in Figure 5. The computer program used to generate these results and those of the next section are contained in Appendix II.

2.3 Pellet Velocity Requirement

As stated previously, the pellet velocity requirement can now be determined, at least in principle, since the conservation laws may be readily solved. Actually, only the point at which the pellet disappears as a solid entity may be determined. The actual pellet velocity requirement will depend upon how far into the tokamak the pellet must go as a solid entity for the ablated pellet to travel the rest of the way to the plasma center. This would require a transport study and is well beyond the capabilities of the KSU computer. Thus, as a matter of practical compromise, the problem is parameterized so that when the penetration requirement on the solid entity is determined, the mechanical requirements of the design will be known.

Recall the definition of \dot{q} from equation (16). This may be solved for $\rho(r)$ as follows:

$$\frac{dE}{dr} = \frac{\rho L(E)}{m} . \quad (16)$$

Upon rearranging the above relationship, one finds:

$$\rho dr = \frac{m dE}{L(E)} . \quad (44)$$

Integrating from the pellet surface, one has:

$$\int_{r_0}^{\infty} \rho dr = \int_{E_0}^{E_{\infty}} \frac{m dE}{L(E)} ; \quad (45)$$

where E_0 is the electron energy at the pellet surface and E_{∞} is the electron energy of the reacting body. We may approximate E_0 as 10 eV or less and thus, for all practical purposes in very hot plasmas, E_0 is roughly zero. Upon dividing both sides of equation (45) by $\rho_0 r_0$, we may define a new parameter which may be readily evaluated from the solution of the conservation equations:

$$f(\xi) \equiv \int_1^{\infty} \beta d\xi = \frac{m}{\rho_0 r_0} \int_{E_0}^{E_{\infty}} \frac{dE}{L(E)} . \quad (46)$$

Solving the above equation for ρ_0 , one discovers:

$$\rho_0 = \frac{m}{r_0} \int_{E_0}^{E_{\infty}} \frac{dE}{L(E)} f(\xi)^{-1} . \quad (47)$$

If we solve the working form of equation (18) for dr_0/dt^{-1} , we find:

$$\frac{dr_0}{dt} = - \left[\frac{\rho_0}{\rho_s} \right] \xi^{-1/3} \left[\frac{\gamma - 1}{2} \frac{e \langle L(E) \rangle}{m} J r_0 \right]^{1/3} .$$

Upon substitution of equation (47) into the above equality, one may observe:

$$\frac{dr_o}{dt} = - \left\{ \frac{m}{\rho_s} \right\} [\xi^{1/3} f(\xi)]^{-1} \frac{1}{r_o} \left\{ \frac{\gamma - 1}{2} \frac{e \langle L(E) \rangle}{m} J r_o \right\}^{1/3} \int_{E_o}^{E_\infty} \frac{dE}{L(E)} \cdot \quad (48)$$

The above relationship may be rearranged to produce the following equality:

$$r_o^{2/3} \frac{dr_o}{dr} \frac{dr}{dt} = - \left\{ \frac{m}{\rho_s} \right\} [\xi^{1/3} f(\xi)]^{-1} \left\{ \frac{\gamma - 1}{2} \frac{e}{m} \right\}^{1/3} \times \left\{ \frac{J}{E_\infty - E_o} \int_{E_o}^{E_\infty} L(E) dE \right\}^{1/3} \int_{E_o}^{E_\infty} \frac{dE}{L(E)} \cdot \quad (49)$$

Here, a new relationship is introduced:

$$\frac{m}{\rho_s} = 2 \frac{m_D}{\rho_D} \cdot$$

Upon substitution of the above relationship into equation (49) and rearranging (and noting that dr/dt is just the pellet velocity, u), we discover:

$$u r_o^{2/3} \frac{dr_o}{dr} = -2 \left\{ \frac{m_D}{\rho_D} \right\} [\xi^{1/3} f(\xi)]^{-1} \left\{ \frac{\gamma - 1}{2} \frac{e}{m_H} \right\}^{1/3} \frac{1}{\mu_s^{1/3}} \times \left\{ \frac{J}{E_\infty - E_o} \int_{E_o}^{E_\infty} L(E) dE \right\}^{1/3} \int_{E_o}^{E_\infty} \frac{dE}{L(E)} \cdot \quad (50)$$

If equation (50) is integrated, one obtains:

$$\frac{3}{5} u r_p^{5/3} = 2 \left\{ \frac{m_D}{\rho_D} \right\} [\xi^{1/3} f(\xi)]^{-1} \left\{ \frac{\gamma - 1}{2} \frac{e}{m_H} \right\}^{1/3} \frac{a}{\mu_s^{1/3}} \times \int_{(2/a)}^1 \left\{ \frac{J}{E_\infty - E_o} \int_{E_o}^{E_\infty} L(E) dE \right\}^{1/3} \int_{E_o}^{E_\infty} \frac{dE}{L(E)} d(r/a) \cdot$$

Solving the above relationship for u , the pellet velocity requirement, we discover:

$$u = \frac{10}{3} \frac{m_D}{\rho_D} [f(\xi) \xi^{1/3}]^{-1} \left\{ \frac{\gamma - 1}{2} \frac{e}{m_H} \right\}^{1/3} \frac{a}{\mu_s^{1/3}} \\ \times \int_{(l/a)}^1 \left(\frac{J}{E_\infty - E_0} \int_{E_0}^{E_\infty} L(E) dE \right)^{1/3} \int_{E_0}^{E_\infty} \frac{dE}{L(E)} d(r/a); \quad (51)$$

where: J is as given by equation (16);

$$L(E) \doteq (2.35 \times 10^{18} + 4 \times 10^{15} E + 2 \times 10^{15} / E^2)^{-1} \text{ eV m}^2;$$

E is in eV;

$f(\xi)$ is given by equation (46); and

ξ is given by equation (20).

Thus, the only things that are now required are the temperature and density distributions of the plasma. These are taken to be¹:

$$E_\infty = \frac{45}{16} \langle T_\infty \rangle [1 - (r/a)^2]^2; \quad (52a)$$

$$n_\infty = \frac{3}{2} \langle n_\infty \rangle [1 - (r/a)^2]. \quad (52b)$$

2.4 Results of the Theoretical Treatment

As stated previously, the equations of conservation were solved for two cases, a reference case ($\xi = 1000$) and the CTHR case. As a means of comparison, for the reference case, the value of $[f(\xi) \xi^{1/3}]^{-1}$ obtained by this study is 1.156 whereas Foster and Milora¹ obtained a value of 1.23. The value of this parameter for the CTHR case was found to be 1.0978. The penetration dependent pellet velocity requirement is plotted in Figure 6.

Earlier, we decided to neglect disassociation and ionization processes. At a distance of 10 pellet radii from the pellet center, we may observe that the temperature increase over the pellet surface temperature is a factor of 26.7 for the reference case and 63.5 for the CTHR case. Assuming that the temperature at the pellet surface is 20 K, we find that for the reference case, the temperature is ~535 K, and for CTHR, the temperature is ~1270 K. This is sufficient justification for neglecting the dissociation and ionization processes.

3.0 Fuel Injector Design

3.1 Gas Gun Design

The light gas gun is the type of design chosen for the fuel-injection system of CTHR, whose pertinent plasma dimensions are given by a minor radius of 1.40 m, an aspect ratio of 4.33, and an elongation of 1.60. The conceptual design is shown in Figs. 7 and 8. It is a system which is easily used in rapid-fire mode, say, one to three pellets per second per gas gun system. Three such systems will be utilized in CTHR. The pellet radius is given by the following:

$$r_p = [(3/(4\pi)) (m_D/\rho_D) (n_\infty V_{pl}) / (\dot{p} \tau)]^{1/3} ,$$

where: $V_{pl} \equiv$ volume of the reacting body,

$\dot{p} \equiv$ rate at which pellets are injected into CTHR, and

$\tau \equiv$ confinement time for CTHR, equal to 1.64 sec.

By substituting appropriate values into the above equation, we find that for a spherical pellet, $r_p \doteq 3.073$ mm.

The final parameter required for this design is the pressure used by the light gas gun. Since the required pellet velocity is a function of the penetration depth and since the required gas pressure depends on the velocity needed, the gas pressure may be expressed as a function of penetration depth.

The gas pressure as a function of injection velocity is given by the following expression:⁹

$$P = \frac{m_{s+p}}{A_s} \frac{u u_{\max}}{2L} \left(\frac{\gamma-1}{\gamma+1} \right) \left[\left(1 - \frac{u}{u_{\max}} \right)^{\frac{\gamma+1}{\gamma-1}} - 1 \right]$$

(assuming that the propellant is an ideal gas, expansion of the propellant is to zero pressure and that the injection system is frictionless). Numerically,

with assumed parameters where P is the gas pressure in (Pa) and u is the required pellet velocity (in m/sec), one has for hydrogen gas:

$$P = 5033.7 u \left[\left(1 - \frac{u}{6578.4} \right)^{-6} - 1 \right],$$

and for D+T gas:

$$P = 3187.7 u \left[\left(1 - \frac{u}{4166.0} \right)^{-6} - 1 \right].$$

The required gas pressure is plotted as a function of the penetration fraction in Figures 9 and 10 for natural hydrogen and a deuterium-tritium mixture, respectively.

This brings up one final question. What type of performance can we expect from a light-gas-gun such as that which is conceptualized here? If the performance of the gun were limited to that already achieved by others,⁷ i.e., a pellet-injection speed of 330 m/sec, the pellet would disappear as a solid entity after penetrating about 25% of the distance to the pressure center (see Fig. 4). The maximum performance to be expected of a light-gas gun employed in the refueling of CTR is determined by assuming that the maximum applicable propellant pressure is 100 GPa. Thus, the performance limit for the D+T propellant is found from Figs. 10 and 6 to result in a penetration of about 45% of the pressure-center distance, achieved with an injection velocity of approximately 3600 m/sec. Based on Figs. 9 and 6, the limit for the H_2 propellant turns out to be somewhat higher, say a maximum penetration of about 50%, corresponding to an injection velocity of approximately 5500 m/sec. The extremely high pressure assumed above is justified when one notes that this value may be decreased by a factor of 5 to 10 were this design modified such that each shell could be used only for a single shot.

3.2 Fuel Handling System

The fuel handling system prepares the D+T pellets and subsequently, makes them available to the fuel injection system. To make fuel handling simpler, the pellets have been given a cylindrical shape in place of the spherical one assumed earlier in this report. The net result of this is a cylinder diameter $\sqrt[3]{16/3}$ times the sphere's radius and a height equal to the diameter. The effects of this geometry shift on ablation are unknown but, obviously, expected to be small. [Note that the ratio of surface area to volume is increased only by about 15 percent.]

The fuel handling system works on the following principles (see Figure 11). Initially, the D+T mixture is cooled by a liquid nitrogen stage. The next stage, which liquifies the fuel, is a liquid helium stage. These discrete steps are suggested since the pellets will be mass produced and liquid helium is more expensive than liquid nitrogen. The very cold D+T mixture (5 to 10 K) is injected into a mold whose temperature is maintained at 4.2 K by the liquid helium until the pellets are frozen. The newly formed pellets are subsequently ejected from the mold by the liquid D+T for the next cycle of pellets. Several of these pellet makers could be in operation simultaneously and quality control may be instituted if and where necessary.

Now that the fuel pellets have been formed, they may be loaded into the fuel injector mechanically. The remaining problem is to implement the system into CTHR.

3.3 Implementation into CTHR

For implementation into CTHR, several criteria must be taken into account:

1. The pellet must take the shortest path possible to the fusion plasma.
2. The amount of radiation shielding used to isolate the system from the environment should be minimized within good safety practices.
3. The fuel injection system must be located at a position in the system to make maintenance feasible.
4. Surrounding systems (and subsystems) must be affected as little as possible by the introduction of the fuel injection system.

These goals are attained as shown in Figures 12 and 13. Figure 12 illustrates the overall system and how each of 3 injector systems may be put into CTHR. In Figure 13, a closeup view is provided, to illustrate maintenance possibilities and to illustrate better radiation-shield requirement and system implementation. As can be seen from Fig. 12, interaction of the fuel injection systems with CTHR is very limited, indeed.

4.0 Conclusions

After rejection of several hot-fueling schemes because of inherent difficulties anticipated in practical implementation, the three possibilities of tokamak refueling with cold fuel were considered: liquid-fuel-jet injection, gas puffing, and frozen-pellet injection. It was concluded that pellet fueling appeared to be the most likely of these three alternatives to succeed as a practical scheme. Several methods by which pellet fueling could be implemented into CTHR were then investigated. This led to the further conclusion that the light-gas-gun approach entailed probably the fewest number of real drawbacks. The ORNL Neutral-Gas-Shielding-Model was used to model pellet ablation in CTHR and a required pellet injection velocity was determined as a function of required penetration-depth in the CTHR plasma. On this basis, a conceptual design for a light-gas-gun fuel-injection system for CTHR was developed and the gas-pressure requirement for the device was determined. Finally, suggested means for implementation into the overall design of CTHR have been discussed.

It must be recalled, however, that the penetration-depth requirement for CTHR has not been firmly established. This penetration-depth requirement will ultimately determine the potential feasibility of the design presented in this report. To this end, such a study is highly recommended.

In conclusion, the findings of this study are as follows:

1. If it is determined that the required penetration depth must be 50% of the pressure center distance or more, the light-gas gun will most probably not succeed in refueling CTHR.
2. If the required penetration is between 45% and 50% of the pressure center distance, the D+T propellant mixture will most likely not

succeed in achieving the required injection velocity, but the light-gas-gun design might be practically feasible with H_2 as the propellant.

3. If the required penetration depth is between 25% and 45% of the pressure center depth, the proposed light-gas-gun injection system should be able to refuel the CTHR, with either propellant.
4. If the required penetration depth is less than 25% of the pressure center distance, the proposed scheme of light-gas gun injection will be able to achieve the necessary penetration velocity based on currently tested and proven state-of-the-art technology.

