

TIME DEPENDENCE OF THE MAGNETIC
FIELD IN A RECTANGULAR TOROID

by

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1.0 INTRODUCTION

The purpose of this thesis is to formulate an equation which describes the time dependence of the flux distribution in a rectangular toroid of ferromagnetic material subject to given boundary conditions and to provide a numerical solution to the equation. This is of interest because it allows one to predict the time necessary to release a no-work magnet in an electro-mechanical system or define the switching time of a toroid used in the core plane of many modern digital computers. An analytical solution also provides a convenient tool for optimizing the many problem variables or studying the effect of changing one or more of the variables on the magnetic performance of a given electromagnetic circuit.

Development of theory and assumptions necessary to derive such an equation for the special case considered, i.e., for the toroid of rectangular cross section as illustrated in Fig. 7, or the electromechanical system as illustrated in Fig. 2, are given in section 2.0. The resulting equation obtained by manipulation of Maxwell's equations was found to resemble the diffusion equation when subject to the assumptions required for the problem of interest. The Hysteretic Diffusion Equation, equation (2-28), expresses the desired relationship for specifying the magnetic field intensity H as a function of position in the core and time. It should be noted that H is written as simply H for convenience and represents $H(x,y,z,t)$.

$$(2-28) \quad \frac{\partial H}{\partial t} = \frac{(C_2 + H)^2}{\sigma C_1 C_2} \left[\frac{\partial^2 H}{\partial x^2} + \frac{\partial^2 H}{\partial y^2} \right]$$

where H is the magnetic field intensity (z component)

σ is the material conductivity

C_1 is a constant specifying material properties

C_2 is a constant specifying material properties.

A discussion of the effects of eddy currents and concepts used to confirm the existence of a flux distribution in the cross section during a transient condition are given to aid the reader in understanding the problem.

The task of evaluating H was accomplished by using the Modified Euler method of numerical integration in which the rectangular cross section was thought of as being divided into a grid. Each grid represented a toroid with a time-variant field intensity. This allowed one to express the Hysteretic Diffusion Equation in its finite difference form and approximate $\partial H / \partial t$ by numerical techniques. Once this quantity was known, Euler's method was used to predict an approximate value for the field intensity of each grid area at a time Δt later, thus an approximate solution was obtained for H as a function of time for each grid area of the cross section. Values of flux density B , were obtained by substitution of H in an equation approximating the B-H relationship for the given ferromagnetic material.

One must remember the numerical process yields only an approximation to the actual values. Accuracy increases as the grid size decreases; however, the process becomes very slow for

very small grid sizes because the maximum time increment allowable to insure convergence of the Modified Euler process is dependent on many of the problem variables as shown by equation (3-44).

$$(3-44) \quad 0 < \Delta t < \frac{\sigma C_1 C_2 h^2}{2(C_2 + H_{\max})^2}$$

Consideration of conditions for convergence and accuracy of the process are given in section 3.4.

The last two sections of the thesis describe the program and various input variables needed to execute the process. Each input variable and its function in the program is discussed in section 4.2B. A sample problem with a typical input data set and the corresponding output data is given in section 5.0. Results of the data indicate a flux distribution and flux decay as predicted by sections 2.3 and 2.4.

2.0 DERIVATION OF THE HYSTERETIC DIFFUSION EQUATION FOR A MAGNETIC BOUNDARY VALUE PROBLEM

2.1 Introduction

A problem of interest to many is the solution of the diffusion equation sometimes called the heat equation. It is useful in describing the temperature distribution in an iron bar as a function of time, or in the magnetic case, the time dependence of the flux distribution in a toroid of magnetic material subject to given boundary conditions. It is the latter case which will be considered in detail in this thesis. The solution of this particular boundary value problem is important in predicting performance objectives and analysis of many components in modern digital computers as well as being of use for analytical evaluation of heat transfer characteristics and other problems in many other fields described by the equation. For example, the memory of most modern computers consists of core planes in which rectangular toroids are affixed at the junction of two write windings as illustrated in Fig. 1.

To magnetize the toroid one-half of the current needed must flow in the proper direction through both write wires threaded through the core; thus the core is magnetized in one particular direction which designates a "1" bit and the opposite direction which designates a "0" bit. During a read out of the memory, the direction of magnetization determines the direction of current flow induced in a read winding and in turn determines the

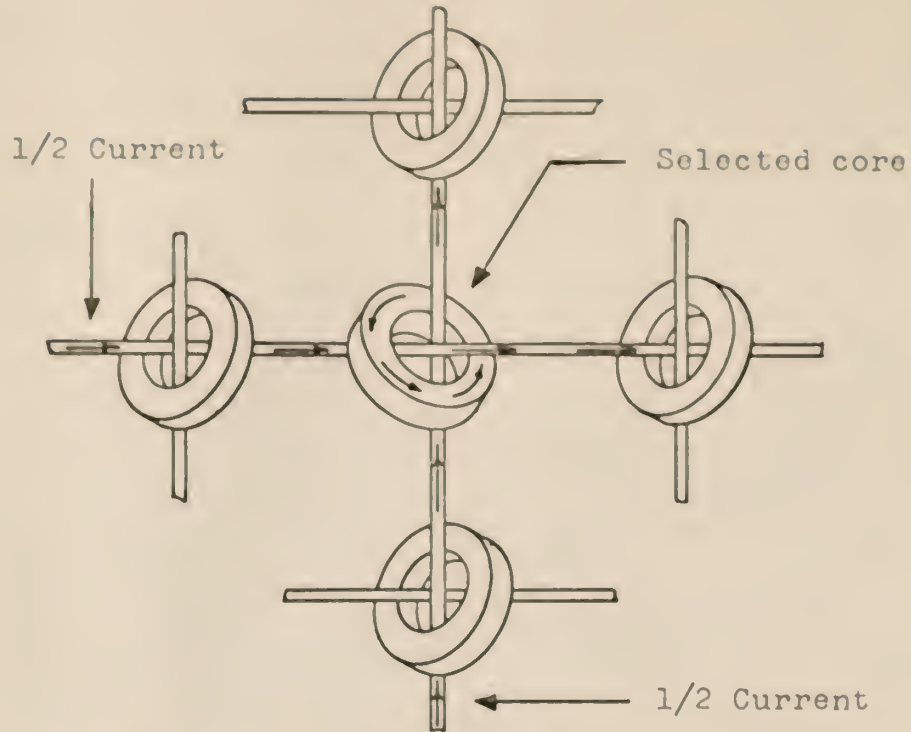


Fig. 1. Section of a memory core plane.

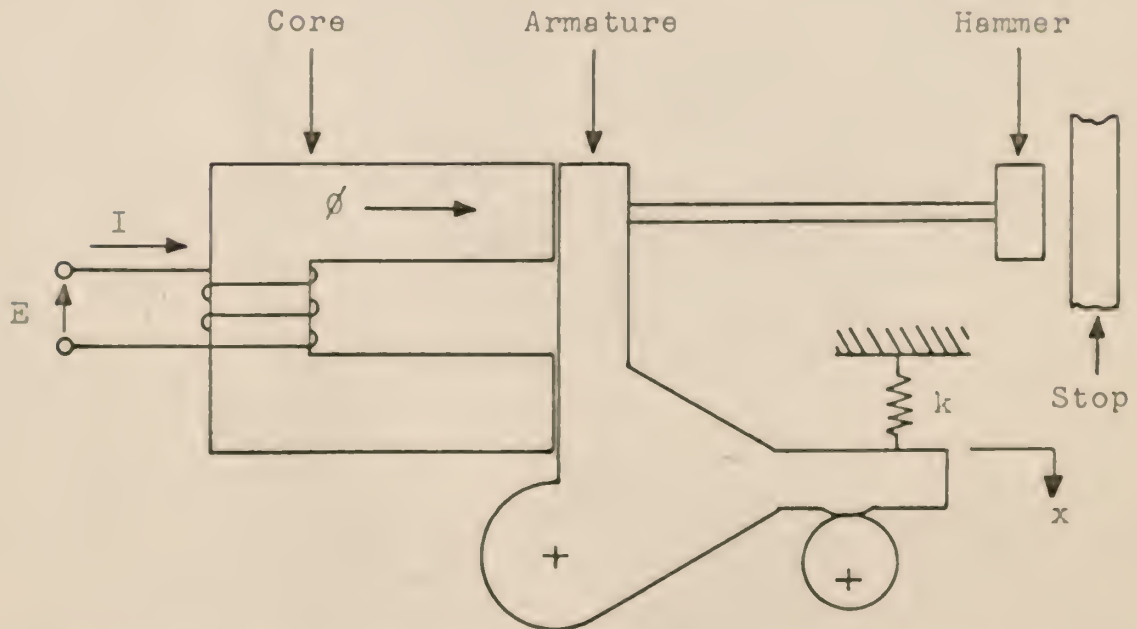


Fig. 2. Electromechanical device (no-work magnet).

presence of a "0" or "1" bit. One can see that it may be desirable to know the time to accomplish flux reversal during the write time. Furthermore, an analytical method of determining the reversal time for various cores of different sizes, shapes, materials, etc., would be convenient since this time could be a significant factor in determining the speed of the computer.

The solution could also provide information useful for improvements in the design of present-day input-output equipment for computation systems. Many high-speed card punches and printers use electromechanical devices in various punching and printing techniques, all of which require elaborate sequential timing and mechanical movement to accomplish the desired result.

Let us consider the electromechanical device in Fig. 2. Suppose a voltage E causes a current I to flow through the coil which produces a flux ϕ sufficient to overcome the force kx and hold the core and armature together. When I is removed (i.e., $E = 0$), the holding force is removed and the spring force allows the hammer to impact the stop. An extension of this principle is used for modern printers which can print at the rate of twelve hundred lines per minute with one hundred twenty characters per line. The time required to release the magnet becomes of prime importance since sequential timing and logic circuits required to release the armature cause a release and hold operation cycle to occur at very high rates of speed. A typical release time might be one millisecond. This same mechanism is used for obtaining punched cards and the same discussion could apply.

Due to effects of mechanical inertia, eddy currents generated by the rapidly changing boundary conditions on the magnetic circuit, and other factors, it may be desirable to provide an analytical solution which would account for changes in the parameters affecting the electromagnetic performance of an electromechanical system as outlined in the previous discussion. Assuming that the junction of the armature and core assembly of Fig. 2 does not provide an additional reluctance to impede the flow of flux across it other than that of the material itself, one can use the results of the following procedure to gain insight on effects that parameter changes produce on the electromagnetic performance.

2.2 Magnetic Circuit Concepts

Magnetic circuit considerations are closely analogous to those of resistive electrical circuits; however, the cause and effect relationship in the magnetic case is nonlinear, i.e., the reluctance of a d-c magnetic circuit depends on the flux in the circuit, while for the d-c electrical case resistance is relatively unaffected by the amount of current.

If one considers a toroid of magnetic material with a coil of wire wound tightly and distributed uniformly around it, a magnetic circuit problem is encountered. (See Fig. 3.)

The voltage E produces a flux ϕ in the magnetic material. The flux lines are perpendicular to a cross section of the toroid and should be uniformly distributed over the cross section in

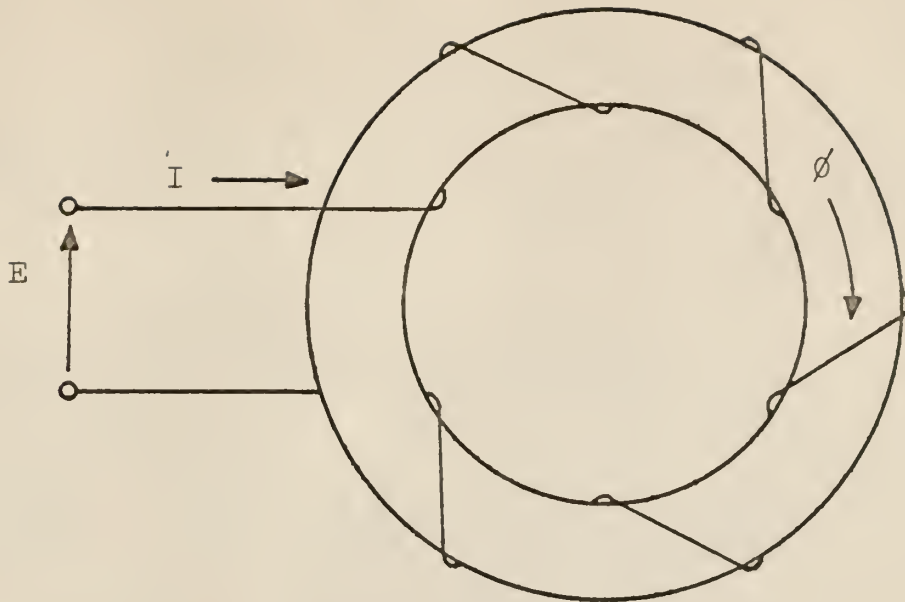


Fig. 3. Toroidal magnetic circuit.

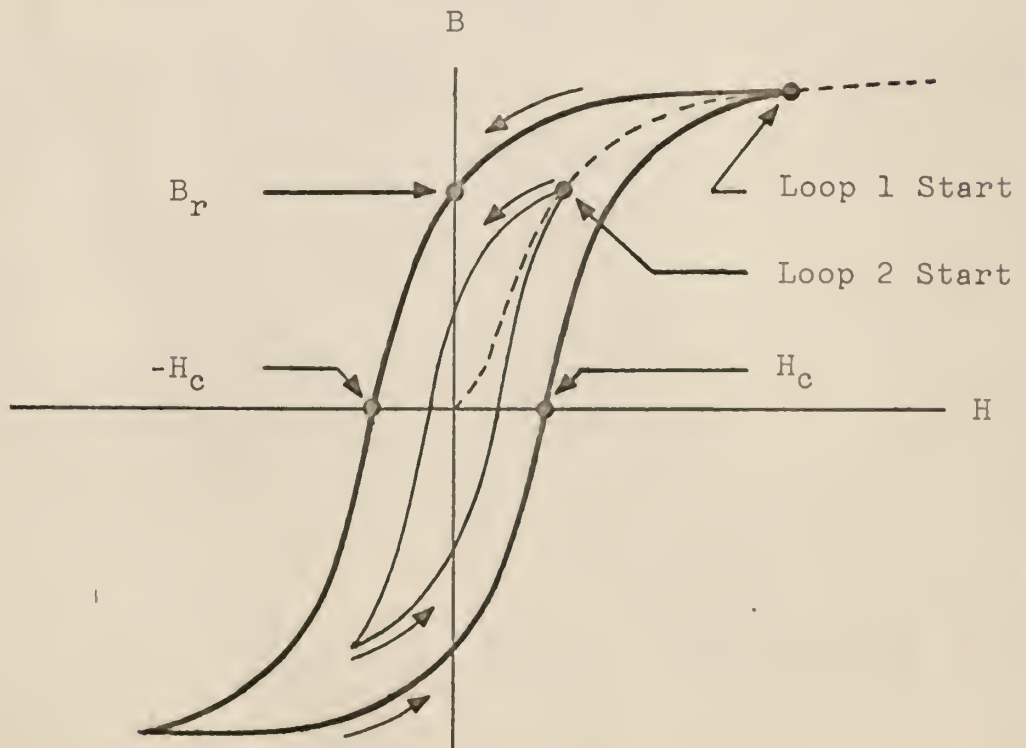


Fig. 4. Hysteresis loops and magnetization curve.

the steady-state case. To insure uniform boundary conditions produced by the current I , assume the cross-sectional diameter of the ring is very small when compared to the inside and outside diameters of the ring. Under these conditions the following relationships hold:

$$(2-1) \quad \mathcal{F} = NI = Hl = R\phi$$

$$(2-2) \quad B = \mu H = \phi/A$$

$$(2-3) \quad R = \mu l/A$$

$$(2-4) \quad \text{Force (F)} = B^2 A / 2\mu_0$$

The above also assumes that no leakage occurs and that one is only considering the steady-state solution; however, the transient operation is of major importance when an electromagnet or rectangular toroid is used in the applications as outlined in section 2.1.

If one desires to know the value of the flux at several discrete points in time during a build-up or decay cycle, one must use a different method of analysis than presented previously. During the transient state a relationship exists between flux and magnetic field intensity as illustrated by the hysteresis curve of Fig. 4. This curve specifies the magnetic properties of the ferromagnetic material.

One can also see that during a flux build-up or decay residual magnetism B_r , and coercive force H_c , depend on the original values of the flux density B and magnetic field intensity H . A typical ferrite material in the memory of a computer exhibits what is known as a square-loop property. This means that

the material's B-H characteristic has a square or rectangular hysteresis loop, which allows an immediate change in the direction of magnetization once the proper value of H is present. In all cases the value of B for a corresponding H follows the upper curve if in the decay portion of the cycle and follows the lower curves if in the build-up portion of the cycle. Furthermore, during either build-up or decay, a flux distribution exists across the cross section of the magnetic core due to eddy currents generated from the changing flux; i.e., the outside of the core reaches the new value of flux density immediately while the center of the core remains at the original value and gradually attains the same value specified by the boundary in the steady-state condition. A more complete discussion of the eddy-current problem and its relation to the flux distribution is given in the next section.

2.3 Eddy Currents

Eddy currents are generated by a change of flux in the magnetic material of a core as shown in Fig. 3, and become a very important factor in the analysis of the transient behavior of high-speed electromagnetic circuits. These currents affect the change in flux by tending to produce an opposing flux, thus changing the flux distribution across the pole face in addition to delaying flux build-up or decay. Flux build-up is not affected nearly so significantly as flux decay because the eddy-current opposition flux is a small part of the total flux

applied. Since a no-work magnet uses flux decay to accomplish its function, the eddy currents play a large part in controlling the release operation; thus they can limit the speed of operation for many of the electromechanical systems used in modern input-output machines of digital computation systems.

One can visualize the effect of eddy currents on the flux distribution by considering a cylindrical core to be made up of concentric shells, each shell being a hollow cylinder of differential thickness. Each shell constitutes a short-circuited turn enclosing part of the core flux. The outside shells link all the flux while the inner shells link only a small part of it. A voltage is induced in the shells upon a change of flux. The magnitude of the induced emf is determined by the amount of flux linked by the coil; thus the outer shells have larger emf's and eddy currents than the inner shells. Induced currents apply a magnetomotive force to the part of the core that lies within it; thus the center of the core is subject to the magnetomotive force of all eddy currents while the surface is subject to none. This accounts for the distribution present in the transient state. Induced currents tend to oppose any change in flux; thus they will tend to sustain a decreasing field and oppose an increasing field. It can be seen that the flux distribution is altered by the eddy currents and is not only a function of time but also of core radius. An approximate flux distribution with radius is shown in Figs. 5 and 6. Figure 5 illustrates the effect of eddy currents during the transient flux build-up condition while Fig. 6 illustrates the effect of eddy currents

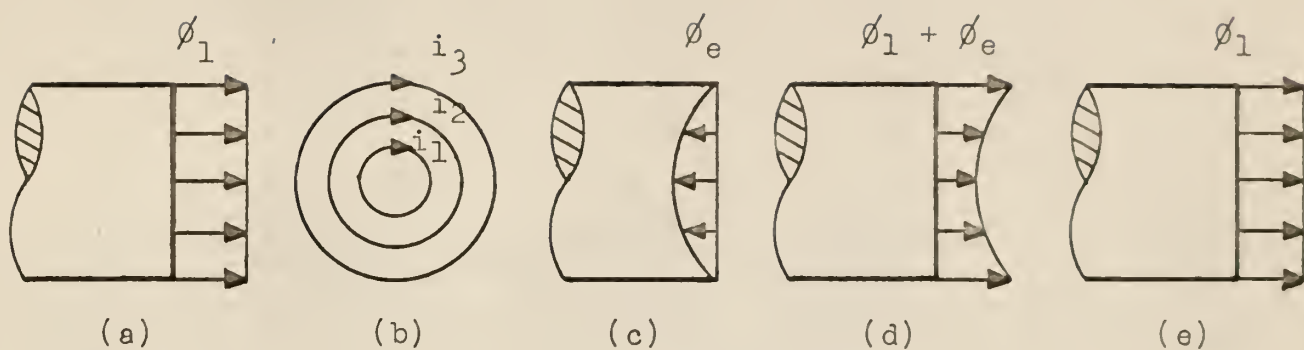


Fig. 5. Effect of eddy currents on flux buildup in a ferromagnetic material.