5.0 Selected References

1. S. L. Milora and C. A. Foster, "ORNL Neutral Gas Shielding Model for Pellet-Plasma Interactions," ORNL/TM-5776(1977).
2. R. G. Mills, "Linear Resonance Acceleration of Pellets," Proceedings of the Princeton Fueling Workshop, (1978) CONF-771129.
3. C. Bruno, "Fueling by Liquid Jets," Proceedings of the Princeton Fueling Workshop, (1978) CONF-771129.
4. F. S. Felber, "Laser Acceleration of Reactor Fuel Pellets," Nuclear Fusion 18 1469, (1978).
5. C. A. Foster and S. L. Milora, "ORNL Pellet Acceleration Program," Proceedings of the Princeton Fueling Workshop, (1978) CONF-771129.
6. L. D. Stewart, "Summary of Fueling by Pellet Injection," Proceedings of the Princeton Fueling Workshop, (1978) CONF-771129.
7. S. L. Milora, C. A. Foster, P. H. Edmonds, and G. L. Schmidt, "Hydrogen Pellet Fueling Experiment on the ISX-A Tokamak," ORNL/TM-6496 (1978).
8. R. F. Flago, "A Review of Gas Gun Technology With Emphasis on Fusion Fueling Applications," Proceedings of the Princeton Fueling Workshop, (1978) CONF-771129.
9. S. L. Milora and C. A. Foster, "Pneumatic Hydrogen Pellet Injection for the ISX-A Tokamak," ORNL/TM-6598 (1978).
10. L. L. Lengyel and W. Riedmüller, "Estimates on the Acceleration of Pellets by Gasdynamic and Electrostatic Means," IPP 4/171 (1978).
11. S. L. Gralnick, "Fuel Injection," PPPL/MATT-1050.
12. S. L. Gralnick, "Solid Deuterium Evaporation in a Fusion Plasma," Nuclear Fusion 13 703, (1973).
13. McAlees, et. al., "Plasma Engineering in EPR," ORNL/TM-5573, pp. 92-108.
14. P. B. Parks, R. J. Turnbull, and C. A. Foster, "A Model for the Ablation Rate of a Solid Hydrogen Pellet in a Plasma," Nuclear Fusion 17 539, (1977)
15. C. A. Foster, et. al., "Solid Hydrogen Pellet Injection into the ORMAK Tokamak," Nuclear Fusion 17 1067, (1977)
16. D. F. Vaslow, "Scaling Law for Ablation of a Hydrogen Pellet in a Plasma," GA-A13659, (1976)
17. L. L. Lengyel, "On Pellet Ablation in Hot Plasmas and the Problem of Magnetic Shielding," IPP 4/160, (1977)

18. L. W. Jørgensen, A. H. Sillesen, and F. Øster, "Ablation of Hydrogen Pellets in Hydrogen and Helium Plasmas," *Plasma Physics* 17, 453, (1974).
19. C. T. Chang, "The Magnetic Shielding Effect of a Re-fuelling Pellet," *Nuclear Fusion* 15, 595, (1975).
20. A. H. Shapiro, The Dynamics and Thermodynamics of Compressible Fluid Flow, vol. 1, (John Wiley and Sons, Inc., 1953).
21. M. Gros, P. Bertrand, and M. R. Feix, "Connection Between Hydrodynamic, Water Bag, and Vlasov Models," *Plasma Physics* 20, 1075, (1978).

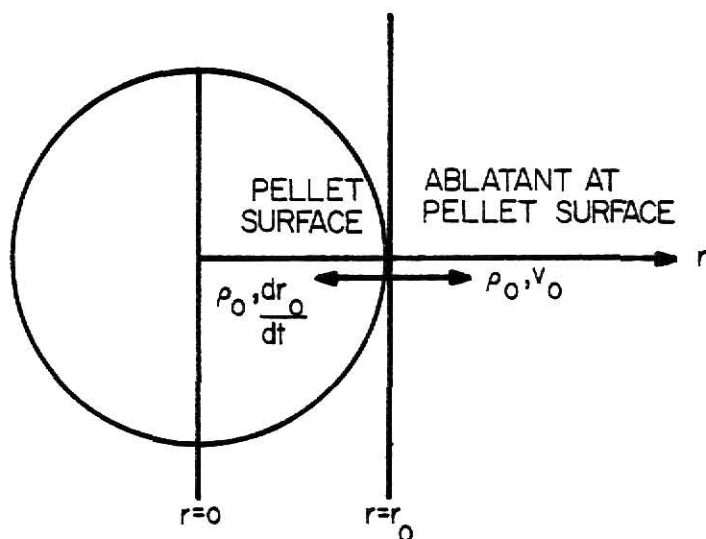


Figure 1. Explanatory diagram of spatial coordinate system and velocity and density boundary conditions at pellet surface. (The origin of the coordinate system moves with the pellet center at the pellet speed, u .)

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DOCUMENT(S) IS OF
POOR LEGIBILITY IN
THE ORIGINAL**

**THIS IS THE BEST
COPY AVAILABLE**

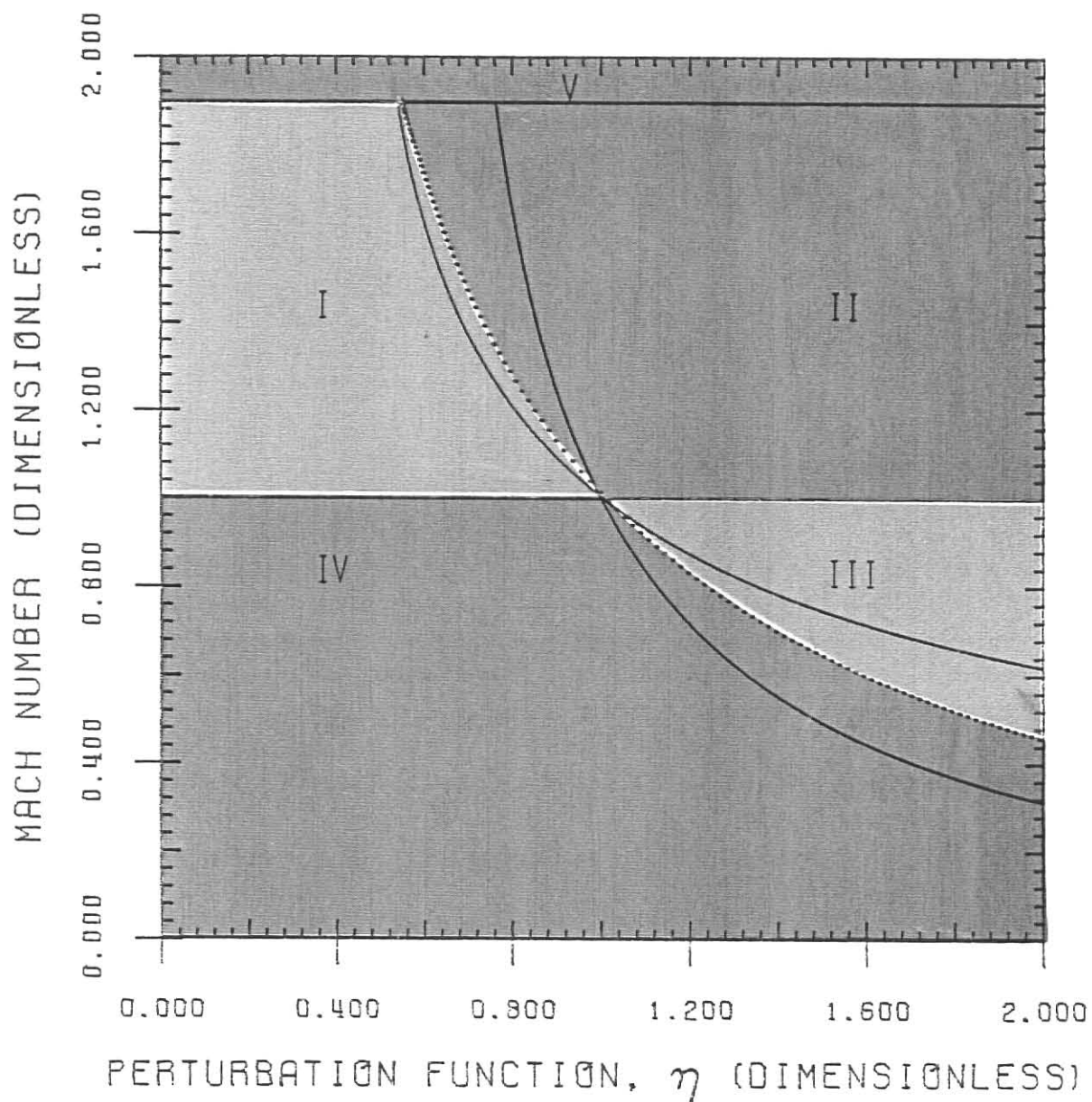


Figure 2. The M, η domain is divided into five regions. In Region V, $M > \sqrt{5/\gamma}$. In Regions II and III, $U < 0$. In Regions I and IV, $U > 0$. In Regions III and IV, $M < 1$. In Regions I and II, $\sqrt{5/\gamma} > M > 1$. Hence, the physically realistic solution must lie in Regions I or III. Furthermore, any solution in Regions II, IV, or V must be physically impossible. The dotted line represents the solution to $U = 0$.

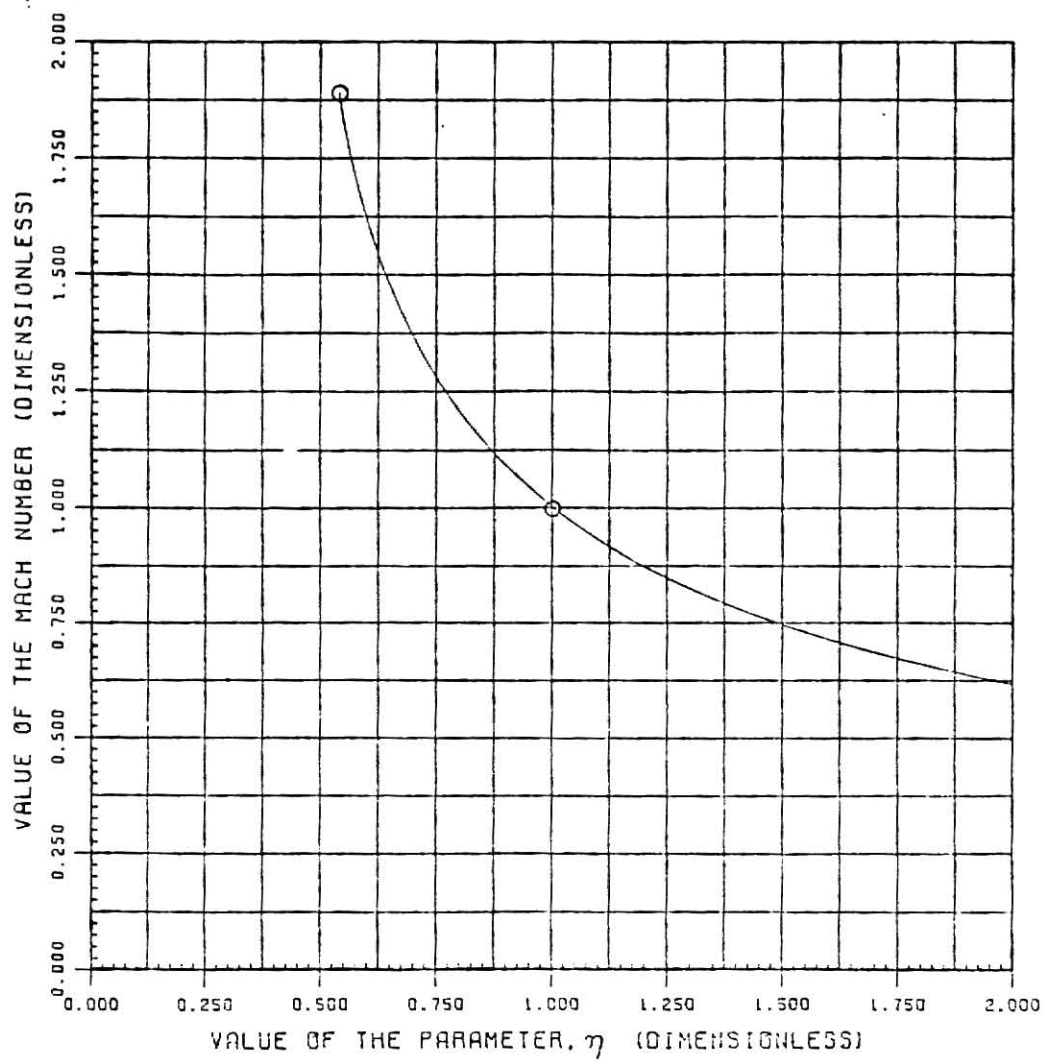


Figure 3. Solution of differential equation (31) for the physically realistic sound-barrier conditions.

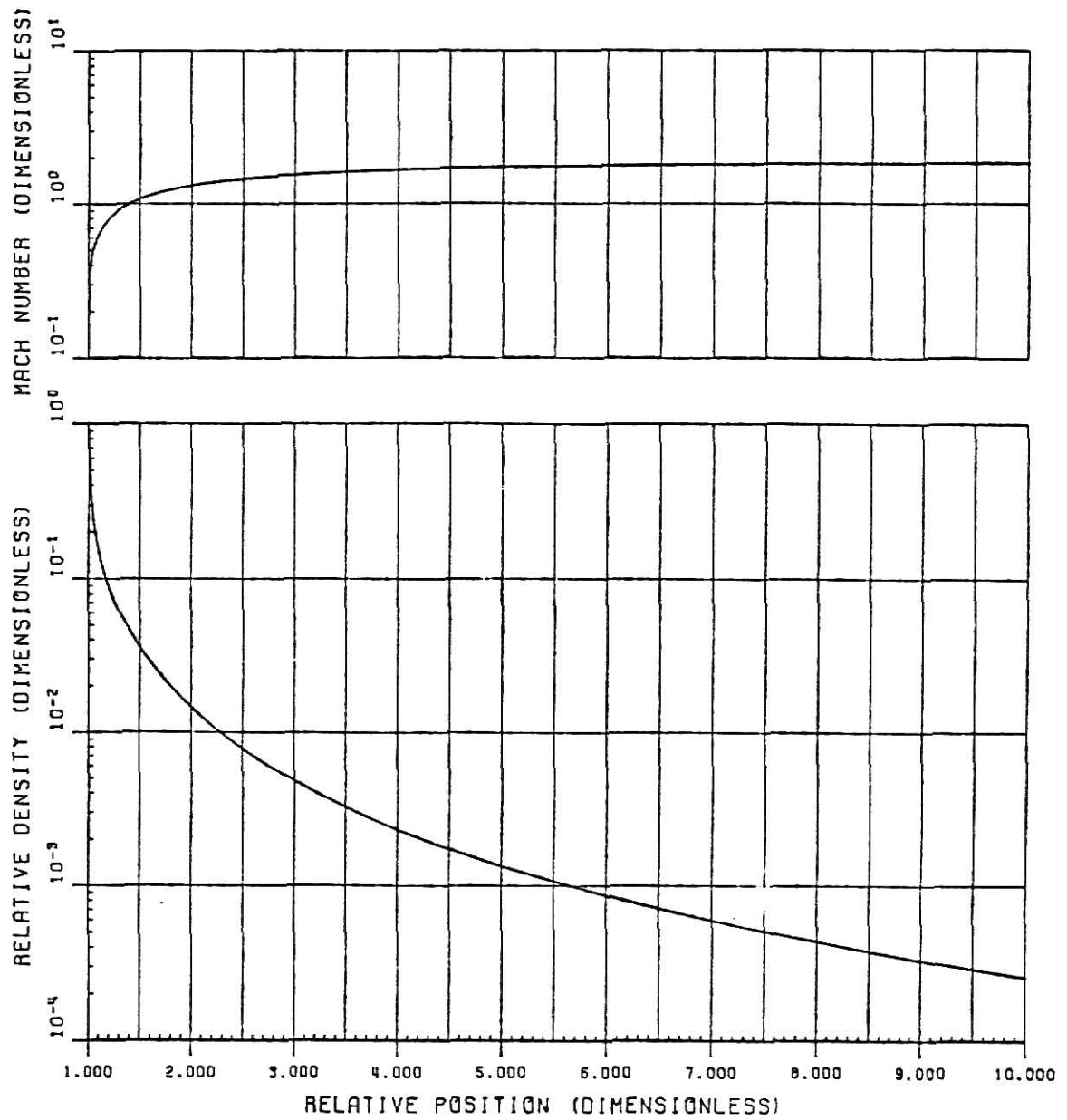


Figure 4. Solution to system of coupled differential equations for case of $\xi = 1000$ (reference case - Milora and Foster¹).