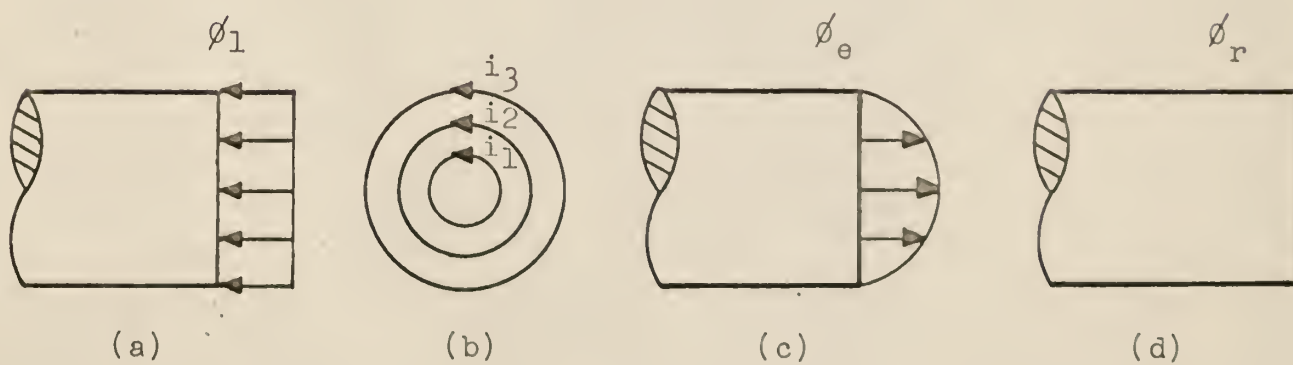


Fig. 6. Effect of eddy currents on flux decay in a ferromagnetic material.

during flux decay.

Let us represent the applied field intensity with vectors in the direction of the applied field. We may then illustrate the flux at $t = 0^+$ by Fig. 5a. Once the field is applied, eddy currents i_1, i_2, i_3, \dots , are set up in the core as illustrated by Fig. 5b. These eddy currents induce an opposing flux ϕ_e shown in Fig. 5c. The resultant flux at $t = 0^+$ is then $\phi_1 + \phi_e$ and is illustrated by Fig. 5d. Since the eddy currents depend on a changing flux they will decrease as time increases, thus ϕ_e will approach zero in the steady state and the steady-state flux will be equal to the applied flux ϕ_1 .

Upon release the same phenomenon occurs; however, the core is originally in the magnetized state and eddy currents will be induced opposite in direction to that of the previous example. This follows from Lenz's Law. The only flux present at $t = 0^+$ is that induced by the eddy currents as shown in Fig. 6b. As time increases, ϕ_e decreases until the flux returns to zero. Actually it will return to a residual value ϕ_r .

2.4 Flux Distribution in the Rectangular Toroid

With the preceding discussion of eddy-current phenomena in mind, one can readily visualize a flux distribution existing in a toroid of rectangular cross section. Furthermore, one knows that the flux distribution is changing with respect to time and values of flux at any point within the core depend on the distance from the outer edge of the core at which the boundary

condition is applied, the boundary condition itself, the material of which the toroid is made, and time. If one considers a rectangular cross section to be divided into many small squares of nearly infinitesimal area, the value of flux density in a given square within the core is a function of time and is different for each area. Time dependence of the flux density may be obtained if one can find a relationship sufficient to specify the value of magnetic field intensity for this area as a function of time. After division into small areas one may then consider each area separately as a rectangular toroid with uniform flux density provided the area is very small in comparison to the total cross-section area. Values of H , B , and ϕ for each area can then be calculated for each value of time according to the relationships given in section 2.7. One might note that once H is determined, it is a simple matter to evaluate B from the hysteresis loop. Multiplication of B and the area corresponding to the value of B calculated will yield a value of flux for that particular area. It then becomes possible to calculate total values of H , B , and ϕ for the entire cross section. If the above calculations are made for each value of time in the transient state, curves relating the time dependence of the flux distribution and total flux can be computed for various boundary conditions and the effect of parameter changes can be analyzed theoretically. Since no literature was found that expressed the magnetic field intensity as a function of position and time for the case in question, equation (2-28) was derived. It will be referred to as the "Hysteretic Diffusion Equation" and its

derivation is included in the next section.

2.5 Derivation of the Hysteretic Diffusion Equation

The derivation presented requires only basic electromagnetic theory and simply applies Maxwell's Equations to the special case being considered. To refresh the reader's memory, the time variant Maxwell Equations are listed below in both point and integral form. Overbarred symbols designate vector quantities in the cartesian co-ordinate system. This system was the most convenient for the rectangular cross section being considered. A co-ordinate transformation of the resulting equation could be used if different cross-section shapes are studied.

Maxwell's Equations

(Point form)

(Integral form)

(2-1) A. $\bar{\nabla} \times \bar{E} = -\partial\bar{B}/\partial t$	B. $\oint \bar{E} \cdot d\bar{l} = (\partial/\partial t) \int_s \bar{B} \cdot d\bar{S}$
(2-2) A. $\bar{\nabla} \times \bar{H} = \bar{i} + \partial\bar{D}/\partial t$	B. $\oint \bar{H} \cdot d\bar{l} = \int_s (\bar{i} + \partial\bar{D}/\partial t) \cdot d\bar{S}$
(2-3) A. $\bar{\nabla} \cdot \bar{B} = 0$	B. $\oint \bar{B} \cdot d\bar{S} = 0$
(2-4) A. $\bar{\nabla} \cdot \bar{D} = \rho$	B. $\oint \bar{D} \cdot d\bar{S} = \int_v \rho \cdot dv$

Consider a toroid with a very large inside diameter in comparison to the bar diameter or an infinite bar of ferromagnetic material which undergoes a change of magnetomotive force on its boundary due to a change of the exciting current I supplied from a voltage E . (See Fig. 7.)

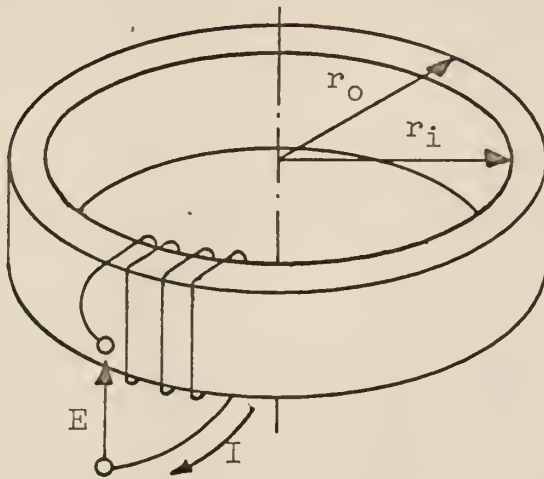


Fig. 7a. Rectangular toroid.

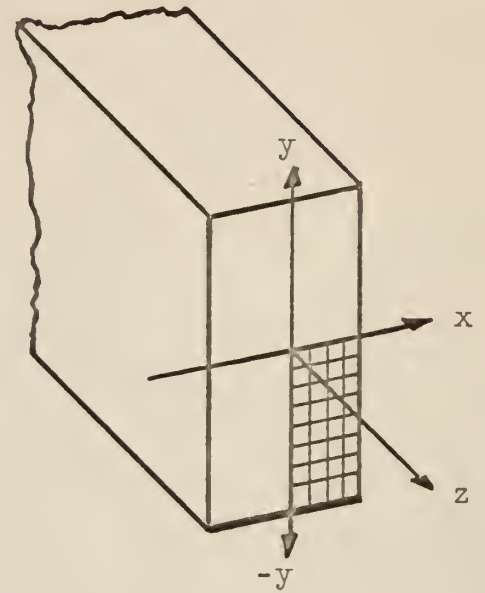


Fig. 7b. Cross section of a rectangular toroid.

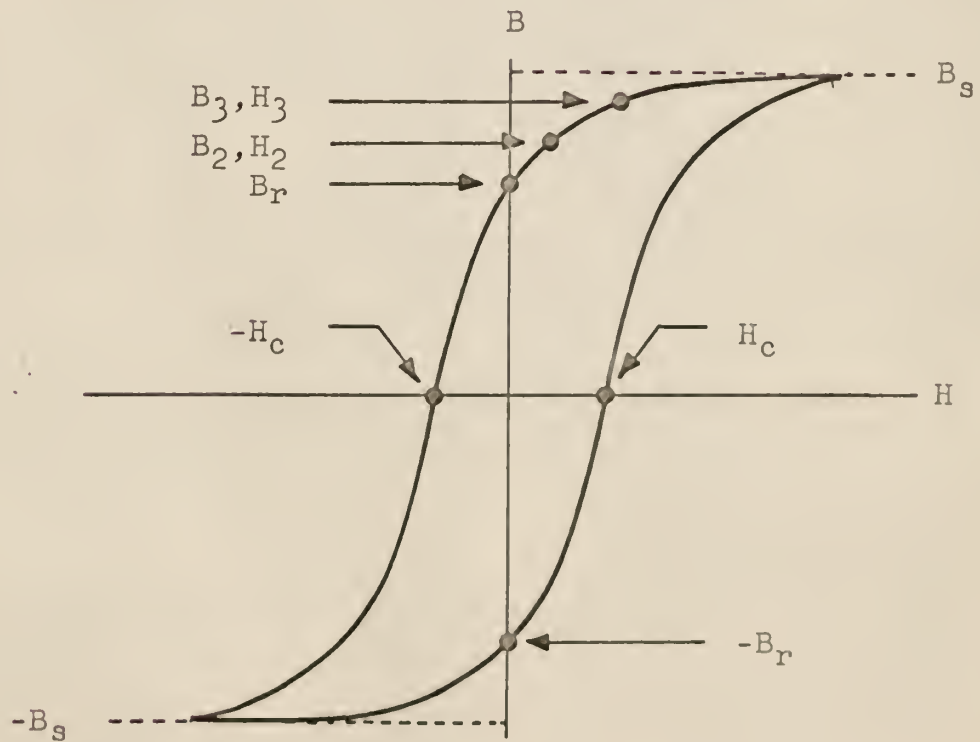


Fig. 8. Hysteresis loop.

It is known with some certainty that a flux distribution exists within the core due to eddy currents while the magnetization of the core is in the transient state, thus steady-state techniques cannot be applied; however, Maxwell's Equations hold in the transient state as well as for the steady state, thus an equation relating the effects of \bar{H} as it varies with x , y , z , and t can be derived.

First we may assume that an mmf has been applied. When applied, a magnetic field \bar{H} , an electric field \bar{E} , and a flux density \bar{B} , exist in the toroid, and all are governed by Maxwell's Equations. If the electric field current density \bar{D} , is neglected, Maxwell's second equation becomes

$$(2-5) \quad \bar{\nabla} \times \bar{H} = \bar{i}$$

Since the current flow in a conductor is in the direction of the applied electric field and perpendicular to an incremental surface \bar{dS} , the total current I is obtained by integrating the current density \bar{i} over the surface; i.e.,

$$(2-6) \quad I = \int_s \bar{i} \cdot \bar{dS} = \int_s (\bar{\nabla} \times \bar{H}) \cdot \bar{dS}$$

These currents are the eddy currents present in the transient state and can be represented by the scalar multiplication

$$(2-7) \quad \bar{i} = \sigma \bar{E}$$

Thus the relationship between the magnetic field intensity \bar{H} , electric field intensity \bar{E} , material conductivity σ , and the conduction current density \bar{i} , is determined since

$$(2-8) \quad \mathbf{I} = \int_S \bar{\mathbf{i}} \cdot d\bar{\mathbf{S}} = \int_S (\nabla \times \bar{\mathbf{H}}) \cdot d\bar{\mathbf{S}} = \sigma \int_S \bar{\mathbf{E}} \cdot d\bar{\mathbf{S}}$$

Therefore one can now show that

$$(2-9) \quad \nabla \times \bar{\mathbf{H}} = \sigma \bar{\mathbf{E}} = \bar{\mathbf{i}}$$

A relationship between the flux density $\bar{\mathbf{B}}$ and magnetic field intensity $\bar{\mathbf{H}}$ can be found by combining equation (2-5) and Maxwell's first equation, equation (2-10).

$$(2-10) \quad \nabla \times \bar{\mathbf{E}} = - \partial \bar{\mathbf{B}} / \partial t$$

Combining equation (2-9) with equation (2-10) results in

$$(2-11) \quad \nabla \times \nabla \times \bar{\mathbf{H}} = - \partial \bar{\mathbf{B}} / \partial t$$

In the general case $\bar{\mathbf{B}}$, $\bar{\mathbf{H}}$, and $\bar{\mathbf{E}}$ are functions of x , y , z , and t and should be written $\bar{\mathbf{B}}(x,y,z,t)$, $\bar{\mathbf{H}}(x,y,z,t)$, and $\bar{\mathbf{E}}(x,y,z,t)$. It is also known that a nonlinear relationship exists between $\bar{\mathbf{B}}$ and $\bar{\mathbf{H}}$; thus one can say that

$$(2-12) \quad \bar{\mathbf{B}} = f(\bar{\mathbf{H}})$$

$$(2-13) \quad \partial \bar{\mathbf{B}} / \partial t = (\partial \bar{\mathbf{B}} / \partial \bar{\mathbf{H}}) (\partial \bar{\mathbf{H}} / \partial t)$$

Further examination of the functions $\bar{\mathbf{B}}(x,y,z,t)$ and $\bar{\mathbf{H}}(x,y,z,t)$ illustrates the fact that total derivatives may be used since

$$(2-14) \quad \frac{d\bar{\mathbf{B}}}{dt} = \frac{\partial \bar{\mathbf{B}}}{\partial x} \frac{dx}{dt} + \frac{\partial \bar{\mathbf{B}}}{\partial y} \frac{dy}{dt} + \frac{\partial \bar{\mathbf{B}}}{\partial z} \frac{dz}{dt} + \frac{\partial \bar{\mathbf{B}}}{\partial t}$$

$$(2-15) \quad \frac{d\bar{H}}{dt} = \frac{\partial \bar{H}}{\partial x} \frac{dx}{dt} + \frac{\partial \bar{H}}{\partial y} \frac{dy}{dt} + \frac{\partial \bar{H}}{\partial z} \frac{dz}{dt} + \frac{\partial \bar{H}}{\partial t}$$

and the derivatives of distance with respect to time are equal to zero if the toroid is stationary with respect to the exciting mmf. Thus we have

$$\frac{\partial \bar{B}}{\partial t} = d\bar{B}/dt$$

$$\frac{\partial \bar{H}}{\partial t} = d\bar{H}/dt$$

$$(2-16) \quad d\bar{B}/dt = (d\bar{B}/d\bar{H})(d\bar{H}/dt) = U(\bar{H})d\bar{H}/dt$$

$$(2-17) \quad \frac{\partial \bar{B}}{\partial \bar{H}} = d\bar{B}/d\bar{H} = U(\bar{H})$$

The quantity $U(\bar{H})$ represents the relationship between \bar{B} and \bar{H} as illustrated by the slope of the hysteresis loop of Fig. 8. Combining equations (2-17), (2-13), and (2-11), we find that the following equation exists.

$$(2-18) \quad \frac{\partial \bar{H}}{\partial t} = -(1/\sigma U(\bar{H})) (\bar{\nabla} \times \bar{\nabla} \times \bar{H})$$

Solution of equation (2-18) will provide an expression for \bar{H} as a function of the position within the rectangular toroid and time; however, a solution is very hard to obtain. If one assumes the only component of field intensity present is along the z axis (see Fig. 7) and that this value of field intensity \bar{H} is the same value at all points on the z axis, equation (2-18) reduces to a form generally recognized as the diffusion equation in two dimensions; i.e.,

$$(2-19) \quad \frac{\partial \bar{H}}{\partial t} = \bar{\nabla}^2 \bar{H} / \sigma U(\bar{H})$$

The above assumptions require that

$$H_x = H_y = \frac{\partial H_z}{\partial z} = 0$$

where in general \bar{H} is given by the vector equation

$$\bar{H} = H_x \bar{i} + H_y \bar{j} + H_z \bar{k} .$$

Thus equation (2-18) may be written as

$$(2-20) \quad \frac{\partial \bar{H}}{\partial t} = - \frac{1}{\sigma U(H)} (\bar{\nabla} \times \bar{\nabla} \times H_z \bar{k})$$

Expanding the expression $(\bar{\nabla} \times \bar{\nabla} \times \bar{H})$ of equation (2-18) with $\bar{H} = H_z \bar{k}$ and noting that

$$\bar{\nabla} \cdot \bar{H} = \left[\frac{\partial}{\partial x} \bar{i} + \frac{\partial}{\partial y} \bar{j} + \frac{\partial}{\partial z} \bar{k} \right] \cdot (H_z \bar{k}) = \frac{\partial H_z}{\partial z} = 0$$

we may evaluate the expression $(\bar{\nabla} \times \bar{\nabla} \times \bar{H})$, i.e.,

$$\begin{aligned} \bar{\nabla} \times \bar{\nabla} \times \bar{H} &= \bar{\nabla} (\bar{\nabla} \cdot \bar{H}) - (\bar{\nabla} \cdot \bar{\nabla}) \bar{H} \\ \bar{\nabla} \times \bar{\nabla} \times \bar{H} &= \bar{\nabla} (0) - \bar{\nabla}^2 \bar{H} \\ (2-21) \quad \bar{\nabla} \times \bar{\nabla} \times \bar{H} &= - \bar{\nabla}^2 H_z \bar{k} \end{aligned}$$

Thus equation (2-18) is now reduced to the diffusion equation for the two-dimensional case, i.e.,

$$(2-22) \quad \frac{\partial H}{\partial t} = \frac{1}{\sigma U(H)} \left[\frac{\partial^2 H_z}{\partial x^2} + \frac{\partial^2 H_z}{\partial y^2} \right]$$

An approximate solution of the above equation may be accomplished through use of numerical integration techniques once a suitable approximation for $U(H)$ is determined. The remainder of

the thesis will deal with the special case outlined above. The vector notation may be deleted since the only component present is H_z and the equation was solved only for this particular case.

The only remaining unknown is the functional relationship existing between B and H . One knows this relationship is expressed by the hysteresis curves of various magnetic materials; thus an equation which approximates the specific curve of interest would be desirable. Development of such an equation is undertaken in the next section.

2.6 Approximation of the B-H Curve for $U(H)$

Because of the flux distribution within the toroidal core, each small area as defined in section 2.4, will have a different value of field intensity and flux density at any given time in the transient condition; thus an approximation would allow ϕ to be calculated directly for each area once H for that area is known. This may be accomplished if an equation can be found which approximates the particular B-H curve of interest. Since this particular curve is determined by the initial values of H for a given material, an approximation of only this B-H curve would be sufficient to compute B once H is known.

A modified Froelich approximation equation as given by equation (2-23) can be used for generation of one-fourth of a given hysteresis loop.

$$(2-23) \quad B = \frac{C_1 H}{C_2 + H} + B_r$$

where B is the flux density

B_r is the residual flux density

H is the field intensity

C_1 constant specifying material properties

C_2 constant specifying material properties.

Values of C_1 and C_2 determine the shape of the curve and may be varied to generate an approximation to most hysteresis loops. They may be calculated from B_r and selection of two values, B_1, H_1 , and B_2, H_2 , taken from near the knee of the decay portion of a given hysteresis loop representative of the operating range in a given material. (See Fig. 8.)

Since both points selected must lie on the curve, a simple simultaneous solution of two equations formed by substituting B_1, H_1 , and B_2, H_2 , in equation (2-23) will be sufficient to specify C_1 and C_2 . Performing this operation yields

$$(2-23a) \quad B_1 = \frac{C_1 H_1}{C_2 + H_1} + B_r$$

$$(2-23b) \quad B_2 = \frac{C_1 H_2}{C_2 + H_2} + B_r$$

Solving for C_1 and C_2 we have the following relationships:

$$(2-24) \quad C_2 = \frac{H_2 H_3 (B_3 - B_2)}{(B_2 - B_r) H_3 - (B_3 - B_r) H_2}$$

$$(2-25) \quad C_1 = \frac{(B_2 - B_r)(C_2 + H_2)}{H_2}$$

After determining C_1 and C_2 , one-fourth of the hysteresis loop can be plotted using the values calculated from equation (2-22). If there is a reasonable correspondence between the original curve and the curve plotted using the approximation equation, one can assume C_1 and C_2 are sufficiently accurate to specify the B-H relationship. Reflection and translation of this quarter section of the hysteresis loop generates the remaining portions of the loop. Calculation of B for values of $H < -H_c$ for flux decay and $H > H_c$ for flux build-up are considered in section 4.3.

This method of approximation is perhaps rather crude; however, it does suffice in this case. More accurate approximations are no doubt possible although maybe not practical since the magnetic properties of a given material will vary and cause larger errors than those due to the approximation.

2.7 Hysteretic Diffusion Equation

If we now use the approximation equation (2-22) for evaluating $U(H)$, we have

$$(2-26) \quad U(H) = \frac{dB}{dH} = \frac{d}{dH} \left[\frac{C_1 H}{C_2 + H} + B_r \right]$$

Performing the differentiation we have

$$(2-27) \quad U(H) = \frac{C_1 C_2}{(C_2 + H)^2}$$

Combining equations (2-22) and (2-27), we obtain the hysteretic form of the diffusion equation.

$$(2-28) \quad \frac{\partial H_z}{\partial t} = \frac{(C_2 + H_z)^2}{\sigma C_1 C_2} \left[\frac{\partial^2 H_z}{\partial x^2} + \frac{\partial^2 H_z}{\partial y^2} \right]$$

It is this equation whose solution will yield the time dependence of the flux distribution pattern in the rectangular toroid subject to boundary conditions as outlined in section 2.5.

The remainder of this thesis is concerned with the numerical methods for obtaining this solution, implementation of these methods, and demonstration of the numerical process by actually obtaining a solution of equation (2-28) with its boundary values.

3.0 NUMERICAL SOLUTIONS WITH THE MODIFIED EULER METHOD

3.1 Euler Method for Solving an Ordinary Differential Equation

Numerous methods of obtaining solutions to ordinary and partial differential equations have become practical since the advent of the high-speed digital computer. The Modified Euler Method was chosen for this problem because of its simplicity, although other methods may provide a more accurate approximation to the solution or take less computation time. Before showing the application of Euler's Method and its modification to equation (2-28), let us consider the solution of an ordinary differential equation of first order. In symbolic form we may write

$$(3-1) \quad \frac{dH}{dt} = f(t, H) = D$$

The integral of equation (3-1) gives H as a function of time; thus we have $H = F(t)$. A graph of $F(t)$ is a curve in the H - t plane which may be approximated by a series of short line segments provided the curve is continuous; thus we have the approximation relation (see Fig. 9)

$$(3-2) \quad \Delta H = \Delta t \tan \theta = \left[\frac{dH}{dt} \right]_0 \Delta t = D_0 \Delta t$$

$$(3-3) \quad H_1 = H_0 + D_0 \Delta t$$

If we let $\Delta t = h = t_{i+1} - t_i$, we can express the approximation

by

$$(3-4) \quad H_{i+1} = H_i + D_i h \quad (i = 0, 1, 2, \dots, l)$$

This is known as Euler's Method. However, if h is taken small enough to yield sufficient accuracy the method is too slow; if h is larger, inaccuracies will cause the approximation to be unsatisfactory; furthermore, if the graph is monotonic, the approximation will diverge from the actual curve for any value of h chosen. A modification of this method tends to eliminate the divergence.

3.2 Modified Euler Method for Solving an Ordinary Differential Equation

Starting with an initial value H_0 one can approximate H_1 in the same manner as before to yield

$$(3-5)^1 \quad H_1^{(1)} = H_0 + D_0 h$$

Substituting $H_1^{(1)}$ into equation (3-1), one obtains an approximation for dH/dt at the end of the first interval, i.e.,

$$(3-6)^1 \quad D_1^{(1)} = f(t_1, H_1^{(1)})$$

An improved value of H is then found by multiplying h by the average of the values of dH/dt at the ends of the interval t_0 to t_1 ; thus we have

¹Note that D_i represents $(dH/dt)_i$. This notation will be used throughout the remainder of this section.

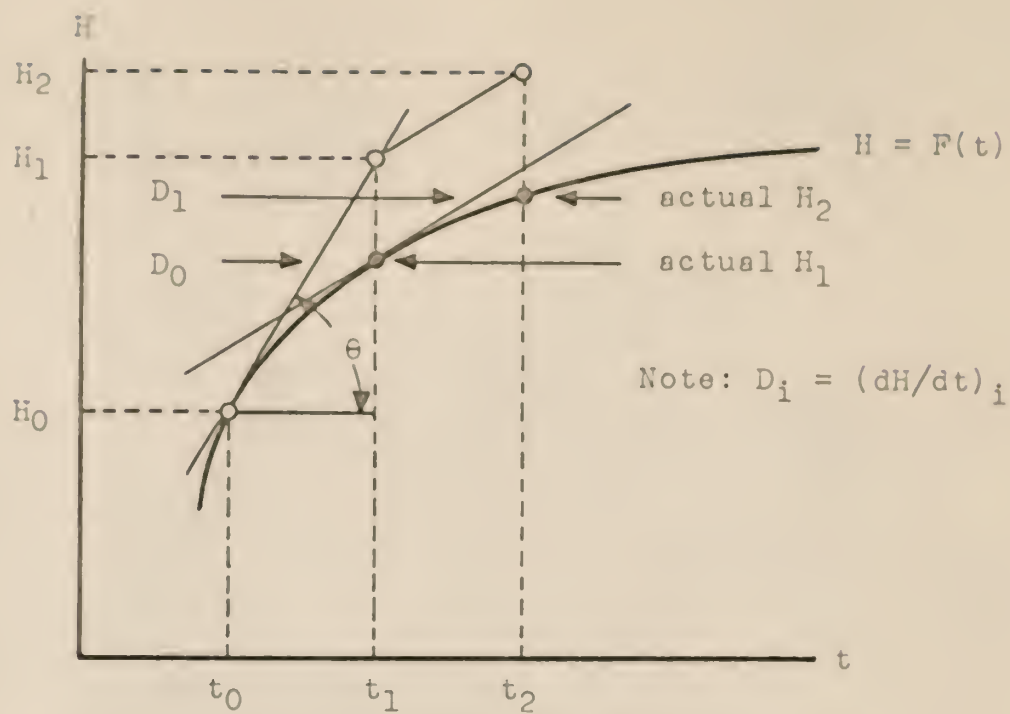


Fig. 9. Approximation with Euler method.

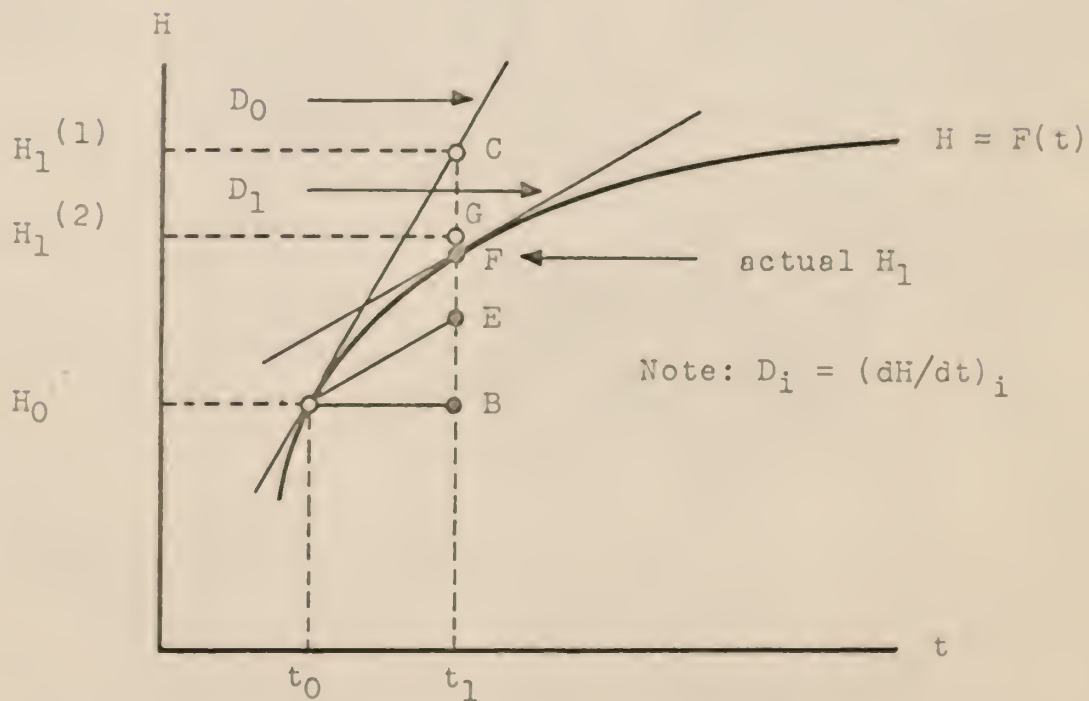


Fig. 10. Approximation with modified Euler method.

$$(3-7) \quad \Delta H = \frac{1}{2} h(D_0 + D_1^{(1)})$$

A more accurate value of H_1 can now be calculated and will be denoted as follows:

$$(3-8)^1 \quad H_1^{(2)} = H_0 + \frac{1}{2} h(D_0 + D_1^{(1)})$$

If we look at Fig. 10, it is evident that this value of $H_1^{(2)}$ is more accurate than $H_1^{(1)}$. $H_1^{(1)}$ is represented by the line $H_0 + BC$ if calculated according to Euler's formula. Substitution of $H_1^{(1)}$ in equation (3-1) gives an approximation of the slope represented by the tangent at point F. If a value of H_1 were calculated using the slope at the end of the interval, we would have $H_1 = H_0 + BE$. When the average of the slopes at the ends of the interval are used in place of D_0 we find that $H_1^{(2)} = H_0 + BG$ which is definitely a better approximation to the real value of H_1 at t_1 than the first value of H_1 computed. This process may be represented symbolically as

$$(3-9) \quad \Delta H = \frac{1}{2} h(D_0 + D_1^{(1)})$$

$$\Delta H = \frac{1}{2} (BE + BC) = \frac{1}{2} (BE + BE + EC)$$

$$(3-10) \quad \Delta H = BE + \frac{1}{2} EC$$

$$(3-11) \quad H_1^{(2)} = H_0 + \Delta H = H_0 + (BE + \frac{1}{2} EC)$$

¹The superscript (k) on $H_1^{(k)}$ indicates the kth value of H_1 where H_1 is the approximation to the actual value of H_1 at a time $t = t_1 = 0 + \Delta t$.

The new value is much closer to the actual value of $H_1 = (H_0 + BF)$ than before. Continuation of this process by again calculating an approximation to the slope at the end point of the interval and by substituting $H_1^{(2)}$ in equation (3-1), will yield a more accurate approximation to the slope at point F. A new value of $H_1^{(3)}$ can then be calculated.

$$(3-12) \quad H_1^{(3)} = H_0 + \frac{1}{2} h(D_0 + D_1^{(2)})$$

$H_1^{(3)}$ will be more accurate than $H_1^{(2)}$ since the approximation of the slope will be improved; i.e., $D_1^{(2)}$ is more accurate than $D_1^{(1)}$. Further continuation of the process will yield successively more accurate approximations to the actual value of H_1 . The process may be continued until the value of $H_1^{(k+1)} = H_1^{(k)}$, where $H_1^{(k)}$ is the approximation of the actual value of H_1 . It must be noted that $H_1^{(k)}$ is only an approximation to the actual H_1 . As successive values of $H_1^{(k)}$ are generated, $H_1^{(k)}$ will converge to a value $H_1^{(n)}$ which will not be the same as the actual value of H_1 . To force the approximation $H_1^{(k)}$ to converge to H_1 we must make h very small. The approximate solution will approach the exact solution as $h \rightarrow 0$; however, as $h \rightarrow 0$ the number of calculations increases and the computation time becomes very large. It then becomes necessary to determine the magnitude of errors permissible in relation to the time available and adjust h accordingly.

To illustrate the fact that the limit of $H_1^{(k)}$ does not approach the actual H_1 as $k \rightarrow \infty$, let us consider the following example. Suppose $f(t, H)$ is expressed by equation (3-13) and

the initial conditions $t_0 = 0.0$ and $H_0 = 1.0$.

$$(3-13) \quad D = f(t, H) = t + H$$

Substituting these values of t_0 and H_0 in equation (3-13) we obtain

$$D_0 = f(t_0, H_0) = t_0 + H_0 = 1.0$$

If we select $h = 0.05$, we may then write

$$\begin{aligned} H_1^{(1)} &= H_0 + D_0 h = 1.05 \\ D_1^{(1)} &= t_1 + H_1^{(1)} = 1.10 \end{aligned}$$

The second approximation to H_1 and D_1 are

$$\begin{aligned} H_1^{(2)} &= H_0 + \frac{1}{2} h(D_0 + D_1^{(1)}) = 1.0525 \\ D_1^{(2)} &= t_1 + H_1^{(2)} = 1.1025 \end{aligned}$$

Continuing we have successive approximations to H_1 and D_1 as follows.

$$\begin{aligned} H_1^{(3)} &= H_0 + \frac{1}{2} h(D_0 + D_1^{(2)}) = 1.05256 \\ D_1^{(3)} &= t_1 + H_1^{(3)} = 1.10256 \\ H_1^{(4)} &= H_0 + \frac{1}{2} h(D_0 + D_1^{(3)}) = 1.05256 \end{aligned}$$

Since $H_1^{(4)} = H_1^{(3)}$ we should now stop the process if we desire $H_1^{(k)}$ to agree with $H_1^{(k+1)}$ only to the fifth decimal place. Further values of $H_1^{(k)}$ for $3 \leq k \leq \infty$ will yield the same value for $H_1^{(k)}$ in the first five decimal positions as $H_1^{(3)}$. This points out the fact that the approximation $H_1^{(k)}$ will not

converge to the actual H_1 as $k \rightarrow \infty$ since H_1 actual = 1.05254.

We now take

$$H_1 = H_1^{(3)} = 1.0526$$

$$D_1 = D_1^{(3)} = 1.1026$$

Continuing, we can calculate the first approximation for H_2 and D_2 .

$$H_2^{(1)} = H_1 + D_1^{(1)}h = 1.1077$$

$$D_2^{(1)} = t_2 + H_2^{(1)} = 1.2077$$

Then we calculate second and third approximations for H_2 and D_2 (i.e., $H_2^{(2)}, H_2^{(3)}$ and $D_2^{(2)}, D_2^{(3)}$) and stop since $D_2^{(3)} = D_2^{(2)}$. Consequently $H_2^{(3)}$ agrees with $H_2^{(2)}$ and we have

$$(3-14) \quad H_2 = H_2^{(3)} = 1.1104$$

$$(3-15) \quad D_2 = D_2^{(3)} = 1.2104$$

Further continuations will yield successive values of H_3, D_3, H_4, D_4 , etc.

To evaluate the accuracy of the method let us compare the approximated values of H_1, H_2 , etc., to those calculated from the solution to $dH/dt = t + H$, which is

$$(3-16) \quad \bar{H} = 2e^t - t - 1$$

These are given in Table 1. Accuracies of H indicated in Table 1 can be improved only by using a smaller value for h .

Table 1. Comparison of approximate and exact solution
of $dH/dt = t + H = D$.

i	t	H_i^{act}	H_i^{approx}	D_i^{act}	D_i^{approx}
0	0.00	1.00000	1.0000	1.00000	1.0000
1	0.05	1.05254	1.0526	1.10254	1.1026
2	0.10	1.11034	1.1104	1.21034	1.2104

In general, execution of the Modified Euler's method for the equation $D = f(t, H)$ is as follows.

Step 1. Obtain the initial conditions

$$t = t_0$$

$$h = \text{constant} = \Delta t$$

$$H = H_0$$

Step 2. Evaluate the approximations to H_1 by generating the approximations

$$D_0, H_1^{(1)}, D_1^{(1)}, H_1^{(2)}, D_1^{(2)} \dots H_1^{(k)}, D_1^{(k)},$$

where

$$(3-17) \quad D_0 = f(t_0, H_0)$$

$$(3-18) \quad H_1^{(1)} = H_0 + D_0 h$$

$$(3-19) \quad D_1^{(k)} = f(t_1, H_1^{(k)})$$

$$(3-20) \quad H_1^{(k)} = H_0 + \frac{1}{2} h (D_0 + D_1^{(k-1)})$$

for $k = (1, 2, 3, \dots, p)$ and p an integer such that

$$(3-21) \quad D_1^{(p)} = D_1^{(p-1)}$$

$$(3-22) \quad H_1^{(p+1)} = H_1^{(p)}$$

When this condition occurs take

$$(3-23) \quad D_1 = D_1^{(p-1)}$$

$$(3-24) \quad H_1 = H_1^{(p)}$$

Then proceed to Step 3.

Step 3. Evaluate the approximations to H_2 from the calculated approximation to H_1 and D_1 obtained as the result of Step 2 and generation of

$$D_2^{(1)}, H_2^{(2)}, D_2^{(2)}, \dots, H_2^{(k)}, D_2^{(k)}$$

where

$$(3-25) \quad H_2^{(k)} = H_1 + D_1 h$$

$$(3-26) \quad D_2^{(k)} = f(t_2, H_2^{(k)})$$

$$(3-27) \quad H_2^{(k)} = H_1 + \frac{1}{2} h (D_1 + D_2^{(k-1)})$$

for $k = 1, 2, 3, \dots, q$ and q an integer where

$$(3-28) \quad D_2^{(q)} = D_2^{(q-1)}$$

$$(3-29) \quad H_2^{(q+1)} = H_2^{(q)}$$

When this condition occurs take

$$(3-30) \quad D_2 = D_2^{(q-1)}$$

$$(3-31) \quad H_2 = H_2^{(q)}$$

Then proceed to Step 4.

Step 4 will generate the approximation to H_3 and D_3 after which we proceed to Step 5 for an approximation to H_4 and D_4 and so on until we have an approximation for H_n and

D_n . The process should stop when the n^{th} approximation has been evaluated. The value chosen for n will depend on the accuracy desired and maximum value of Δt to be used. A flow chart of the process is given in Fig. 11 for the function $D = f(t, H)$.

If the function $f(t, H)$ is specified by the Hysteretic Diffusion Equation and we apply the Modified Euler Method as before, we can solve for the value of $H(x, y, z, t)$ as specified by the partial differential equation (2-28).

3.3 Numerical Solution of the Hysteretic Diffusion Equation with the Modified Euler Method

When applying the Modified Euler Method to the Hysteretic Diffusion Equation, equation (2-28), we find that evaluating successive values of $\partial H/\partial t$ becomes somewhat more complex since $f(t, H)$ now depends on many variables; i.e.,

$$(3-32) \quad f(t, H) = \frac{(C_2 + H_z)^2}{\sigma C_1 C_2} \left[\frac{\partial^2 H_z}{\partial x^2} + \frac{\partial^2 H_z}{\partial y^2} \right]$$

Because of this it becomes more difficult to visualize the physical significance of $\partial H/\partial t$ in terms of equation variables; however, no problem should exist in evaluating $\partial H/\partial t$ if we express the equation in its finite difference form and establish initial and boundary conditions. The quantity $\partial H/\partial t$ may now be evaluated simply by a series of numerical operations. Successive values of $\left. \partial H/\partial t \right|_{t=t_i}$ can be found through the recurrence relations as outlined in the flow chart of the Modified Euler

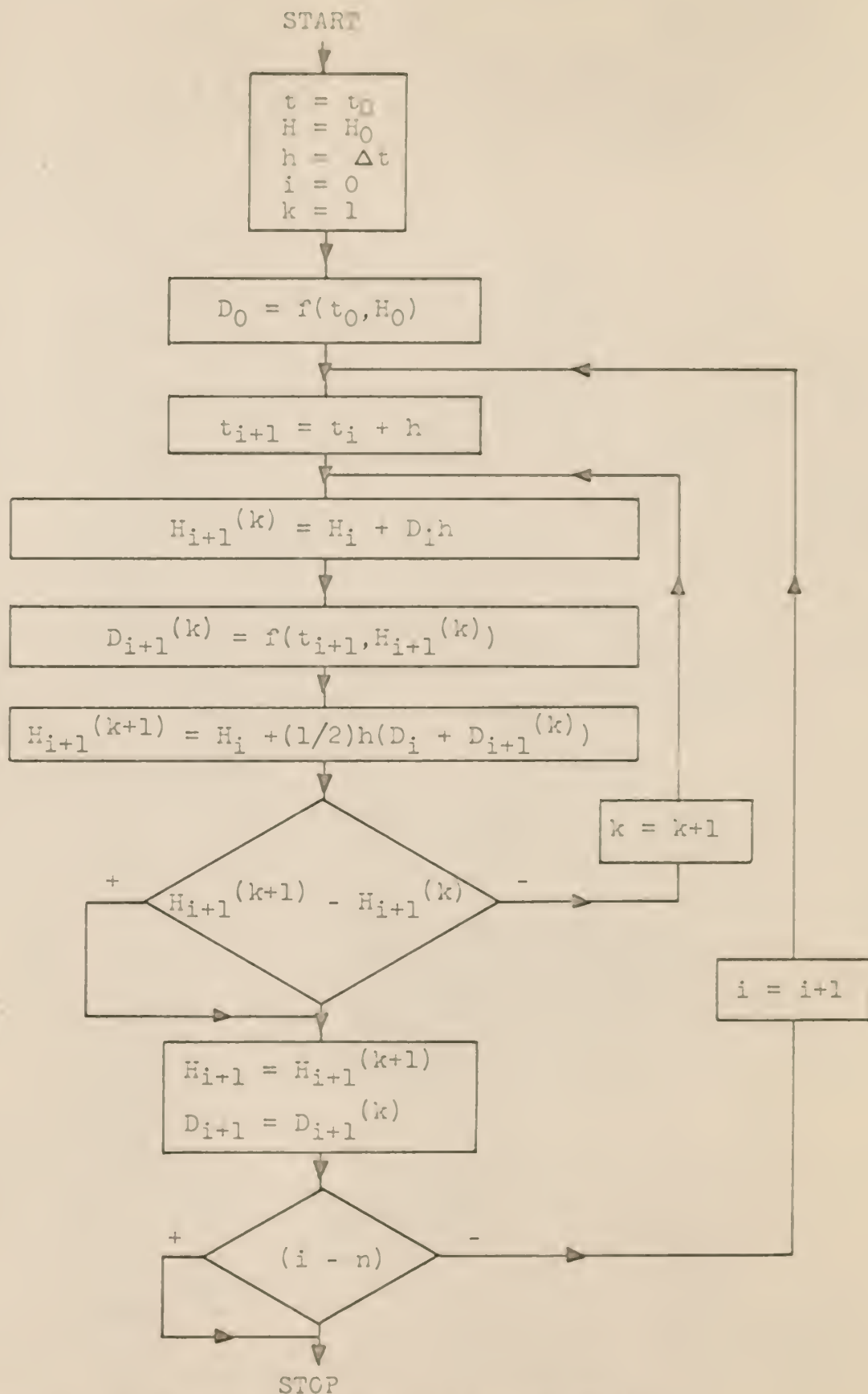


Fig. 11. Flow chart of modified Euler method for $dH/dt = f(t, H) = D$.