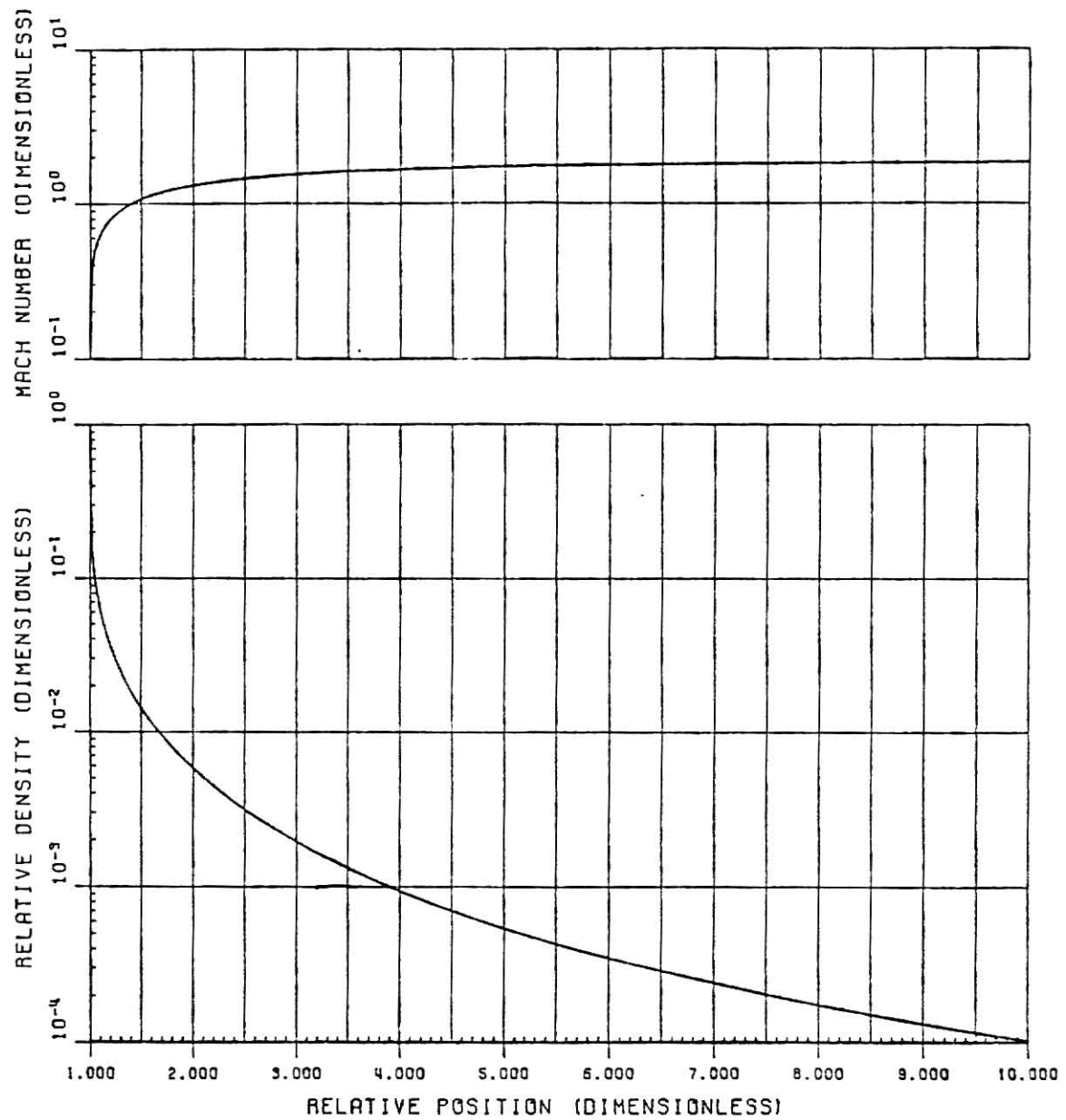


Figure 5. Solution to system of coupled differential equations for CTHR case ($\xi = 15,301$).

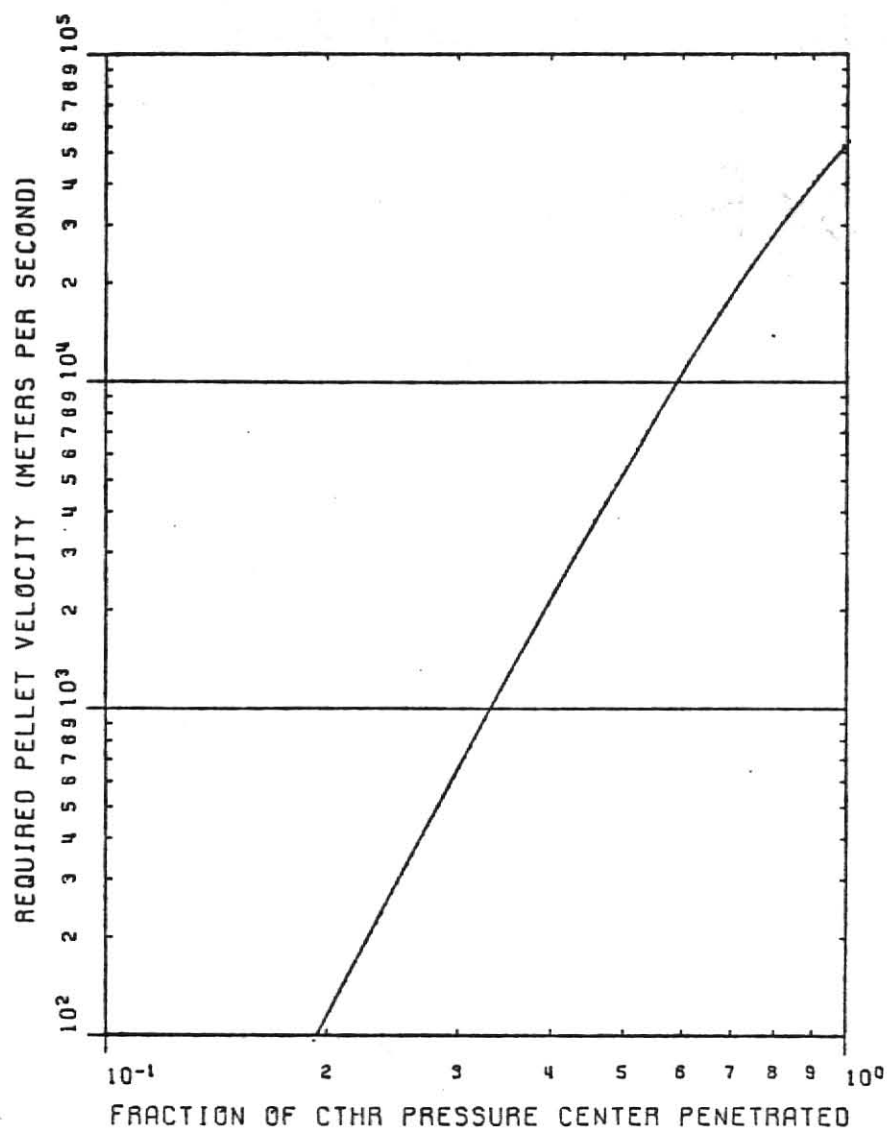


Figure 6. Required pellet velocity is plotted as a function of the point at which the pellet will be totally ablated, which is known as the penetration fraction or depth.

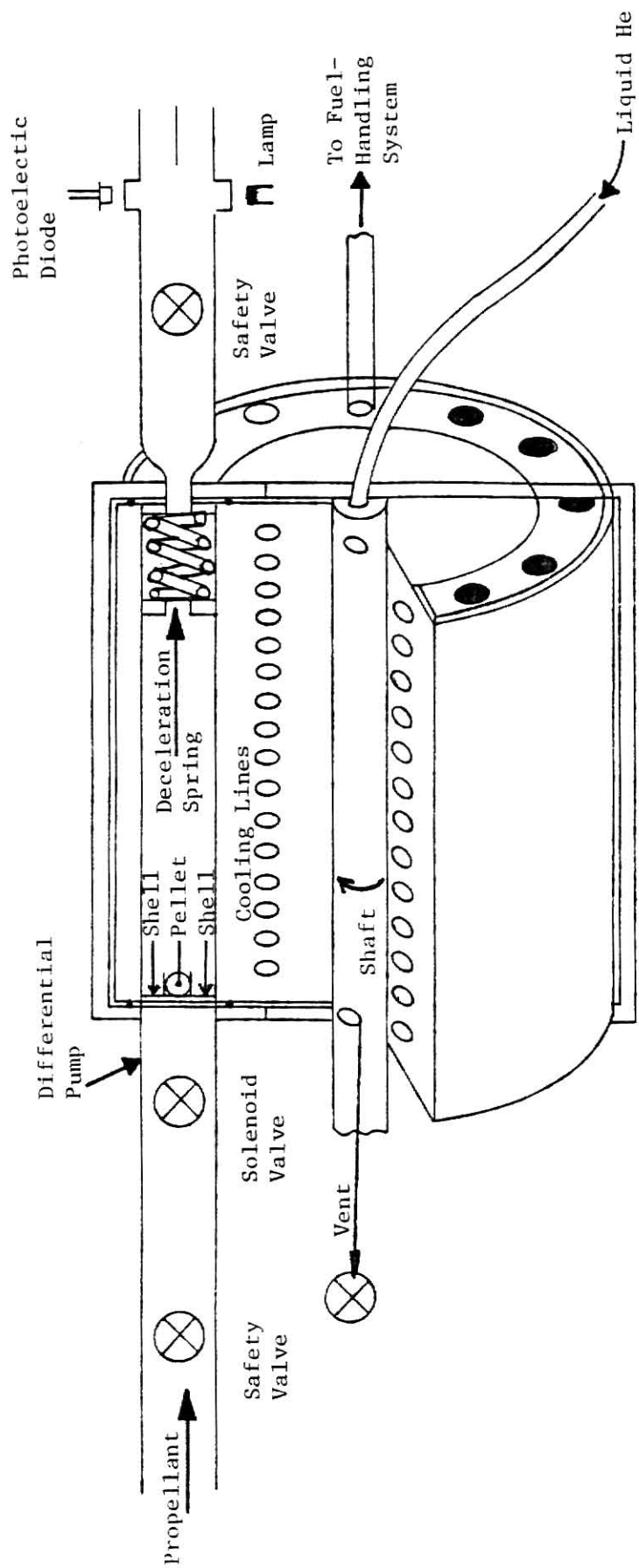


Figure 7. Fuel injection system for CTHR

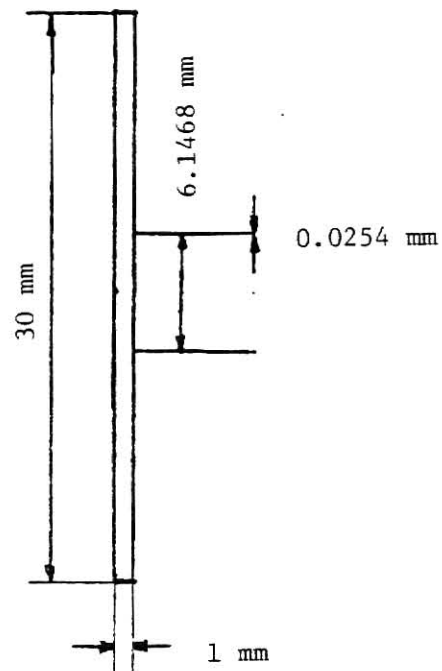


Figure 8. The shell for the fuel injector is a 1mm by 30mm dia disk intersected by a cylindrical shell whose thickness is 0.0254 mm and whose inner diameter is 6.1648 mm.

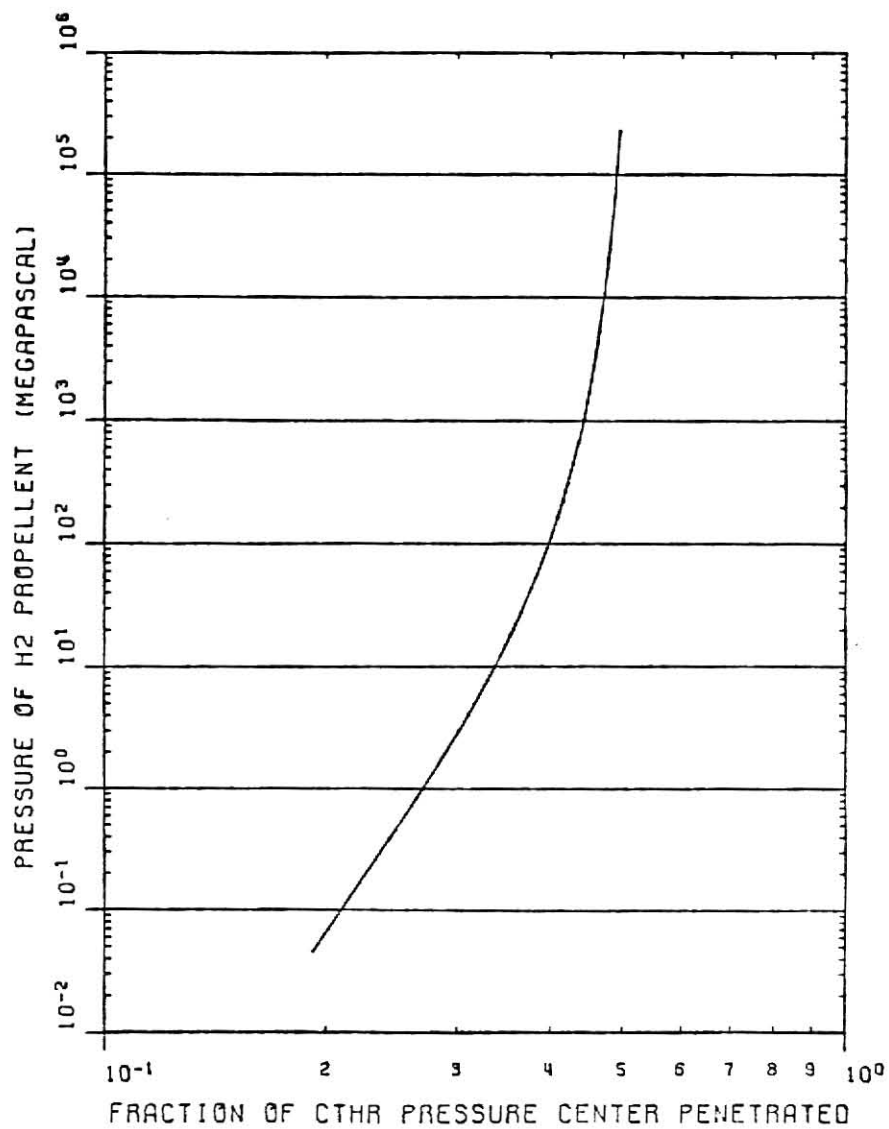


Figure 9. The required gas pressure, utilizing H₂ as propellant, is plotted as a function of the penetration fraction.