Process. Although some care must be exercised to limit the size of Δt to insure that the

$$(3-33) \quad \lim_{k \rightarrow \infty} (H_i^{(k+1)} - H_i^{(k)}) \rightarrow 0$$

the problem is a straightforward procedure as outlined by the flow chart of Fig. 11.

In the following discussion, the magnetic field intensity for the I, J^{th} point in the grid of Fig. 12 is specified by $H(I, J)$; the i^{th} approximation to H at the I, J^{th} point is denoted by $H(I, J)_i$; and the k^{th} value of the i^{th} approximation of H for the I, J^{th} point is $H(I, J)_i^{(k)}$. Note that the superscript does not indicate that $H(I, J)_i$ is raised to a power or that it is the k^{th} derivative of H . Let us also denote $(\partial H / \partial t)_i$ by P_i .

3.4 Conditions for Convergence

Conditions exist under which the process will not yield a solution. This is due to the fact that $\lim_{k \rightarrow \infty} (H_i^{(k+1)} - H_i^{(k)})$ will not approach zero; however, we can establish a relationship between the time increment and other parameters such as grid size, C_1, C_2, σ , and H_{max} to insure convergence of the process and existence of a solution.

Let us divide a rectangular cross section of the toroid considered in section 3.0 into a rectangular grid network as shown in Fig. 12. Each grid is of equal area h^2 and $H(I, J)$

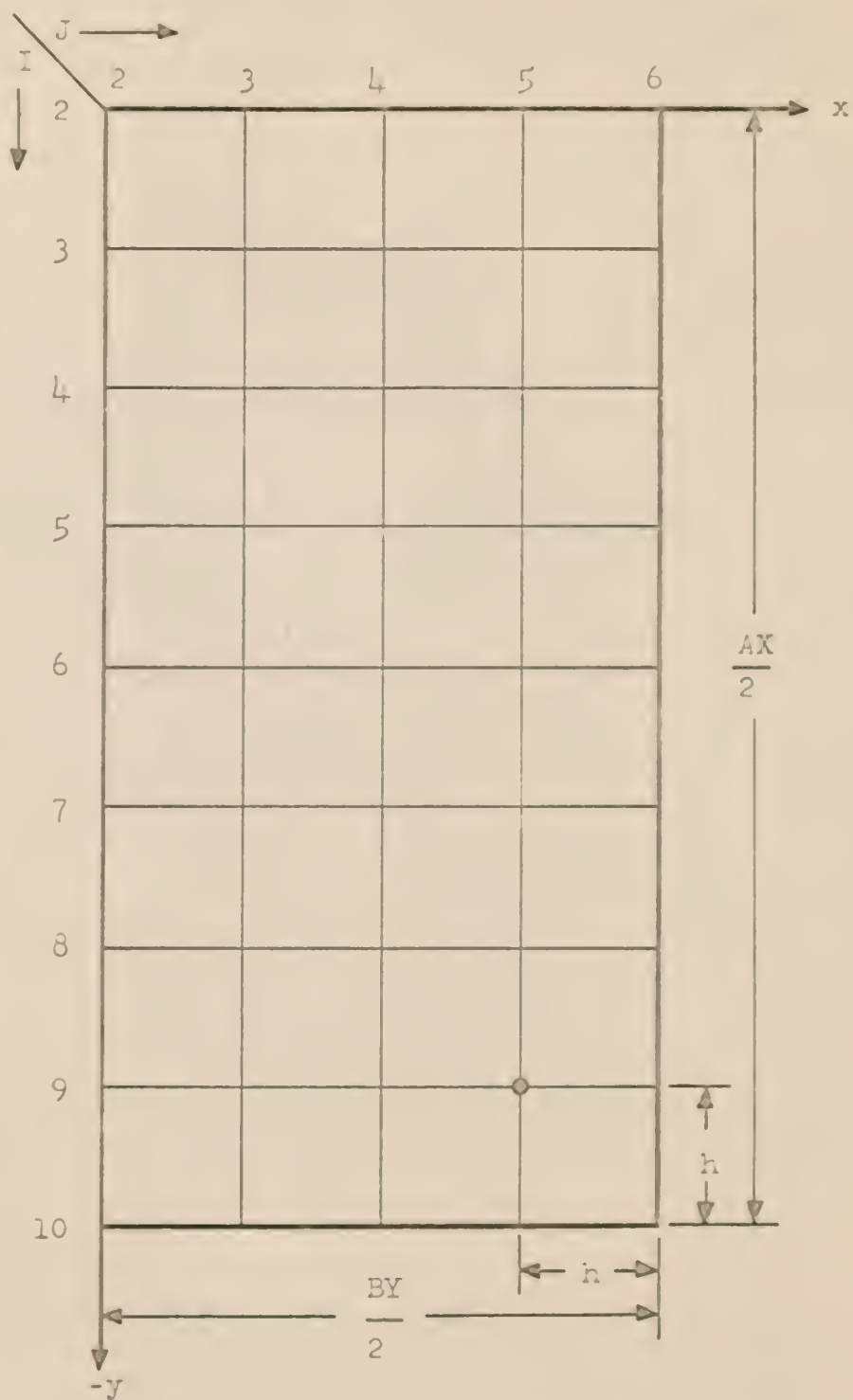


Fig. 12. One-fourth of rectangular toroid cross section for sample problem given in section 5.0. (See Fig. 7b.)
 Note: $NI = 4$, $NJ = 8$, $NY = 6$, $NX = 10$.

is the field intensity at the intersection of the I^{th} row and J^{th} column of the grid.

The Hysteretic Diffusion Equation for the I, J^{th} point is given by equation (3-34).

$$(3-34)^1 \quad P(I, J) = \frac{(C_2 + H(I, J))^2}{\sigma C_1 C_2} \left[\frac{\partial^2 H(I, J)}{\partial x^2} + \frac{\partial^2 H(I, J)}{\partial y^2} \right]$$

Expressed in its finite difference form the equation becomes

$$V(I, J) = H(I + 1, J) - 2H(I, J) + H(I - 1, J)$$

$$W(I, J) = H(I, J + 1) - 2H(I, J) + H(I, J - 1)$$

$$(3-35)^1 \quad P(I, J) = \frac{(C_2 + H(I, J))^2}{\sigma C_1 C_2} \left[\frac{W(I, J)}{h^2} + \frac{V(I, J)}{h^2} \right]$$

If we wish to evaluate $P(I, J)$ at $I = 5, J = 3$, we obtain

$$V(5, 3) = H(6, 3) - 2H(5, 3) + H(4, 3)$$

$$W(5, 3) = H(5, 4) - 2H(5, 3) + H(5, 2)$$

$$(3-36)^1 \quad P(I, J) = \frac{(C_2 + H(5, 3))^2}{\sigma C_1 C_2} \left[\frac{W(5, 3)}{h^2} + \frac{V(5, 3)}{h^2} \right]$$

3.5 Relationship Between Δt and Other Parameters to Insure Convergence

To investigate the relationship between Δt and other parameters of the problem let us consider a sample calculation of the

¹Note that P_i represents $(\partial H / \partial t)_i$.

approximation to H at a corner position of the grid not on the boundary, i.e., the position $I = NY - 1, J = NX - 1$. (See Fig. 7 and Fig. 12, position $(I, J) = (5, 9)$.)

Let us apply initial conditions and boundary conditions when all points inside the boundary are at $H_0 = H_b$ and all points on the boundary are at $H_0 = -mH_b$ where m is some real constant.

If we examine the Hysteretic Diffusion Equation expressed in its finite difference form (equation 3-48), we find that the largest value of P will occur at the $I = NY - 1, J = NX - 1$ position. To approximate a value of H at this position after a time Δt , we use the equation

$$(3-37) \quad H(NY - 1, NX - 1)_1^{(1)} = H(NY - 1, NX - 1)_0 + P_0(NY - 1, NX - 1) \Delta t$$

Since we are considering only one position in the grid we may drop the subscripts; thus

$$H_1^{(1)} = H_b + P_0 \Delta t$$

Evaluating P_0 we find that

$$P_0 = \frac{(C_2 + H_{\max})^2}{\sigma C_1 C_2} \left[\frac{\partial^2 H}{\partial x^2} + \frac{\partial^2 H}{\partial y^2} \right]$$

since $\frac{\partial^2 H}{\partial x^2} = \frac{\partial^2 H}{\partial y^2}$ at this position we have

$$(3-38) \quad P_0 = 2x_0(H(NY, NX - 1) - 2H(NY - 1, NX - 1) + H(NY - 2, NX - 1))$$

$$P_0 = 2x_0(-mH_b - 2H_b + H_b)$$

$$P_0 = -2x_0(m + 1) H_b$$

$$(3-39) \quad H_1^{(1)} = H_b - 2x_0(m + 1) H_b \Delta t$$

We know the approximation H_1 is between $-mH_b$ and $+H_b$, thus $2x_0(m + 1) H_b \Delta t$ has to be chosen such that

$$(3-40) \quad -mH_b \leq (H_b - 2x_0(m + 1)H_b \Delta t) \leq H_b .$$

We can therefore limit the value of Δt to satisfy this condition; i.e.,

$$(3-41) \quad 0 \leq 2x_0(m + 1)H_b \Delta t \leq (m + 1)H_b .$$

Thus the possible values of Δt become

$$(3-42) \quad 0 \leq \Delta t \leq 1/2x_0$$

where

$$x_0 = \frac{(C_2 + H_{\max})^2}{\sigma C_1 C_2 h^2}$$

Let us now investigate a few special cases.

Case I. For $\Delta t = 1/2x_0$ we have

$$P_0 = -2x_0(m + 1)H_b$$

$$H_1^{(1)} = H_b - 2x_0(m + 1)H_b \Delta t = -mH_b$$

$$P_1^{(1)} = +2x_0(m + 1)H_b$$

$$H_1^{(2)} = H_b$$

$$P_1^{(2)} = -2x_0(m + 1)H_b$$

$$H_1^{(3)} = -mH_b$$

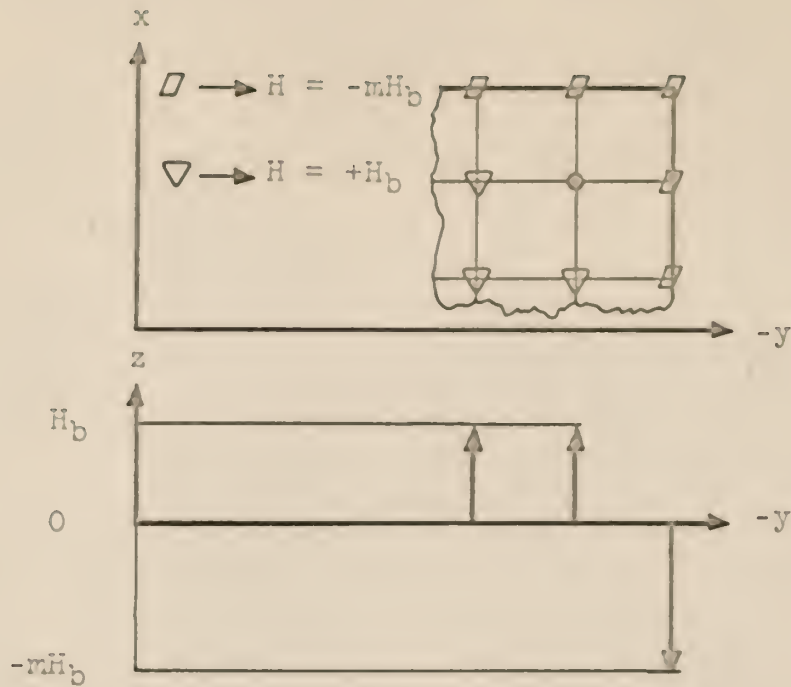


Fig. 13. Grid section of Fig. 12. Indicates values of H_z adjacent to $H_z(9,5)$ at $T = 0^+$.

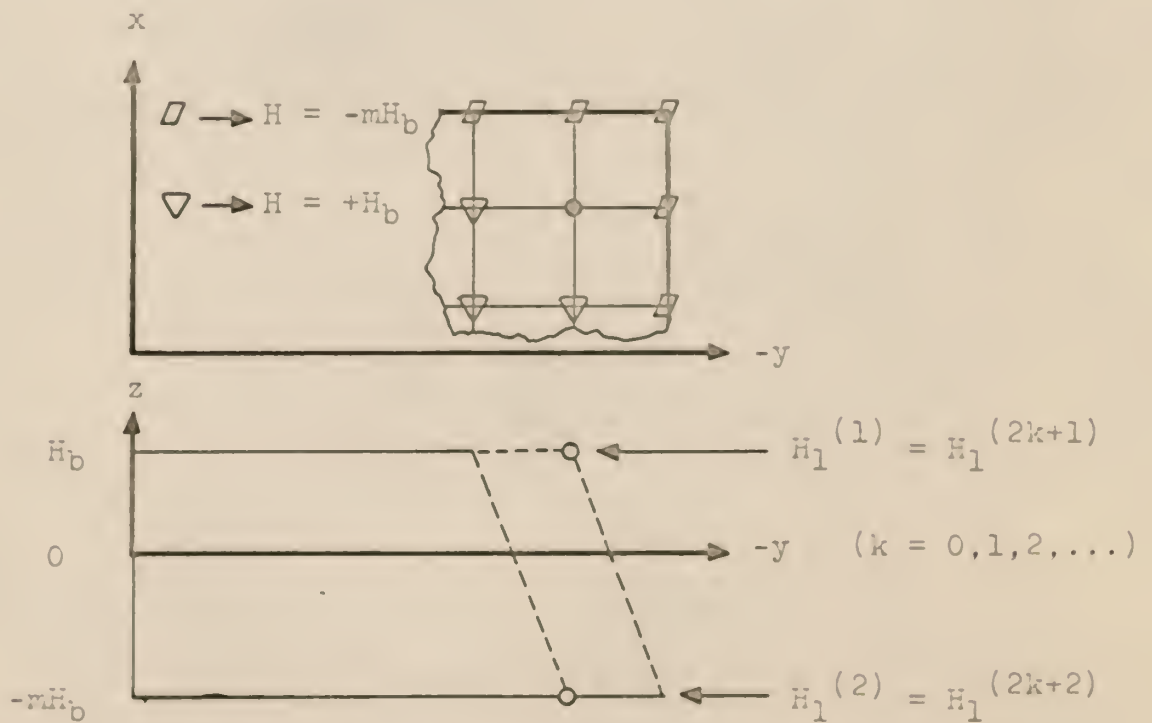


Fig. 14. Grid section of Fig. 12. Indicating values of $H_z(9,5)_1^{(k)}$, $k = 1, 2, 3, \dots, p$; for $H_{\text{boundary}} = -mH_b$, $\Delta t = 1/2x_0$.

$$P_1^{(3)} = +2x_0(m+1)H_b$$

$$H_1^{(4)} = H_b$$

Thus we can see that an oscillation between H_b and $-mH_b$ occurs and the approximation to H_1 can never be found since the $\lim_{k \rightarrow \infty} (H_1^{(k+1)} - H_1^{(k)}) \rightarrow 0$ is not satisfied.

Case II. For $\Delta t = 1/4x_0$ we have

$$P_0 = -2x_0(m+1)H_b$$

$$H_1^{(1)} = H_b - 2x_0(m+1)H_b \Delta t = \frac{H_b}{2} (1 - m)$$

$$P_1^{(1)} = 0$$

$$H_1^{(2)} = H_b - x_0(m+1)H_b \Delta t = \frac{H_b}{4} (3 - m)$$

$$P_1^{(2)} = -x_0(m+1)H_b$$

$$H_1^{(3)} = H_b - (3/2)x_0(m+1)H_b \Delta t = \frac{H_b}{8} (5 - 3m)$$

$$P_1^{(3)} = -x_0/2(m+1)H_b$$

$$H_1^{(4)} = H_b - (5/4)x_0(m+1)H_b \Delta t = \frac{H_b}{16} (11 - 5m)$$

Thus the process is converging to some value $H_1^{(n)}$ and $\lim_{k \rightarrow \infty} (H_1^{(k+1)} - H_1^{(k)}) \rightarrow 0$ is satisfied. One can easily

visualize convergence by considering H to be zero on the boundary (i.e., $m = 0$), and establishing approximations $H_1^{(k)}$ as k increases. This process is shown in Fig. 16.

In order to have convergence we must not include zero or $1/2x_0$ in the permissible values; thus

$$(3-43) \quad 0 < \Delta t < 1/2x_0$$

Selection of the Δt to be used is dependent on the computation time available and accuracies desired. As before, large

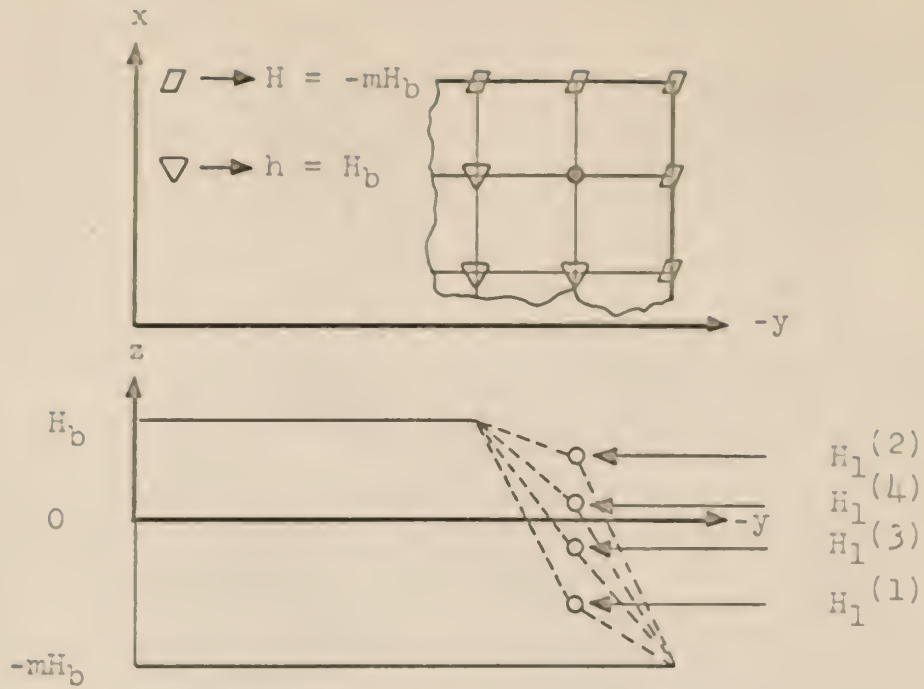


Fig. 15. Grid section of Fig. 12. Indicating values of $H_2(9,5)_1^{(k)}$, $k = 1, 2, 3, 4$, for $H_{\text{boundary}} = -mH_b$, $\Delta t = 1/4x_0$.

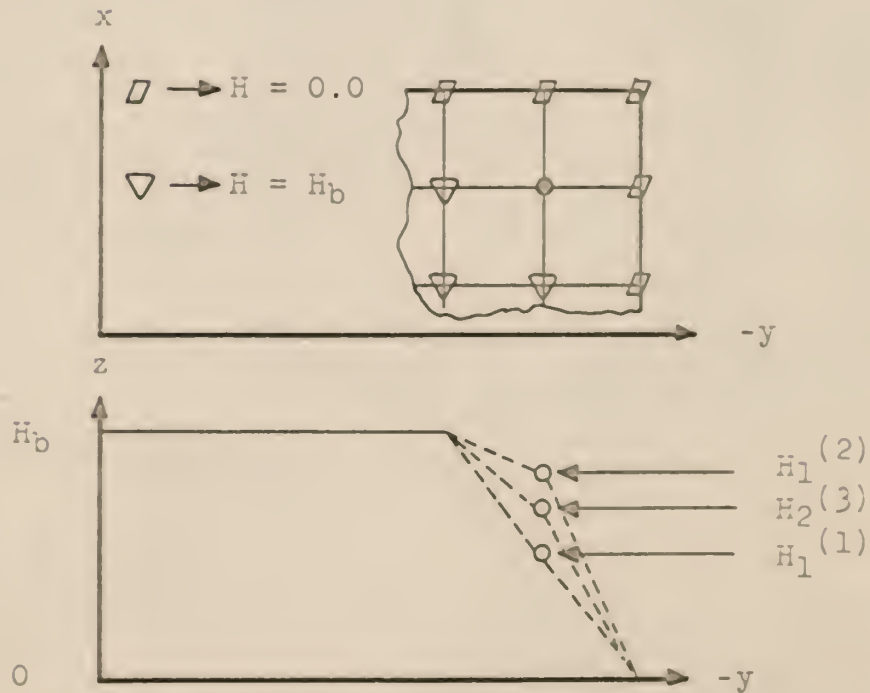


Fig. 16. Grid section of Fig. 12. Indicating values of $H_2(9,5)_1^{(k)}$, $k = 1, 2, 3$, for $H_{\text{boundary}} = 0.0$, $\Delta t = 1/4x_0$.

Δt 's introduce inaccuracies; however, they may allow evaluation of the maximum value of H more rapidly. This is not necessarily the case since more iterations are necessary to evaluate approximations when Δt is large, and it becomes possible to use more time in evaluating these approximations than to evaluate H in smaller steps of Δt . Selection of a permissible Δt for convergence requires that

$$(3-44) \quad 0 < \Delta t < \frac{\sigma C_1 C_2 h^2}{2(C_2 + H_{\max})^2} .$$

If we force the corner position to converge all other positions along the boundary will converge since they only require that

$$(3-45) \quad 0 < \Delta t < \frac{\sigma C_1 C_2 h^2}{(C_2 + H_{\max})^2} .$$

To insure convergence we may write

$$(3-46) \quad 0 < \Delta t < C_0 \frac{\sigma C_1 C_2 h^2}{(C_2 + H_{\max})^2}$$

where the parameters of equation (3-46) are defined as follows:

C_0 = constant and is always < 0.5

σ = material conductivity

C_1 = constant for B-H approximation (see section 2-6)

C_2 = constant for B-H approximation (see section 2-6)

h = the grid size selected (see Fig. 12)

H_{\max} = the magnitude of the maximum value of $H(I,J)$
specified by the initial conditions of
boundary conditions at time $t = 0^+$.

It is now evident that Δt has to be selected in accordance with the conditions of equation (3-44) and is not an arbitrary variable. It might also be noted that smaller grid areas h^2 will require smaller time increments. This condition specifies the accuracy of the method in the same manner as limiting the value of Δt to be small in the example of section 3.2.

With the various limitations on Δt now specified we can proceed to implement the solution to equation (2-28) by constructing a program to perform the almost endless number of numerical calculations necessary. Discussion of this task is given in the next section.

4.0 PROGRAM FOR SOLUTION OF THE HYSTERETIC DIFFUSION EQUATION

4.1 Introduction

The basic function of the program is to execute the Modified Euler Method of numerical integration to solve the Hysteretic Diffusion Equation and to obtain the magnetic field intensity distribution pattern in the cross section of a rectangular toroid subject to given initial time and boundary conditions. Before this process may be executed we must first define the problem in terms of variable names in accordance with the FORTRAN IV programming language, read in the input data, and establish output data formats. A method of completing this task is given in the next section. It is assumed that the reader is familiar with the FORTRAN IV language.

Discussion of the program is given in various sections with card numbers indicating locations of the sections in the program listing of Appendix C. Notation and symbolism used in the program follow closely that used in previous discussion, although some modifications were necessary to satisfy programming requirements.

4.2 Main Program EMEX

This program controls all numerical operations and may be subdivided to indicate the particular functions performed.

4.2A Comments (MG90000 - MG90106)

The cards serve as a partial list of variable names in the program and specify various numerical values to provide given print-out formats.

4.2B Input Data (MG90122 - MG90138)
and (MG90262)

The first five cards of each data set give all data necessary to execute a given problem except the initial and boundary conditions which are contained on the remaining cards 6 through n.

Any variable name beginning with I, J, K, L, M, or N is fixed point and read according to an I5 format. Variable names beginning with other letters are floating point quantities and are read according to an E16.8 format. The E16.8 format allows any of the common systems of units to be used without alteration of the input-output format specifications. Input variables are read in the sequential order as specified below.

Card 1	EO, HSO, BR, HMAX, CUO
Card 2	TO, TKM, COND, AX, BY
Card 3	CDT, KZ, NI, NJ, K, KO, KI, KG
Card 4	B2, H2, B3, H3
Card 5	H10, DHX, I3, I4, I5
Card 6	.
.	.
.	((HO(I,J), J = 2,NX), I = 2,NY)
¹ Card n	.

¹The value of $n = 5 + (NI + 1)(NJ + 1)/5$, $NX = NJ + 2$, $NY = NI + 2$.

Definition of the variables and methods of obtaining numerical values for them are given as follows.

1. BR - Residual flux density specified as the flux density for the magnetic field intensity $H = 0.0$.
2. B2, H2 - These are values of flux density and magnetic field intensity taken near the knee of the decay portion of the hysteresis curve. (See section 2.6.)
B3, H3
3. H10, DHX - Variables used for control information which specify the number of points on the approximate hysteresis loop. For a symmetrical hysteresis loop as shown in Fig. 22, choose $I3 = 1 + 2(H10/DHX)$.
I3
 - a. H10 - Specifies the maximum value of H to be plotted on the hysteresis loop approximation.
 - b. DHX - Specifies the increments of H plotted on the hysteresis loop approximation.
 - c. I3 - Specifies the number of points on the decay portion of the hysteresis loop approximation. The maximum value is 100.
4. HMAX - Specifies the maximum value of field intensity applied where $H = \mathcal{F}/l = NI/l$. (See equations 2-1, 2, 3, 4.)
5. CUO - Constant to allow for units conversion when computing values of force. $\text{Force} = B^2 A / 2\mu_0$ where $\mu_0 = 4\pi\text{CUO}$.
6. COND - Conductivity of ferromagnetic material.
7. AX - "X" dimension of the rectangular cross section as shown in Fig. 12.
8. BY - "Y" dimension of the rectangular cross section as shown in Fig. 12.
9. NI - Number of divisions of width CX in the "X" direction. (See Fig. 12, $h = CX$)
10. NJ - Number of divisions of width CX in the "Y" direction. (See Fig. 12, $h = CX$.)
11. TO - Specifies the initial value of time $T(K)$.

12. K - Subscript indicating iterations. Appears as $T(K)$.
13. I5 - Specifies execution for flux decay or flux buildup run. (See Appendix C for values.)
14. CDT - Specifies length of time increment. Same as C_0 of section 3.5. Use $CDT < 0.5$ to insure convergence, smaller for higher accuracy.
15. KZ - Specifies physical position of data set in relation to the last set; e.g., if four data sets are being run, $KZ = 1$ for the last set to be run, $KZ = 4$ for the first set to be run, etc.
16. KI - Specifies increments of time for which data is to be plotted; e.g., if we have 4000 values of total force or flux the printer need not plot every point; thus it plots every KI^{th} point. (Use KI such that $K_0/KI < 800$.)
17. KG - Specifies increments of time for which a print-out of the flux distribution pattern is desired; i.e., a print-out for $T(0)$, $T(KG)$, $T(2KG)$, etc.
18. EO - Specifies maximum error norm E as given by equation (4-1) to be less than EO ; i.e., $E \leq EO$.
19. HSO
TKM
K0 - Specify halt condition. A halt will occur immediately after any one of the conditions given below are met. One should also note that $T(K) = K \Delta t$.
- a. HSO - HS is the sum of $H(I,J)$ for all points in one-fourth of the lattice. Halt occurs for $HS \geq HSO$, buildup; $HS \leq HSO$, decay.
- b. TKM - Specifies maximum time $T(K)$ to be considered. Halt occurs when $T(K) \geq TKM$.
- c. K0 - Specifies maximum number of increments desired. Halt occurs when $K \geq K0$.
20. I4 - Specifies data selected for output and format specifications to be used. A detailed discussion of I4 is given in section 4.2G.
21. HO(I,J) - Specifies numerical value of magnetic field intensity at every point in one-fourth of the lattice at $T(0) = 0^+$. This applies the boundary conditions and initial conditions at $T(0) = 0$. Quantities are "read" according to the instruction $((HO(I,J), J = 2, NX), I = 2, NY)$.

This completes the definition of data set constituents. Further information may be obtained by studying the sample problem given in section 5.0.

4.2C Evaluation of the Hysteresis Loop Approximation Formula (MG90140 - MG90250)

This section evaluates the constants C_1 and C_2 as given equations (2-24) and (2-25), then calculates the B-H relationship as illustrated by the hysteresis loop of Fig. 22. A graphical and printed output of an approximation to the material hysteresis loop is available at the beginning of each set of output data. To accomplish this task a duplication of the subroutine FLUD was included; however, the variables were changed and an output routine was added. Refer to section 4.3 for discussion of the procedure.

4.2D Calculation of Constants and Page Heading Routine

These cards are used to place a heading at the beginning of each output data set which defines calculated variables specifying operations to be performed; provide a print-out of data read from the input data set; check for input data errors; etc. Page 76 of section 5.2 was generated by cards contained in sections 4.2D and 4.2E.

4.2E Initial Conditions Print-out Routine (MG90724 - MG90874)

This section provides a print-out of the initial input data and corresponding initial values of flux density and force at $T(0) = 0^+$. Cards MG90724 - MG90784 provide the output in a list form as shown for Tape 9. The remaining cards give a matrix output format as shown for Tape 15.

In addition to the print-out instructions total values of flux (FLUXT) and theoretical force (FORCT) are calculated and printed following listing of $B(I,J)$, $H(I,J)$, and $FORCE(I,J)$.

4.2F Programming the Modified Euler Process for the Hysteretic Diffusion Equation (MG90876 - MG91060)

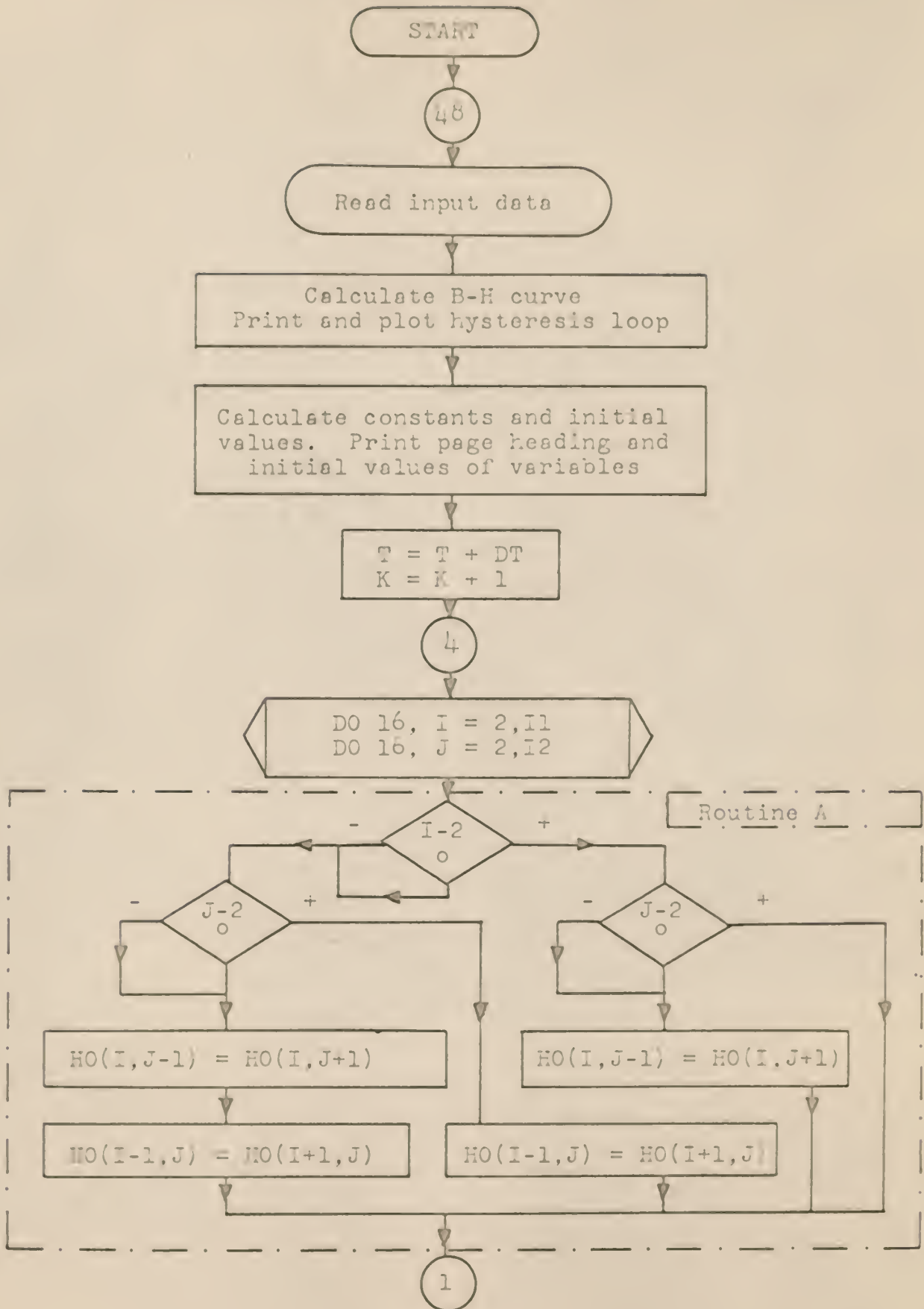
Execution of the process is nearly the same as described in section 3.2 and as outlined by the flow chart of Fig. 11; however, the condition $(r_o - (r_o/r_i)) \gg 1$ (see Fig. 7a) allows use of one-fourth of the cross-sectional area because symmetrical boundary conditions are present when the magnetic field is applied.

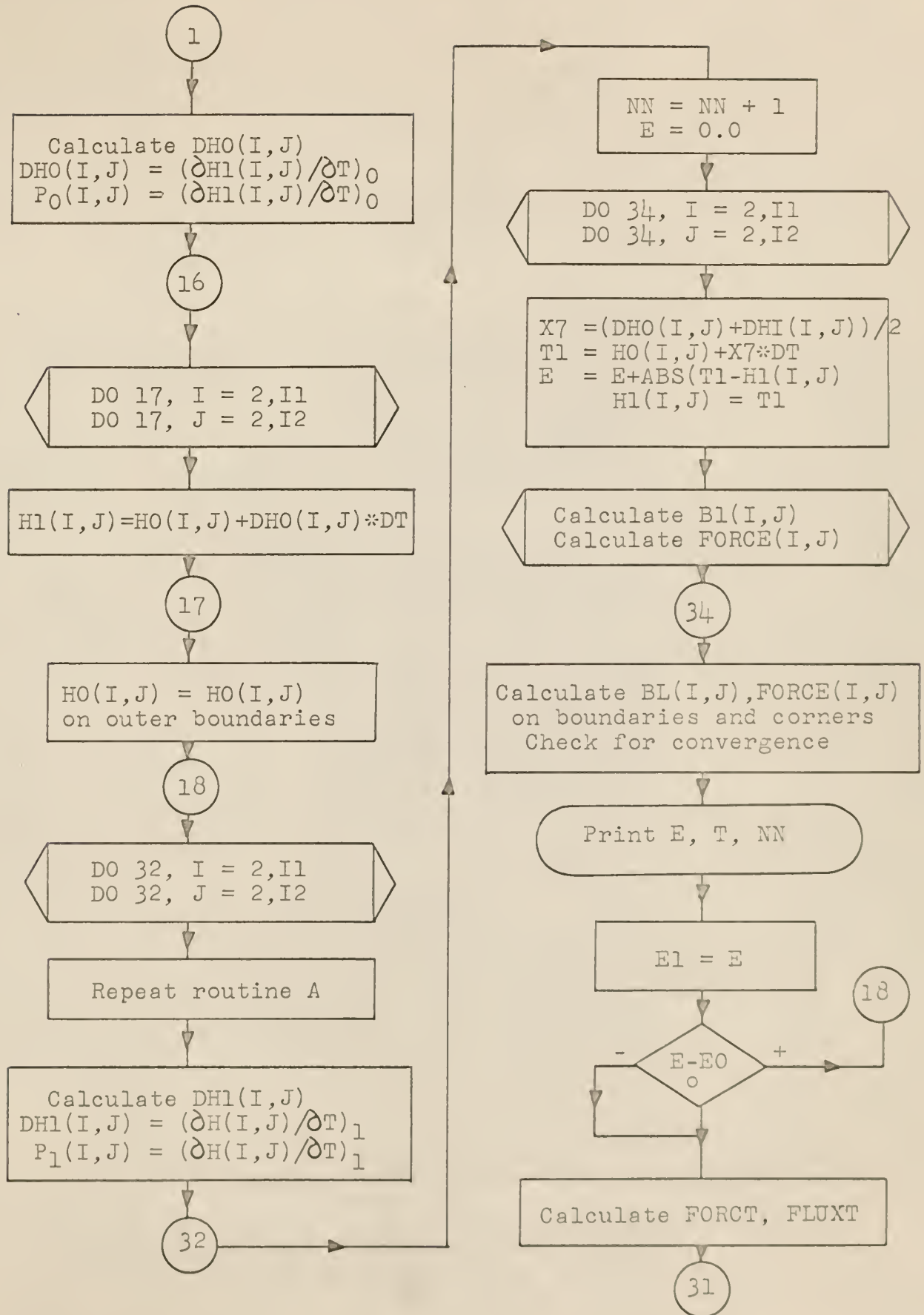
The finite difference representation of the Hysteretic Diffusion Equation (equation 3-36) requires that H for all points in the lattice adjacent to the point under consideration be known. We must compute H only at all points inside the outer boundary since the boundary conditions require flux on the boundary to remain constant. Because of symmetry, values of H immediately to the left of the Y-axis are equal to those

immediately to the right and values of H immediately below the X-axis are equal to those directly above. (See Fig. 12.) Cards MG90884 - MG90900 and MG90946 - MG90962 were coded to accomplish this task when values of I and J placed the point in question on the inner boundaries. The desired values of P_0 were computed for every I and J as were the approximations $H_1^{(1)}$, $P_1^{(1)}$, $H_2^{(2)}$, . . ., etc. While computing successive approximations to H, an accumulative error norm, E was generated. This error is expressed by equation (4-1).

$$(4-1) \quad E = \sum_{I=2}^{NX-1} \sum_{J=2}^{NY-1} \left| H(I,J)_i^{(k+1)} - H(I,J)_i^{(k)} \right|$$

This method was chosen in order that new approximations for the surrounding points would be considered when calculating successive approximations. The maximum error is specified by E0 and should be small if a high degree of accuracy is desired. Successive approximations for H will continue until the matrix error norm, $E < E0$. When a convergence condition is satisfied, NN (variable specifying number of successive approximations to H, i.e., the same as (k) in $H_i^{(k)}$) is returned to a value of 1; values of flux density, total force, and total flux are computed; and a new approximation of $H = H_{i+1}$ is begun. If the error is not converging rapidly enough to yield a solution in a reasonable amount of time or if the error is diverging, an error message will alert the operator and a machine halt will occur. A flow diagram of the Modified Euler Method for the special problem considered is given in Fig. 17.