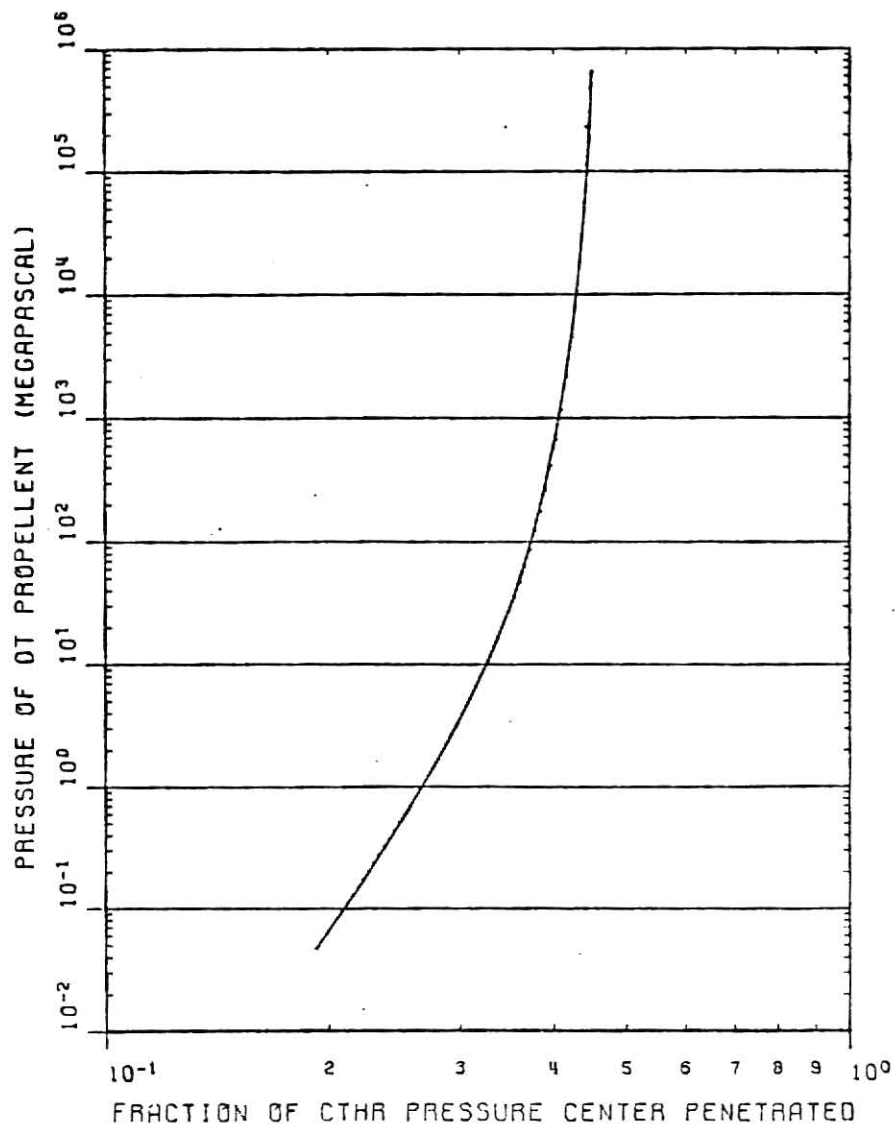


Figure 10. The required gas pressure, utilizing DT as propellant, is plotted as a function of the penetration fraction.

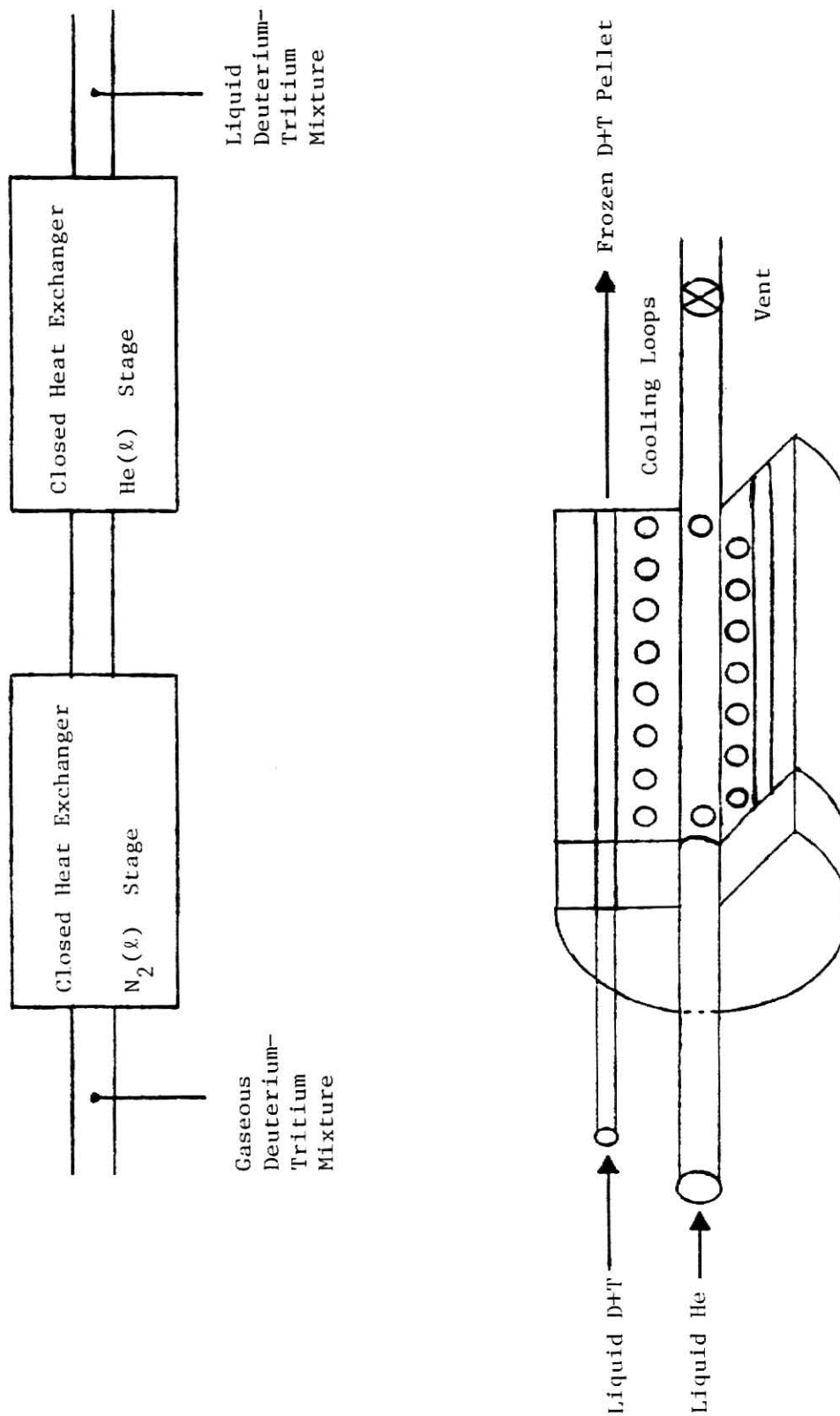


Figure 11. The fuel handling system is as diagrammed above. The liquid DT mixture enters the cylindrical chamber (above), is frozen by the liquid He, and on the next cycle is forced out as a pellet by the entering DT(ℓ).

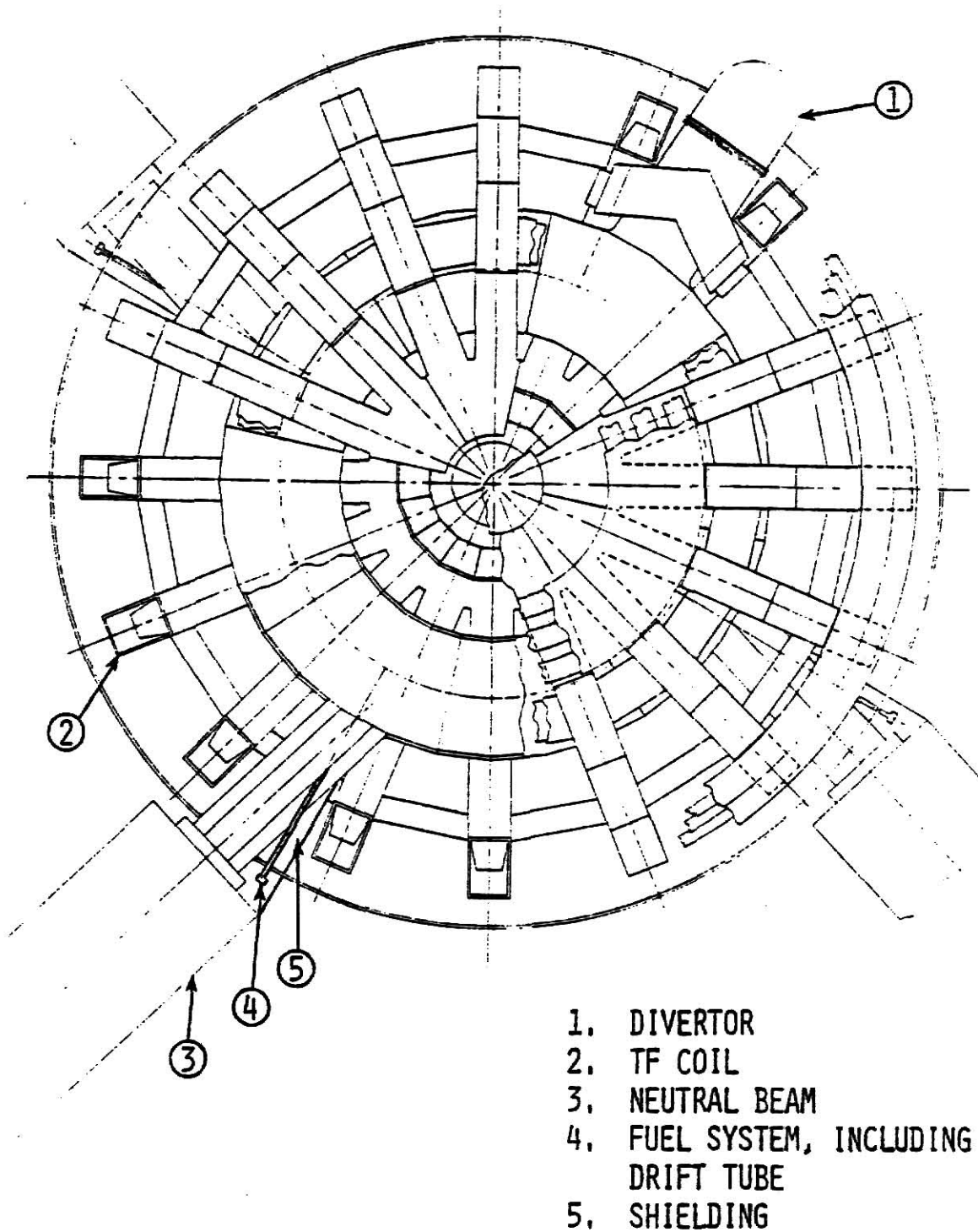


Figure 12. Incorporation of fueling systems into CTHR.

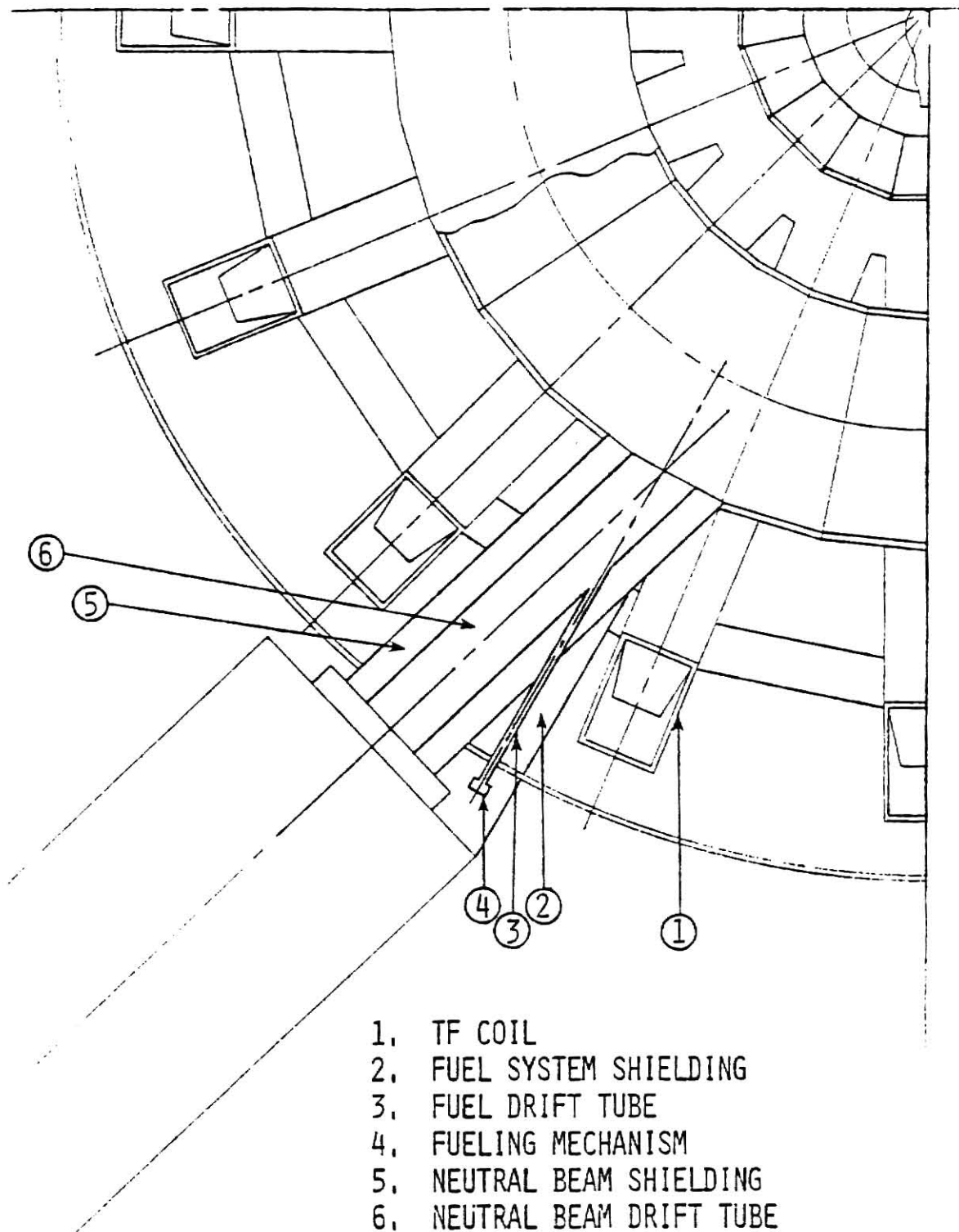


Figure 13. Closeup view of incorporation of pellet injector into CTHR.

Appendix I. Computer Program to Determine M_o .

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C*****
C****                                     ****ML001
C****                                     ****ML002
C****                                     ****ML003
C**** PROGRAM AND SUBROUTINES BY KENNETH DALE MATNEY. ****ML004
C****                                     ****ML005
C**** THE PURPOSE OF THIS PROGRAM IS TO CALCULATE THE DEPENDENCE OF THE ****ML006
C**** MACH NUMBER ON THE FUNCTION, ETA, ASSUMING THAT LAMBDA-1 IS THE ****ML007
C**** APPROPRIATE ROOT OF THE QUADRATIC EQUATION. ****ML008
C****                                     ****ML009
C****                                     ****ML010
C*****
C*****
ISN 0002      IMPLICIT REAL*8(A-H,O-Z)                                ML011
ISN 0003      REAL ERR,E,ER                                           ML012
ISN 0004      REAL*8 K1,K2,K3,K4                                       ML013
ISN 0005      DIMENSION X(1005),Y(1005,1),K1(11),K2(11),K3(11),K4(11),E(2), ML014
                                <DMDF(1005),YQ(2),YP(8),YC(8),ER(8)      ML015
                                COMMON H                                ML016
ISN 0006      COMMON H                                                  ML017
ISN 0007      COMMON/HANDAT/CRIT,PSI,GAMMA,IFRKG                       ML018
ISN 0008      1 FORMAT(09.2,1X,023.16,7X,11)                            ML019
ISN 0009      2 FORMAT(1PD22.15,8X,022.15,12X,'L1 N,M DATA ',I4)      ML020
ISN 0010      10 FORMAT(/,11X,'**THE VALUE OF THE LAST INCREMENT IS:',1PD12.2,'**') ML021
ISN 0011      11 FORMAT(1H1,112X,'PAGE ',I5)                           ML022
ISN 0012      12 FORMAT(11X,'THE MACH NUMBER AND DM/DN AS A FUNCTION OF F FOR H = ' ML023
                                <,1PD9.2,'.')                        ML024
ISN 0013      13 FORMAT(11X,'THE MACH NUMBER AND DM/DN AS A FUNCTION OF F FOR H = ' ML025
                                <,1PD9.2,' — CONTINUED.')          ML026
ISN 0014      14 FORMAT(/,22X,'N',27X,'M',25X,'DM/DN'/6X,3(5X,23('—')) ML027
ISN 0015      15 FORMAT(/,11X,1PD23.16,5X,023.16,5X,023.16)          ML028
ISN 0016      22 FORMAT(11X,'N AND DN/DM AS A FUNCTION OF THE MACH NUMBER FOR H = ' ML029
                                <,1PD9.2,'.')                        ML030
ISN 0017      23 FORMAT(11X,'N AND DN/DM AS A FUNCTION OF THE MACH NUMBER FOR H = ' ML031
                                <,1PD9.2,' — CONTINUED.')          ML032
ISN 0018      24 FORMAT(/,22X,'M',27X,'N',25X,'DN/DM'/6X,3(5X,23('—')) ML033
ISN 0019      ND=1005                                                  ML034
ISN 0020      GAMMA=1.400                                              ML035
ISN 0021      IPAGE=1                                                  ML036
ISN 0022      99 CONTINUE                                              ML037
ISN 0023      IFRKG=-1                                                ML038
ISN 0024      WRITE(6,11)IPAGE                                         ML039
ISN 0025      IPAGE=IPAGE+1                                             ML040
ISN 0026      X(1)=1.00                                                ML041
ISN 0027      Y(1,1)=1.00                                              ML042
ISN 0028      N=1                                                      ML043
C****                                     ****ML044
C****                                     ****ML045
C**** READ THE VALUE OF THE INCREMENT TO BE TAKEN, THE VALUE OF THE ****ML046
C**** HEATING PARAMETER, AND IPUNCH. IF IPUNCH IS SET EQUAL TO ZERO, ****ML047
C**** THE PUNCHING OF THE OUTPUT WILL BE SUPPRESSED. ****ML048
C****                                     ****ML049
C****                                     ****ML050
ISN 0029      READ(5,1,END=9999)H,PSI,IPUNCH                          ML051
ISN 0030      D=DABS(H)                                                ML052
ISN 0031      ITAG=M/2.YO-1/DABS(H)                                    ML053
ISN 0032      HSAV=H                                                    ML054
ISN 0033      IF(H.LT.0.)WRITE(6,22)HSAV                               ML055
ISN 0035      IF(H.GT.0.)WRITE(6,12)HSAV                               ML056

```


ISN 0037	IF(H.LT.0.)WRITE(6,24)	ML057
ISN 0039	IF(H.GT.0.)WRITE(6,14)	ML058
ISN 0041	NINC=5.0-02/DABS(H)+0.100	ML059
ISN 0042	IF(NINC.LT.1)NINC=1	ML060
ISN 0044	NINC1=1	ML061
ISN 0045	K0=0	ML062
ISN 0046	XSAV=PSI*(1.00/3.00)	ML063
ISN 0047	IF(H.LT.0.00)XSAV=DSJKT(5.00/GAMMA)	ML064
ISN 0049	NTOT=(XSAV-1.00)/DABS(H)+1.100	ML065
ISN 0050	NSAV=NTOT	ML066
ISN 0051	100 CONTINUE	ML067
ISN 0052	NINTER=NTGT	ML068
ISN 0053	IF(NINTER.GT.1000)NINTER=1003	ML069
ISN 0055	IF(NTOT.LE.1000)NTOT=0	ML070
ISN 0057	IF(NTOT.GT.1000)NTOT=NTOT-1000	ML071
ISN 0059	X0=1.00	ML072
ISN 0060	Y0(1)=1.00	ML073
ISN 0061	NC=1	ML074
ISN 0062	A1=0.00	ML075
ISN 0063	A2=0.00	ML076
ISN 0064	ERR=1.E-14	ML077
ISN 0065	CALL HAMMIN(X0,Y0,NC,NINTER,D,A1,A2,ERR,X,Y,IER,YP,YC,K1,K2,K3,K4, <E,ER,ND)	ML078
ISN 0066	IFRKG=IFRKG+1	ML079
ISN 0067	IF(NTOT.EQ.0)D=XSAV-X(NINTER-1)	ML080
ISN 0069	IF(NTOT.EQ.0)CALL RKGSQ(NINTER,X,Y,NC,D,E,K1,K2,K3,K4,ND)	ML081
ISN 0071	DO 200 J=1,NINTER,1	ML082
ISN 0072	IF(H.GT.0.)DMDF(J)=F(L,X,Y,J)	ML083
ISN 0074	IF(H.LT.0.)DMDF(J)=F(L,X,Y,J)	ML084
ISN 0076	200 CONTINUE	ML085
ISN 0077	IF(I.PUNCH.EQ.0)GO TO 251	ML086
ISN 0079	DO 250 J=1,1000,1	ML087
ISN 0080	IF(J.GT.NINTER)GO TO 251	ML088
ISN 0082	IF(H.LT.0.00)PUNCH 2,Y(J,1),X(J),ITAG	ML089
ISN 0084	IF(H.GT.0.00)PUNCH 2,X(J),Y(J,1),ITAG	ML090
ISN 0086	ITAG=ITAG+H/9.90-1/DABS(H)	ML091
ISN 0087	250 CONTINUE	ML092
ISN 0088	251 CONTINUE	ML093
ISN 0089	DO 300 J=NINC1,NINTER,NINC	ML094
ISN 0090	IF(NINTER.LT.1000.AND.NINTER.LT.NINC1)GO TO 299	ML095
ISN 0092	WRITE(6,15)X(J),Y(J,1),DMDF(J)	ML096
ISN 0093	JINTER=J	ML097
ISN 0094	K0=K0+1	ML098
ISN 0095	IF(K0.NE.25)GO TO 300	ML099
ISN 0097	K0=0	ML100
ISN 0098	WRITE(6,11)IPAGE	ML101
ISN 0099	IPAGE=IPAGE+1	ML102
ISN 0100	IF(H.LT.0.)WRITE(6,23)HSAV	ML103
ISN 0102	IF(H.GT.0.)WRITE(6,13)HSAV	ML104
ISN 0104	IF(H.LT.0.)WRITE(6,24)	ML105
ISN 0106	IF(H.GT.0.)WRITE(6,14)	ML106
ISN 0108	GO TO 300	ML107
ISN 0109	299 CONTINUE	ML108
ISN 0110	WRITE(6,15)X(NINTER),Y(NINTER,1),DMDF(NINTER)	ML109
ISN 0111	GO TO 301	ML110
ISN 0112	300 CONTINUE	ML111
ISN 0113	301 CONTINUE	ML112
ISN 0114	IF(NTOT.NE.0)GO TO 350	ML113
		ML114

ISN 0116	WRITE(6,10)D	ML115
ISN 0117	GO TO 99	ML116
ISN 0118	350 CONTINUE	ML117
ISN 0119	DO 400 J=1,4,1	ML118
ISN 0120	X(J)=X(NINTER-4+J)	ML119
ISN 0121	Y(J,1)=Y(NINTER-4+J,1)	ML120
ISN 0122	400 CONTINUE	ML121
ISN 0123	NINC1=NINC+1	ML122
ISN 0124	GO TO 100	ML123
ISN 0125	9999 CONTINUE	ML124
ISN 0126	STOP	ML125
ISN 0127	END	ML126

```

C*****PC005
C***PC010
C***PC015
C*** THE PURPOSE OF HAMMIN IS TO SOLVE A SET OF SIMULTANEOUS FIRST PC020
C*** ORDER FUNCTIONALS OF Y WITH RESPECT TO X USING HAMMING'S FIFTH PC025
C*** ORDER PREDICTOR-CORRECTOR. FOURTH ORDER RUNGA-KUTTA-3ILL PC030
C*** PROVIDES STARTING VALUES WHICH ARE ITERATED UPON BY THREE OTHER PC035
C*** INTEGRATION FORMULAE UNTIL A CONSISTENT STARTER IS OBTAINED. PC040
C***PC045
C***PC050
C*** X0 = INITIAL VALUE OF X. PC055
C***PC060
C*** Y0 = ARRAY OF INITIAL Y-VALUES. PC065
C***PC070
C*** NC = NUMBER OF COUPLED EQUATIONS. PC075
C***PC080
C*** H = INTEGRATION INCREMENT PC085
C***PC090
C*** N = NUMBER OF INTEGRATIONS TO BE PERFORMED. PC095
C***PC100
C*** A1,A2 = PREDICTOR-CORRECTOR PARAMETERS TO BE CHOSEN BY USER. PC105
C***PC110
C*** ERR = CONVERGENCE PARAMETER FOR STARTER. PC115
C***PC120
C*** X = ARRAY OF X-VALUES PC125
C***PC130
C*** Y = 2-DIMENSIONAL ARRAY OF Y-VALUES. PC135
C***PC140
C*** IER = ERROR PARAMETER. IF IER=0, INTEGRAL INSIDE OF ASYMPTOTES. PC145
C*** IF IER=1, STARTER FAILED TO CONVERGE. PC150
C*** IF IER=2, INTEGRAL DIVERGES. PC155
C***PC160
C*** YC,W1,W2,W3,W4,E = WORK ARRAYS OF LENGTH NC. PC165
C***PC170
C*** YP,ER = WORK ARRAY OF MINIMUM SIZE (4,NC). PC175
C***PC180
C*** ND = DIMENSION SIZE OF Y IN MAIN PROGRAM. PC185
C***PC190
C*** TO GET DOUBLE PRECISION VERSION, REMOVE C'S FROM PC240 AND PC245. PC195
C***PC200
C***PC205
C*****PC210
ISN 0002 SUBROUTINE HAMMIN(X0,Y0,NL,N,H,A1,A2,ERR,X,Y,IER,YP,YC,W1,W2,W3,W4 PC215
          $,E,ER,ND) PC220
ISN 0003 REAL*8 X0,Y0,H,A1,A2,X,Y,YP,YC,W1,W2,W3,W4,C1,C2,C3,C4,C5,C6,C7,EP PC225
          $,EC,A0,B0,B1,B2,B3,BCN1,ECO,BC1,BC2,CRIT,F,PSI,GAMMA PC230
ISN 0004 COMMON/HAMDAT/CRIT,PSI,GAMMA,IFRKG PC235
ISN 0005 DIMENSION Y0(1),X(1),Y(ND,1),YP( 4,1),YC(1),E(1),W1(1),W2(1),W3(1) PC240
          $,W4(1),ER( 4,1) PC245
ISN 0006 IER=0 PC250
ISN 0007 C1=9.0+00 PC255
ISN 0008 C2=19.0+00 PC260
ISN 0009 C3=5.0+00 PC265
ISN 0010 C4=24.0+00 PC270
ISN 0011 C5=4.0+00 PC275
ISN 0012 C6=3.0+00 PC280