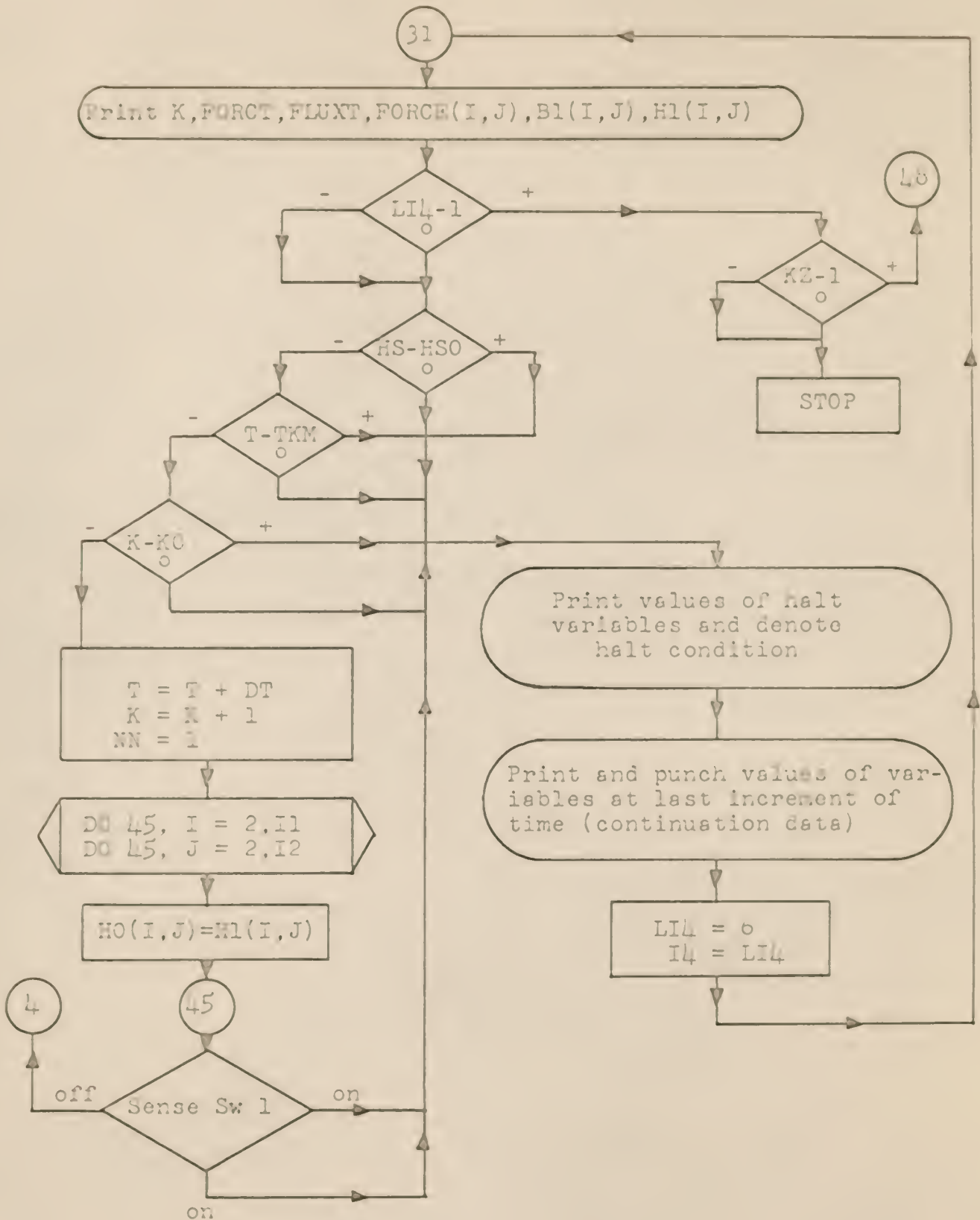


Fig. 17.
Flow chart of EMEX.

4.2G Output Data (MG91078 - MG91410)

The remaining cards are related to output data and appropriate headings for identification. Other special control functions are performed to specify the halt condition and to print-out all existing data at halt time.

Output data are available on Tapes 6, 9, 15, and 16. A sample of the output is given for Data Set 3 in section 5.2. Each tape has an output data heading for identification and prints initial values of H, B, FORCE, total force, total flux, and boundary values. Output data contained on the various tapes are as given below.

1. Tape 6 - Contains approximation to hysteresis loop and lists total force, total flux, and time. (See section 5.2A, Output Tape 6.)
2. Tape 9 - Lists all values of FORCE, B, and H with the corresponding time. The number of iterations required to satisfy the condition for convergence is also listed with the associated error norm E, i.e., lists values of E, T(K), and NN.
3. Tape 15 - Provides the same information as Tape 9 except that the output appears in a matrix form as shown in section 5.2C. Variables to be printed are selected by the value of I4 chosen. (See Table 2.)
4. Tape 16 - Data on this tape is used to punch a deck which may be used as input to a new run continuing from the existing values of variables at the completion of the first run. A print-out of the same data occurs on the output of Tape 6. One must change the value of TKM, HSO, and KO to continue.

Desired output data can be selected by specifying different values of I4. These data are not written on the tapes if not desired since the write time is very large when compared to the

computation time. I_4 may range from 1-7 and the output can be determined from Table 2. If $I_4 = 7$, only Tape 6 will have a useful output.

Table 2. Values of I_4 and related output data.

I_4	Tape 9 (list)			:	Tape 15 (matrix)			
	H(I,J)	B(I,J)	FORCE(I,J)	NN	H(I,J)	B(I,J)	FORCE(I,J)	
1	x	x	x	x	x	x	x	x
2	x	x	x	x				
3				x		x		
4				x	x			
5				x				x
6				x	x	x	x	x
7				x				

4.3 Subprogram FLUD

Subprogram FLUD computes a value of B for a corresponding H. The approximation is given by equation (2-22); however, the value of B is computed by four different methods depending on the value of H and whether the flux decay or flux buildup problem is being considered. We assume the hysteresis loop to be approximated by reflecting the decay portion of the loop and translating the entire decay loop to the right by $2H_c$. (See Fig. 18.) Let us consider the following cases to illustrate the regions and four different methods of computing B. Figure 18 illustrates the regions considered.

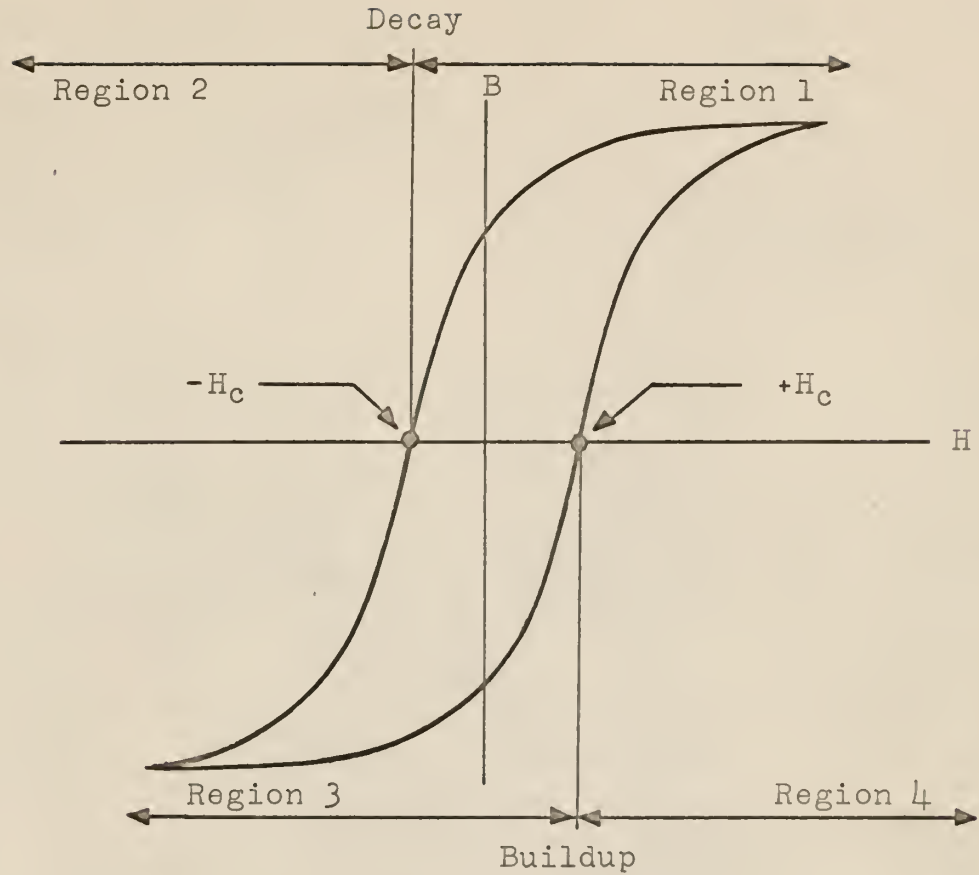


Fig. 18. Approximating the hysteresis loop.

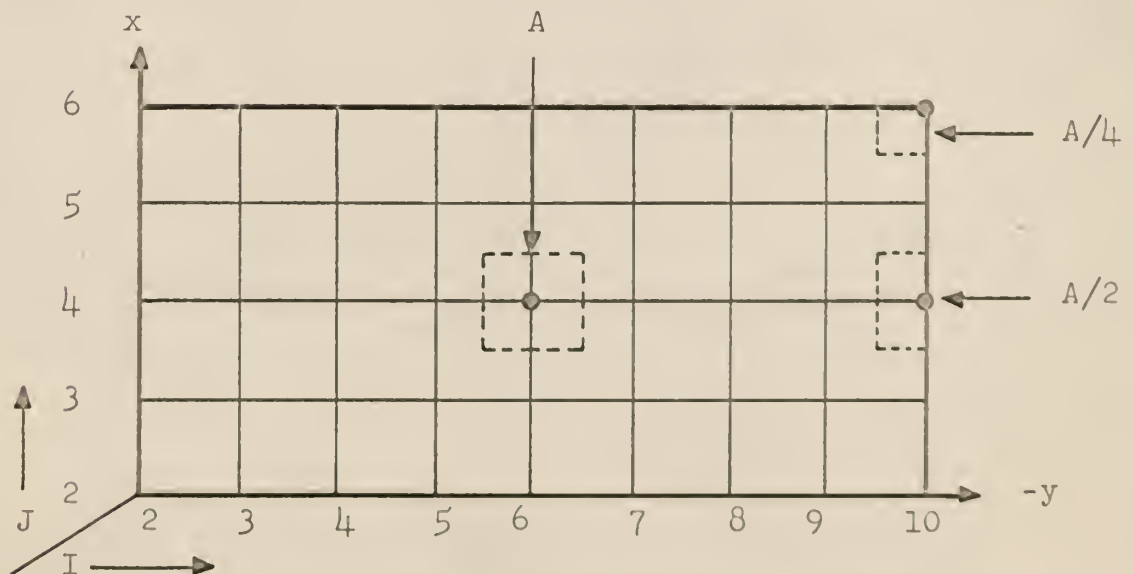


Fig. 19. Correction of area for corners and boundaries.

Region 1.¹ (Decay run, $-H_c < H < \infty$)

Flux density B is computed by a simple substitution of H in equation (2-23).

Region 2. (Decay run, $-\infty < H < -H_c$)

Since Region II of the hysteresis loop is a reflection of Region I about the H axis, we may compute B by evaluating the difference between H and H_c , forming a new variable equal to the sum of $-H_c$ and $|H - H_c|$ to replace H in equation (2-23). Then compute B with H replaced by $|H - H_c| - H_c$. A minus sign is then assigned to B in order to obtain the reflection characteristic.

Region 3. (Buildup run, $-\infty < H < H_c$)

Values of B in this region may be obtained by adding the quantity $2H_c$ to H and computing the negative of B with H replaced by the sum $H + 2H_c$.

Region 4. (Buildup run, $H_c < H < \infty$)

Values of B in this region are computed by subtracting the quantity $2H_c$ from H and computing B with H replaced by the difference $H - 2H_c$.

Subroutine FLUD determines the particular region in which H falls, and computes the corresponding B according to the rules discussed above. A flow diagram of the subroutine is given in Fig. 20.

4.4 Subprogram FORCX

This subroutine computes a fictitious force for each point in the lattice where $\text{Force} = B^2 A / 2\mu_0$. The reason for calculating forces is to obtain a force time relation for the electromechanical system shown in Fig. 2. The rectangular toroid was used

¹The value H_c is a positive real number.

SUBROUTINE START

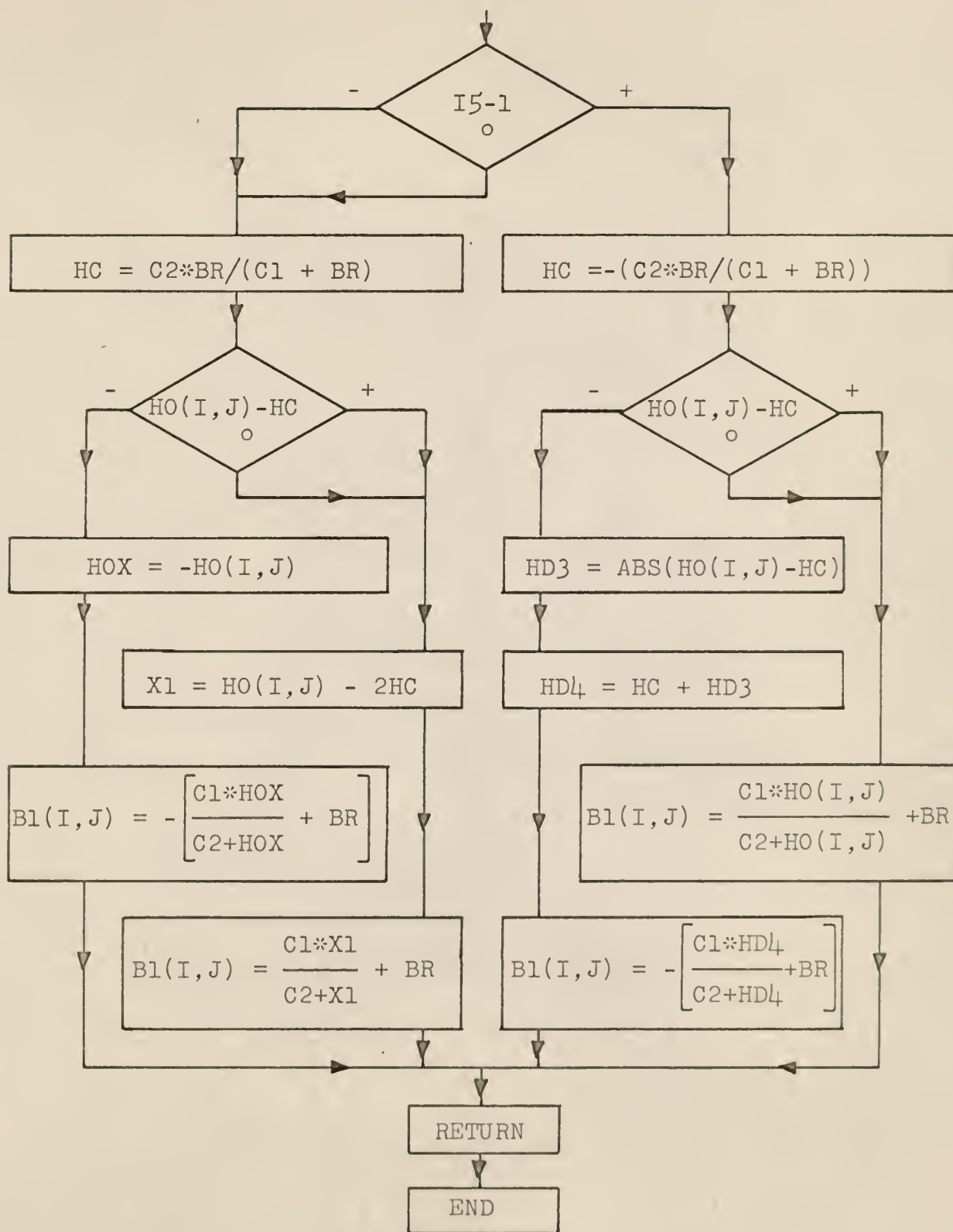


Fig. 20. Flow chart of FLUD.

as an approximation to the system with the assumption that the gap would not increase the reluctance of the flux path.

Total force in the toroid is the summation of forces at each grid point; i.e.,

$$(4-2) \quad \text{Total force (FORCT)} = \sum_{I=2}^{NY-1} \sum_{J=2}^{NX-1} \text{FORCE}(I,J)$$

Computation of the force may be approximated by the formula given if the area A is adjusted as follows. Since approximate values of B are known only for points at intersecting grid lines, areas used for computing force values along all boundaries on the lattice should only be one-half of those for the inside points and areas used for the corners should be only one-fourth of the inside areas. (See Fig. 19.)

Cards (MG91466 - MG91480) select the correct formula for computing force depending on values of I and J ; (i.e., replace A by $A/2$ on all boundaries and A by $A/4$ for all corners). A flow chart of the subroutine is given in Fig. 21.

4.5 Subprogram FLUX

Subroutine FLUX evaluates total flux by forming the product B^2A . Operations performed closely resemble those of subroutine FORCX except only total flux is computed and the flux for each point in the lattice is not. Once B is determined and subroutine FLUX is called, B is multiplied by the appropriate area and added to the existing value of total flux. When all points

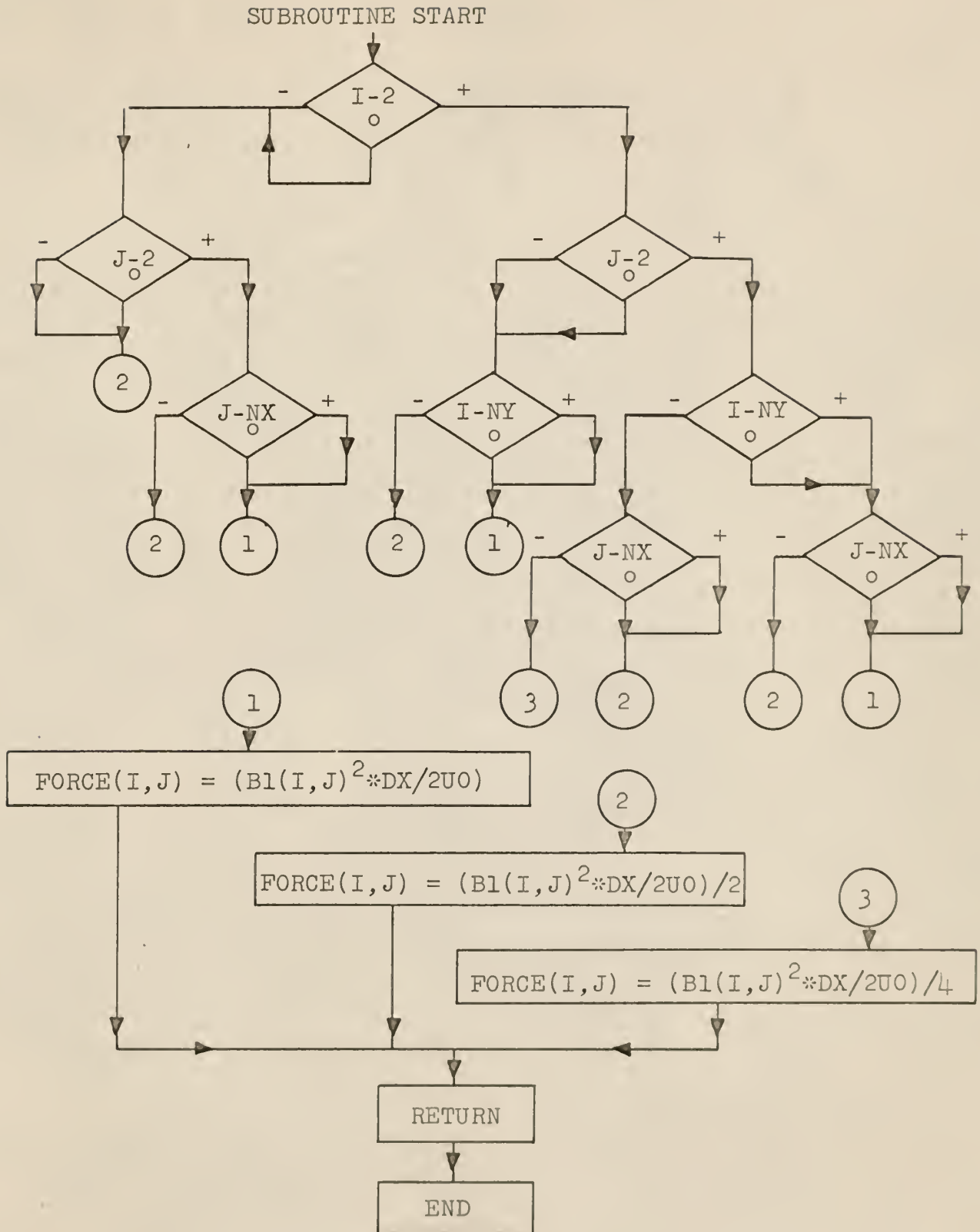


Fig. 21. Flow chart of FORCX.

have been summed, the total flux is known; however, no matrix print-out of the individual flux values can occur since the individual values are not stored. A flow diagram would be a duplication of that for subroutine FORCX with a change of variables.

A listing of the main program and associated subprograms with four sample data sets is given in Appendix C. Further information concerning their use will be considered next.

5.0 USING THE PROGRAM ON A SAMPLE PROBLEM

5.1 Preparing the Input Data

Perhaps the best method of discussing the use of the program is to consider a sample problem of the type outlined in section 2.0. Formulation of an input data set specifying all input parameters discussed in section 4.2B must be obtained before execution. Let us assume we are given a rectangular toroid of ferromagnetic material as illustrated by Fig. 7, with an initial field intensity of H_0 . We then reverse the applied field such that the boundary at time = 0^+ is at $-H_0$. If H_0 is sufficient to saturate the core a flux reversal will occur and a flux distribution will exist in the core during the transient state. We can determine the distribution pattern as a function of time by solving for $H(I,J)$, $B(I,J)$, and $FORCE(I,J)$ as a function of time. Let us assume that we desire to investigate the time dependence of the distribution pattern and the relationship between total flux (FLUXT), total force (FORCT), and time. Typical numerical values may be as given below.

- Given:
- A. Material is 2.5 per cent silicon iron with a hysteresis loop given by Fig. 22.
 - B. Material conductivity (COND) - $COND = 2.5 \times 10^{-7}$ mhos
 - C. Material dimensions (see Fig. 12)
 - $AX = 0.004$ meter
 - $BY = 0.008$ meter
 - D. Initial conditions
 - $H_0 = 250$ amp-t/m
 - $T_0 = 0.0$ sec

- G. Maximum field intensity - HMAX = 250 amp-t/m
(see section 3.4)
- H. Boundary conditions
 1. HO(I,NX), I = 2, NY - HO = -250 amp-t/m
 2. HO(NY,J), J = 2, NX

The above information is sufficient to calculate numerical values for the remaining input data variables. (See section 4.2B.) We will use the mks system of units; thus the constant $\mu_0 = 4\pi \times 10^{-7}$ webers/amp-t/m = $4\pi\text{CUO}$; hence

$$\text{CUO} = 1.0 \times 10^{-7} \text{ weber/amp-t/meter.}$$

We obtain BR, B2, H2, B3, and H3 from the hysteresis loop for the given material. An approximation to this loop is shown by Fig. 22. Typical values for 2.5 per cent silicon iron are given below.

$$\begin{aligned} \text{BR} &= 0.71 \text{ weber/m}^2 \\ \text{H2} &= 11.94 \text{ amp-t/m} \\ \text{B2} &= 0.80 \text{ weber/m}^2 \\ \text{H3} &= 103.50 \text{ amp-t/m} \\ \text{B3} &= 1.20 \text{ webers/m}^2 \end{aligned}$$

Values of H10, DHX, and I3 to yield a symmetrical hysteresis loop approximation with a range of H from +400 amp-t/m to -400 amp-t/m are

$$\begin{aligned} \text{H10} &= 400 \\ \text{DHX} &= 10 \\ \text{I3} &= 81 \end{aligned}$$

Note: Maximum value of I3 = 100
 For symmetrical loop
 $\text{I3} = 1 + 2(\text{H10}/\text{DHX})$

Grid size is selected by specifying NI, the number of divisions desired in the X-direction. (See Fig. 12.) This also specifies the quantity NJ since the grid is square and $NJ = (BY)(AX)/4NI$; hence we shall use $NI = 4$ and $NJ = 8$. The quantity CDT must be less than 0.5 to insure convergence, hence $CDT = 0.2$. Remaining variables of the first five cards control the program as discussed in section 4.2B. We desire to begin at $T(K) = 0$, hence let $K = 0$; to print-out all data available in both formats, hence $I4 = 1$; to investigate a flux decay, hence $I5 = 2$; to plot every value of FORCT and FLUXT versus time $T(K)$, hence $KI = 1$; to print-out the distribution pattern for all variables $H(I,J)$, $B(I,J)$, and $FORCE(I,J)$ for every fifth value of $T(K)$, hence $KG = 5$; and to execute this input data set third from the last, hence $KZ = 3$. Collecting the above quantities in a list we have

CDT = 0.2	K = 0	KI = 1
NI = 4	I4 = 1	KG = 5
NJ = 8	I5 = 2	KZ = 3

Variable EO is the maximum error norm E, thus if we wish the total error norm for one-fourth of the lattice to be less than or equal to 1.0, we require $EO < 1.0$. Variables HSO, TKM, and KO halt execution. (See section 4.2B.) Let us halt when $T(K)$ becomes ≥ 20.0 milliseconds, or when $K \geq 100$, and not stop for the condition $HS \leq HSO$, hence we have

$$EO = 1.0$$

$$HSO = -1.0 \times 10^{-7}$$

$$TKM = 20.0$$

$$KO = 100$$

$$\text{Note: } 1. T(K) = K \Delta t$$

$$2. \Delta t = CDT \frac{\sigma C_1 C_2 h^2}{(C_2 + HMAX)^2}$$

This completes formulation of a typical input data set.

If we collect these values, sequence them according to the "read" instructions (see section 4.2B) and express the numbers in accordance with the format specifications, we have a data set as given by Table 3.

5.2 Output Data

Output data are available on Tapes 6, 9, and 15. The output data set included was generated by execution of the program with input data set 3. $I_4 = 1$ was used to print-out all available data. Due to the voluminous amount of data available only a small number of distribution patterns are included; however, patterns for all increments of time considered may be obtained by specifying $KG = 1$. Duplications of page headings, etc., were deleted. When the execution is stopped, the halt condition is indicated (see page 101) and all data for the last value of time are printed-out on all tapes.

5.2A Output Tape 6 (Pages 71 - 82)

Tape 6 lists values for the approximate hysteresis loop (pages 71 - 74) and provides a graphical representation of the

approximation (page 75). A listing of the total force, total flux, and the corresponding time then follows, (pages 78 and 79). Next is a graphical output illustrating the time dependence of total force and total flux, (pages 80 and 81, respectively). At completion of the run, a print-out of the continuation data available as a punched card output of Tape 16 occurs.

5.2B Output Tape 9 (Pages 83 - 92)

The first five pages of print-out for this tape duplicate pages 71 - 74 and 76 of Tape 6 output and are not included; however, the following page lists initial and boundary conditions at time $T(K) = 0^+$, (page 83). A listing of corresponding values of flux density, force, and time are also given at $T(K) = 0^+$. As time increases in increments of Δt , a complete list specifying distribution patterns at $T(K)$ ($K = 0, KG, 2KG, 3KG, \dots$) is listed according to format specifications as shown by pages 83 - 92. This tape also contains a listing of iterations and their associated error norm E (pages 84 and 86).

5.2C Output Tape 15 (Pages 93 - 100)

Tape 15 gives a more convenient output format since data are given in a matrix format which eases interpretation by locating values calculated for specific points in the same physical location as they would occupy in the matrix of Fig. 12. All

data available on Tape 9 except error norm values are available on Tape 15. Total values of force and flux are listed following each matrix output. Duplicates of pages 71-74 and 76 of Tape 9 precede the output and are not included.

0	14	725388E	01	-0	0	27000000E	03	94
0	14	625316E	01	-0	0	26000000E	03	95
0	14	520148E	01	-0	0	25000000E	03	96
0	14	05484E	01	-0	0	23999999E	03	97
0	14	292881E	01	-0	0	23000000E	03	98
0	14	169847E	01	-0	0	22000000E	03	99
0	14	039937E	01	-0	0	20999999E	03	100
0	13	902239E	01	-0	0	20000000E	03	101
0	13	756369E	01	-0	0	19000000E	03	102
0	13	601458E	01	-0	0	18000000E	03	103
0	13	436639E	01	-0	0	16999999E	03	104
0	13	260929E	01	-0	0	16000000E	03	105
0	13	073211E	01	-0	0	15000000E	03	106
0	12	872214E	01	-0	0	13999999E	03	107
0	12	656473E	01	-0	0	13000000E	03	108
0	12	424308E	01	-0	0	12000000E	03	109
0	12	173766E	01	-0	0	11000000E	03	110
0	11	902577E	01	-0	0	09999999E	03	111
0	11	608078E	01	-0	0	90000000E	02	112
0	11	287129E	01	-0	0	80000000E	02	113
0	10	935999E	01	-0	0	70000000E	02	114
0	10	550223E	01	-0	0	59999999E	02	115
0	10	124402E	01	-0	0	50000000E	02	116
0	9	6519651E	00	-0	0	40000000E	02	117
0	8	5228485E	00	-0	0	30000000E	02	118
0	7	8633461E	00	-0	0	20000000E	02	119
0	7	0999999E	00	-0	0	09999999E	02	120
0	6	215918E	00	0	0	99999999E	02	121
0	5	1999857E	00	0	0	20000000E	02	122
0	3	9970580E	-00	0	0	30000000E	02	123
0	2	598342E	-00	0	0	40000000E	02	124
0	8	1247769E	-01	0	0	50000000E	02	125
0	1	1858648E	-00	0	0	59999999E	02	126
0	2	8646551E	-00	0	0	70000000E	02	127
0	4	2506013E	-00	0	0	80000000E	02	128
0	6	4049356E	00	0	0	90000000E	02	129
0	7	2587077E	00	0	0	09999999E	03	130
0	8	020694E	00	0	0	10000000E	03	131
0	8	6551375E	00	0	0	12000000E	03	132
0	9	2334213E	00	0	0	13000000E	03	133
0	9	7490710E	00	0	0	13999999E	03	134
0	10	211739E	01	0	0	15000000E	03	135
0	10	629193E	01	0	0	16000000E	03	136
0	11	007752E	01	0	0	16999999E	03	137
0	11	352608E	01	0	0	18000000E	03	138
0	11	668073E	01	0	0	18000000E	03	139
0	11	957748E	01	0	0	19000000E	03	140
0	12	224673E	01	0	0	20000000E	03	141
0	12	471426E	01	0	0	20999999E	03	142
0	12	471426E	01	0	0	22000000E	03	143
0	12	471426E	01	0	0	23000000E	03	144

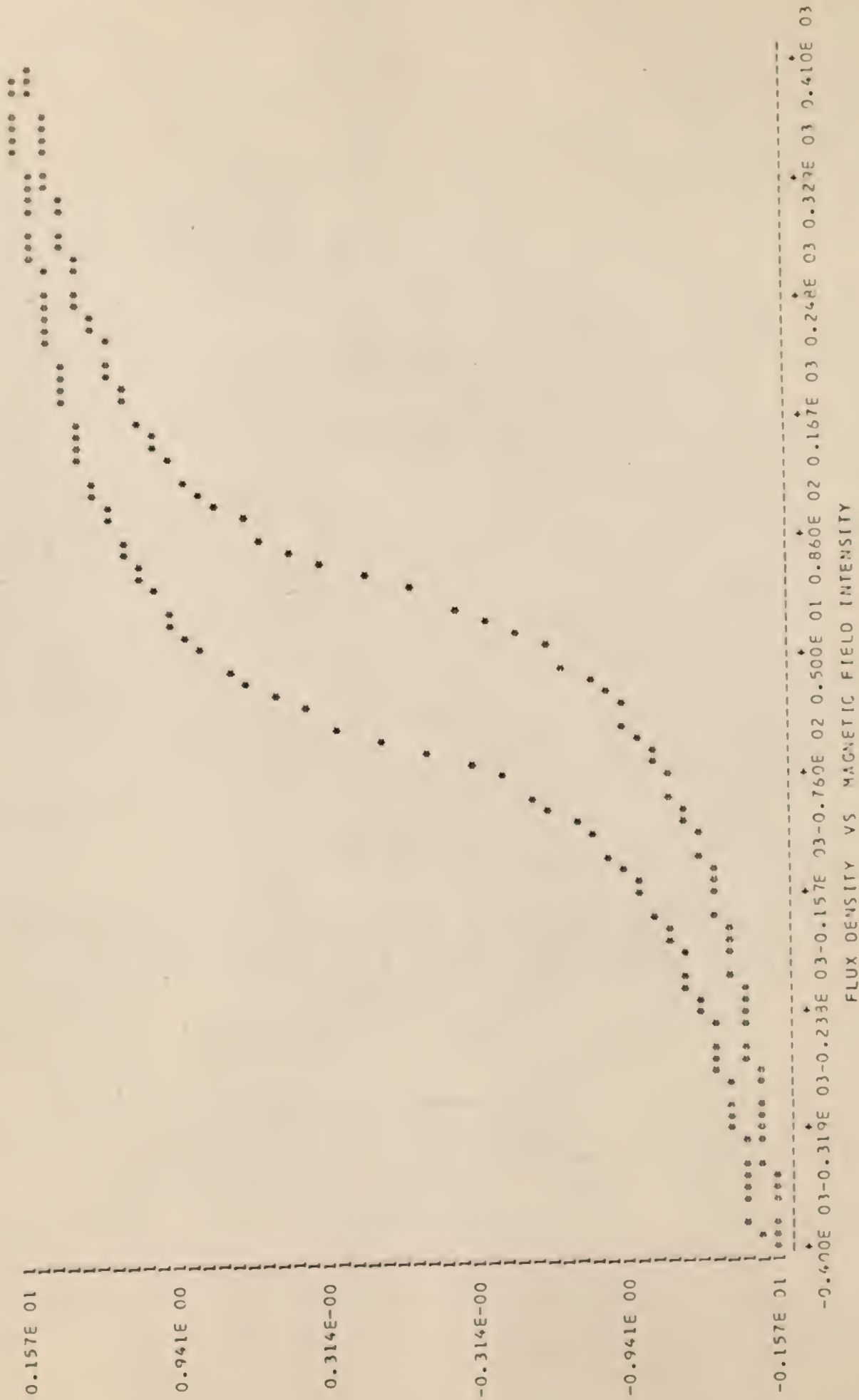


Fig. 22. Printer-plot of approximate hysteresis loop.

THIS IS A DECAY RUN
 LATTICE COVERS ONLY 1/4 OF THE CROSS SECTION AREA
 LATTICE SIZE = NI BY NJ = 4 BY 8 (MAX = 48 BY 48)
 NI = 4 NJ = 8 ANJ = 0.8000000E 01 CDT = 0.2000000E-00

CX	0.499999999E-03	0.11655107E 01	0.14268442E 03	0.13480794E-03	DT
EO		HSO			HMAX
BR	0.099999999E 01	-0.099999999E 08	0.709999999E 00	0.25000000E 03	CUO
HMAX		TKM			BY
CUO		0.20000000E-01	0.25000000E 07	0.399999999E-02	0.799999999E-02
TO		KZ	KI	KG	
		NI	KO		
		NJ	0		
		4	100		
CDT	0.20000000E-00	H2	0		
B2		DHX	13		
0.80000000E 00	0.11939999E 02	0.12000000E 01	14		
H10			15		
			1		
0.40000000E 03	0.09999999E 02	81	2		

EO = MAXIMUM MATRIX ERROR
 HSO = MINIMUM SUM OF HI(I,J)
 BR = RESIDUAL FLUX DENSITY
 HMAX = MAXIMUM FIELD INTENSITY APPLIED
 CUO = CONSTANT USED TO ALLOW ALL UNITS
 TO = INITIAL VALUE OF TIME IN SECONDS
 TKM = SPECIFIES MAXIMUM TIME
 COND = MATERIAL CONDUCTIVITY
 AX = CORE WIDTH
 BY = CORE HEIGHT
 KZ = NO OF SETS OF DATA
 NI = CONSTANT MULTIPLIER FOR THE TIME INCREMENT
 NJ = NO OF DIVISIONS IN THE X DIRECTION
 K0 = NO OF DIVISIONS IN THE Y DIRECTION
 KI = NO OF INCREMENTS OR TIME SUBSCRIPT
 H2 = MAXIMUM NO OF INCREMENTS PER RUN
 H3 = SPECIFIES INCREMENT FOR FORCT AND FLUX PLOTS
 H10 = SPECIFIES INCREMENT FOR OUTPUT H1, B1, FORCE
 DHX = FIELD INTENSITY CORRESPONDING TO B2 VALUE
 I3 = FLUX DENSITY FROM HYSTERESIS CURVE
 I4 = FIELD INTENSITY CORRESPONDING TO B3 VALUE
 I5 = MAXIMUM VALUE OF FIELD INTENSITY FOR B VS H
 I4 = INCREMENT NO OF POINTS ON B VS H CURVE
 I5 = USED TO SPECIFY OUTPUT FORMAT
 USE 1 FOR OUTPUT TAPES 9 AND 15 (ALL DATA)
 USE 2 FOR OUTPUT TAPE 9 ONLY (ALL DATA)
 USE 3 FOR OUTPUT TAPE 15 ONLY (H1)
 USE 4 FOR OUTPUT TAPE 15 ONLY (B1)
 USE 5 FOR OUTPUT TAPE 15 ONLY (FORCE)
 USE 6 FOR OUTPUT TAPE 15 ONLY (H1, B1, FORCE)
 USE 7 FOR OUTPUT TAPE 15 ONLY (ITERATIONS ON 9)
 I5 = VALUE TO SPECIFY BUILDUP OR DECAY

CX USE 0 OR 1 FOR BUILDUP
 CI USE 2 FOR DECAY
 C2 DIVISION WIDTHS
 DT = CONSTANT IN MODIFIED FROELICH APPROXIMATION
 = CONSTANT IN MODIFIED FROELICH APPROXIMATION
 = TIME INCREMENT IN SECONDS

K	0	1	2	3	4	5	6	7	8	9	10	11	12	13	14	15	16	17	18	19	20	21	22	23	24	25	26	27	28	29	30	31	32	33	34	35	36	37	38	39	40	41	42	43	44	45	46	47	48	49	50
TOTAL FORCE (FORCT)	0.64588968E 01	0.61366858E 01	0.58863641E 01	0.56693082E 01	0.54746731E 01	0.52970427E 01	0.51331148E 01	0.49807078E 01	0.483382594E 01	0.47045754E 01	0.45787790E 01	0.44601419E 01	0.43480699E 01	0.42420702E 01	0.41417268E 01	0.40456684E 01	0.395566297E 01	0.38711483E 01	0.37897687E 01	0.37121330E 01	0.36379664E 01	0.35670491E 01	0.34992010E 01	0.34342711E 01	0.33721296E 01	0.33126634E 01	0.32557724E 01	0.32013676E 01	0.31493678E 01	0.30997001E 01	0.30522969E 01	0.30070964E 01	0.29640412E 01	0.29230534E 01	0.28840140E 01	0.28468167E 01	0.28113715E 01	0.27776012E 01	0.27454463E 01	0.27148397E 01	0.26857109E 01	0.26579735E 01	0.26315439E 01	0.26063437E 01	0.25822925E 01	0.25593032E 01	0.25372707E 01	0.25160980E 01	0.24957021E 01	0.24760111E 01	0.24569634E 01
TIME (K)	-0.13480794E-03	0.26961589E-03	0.40442384E-03	0.53923178E-03	0.67403973E-03	0.80884767E-03	0.94365562E-03	0.10784635E-02	0.12132715E-02	0.13480794E-02	0.14828873E-02	0.16176953E-02	0.17525032E-02	0.18873112E-02	0.20221191E-02	0.21569270E-02	0.22917350E-02	0.24265429E-02	0.25613508E-02	0.26961587E-02	0.28309666E-02	0.29657745E-02	0.31005825E-02	0.32353904E-02	0.33701983E-02	0.35050063E-02	0.36398142E-02	0.37746221E-02	0.39094301E-02	0.40442379E-02	0.41790458E-02	0.43138537E-02	0.44486616E-02	0.45834695E-02	0.47182774E-02	0.48530853E-02	0.49878932E-02	0.51227011E-02	0.52575090E-02	0.53923169E-02	0.55271248E-02	0.56619327E-02	0.57967406E-02	0.59315485E-02	0.60663564E-02	0.62011642E-02	0.63359721E-02	0.64707800E-02	0.66055879E-02	0.67403958E-02	
TOTAL FLUX (FLUXT)	0.76726136E-05	0.73824580E-05	0.71434717E-05	0.69226922E-05	0.67250592E-05	0.65340440E-05	0.63515279E-05	0.61759755E-05	0.60062890E-05	0.58416083E-05	0.56813462E-05	0.55249920E-05	0.53721312E-05	0.52224213E-05	0.50755748E-05	0.49313471E-05	0.47897567E-05	0.46521715E-05	0.45183534E-05	0.43878131E-05	0.42601639E-05	0.41350934E-05	0.40123452E-05	0.38917053E-05	0.37729935E-05	0.36560558E-05	0.35407588E-05	0.342698867E-05	0.33146381E-05	0.32036224E-05	0.30938599E-05	0.29852786E-05	0.28778397E-05	0.27722174E-05	0.26686006E-05	0.256667510E-05	0.246667510E-05	0.23682502E-05	0.22712404E-05	0.21756566E-05	0.20819284E-05	0.19901370E-05	0.19002638E-05	0.18123583E-05	0.17264927E-05	0.16430592E-05	0.15622683E-05	0.14839365E-05	0.14078944E-05	0.13339885E-05	0.12620801E-05