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ISN 0013	C7=8.D+00	PC285
ISN 0014	EP=(2.51D+02-C2*A1-C7*A2)/6.D+00	PC290
ISN 0015	EC=(-C2+11.D+00*A1-C7*A2)/6.D+00	PC295
ISN 0016	A7=1.D+00-A1-A2	PC300
ISN 0017	B0=(55.D+00+C1*A1+C7*A2)/C4	PC305
ISN 0018	B1=(-59.D+00+C2*A1+32.D+00*A2)/C4	PC310
ISN 0019	B2=(37.D+00-C3*A1+C7*A2)/C4	PC315
ISN 0020	B3=(A1-C1)/C4	PC320
ISN 0021	BCN1=(C1-A1)/C4	PC325
ISN 0022	BC0=(C2+13.D+00*A1+C7*A2)/C4	PC330
ISN 0023	BC1=(-C3+13.D+00*A1+32.D+00*A2)/C4	PC335
ISN 0024	BC2=(1.D+00-A1+C7*A2)/C4	PC340
ISN 0025	IF(IFRKG.GE.0)GO TO 155	PC345
ISN 0027	X(1)=X0	PC350
ISN 0028	DO 100 L=1,NC,1	PC355
ISN 0029	Y(1,L)=Y0(L)	PC360
ISN 0030	100 CONTINUE	PC365
ISN 0031	DO 115 J=2,4,1	PC370
ISN 0032	CALL RKGSDQ(J,X,Y,NC,H,E,W1,W2,W3,W4,ND)	PC375
ISN 0033	115 CONTINUE	PC380
ISN 0034	JCHECK=0	PC385
ISN 0035	120 CONTINUE	PC390
ISN 0036	JCHECK=JCHECK+1	PC395
ISN 0037	DO 130 L=1,NC,1	PC400
ISN 0038	YP(2,L)=Y(1,L) +H*(C1*F(L,X,Y,1)+C2*F(L,X,Y,2)-C3*F(L,X,Y,3)+F(L,X,Y,4))/C4	PC405
ISN 0039	YP(3,L)=Y(1,L) +H*(F(L,X,Y,1)+C5*F(L,X,Y,2)+F(L,X,Y,3))/C6	PC410
ISN 0040	YP(4,L)=Y(1,L) +H*(F(L,X,Y,1)+C6*(F(L,X,Y,2)+F(L,X,Y,3))+F(L,X,Y,4))/C7C6	PC415
ISN 0041	DO 125 J=2,4,1	PC420
ISN 0042	ER(J,L)=(Y(J,L)-YP(J,L))/Y(J,L)	PC425
ISN 0043	ER(J,L)=ABS(ER(J,L))	PC430
ISN 0044	125 CONTINUE	PC435
ISN 0045	130 CONTINUE	PC440
ISN 0046	ITEST=0	PC445
ISN 0047	DO 140 L=1,NC,1	PC450
ISN 0048	DO 135 J=2,4,1	PC455
ISN 0049	IF(ER(J,L).GT.ERR)ITEST=1	PC460
ISN 0051	Y(J,L)=YP(J,L)	PC465
ISN 0052	135 CONTINUE	PC470
ISN 0053	140 CONTINUE	PC475
ISN 0054	IF(JCHECK.GE.500)GO TO 145	PC480
ISN 0056	IF(ITEST.EQ.1)GO TO 120	PC485
ISN 0058	145 DO 150 L=1,NC,1	PC490
ISN 0059	YP(1,L)=0.D+00	PC495
ISN 0060	YC(L)=0.D+00	PC500
ISN 0061	150 CONTINUE	PC505
ISN 0062	155 CONTINUE	PC510
ISN 0063	DO 300 J=4,N,1	PC515
ISN 0064	X(J+1)=X(J)+H	PC520
ISN 0065	DO 175 L=1,NC,1	PC525
ISN 0066	YP(2,L)=YP(1,L)	PC530
ISN 0067	175 CONTINUE	PC535
ISN 0068	DO 200 L=1,NC,1	PC540
ISN 0069	YP(1,L)=A0*Y(J,L)+A1*Y(J-1,L)+A2*Y(J-2,L)+H*(B0*F(L,X,Y,J)+B1*F(L,X,Y,J-1)+B2*F(L,X,Y,J-2)+B3*F(L,X,Y,J-3))	PC545
ISN 0070	200 CONTINUE	PC550
ISN 0071	DO 225 L=1,NC,1	PC555
		PC560
		PC565
		PC570

ISN 0072	Y(J+1,L)=YP(1,L)-(EP/(EP-EC))*(YP(2,L)-YC(L))	PC575
ISN 0073	225 CONTINUE	PC580
ISN 0074	DO 250 L=1,NC,1	PC585
ISN 0075	YC(L)=A0*Y(J,L)+A1*Y(J-1,L)+A2*Y(J-2,L)+H*(BCN1*F(L,X,Y,J+1)+BC0*F(L,X,Y,J)+BC1*F(L,X,Y,J-1)+BC2*F(L,X,Y,J-2))	PC590
		PC595
ISN 0076	250 CONTINUE	PC600
ISN 0077	DO 275 L=1,NC,1	PC605
ISN 0078	Y(J+1,L)=YC(L)-(EC/(EP-EC))*(YP(1,L)-YC(L))	PC610
ISN 0079	275 CONTINUE	PC615
ISN 0080	IF(J-N)285,300,400	PC620
ISN 0081	285 CONTINUE	PC625
	C IF(Y(J+1,1).GT.1.D+00.OR.Y(J+1,1).LT.Y(J,1))GC TC 500	PC630
ISN 0082	300 CONTINUE	PC635
ISN 0083	IER=ITEST	PC640
ISN 0084	400 N=N+1	PC645
ISN 0085	RETURN	PC650
ISN 0086	END	PC670

```

C*****
C***
C***
C*** THE PURPOSE OF RKGSDQ IS TO SOLVE, USING RUNGA-KUTTA-GILL, A SET
C*** OF SIMULTANEGUS DIFFERENTIAL EQUATIONS, WHICH ARE FIRST OKDER IN
C*** A COMMON (DUMMY) VARIABLE. THE QUADRATURE IS FOURTH ORDER.
C***
C***
C*** J IS THE ITERATE OF Y BEING SOLVED.
C***
C***
C*** X(J) IS THE DUMMY VARIABLE.
C***
C*** Y(L,J) IS THE SOLUTION TO THE L'TH DIFFERENTIAL EQUATION.
C***
C*** N IS THE NUMBER OF SIMULTANEGUS DIFFERENTIAL EQUATIONS.
C***
C*** H IS THE INCREMENT TO BE ADDED TO X(J-1) TO GET X(J).
C***
C*** E IS A PARAMETER USED IN THE RULE OF COLLATZ, OF LENGTH,N.
C***
C*** K1 IS A REAL WORK ARRAY OF LENGTH, N.
C***
C*** K2 IS A REAL WORK ARRAY OF LENGTH, N.
C***
C*** K3 IS A REAL WORK ARRAY OF LENGTH, N.
C***
C*** K4 IS A REAL WORK ARRAY OF LENGTH, N.
C***
C***
C*****
ISN 0002 SUBROUTINE RKGSDQ(J,X,Y,N,H,E,K1,K2,K3,K4,M) RK 31
ISN 0003 REAL*6 X,Y,H,K1,K2,K3,K4,AMD,ADD,AMS,AM,AD,AX,F RK 32
C REAL K1,K2,K3,K4 RK 33
C DIMENSION X(1),Y(M,1),E(1),K1(1),K2(1),K3(1),K4(1) RK 34
AMD=2.0+00 RK 35
ADD=1.0+00+DSQRT(5.0-01) RK 36
AMS=5.0-01-DSQRT(5.0-01) RK 37
AM=2.0+00-DSQRT(2.0+00) RK 38
AD=2.0+00+DSQRT(2.0+00) RK 39
AX=0.0+00 RK 40
ISN 0010 X(J)=X(J-1) RK 41
ISN 0011 DO 100 L=1,N,1 RK 42
ISN 0012 Y(J,L)=Y(J-1,L) RK 43
ISN 0013 100 CONTINUE RK 44
ISN 0015 DO 110 L=1,N,1 RK 45
ISN 0016 K1(L)=F(L,X,Y,J)*H RK 46
ISN 0017 110 CONTINUE RK 47
ISN 0018 DO 120 L=1,N,1 RK 48
ISN 0019 Y(J,L)=Y(J,L)+(K1(L)/AMD) RK 49
ISN 0020 120 CONTINUE RK 50
ISN 0021 X(J)=X(J)+(H/AMD) RK 51
ISN 0022 DO 130 L=1,N,1 RK 52
ISN 0023 K2(L)=F(L,X,Y,J)*H RK 53
ISN 0024 130 CONTINUE RK 54
ISN 0025 DO 140 L=1,N,1 RK 55
ISN 0026 Y(J,L)=Y(J,L)+(K2(L)-K1(L))*(AMD-ADD) RK 56

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ISN 0027	140	CONTINUE	RK 57
ISN 0028	DO 150	L=1,N,1	RK 58
ISN 0029		K3(L)=F(L,X,Y,J)*H	RK 59
ISN 0030	150	CONTINUE	RK 60
ISN 0031	DO 160	L=1,N,1	RK 61
ISN 0032		Y(J,L)=Y(J,L)+AMS*K1(L)-K2(L)+ADD*K3(L)	RK 62
ISN 0033	160	CONTINUE	RK 63
ISN 0034		X(J)=X(J-1)+H	RK 64
ISN 0035	DO 170	L=1,N,1	RK 65
ISN 0036		K4(L)=F(L,X,Y,J)*H	RK 66
ISN 0037	170	CONTINUE	RK 67
ISN 0038	DO 180	L=1,N,1	RK 68
ISN 0039		Y(J,L)=Y(J-1,L)+(K1(L)+AM*K2(L)+AD*K3(L)+K4(L))/AX	RK 69
ISN 0040	180	CONTINUE	RK 72
ISN 0041		RETURN	RK 73
ISN 0042		END	RK 74

CS/360 FORTRAN H

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C*****M,N01
C****M,N02
C****M,N03
C**** THE PURPOSE OF THIS FUNCTION IS TO PROVIDE DERIVATIVE FUNCTIONS ****M,N04
C**** FOR SUBROUTINES RKGSDQ AND HAMMIN. ****M,N05
C****M,N06
C****M,N07
C*****M,N08
ISN 0002      FUNCTION F(L,X,Y,J)M,N09
ISN 0003      REAL*8 X,Y,S,G,M,NUM,DEN,F3,L1,L2,F,CRIT,FNEW,HM,N10
ISN 0004      COMMON HM,N11
ISN 0005      COMMON/HAMDAT/CRIT,S,G,IFRKG M,N12
ISN 0006      DIMENSION X(1),Y(1005,1)M,N13
ISN 0007      M=Y(J,1)M,N14
ISN 0008      IF(H.LT.0.00)M=X(J)M,N15
ISN 0010      F3=X(J)*X(J)*X(J)M,N16
ISN 0011      IF(H.LT.0.00)F3=Y(J,1)*Y(J,1)*Y(J,1)M,N17
ISN 0013      FNEW=X(J)M,N18
ISN 0014      IF(H.LT.0.00)FNEW=Y(J,1)M,N19
ISN 0016      NUM=-3.00*(2.00-M*M*(1.00+F3)-G*(1.00-M*M*F3))*MM,N20
ISN 0017      DEN=(7.00-M*M*(1.00+6.00*F3))*FNEWM,N21
ISN 0018      IF(M.EQ.1.00.AND.FNEW.EQ.1.00)GO TO 00M,N22
ISN 0020      IF(H.LT.0.00)GO TO 100M,N23
ISN 0022      F=NUM/DENM,N24
ISN 0023      RETURNM,N25
ISN 0024      100 CONTINUEM,N26
ISN 0025      F=DEN/NUMM,N27
ISN 0026      RETURNM,N28
ISN 0027      200 CONTINUEM,N29
ISN 0028      L1=3.00*((G-1.00)-DSQRT((2.00-G)*(2.00-G)+7.00))M,N31
ISN 0029      L2=3.00*((G-1.00)+DSQRT((2.00-G)*(2.00-G)+7.00))M,N30
ISN 0030      IF(H.LT.0.00)GO TO 300M,N32
ISN 0032      F=-(L1+1.801)/1.401M,N33
ISN 0033      RETURNM,N34
ISN 0034      300 CONTINUEM,N35
ISN 0035      F=-1.401/(L1+1.801)M,N36
ISN 0036      RETURNM,N37
ISN 0037      ENDM,N38

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Appendix II. Computer Program to Solve Equations of Conservation and
Determine the Pellet Injection Velocity Requirement

CS/360 FORTRAN H

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C*****A0001
C***A0002
C***A0003
C*** PROGRAM BY KENNETH D. MATNEY. A0004
C***A0005
C***A0006
C*** THE ULTIMATE PURPOSE OF THIS PROGRAM IS TO DETERMINE FUEL PELLET A0007
C*** VELOCITY REQUIREMENTS FOR TOKAMAK TYPE REACTORS USING AN EQUAL A0008
C*** PARTS MIXTURE OF DEUTERIUM AND TRITIUM. A0009
C***A0010
C***A0011
C*****A0012
ISN 0002      IMPLICIT REAL*8(A-H,O-Z) A0013
ISN 0003      REAL EW(8),D,ER,MACH(2005),DENS(2005),RPCS(2005) A0014
ISN 0004      COMMON/FDATA/PSI,GAMMA,T,A,B,C,DARG A0015
ISN 0005      COMMON/HAMDAT/CRIT,IFRKG,LG A0016
ISN 0006      COMMON R(1005),Y(1005,2),D(2),XINIT A0017
ISN 0007      DIMENSION W1(2),W2(2),W3(2),W4(2),W(82),X(82),Y0(2),Y1(2),Y2(8) A0018
ISN 0008      2 FORMAT(2(D34.27,6X)) A0019
ISN 0009      3 FORMAT(1PD22.15,8X,D22.15,13X,'MACH DATA ',15) A0020
ISN 0010      4 FORMAT(1PD22.15,8X,D22.15,13X,'DENS DATA ',15) A0021
ISN 0011      5 FORMAT(4(D16.5,4X)) A0022
ISN 0012      6 FORMAT(16,4X,D21.14) A0023
ISN 0013      10 FORMAT(///) A0024
ISN 0014      11 FORMAT(1H1,115X,'PAGE ',15) A0025
ISN 0015      12 FORMAT(1H0,5X,'THE VALUE OF THE LOSS FUNCTION IS ',1PD25.18) A0026
ISN 0016      13 FORMAT(1H0,5X,'THE VALUE OF THE PELLET RADIUS IS ',1PD25.18/6X, A0027
        $,'THE VALUE OF THE ABLATANT VELOCITY AT THE BOUNDARY IS ',D25.18) A0028
ISN 0017      14 FORMAT(1H0,5X,'THE VALUE OF THE HEATING PARAMETER IS ',1PD25.18) A0029
ISN 0018      15 FORMAT(1H0,5X,'THE STABILITY POINTS ARE ',1PD25.18,' AND ',D25.18) A0030
ISN 0019      17 FORMAT(1H0,5X,'THE VALUE OF THE DESIRED CONSTANT IS ',1PD25.18) A0031
ISN 0020      18 FORMAT(1H0,5X,'THE BOUNDARY MACH NUMBER IS ',1PD25.18/ A0032
        $6X,'THE CONVERGENCE EFFICIENCY IS ',D25.18) A0033
ISN 0021      19 FORMAT(1H0,5X,14,7X,1PD25.18,2(10X,D25.18)) A0034
ISN 0022      20 FORMAT( 6X,'TABLE',14,'.  INTEGRATION DATA.'/4X,29(1H-)) A0035
ISN 0023      21 FORMAT( 6X,'TABLE',14,'.  INTEGRATION DATA -- CONTINUED.'/6X,42(1H A0036
        $-)) A0037
ISN 0024      22 FORMAT(1H0,5X,'R(',16,') = ',1PD25.18,', WHEN M = ',D25.13,', AND' A0038
        $',/6X,'RHO = ',D25.18,', WHERE DM/DR' = ',D25.18) A0039
ISN 0025      23 FORMAT(/1H0/1H0,5X,'*F(PSI) = ',1PD25.18) A0040
ISN 0026      24 FORMAT(/1H0,5X,'STEP',5X,'VALUE OF THE RELATIVE POSITION',5X,'VALU A0041
        $E OF THE RELATIVE DENSITY ',5X,'VALUE OF ABLATANT MACH NUMBER '/ A0042
        $6X,4(1H-),3(5X,30(1H-))) A0043
ISN 0027      25 FORMAT(/6X,'THE NUMBER OF INTEGRATIONS PERFORMED THIS STEP = ',16) A0044
ISN 0028      26 FORMAT(6X,'FOR THE PARAMETER L/A = ',1PD22.15,':') A0045
ISN 0029      27 FORMAT(1H0,5X,'THE VALUE OF THE PELLET VELOCITY MUST BE ',1PD22.15 A0046
        $,' TIMES THE PRESSURE CENTER DISTANCE.') A0047
ISN 0030      28 FORMAT(1PD22.15,8X,D22.15,11X,'PELLET SPEED',15) A0048
ISN 0031      FL(E)=1.0+00/(2.350+18+4.D+15+E+2.D+21/E/E) A0049
ISN 0032      ND=1005 A0050
ISN 0033      LG=0 A0051
ISN 0034      IPAGE=1 A0052
ISN 0035      ITABLE=1 A0053
ISN 0036      K=0 A0054
ISN 0037      NSTART=1 A0055
ISN 0038      WRITE(6,11)IPAGE A0056

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ISN 0039	IPAGE=IPAGE+1	A0057
ISN 0040	READ(5,6)N	A0058
ISN 0041	DO 100 J=1,N,1	A0059
ISN 0042	READ(5,2)X(J),W(J)	A0060
ISN 0043	X(J+N)=(1.0+00+X(J))/2.0+00	A0061
ISN 0044	X(J)=(1.0+00-X(J))/2.0+00	A0062
ISN 0045	W(J)=W(J)/2.0+00	A0063
ISN 0046	W(J+N)=W(J)	A0064
ISN 0047	100 CCNTINUE	A0065
ISN 0048	NSAVE=N+N	A0066
ISN 0049	N=NSAVE	A0067
ISN 0050	READ(5,5)EMAX,DN,VO,RP	A0068
ISN 0051	E=1.6021892D-19	A0069
ISN 0052	XM=(2.01410222D+00+3.01604972D+00)*1.660571D-27	A0070
ISN 0053	PI=3.141592653589793D+00	A0071
ISN 0054	GAMMA=1.4D+00	A0072
ISN 0055	PSI=0.0+00	A0073
ISN 0056	DO 105 J=1,N,1	A0074
ISN 0057	PSI=PSI+W(J)*FL(EMAX*X(J))	A0075
ISN 0058	105 CCNTINUE	A0076
ISN 0059	WRITE(6,12)PSI	A0077
ISN 0060	WRITE(6,13)RP,VO	A0078
ISN 0061	PL:=PSI*(GAMMA-1.0+00)/2.0+00+E/XM*RP/VO**3*DSQRT(E*EMAX/3.0+00/PI	A0079
	\$/1.660571D-27/548.579D-06)*DN	A0080
ISN 0062	T= MAX*2.0+00/3.0+00	A0081
ISN 0063	READ(5,2)A1SAVE,A2SAVE	A0082
ISN 0064	S=-1.0+00/3.0+00	A0083
ISN 0065	CRIT=DSQRT(5.0+00/GAMMA)	A0084
ISN 0066	READ(5,6)NSAV,HSAY	A0085
ISN 0067	READ(5,6)NTABLE	A0086
ISN 0068	110 CONTINUE	A0087
ISN 0069	READ(5,2,END=999)XMACH,DELTA	A0088
	C****	****A0N01
	C****	****A0N02
	C**** THE PURPOSE OF THIS, AN OBVIOUS ADDITION TO THE MAIN PROGRAM, IS	****A0N03
	C**** TO REDUCE THE CPU BUT MAINTAIN THE INTEGRITY IN COMPUTATION.	****A0N04
	C**** SINCE THE FOLLOWING SECTION ONLY COMPUTES THE VALUE OF THE	****A0N05
	C**** VARIABLE, CONST; IT IS SUBMITTED HERE AS COMPUTED IN EARLIER RUNS	****A0N06
	C**** IN THE FORM OF A DATA STATEMENT. THE SECTION OF PROGRAM WHICH	****A0N07
	C**** FOLLOWS IS THEREFORE BYPASSED, YIELDING A CPU SAVINGS OF ABOUT	****A0N08
	C**** 90%.	****A0N09
	C****	****A0N10
	C****	****A0N11
ISN 0070	DATA CONST/1.097763455886156D0/	A0N12
ISN 0071	IF(IPAGE.GT.0)GO TO 337	A0N13
ISN 0073	ROSAV=1.0+00	A0089
ISN 0074	RHOSAV=1.0+00	A0090
ISN 0075	NINC=5.0-3/HSAY+5.0-3	A0091
ISN 0076	IF(NINC.LT.1)NINC=1	A0092
ISN 0078	NEWTAB=0	A0093
ISN 0079	INITIL=1	A0094
ISN 0080	H=HSAY	A0095
ISN 0081	RO=ROSAV	A0096
ISN 0082	YO(1)=XMACH	A0097
ISN 0083	YO(2)=RHOSAV	A0098
ISN 0084	FPSI=0.0+00	A0099
ISN 0085	NC=2	A0100
ISN 0086	A1=A1SAVE	A0101

ISN 0087	A2=A2SAVE	A0102
ISN 0088	ER=1.E-14	A0103
ISN 0089	IFRKG=-1	A0104
ISN 0090	WRITE(6,14)PSI	A0105
ISN 0091	WRITE(6,15)A1SAVE,A2SAVE	A0106
ISN 0092	WRITE(6,19)INITIL,ROSAV,RHCSAV,XMACH	A0107
ISN 0093	WRITE(6,23)FPSI	A0108
ISN 0094	WRITE(6,11)IPAGE	A0109
ISN 0095	IPAGE=IPAGE+1	A0110
ISN 0096	NSTORE=0	A0111
ISN 0097	IF(NEWTAB.EQ.0)WRITE(6,20)ITABLE	A0112
ISN 0099	IF(NEWTAB.EQ.1)WRITE(6,21)ITABLE	A0113
ISN 0101	WRITE(6,24)	A0114
ISN 0102	NEWTAB=1	A0115
ISN 0103	LINES=0	A0116
ISN 0104	120 CONTINUE	A0117
ISN 0105	N=1003	A0118
ISN 0106	CALL HAMMIN(RO,YO,NC,N,H,A1,A2,ER,R,Y,I,Y2,Y1,W1,W2,W3,W4,D,EW,ND)	A0119
ISN 0107	NSTORE=NSTORE+1000	A0120
ISN 0108	N=1001	A0121
ISN 0109	DO 200 J=3,N,2	A0122
ISN 0110	FPSI=FPSI+(Y(J-2,2)+4.0+00*Y(J-1,2)+Y(J,2))*CABS(H)/3.0+00	A0123
ISN 0111	200 CONTINUE	A0124
ISN 0112	DO 300 J=NSTART,N,INC	A0125
ISN 0113	WRITE(6,19)J,R(J),Y(J,2),Y(J,1)	A0126
ISN 0114	K=K+1	A0127
ISN 0115	RPOS(K)=R(J)	A0128
ISN 0116	MACH(K)=Y(J,1)	A0129
ISN 0117	DENS(K)=Y(J,2)	A0130
ISN 0118	LINES=LINES+1	A0131
ISN 0119	IF(LINES-20)300,250,250	A0132
ISN 0120	250 CONTINUE	A0133
ISN 0121	IF(J+125 -N)275,300,300	A0134
ISN 0122	275 CONTINUE	A0135
ISN 0123	WRITE(6,11)IPAGE	A0136
ISN 0124	IPAGE=IPAGE+1	A0137
ISN 0125	WRITE(6,21)ITABLE	A0138
ISN 0126	WRITE(6,24)	A0139
ISN 0127	LINES=0	A0140
ISN 0128	300 CONTINUE	A0141
ISN 0129	N=N-1	A0142
ISN 0130	IFRKG=IFRKG+1	A0143
ISN 0131	DO 320 J=1,4,1	A0144
ISN 0132	DO 310 L=1,NC,1	A0145
ISN 0133	Y(J,L)=Y(J+N,L)	A0146
ISN 0134	310 CONTINUE	A0147
ISN 0135	R(J)=R(J+N)	A0148
ISN 0136	320 CONTINUE	A0149
ISN 0137	NSTART=NINC+1	A0150
ISN 0138	IF(NSTORE.GE.NSAV)GO TO 330	A0151
ISN 0140	GO TO 120	A0152
ISN 0141	330 CONTINUE	A0153
ISN 0142	WRITE(6,11)IPAGE	A0154
ISN 0143	IPAGE=IPAGE+1	A0155
ISN 0144	YNEFF=Y(N,1)/CRIT*1.02	A0156
ISN 0145	WRITE(6,18)XMACH,YNEFF	A0157
ISN 0146	WRITE(6,25)NSTORE	A0158
ISN 0147	UMDR=F(1,R,Y,N)	A0159

ISN 0148	WRITE(6,22)NSTORE,R(N),Y(N,1),Y(N,2),DMOR	A0160
ISN 0149	VALUE=3.00/4.00/((2.00+(GAMMA-1.00)*Y(N,1)**2)/PSI/Y(N,1)**2/(1.00	A0161
	+GAMMA*Y(N,1)**2))**S	A0162
ISN 0150	FPSI=FPSI+R(N)**(4.00*S)*VALUE	A0163
ISN 0151	CONST=PSI**S/FPSI	A0164
ISN 0152	WRITE(6,23)FPSI	A0165
ISN 0153	WRITE(6,17)CONST	A0166
ISN 0154	IF(YNEFF.LT.90.)GO TO 345	A0167
ISN 0156	DO 333 IP=1,K,1	A0168
ISN 0157	PUNCH 3,RPOS(IP),MACH(IP),IP	A0169
ISN 0158	333 CONTINUE	A0170
ISN 0159	DO 336 IP=1,K,1	A0171
ISN 0160	PUNCH 4,RPOS(IP),DENS(IP),IP	A0172
ISN 0161	336 CCNTINUE	A0173
ISN 0162	337 CONTINUE	A0174
ISN 0163	WRITE(6,11)IPAGE	A0175
ISN 0164	IPAGE=IPAGE+1	A0176
ISN 0165	XINIT=0.000	A0177
ISN 0166	ICOUNT=0	A0178
ISN 0167	XMD=2.0141022200*1.660571D-27	A0179
ISN 0168	XME=548.5790-6*1.660571D-27	A0180
ISN 0169	XMH=1.00866522D0*1.660571D-27	A0181
ISN 0170	A=CONST/(PI*XME)**(1.00/6.00)*(XMD/2.05902)**(4.00/9.00)*DSQRT(E)*	A0182
	S(2.35D19)**(2.00/3.00)*(4.00*PI/3.00)**(5.00/9.00)*(1.500/1.600)**	A0183
	S(7.00/6.00)*(3.00/XMH)**(1.00/3.00)/T**(1.00/6.00)	A0184
ISN 0171	B=9.00/3.7603	A0185
ISN 0172	C=-4.09602/3.807D0	A0186
ISN 0173	DARG=2.8208D1/1.0575D0	A0187
ISN 0174	FVOL=4.00*PI/3.00*RP*RP*RP*2.05902/DN/XMD	A0188
ISN 0175	338 CONTINUE	A0189
ISN 0176	WRITE(6,26)XINIT	A0190
ISN 0177	H=1.D-2	A0191
ISN 0178	V1=G(T,H)/(XM/XMH)**(1.00/3.00)/FVOL**((5.00/9.00)/DN**((2.00/9.00)	A0192
ISN 0179	WRITE(6,27)V1	A0193
ISN 0180	H=1.D-3	A0194
ISN 0181	V2=G(T,H)/(XM/XMH)**(1.00/3.00)/FVOL**((5.00/9.00)/DN**((2.00/9.00)	A0195
ISN 0182	WRITE(6,27)V2	A0196
ISN 0183	XPENET=1.00-XINIT	A0197
ISN 0184	ICOUNT=ICOUNT+1	A0198
ISN 0185	PUNCH 28,XPENET,V2,ICOUNT	A0199
ISN 0186	XINIT=1.00-XPENET*2.D-1**1.D-2	A0200
ISN 0187	IF(MCD(ICOUNT,6).EQ.0)GO TO 339	A0201
ISN 0189	WRITE(6,10)	A0202
ISN 0190	GO TO 340	A0203
ISN 0191	339 CONTINUE	A0204
ISN 0192	WRITE(6,11)IPAGE	A0205
ISN 0193	IPAGE=IPAGE+1	A0206
ISN 0194	340 CONTINUE	A0207
ISN 0195	IF(V2.GT.1.02)GO TO 338	A0208
ISN 0197	345 CONTINUE	A0209
ISN 0198	WRITE(6,11)IPAGE	A0210
ISN 0199	IPAGE=IPAGE+1	A0211
ISN 0200	ITABLE=ITABLE+1	A0212
ISN 0201	GO TO 110	A0213
ISN 0202	999 CCNTINUE	A0214
ISN 0203	STOP	A0215
ISN 0204	END	A0216

```
C*****  
C***  
C***  
C*** THE PURPOSE OF F IS TO PROVIDE THE COUPLED DIFFERENTIAL  
C*** EQUATIONS REQUIRED BY RKGSDQ AND HAMMIN. COMMON BLOCK FDATA IS  
C*** REQUIRED. S IS THE VALUE OF PSI AND G IS THE VALUE OF GAMMA.  
C*** R IS THE DISTANCE FROM PELLET CENTER, RELATIVE TO PELLET RADIUS.  
C***  
C*** STATEMENT NO.      PURPOSE OF STATEMENT  
C*** -----  
C***  
C***          10      DIFFERENTIAL EQUATION FOR MACH NUMBER, M  
C***          20      DIFFERENTIAL EQUATION FOR RELATIVE DENSITY, P  
C***          30      TEST FOR CONTINUITY AT M = 1.D+00  
C***          40      EQUATION FOR G(T,H)  
C***  
C***  
C*****  
ISN 0002    FUNCTION F(L,X,Y,J) FUN18  
ISN 0003      I=ICIT REAL*8(A-H,Q-Z) FUN19  
ISN 0004      REAL*8 M FUN20  
ISN 0005      COMMON/FDATA/S,G,T,A,B,C,D FUN21  
ISN 0006      COMMON/HAMCAT/CRIT,IFRKG,LG FUN22  
ISN 0007      DIMENSION X(1),Y(1005,1) FUN23  
ISN 0008      CAPX(T,RA)=(1.D0-(RA)**2)**2*T FUN24  
ISN 0009      IF(LG.EQ.4)GO TO 40 FUN25  
ISN 0010      A=1.D+00 FUN26  
ISN 0011      B=2.D+00 FUN27  
ISN 0012      M=Y(J,1) FUN28  
ISN 0013      P=Y(J,2) FUN29  
ISN 0014      R=X(J) FUN30  
ISN 0015      IF(L-1)1,10,20 FUN31  
ISN 0016      1 WRITE(6,6)J,L FUN32  
ISN 0017      6 FORMAT(1H0,2X,'...ERROR...AT STEP = ',I3,', DIFF EQ CALLED = ',I3) FUN33  
ISN 0018      F=0.D+00 FUN34  
ISN 0019      RETURN FUN35  
ISN 0020      10 F=(S*(P*M*R*R)**3*(A+G*M*M)-M/R*(B+(G-A)*M*M))/(A-M*M) FUN36  
ISN 0021      RETURN FUN37  
ISN 0022      20 CONTINUE FUN38  
ISN 0023      IF(L-3)25,30,1 FUN39  
ISN 0024      25 F=B*P*M*M/(A-M*M)*(A/R-S*(P*R*R)**3) FUN40  
ISN 0025      RETURN FUN41  
ISN 0026      30 CONTINUE FUN42  
ISN 0027      F=(S*(P*M*R*R)**3*(A+G*M*M)-M/R*(B+(G-A)*M*M)) FUN43  
ISN 0028      RETURN FUN44  
ISN 0029      40 CONTINUE FUN45  
ISN 0030      RA=X(J) FUN46  
ISN 0031      IF(RA.GT.9.99999999D-01)GO TO 45 FUN47  
ISN 0032      DNUM=CAPX(T,RA)+B*CAPX(T,RA)**2+C/CAPX(T,RA)+D FUN48  
ISN 0033      DEND=(1.D0/CAPX(T,RA)+E-4.D0*C/CAPX(T,RA)**3)**(1.D0/3.D0) FUN49  
ISN 0034      F=DNUM/DEND*A FUN50  
ISN 0035      RETURN FUN51  
ISN 0036      45 CONTINUE FUN52  
ISN 0037      F=0.D0 FUN53  
ISN 0038      RETURN FUN54  
ISN 0039      END FUN55  
ISN 0040
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C*****GTH01
C***GTH02
C***GTH03
C*** THE PURPOSE OF THIS FUNCTION IS TO AID IN THE CALCULATION OF THE ***GTH04
C*** REQUIRED PELLET VELOCITY BY CORRECTING FOR TEMPERATURE AND ***GTH05
C*** DENSITY DISTRIBUTION EFFECTS. ***GTH06
C***GTH07
C***GTH08
C*****GTH09
      FUNCTION G(T,H)GTH10
      IMPLICIT REAL*8(A-D,F-H,G-Z)GTH11
      REAL*8 K1(1),K2(1),K3(1),K4(1)GTH12
      COMMON/HAMDAT/CRIT,IFRKG,LGTH13
      COMMON X(1005),Y(1005,2),E(2),XINITGTH14
      M=1005GTH15
      NC=1GTH16
      L=4GTH17
      X(1)=XINITGTH18
      Y(1,1)=0.00GTH19
      NSAV=(1.00-X(1))/H+1.100GTH20
50 CONTINUEGTH21
      N=NSAVGTH22
      IF(NSAV.GT.1000)N=1000GTH23
      DO 100 J=2,N,1GTH24
      CALL RKGSQ(J,X,Y,NC,H,E,K1,K2,K3,K4,M)GTH25
100 CONTINUEGTH26
      IF(NSAV.LE.1000)GO TO 150GTH27
      NSAV=NSAV-999GTH28
      X(1)=X(N)GTH29
      Y(1,1)=Y(N,1)GTH30
      GO TO 50GTH31
150 CONTINUEGTH32
      HNEW=1.00-X(N)GTH33
      N=N+1GTH34
      CALL RKGSQ(N,X,Y,NC,HNEW,E,K1,K2,K3,K4,M)GTH35
      G=Y(N,1)GTH36
      RETURNGTH37
      ENDGTH38

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OS/360 FORTRAN H

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C*****PC005
C***PC010
C***PC015
C*** THE PURPOSE OF HAMMIN IS TO SOLVE A SET OF SIMULTANEOUS FIRST PC020
C*** ORDER FUNCTIONALS OF Y WITH RESPECT TO X USING HAMMING'S FIFTH PC025
C*** ORDER PREDICTOR-CORRECTOR. FOURTH ORDER RUNGA-KUTTA-GILL PC030
C*** PROVIDES STARTING VALUES WHICH ARE ITERATED UPON BY THREE OTHER PC035
C*** INTEGRATION FORMULAE UNTIL A CONSISTENT STARTER IS OBTAINED. PC040
C***PC045
C***PC050
C*** X0 = INITIAL VALUE OF X. PC055
C***PC060
C*** Y0 = ARRAY OF INITIAL Y-VALUES. PC065
C***PC070
C*** NC = NUMBER OF COUPLED EQUATIONS. PC075
C***PC080
C*** H = INTEGRATION INCREMENT PC085
C***PC090
C*** N = NUMBER OF INTEGRATIONS TO BE PERFORMED. PC095
C***PC100
C*** A1,A2 = PREDICTOR-CORRECTOR PARAMETERS TO BE CHOSEN BY USER. PC105
C***PC110
C*** ERR = CONVERGENCE PARAMETER FOR STARTER. PC115
C***PC120
C*** X = ARRAY OF X-VALUES PC125
C***PC130
C*** Y = 2-DIMENSIONAL ARRAY OF Y-VALUES. PC135
C***PC140
C*** IER = ERROR PARAMETER. IF IER=0, INTEGRAL INSIDE OF ASYMPTOTES. PC145
C*** IF IER=1, STARTER FAILED TO CONVERGE. PC150
C*** IF IER=2, INTEGRAL DIVERGES. PC155
C***PC160
C*** YC,W1,W2,W3,W4,E = WORK ARRAYS OF LENGTH NC. PC165
C***PC170
C*** YP,ER = WORK ARRAY OF MINIMUM SIZE (4,NC). PC175
C***PC180
C*** ND = DIMENSION SIZE OF Y IN MAIN PROGRAM. PC185
C***PC190
C*** TO GET DOUBLE PRECISION VERSION, REMOVE C'S FROM PC240 AND PC245. PC195
C***PC200
C***PC205
C*****PC210
ISN 0002 SUBROUTINE HAMMIN(X0,Y0,NC,N,H,A1,A2,ERR,X,Y,IER,YP,YC,W1,W2,W3,W4 PC215
      S,E,ER,ND) PC220
ISN 0003 REAL*8 X0,Y0,H,A1,A2,X,Y,YP,YC,W1,W2,W3,W4,C1,C2,C3,C4,C5,C6,C7,EP PC225
      S,EC,A0,B0,B1,B2,B3,BCN1,BC0,8C1,8C2,CRIT,F PC230
ISN 0004 COMMON/HAMCAT/CRIT,IFRKG,LG333 PC235
ISN 0005 DIMENSION Y0(1),X(1),Y(ND,1),YP( 4,1),YC(1),E(1),W1(1),W2(1),W3(1) PC240
      S,W4(1),ER( 4,1) PC245
ISN 0006 IER=0 PC250
ISN 0007 C1=9.0+00 PC255
ISN 0008 C2=19.0+00 PC260
ISN 0009 C3=5.0+00 PC265
ISN 0010 C4=24.0+00 PC270
ISN 0011 C5=4.0+00 PC275
ISN 0012 C6=3.0+00 PC280

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ISN 0013	C7=8.D+00	PC285
ISN 0014	EP=(2.51D+02-C2*A1-C7*A2)/6.D+00	PC290
ISN 0015	EC=(-C2+11.D+00*A1-C7*A2)/6.D+00	PC295
ISN 0016	A0=1.D+00-A1-A2	PC300
ISN 0017	B0=(55.D+00+C1*A1+C7*A2)/C4	PC305
ISN 0018	B1=(-59.D+00+C2*A1+32.D+00*A2)/C4	PC310
ISN 0019	B2=(37.D+00-C3*A1+C7*A2)/C4	PC315
ISN 0020	B3=(A1-C1)/C4	PC320
ISN 0021	BCN1=(C1-A1)/C4	PC325
ISN 0022	BC0=(C2+13.D+00*A1+C7*A2)/C4	PC330
ISN 0023	BC1=(-C3+13.D+00*A1+32.D+00*A2)/C4	PC335
ISN 0024	BC2=(1.D+00-A1+C7*A2)/C4	PC340
ISN 0025	IF(IFRKG.GE.0)GO TO 155	PC345
ISN 0027	X(1)=X0	PC350
ISN 0028	DO 100 L=1,NC,1	PC355
ISN 0029	Y(1,L)=Y0(L)	PC360
ISN 0030	100 CONTINUE	PC365
ISN 0031	DO 115 J=2,4,1	PC370
ISN 0032	CALL RKGSQ(J,X,Y,NC,H,E,W1,W2,W3,W4,ND)	PC375
ISN 0033	115 CONTINUE	PC380
ISN 0034	JC ECK=0	PC385
ISN 0035	120 CONTINUE	PC390
ISN 0036	JC ECK=JCHECK+1	PC395
ISN 0037	DO 130 L=1,NC,1	PC400
ISN 0038	YP(2,L)=Y(1,L) +H*(C1*F(L,X,Y,1)+C2*F(L,X,Y,2)-C3*F(L,X,Y,3)+F(L,X,Y,4))/C4	PC405
ISN 0039	YP(3,L)=Y(1,L) +H*(F(L,X,Y,1)+C5*F(L,X,Y,2)+F(L,X,Y,3))/C6	PC410
ISN 0040	YP(4,L)=Y(1,L) +H*(F(L,X,Y,1)+C6*(F(L,X,Y,2)+F(L,X,Y,3))+F(L,X,Y,4))/C7*C6	PC415
ISN 0041	DO 125 J=2,4,1	PC420
ISN 0042	ER(J,L)=(Y(J,L)-YP(J,L))/Y(J,L)	PC425
ISN 0043	ER(J,L)=ABS(ER(J,L))	PC430
ISN 0044	125 CONTINUE	PC435
ISN 0045	130 CONTINUE	PC440
ISN 0046	ITEST=0	PC445
ISN 0047	DO 140 L=1,NC,1	PC450
ISN 0048	DO 135 J=2,4,1	PC455
ISN 0049	IF(ER(J,L).GT.ERR)ITEST=1	PC460
ISN 0051	Y(J,L)=YP(J,L)	PC465
ISN 0052	135 CONTINUE	PC470
ISN 0053	140 CONTINUE	PC475
ISN 0054	IF(JCHECK.GE.500)GO TO 145	PC480
ISN 0056	IF(ITEST.EQ.1)GO TO 120	PC485
ISN 0058	145 DO 150 L=1,NC,1	PC490
ISN 0059	YP(1,L)=0.D+00	PC495
ISN 0060	YC(L)=0.D+00	PC500
ISN 0061	150 CONTINUE	PC505
ISN 0062	155 CONTINUE	PC510
ISN 0063	DO 300 J=4,N,1	PC515
ISN 0064	X(J+1)=X(J)+H	PC520
ISN 0065	DO 175 L=1,NC,1	PC525
ISN 0066	YP(2,L)=YP(1,L)	PC530
ISN 0067	175 CONTINUE	PC535
ISN 0068	DO 200 L=1,NC,1	PC540
ISN 0069	YP(1,L)=A0*Y(J,L)+A1*Y(J-1,L)+A2*Y(J-2,L)+H*(B0*F(L,X,Y,J)+B1*F(L,X,Y,J-1)+B2*F(L,X,Y,J-2)+B3*F(L,X,Y,J-3))	PC545
ISN 0070	200 CONTINUE	PC550
ISN 0071	DO 225 L=1,NC,1	PC555
		PC560
		PC565
		PC570

ISN 0072	Y(J+1,L)=YP(1,L)-(EP/(EP-EC))*(YP(2,L)-YC(L))	PC575
ISN 0073	225 CONTINUE	PC580
ISN 0074	DO 250 L=1,NC,1	PC585
ISN 0075	YC(L)=A0*Y(J,L)+A1*Y(J-1,L)+A2*Y(J-2,L)+H*(BCN1*F(L,X,Y,J+1)+BC0*F	PC590
	S(L,X,Y,J)+BC1*F(L,X,Y,J-1)+BC2*F(L,X,Y,J-2))	PC595
ISN 0076	250 CONTINUE	PC600
ISN 0077	DO 275 L=1,NC,1	PC605
ISN 0078	Y(J+1,L)=YC(L)-(EC/(EP-EC))*(YP(1,L)-YC(L))	PC610
ISN 0079	275 CONTINUE	PC615
ISN 0080	IF(J-N)285,300,400	PC620
ISN 0081	285 CONTINUE	PC625
ISN 0082	300 CONTINUE	PC635
ISN 0083	IER=ITEST	PC640
ISN 0084	400 N=N+1	PC645
ISN 0085	RETURN	PC650
ISN 0086	END	PC670

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ISN 0027	140	CONTINUE	RK 57
ISN 0028		DO 150 L=1,N,1	RK 58
ISN 0029		K3(L)=F(L,X,Y,J)*H	RK 59
ISN 0030	150	CONTINUE	RK 60
ISN 0031		DO 160 L=1,N,1	RK 61
ISN 0032		Y(J,L)=Y(J,L)+AMS*K1(L)-K2(L)+ADD*K3(L)	RK 62
ISN 0033	160	CONTINUE	RK 63
ISN 0034		X(J)=X(J-1)+H	RK 64
ISN 0035		DO 170 L=1,N,1	RK 65
ISN 0036		K4(L)=F(L,X,Y,J)*H	RK 66
ISN 0037	170	CONTINUE	RK 67
ISN 0038		DO 180 L=1,N,1	RK 68
ISN 0039		Y(J,L)=Y(J-1,L)+(K1(L)+AM*K2(L)+AD*K3(L)+K4(L))/AX	RK 69
	C	E(L)=(K2(L)-K3(L))/(K2(L)-K1(L))	RK 70
	C	E(L)=ABS(E(L))	RK 71
ISN 0040	180	CONTINUE	RK 72
ISN 0041		RETURN	RK 73
ISN 0042		END	RK 74

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CONCEPTUAL DESIGN OF A COMMERCIAL-
TOKAMAK-HYBRID-REACTOR FUELING SYSTEM

by

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AN ABSTRACT OF A MASTER'S THESIS

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1980

ABSTRACT

A conceptual design of a fuel injection system for CTHR (Commercial Tokamak Hybrid Reactor) is discussed. Initially, relative merits of the cold-fueling concept are compared with those of the hot-fueling concept; that is, fueling where the electron temperature is below 1 eV is compared with fueling where the electron temperature exceeds 100 eV. It is concluded that cold fueling seems to be somewhat more free of drawbacks than hot fueling. Possible implementation of the cold-fueling concept is exploited via frozen-pellet injection. Several methods of achieving frozen-pellet injection are discussed and the light-gas-gun approach is chosen from these possibilities. A modified version of the ORNL Neutral Gas Shielding Model is used to simulate the pellet injection process. From this simulation, the penetration-depth dependent velocity requirement is determined. Finally, with the velocity requirement known, a gas-pressure requirement for the proposed conceptual design is established. The cryogenic fuel-injection and fuel-handling systems are discussed. A possible way to implement the conceptual device is examined along with the attendant effects on the total system.