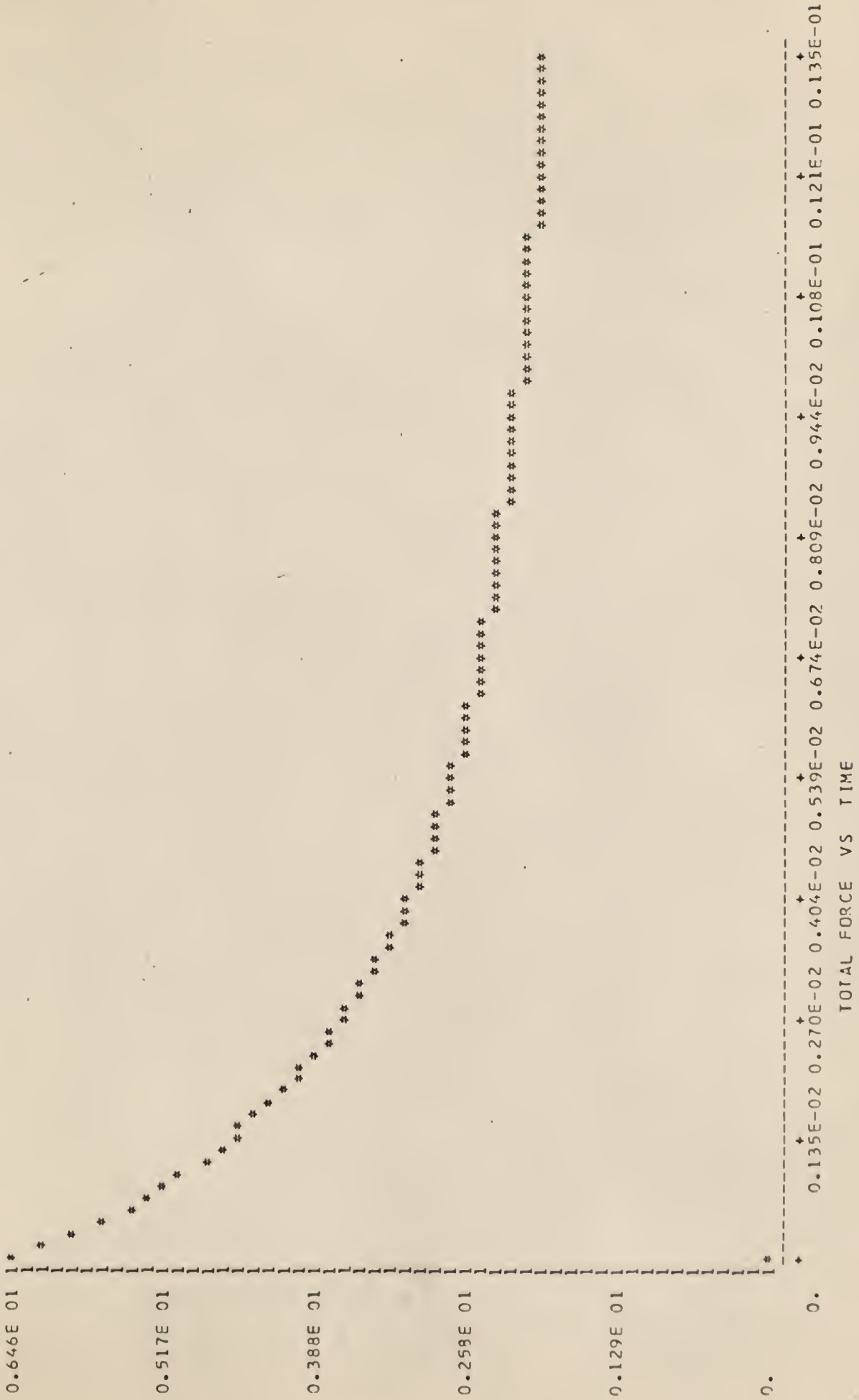


Fig. 23. Printer-plot of total force (newtons) vs time (sec).

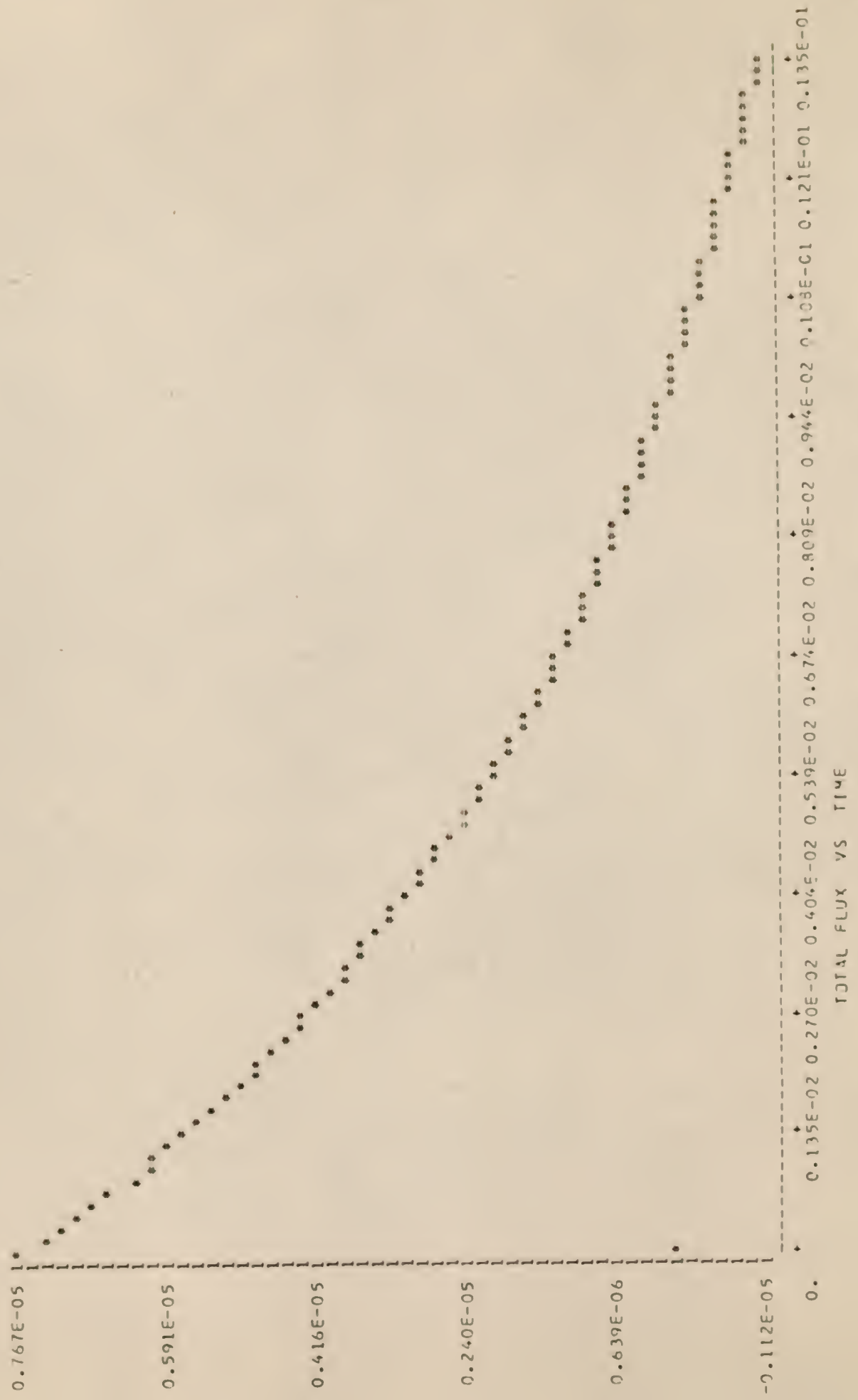


Fig. 24. Printer-plot of total flux (webers) vs time (sec).

THIS IS OUTPUT DATA TO BE USED FOR A CONTINUATION RUN. CHANGE K TO K = 0

0.09999999E 01	-0.09999999E 08	0.70999999E 00	0.25000000E 03	0.99999999E-07
0.13480790E-01	0.20000000E-01	0.25000000E 07	0.39999999E-02	0.79999999E-02
0.20000000E-00	0.3	100	5	
0.80000000E 00	0.11939999E 02	0.12000000E 01	0.10349999E 03	
0.40000000E 03	0.09999999E 02	81		
0.21705081E 02	0.11056926E 02	-0.22811322E 02	-0.88754325E 02	-0.25000000E 03
0.21169190E 02	0.10602032E 02	-0.23092266E 02	-0.88822591E 02	-0.25000000E 03
0.19324502E 02	0.89978344E 01	-0.24069635E 02	-0.89066953E 02	-0.25000000E 03
0.15366369E 02	0.55395380E 01	-0.26207998E 02	-0.89625838E 02	-0.25000000E 03
0.75987671E 01	-0.12978579E 02	-0.30529381E 02	-0.90811701E 02	-0.25000000E 03
-0.72812592E 01	0.14556853E 02	-0.39210112E 02	-0.93332411E 02	-0.25000000E 03
-0.35769892E 02	-0.40476813E 02	-0.57302518E 02	-0.99157947E 02	-0.25000000E 03
-0.92339200E 02	-0.93752974E 02	-0.99211878E 02	-0.116933345E 03	-0.25000000E 03
-0.25000000E 03	-0.25000000E 03	-0.25000000E 03	-0.25000000E 03	-0.25000000E 03

(KZ = 3 PREVIOUS DATA FOR DATA SET NO. 3)

FORCE(I, J)	B(I, J)	H(I, J)	TIME(S)	(I, J)
0.51868869E-01	0.14442192E	0.24290222E	0.67403973E-03	2
0.94222475E-00	0.143336488E	0.23367802E	0.67403973E-03	2
0.94765411E-01	0.138033514E	0.19316754E	0.67403973E-03	2
0.62529676E-01	0.1212639E	0.77800727E	0.67403973E-03	2
0.41465758E-01	0.12912922E	-0.25000000E	0.67403973E-03	2
0.10373606E-00	0.14442076E	0.24289179E	0.67403973E-03	3
0.20444663E-00	0.143336387E	0.23366942E	0.67403973E-03	3
0.18952918E-00	0.138033454E	0.19316351E	0.67403973E-03	3
0.12505901E-00	0.11212623E	0.77800281E	0.67403973E-03	3
0.82931516E-01	0.12912922E	-0.25000000E	0.67403973E-03	3
0.10372317E-00	0.14441179E	0.24281161E	0.67403973E-03	4
0.20442403E-00	0.143335594E	0.23360194E	0.67403973E-03	4
0.18951543E-00	0.13802953E	0.19312953E	0.67403973E-03	4
0.12505569E-00	0.11212474E	0.77795920E	0.67403973E-03	4
0.82931516E-01	0.12912922E	-0.25000000E	0.67403973E-03	4
0.10364556E-00	0.14435775E	0.24232931E	0.67403973E-03	5
0.20428472E-00	0.14330709E	0.23318644E	0.67403973E-03	5
0.18942284E-00	0.13799581E	0.19290102E	0.67403973E-03	5
0.12502786E-00	0.11211227E	0.77759462E	0.67403973E-03	5
0.82931516E-01	0.12912922E	-0.25000000E	0.67403973E-03	5
0.10323692E-00	0.14407289E	0.23980681E	0.67403973E-03	6
0.20335325E-00	0.14304301E	0.23095629E	0.67403973E-03	6
0.18887325E-00	0.13779548E	0.19154983E	0.67403973E-03	6
0.12481143E-00	0.11201519E	0.77476157E	0.67403973E-03	6
0.82931516E-01	0.12912922E	-0.25000000E	0.67403973E-03	6
0.10140903E-00	0.14278539E	0.22880605E	0.67403973E-03	7
0.20006448E-00	0.14181910E	0.22095668E	0.67403973E-03	7
0.18607759E-00	0.13677187E	0.18481231E	0.67403973E-03	7
0.12331796E-00	0.1134300E	0.75534222E	0.67403973E-03	7
0.82931516E-01	0.12912922E	-0.25000000E	0.67403973E-03	7
0.94122320E-01	0.13756634E	0.19001767E	0.67403973E-03	8
0.18599571E-00	0.13674177E	0.18461831E	0.67403973E-03	8
0.17366058E-00	0.13212968E	0.15738061E	0.67403973E-03	8
0.11450005E-00	0.10728835E	0.64510448E	0.67403973E-03	8
0.82931516E-01	0.12912922E	-0.25000000E	0.67403973E-03	8
0.62283787E-01	0.11190571E	0.77157515E	0.67403973E-03	9
0.12328993E-00	0.11133034E	0.75497995E	0.67403973E-03	9
0.11449324E-00	0.10728751E	0.64508279E	0.67403973E-03	9
0.64406883E-01	0.80666644E	0.12661381E	0.67403973E-03	9
0.82931516E-01	0.12912922E	-0.25000000E	0.67403973E-03	10
0.41465758E-01	0.12912922E	-0.25000000E	0.67403973E-03	10
0.82931516E-01	0.12912922E	-0.25000000E	0.67403973E-03	10
0.82931516E-01	0.12912922E	-0.25000000E	0.67403973E-03	10
0.41465758E-01	0.12912922E	-0.25000000E	0.67403973E-03	10

TOTAL FORCE (FORCI) = 0.52970427E 01
TOTAL FLUX (FLUKT) = 0.65340440E-05
TIME(S) = 0.67403973E-03
TIME(S) = 0.67403973E-03

NN =
 NN =
 NN =
 NN =
 TIME(6) = 0.80884767E-03
 E = 0.22873706E 02
 E = 0.21933233E 01
 E = 0.23584962E-00

NN =
 NN =
 NN =
 NN =
 TIME(7) = 0.94365562E-03
 E = 0.17104258E 02
 E = 0.15590467E 01
 E = 0.13956106E-00

NN =
 NN =
 NN =
 NN =
 TIME(8) = 0.10784635E-02
 E = 0.13121015E 02
 E = 0.11226968E 01
 E = 0.95485449E-01

NN =
 NN =
 NN =
 TIME(9) = 0.12132715E-02
 E = 0.10537402E 02
 E = 0.81881452E 00

NN =
 NN =
 NN =
 TIME(10) = 0.13480794E-02
 E = 0.87489747E 01
 E = 0.60457301E 00

FORC (I, J)	BI (I, J)	HI (I, J)	TIME (I, J)	(I, J)
0.49851105E-01	0.14158496E	0.21910437E	0.13480794E-02	2
0.96734820E-01	0.13949812E	0.20339263E	0.13480794E-02	2
0.84842705E-01	0.13060828E	0.14936348E	0.13480794E-02	2
0.42267073E-01	0.92186183E	0.31698651E	0.13480794E-02	2
0.41465758E-01	-0.12912922E	-0.25000000E	0.13480794E-02	2
0.99680693E-01	0.14154968E	0.21893415E	0.13480794E-02	3
0.19352264E-00	0.13948479E	0.20322663E	0.13480794E-02	3
0.16966254E-00	0.92183325E	0.31693427E	0.13480794E-02	3
0.84528904E-01	-0.12912922E	-0.25000000E	0.13480794E-02	3
0.82931516E-01	0.14149047E	0.21836215E	0.13480794E-02	4
0.99569172E-01	0.13941433E	0.20279021E	0.13480794E-02	4
0.19333717E-00	0.13055233E	0.14907693E	0.13480794E-02	4
0.16953912E-00	0.92164657E	0.31659298E	0.13480794E-02	4
0.84494670E-01	-0.12912922E	-0.25000000E	0.13480794E-02	4
0.82931516E-01	0.14119884E	0.21609063E	0.13480794E-02	5
0.99159152E-01	0.13914831E	0.20089148E	0.13480794E-02	5
0.19260000E-00	0.13035659E	0.1486394E	0.13480794E-02	5
0.16903107E-00	0.92069565E	0.31486394E	0.13480794E-02	5
0.84321139E-01	-0.12912922E	-0.25000000E	0.13480794E-02	5
0.82931516E-01	0.14025009E	0.20889440E	0.13480794E-02	6
0.97831076E-01	0.13826019E	0.19470092E	0.13480794E-02	6
0.19014932E-00	0.12964116E	0.14448617E	0.13480794E-02	6
0.16718078E-00	0.91630560E	0.30688532E	0.13480794E-02	6
0.83518212E-01	-0.12912922E	-0.25000000E	0.13480794E-02	6
0.82931516E-01	0.13741264E	0.18899771E	0.13480794E-02	7
0.93912602E-01	0.13553961E	0.17705328E	0.13480794E-02	7
0.18273975E-00	0.12725488E	0.13312110E	0.13480794E-02	7
0.16108289E-00	0.89759270E	0.27370962E	0.13480794E-02	7
0.80141807E-01	-0.12912922E	-0.25000000E	0.13480794E-02	7
0.82931516E-01	0.12878161E	0.14023606E	0.13480794E-02	8
0.82485616E-01	0.12709592E	0.13239589E	0.13480794E-02	8
0.16068070E-00	0.11929654E	0.10096276E	0.13480794E-02	8
0.14156511E-00	0.82936305E	0.14924062E	0.13480794E-02	8
0.66944157E-01	-0.12912922E	-0.25000000E	0.13480794E-02	8
0.82931516E-01	0.91105458E	0.25000000E	0.13480794E-02	9
0.41281864E-01	0.89668476E	0.27213221E	0.13480794E-02	9
0.79979756E-01	0.82028227E	0.14911997E	0.13480794E-02	9
0.66930976E-01	0.39399585E	0.30434300E	0.13480794E-02	9
0.15441298E-01	-0.12912922E	-0.25000000E	0.13480794E-02	9
0.82931516E-01	0.12912922E	0.25000000E	0.13480794E-02	10
0.41465758E-01	-0.12912922E	-0.25000000E	0.13480794E-02	10
0.82931516E-01	0.12912922E	0.25000000E	0.13480794E-02	10
0.82931516E-01	-0.12912922E	-0.25000000E	0.13480794E-02	10
0.41465758E-01	-0.12912922E	-0.25000000E	0.13480794E-02	10

TOTAL FORCE (FORCF) = 0.45787790E 01 TIME(I0) = 0.13480794E-02
TOTAL FLUX (FLUXF) = 0.56813462E-05 TIME(I0) = 0.13480794E-02

FORCE(I, J)	BI(I, J)	HI(I, J)	TIME(20)	(I, J)
0.44763909E-01	1.3410639E	1.6882850E	0.26961587E-02	2 2
0.85440540E-01	1.3106801E	1.5174057E	0.26961587E-02	2 3
0.69391819E-01	1.1811876E	0.95822975E	0.26961587E-02	2 4
0.18530010E-01	0.61038322E	-0.112335042E	0.26961587E-02	2 5
0.41465758E-01	1.29129222E	-0.25000000E	0.26961587E-02	2 6
0.89430710E-01	1.34009360E	0.15138979E	0.26961587E-02	2 3
0.17070544E-00	1.3100064E	0.15138979E	0.26961587E-02	3 3
0.37866169E-00	1.1806685E	0.96650832E	0.26961587E-02	3 4
0.37026459E-01	1.29129222E	-0.112335042E	0.26961587E-02	3 5
0.82931516E-01	1.29129222E	-0.112335042E	0.26961587E-02	3 6
0.89041477E-01	1.3380148E	0.16671370E	0.26961587E-02	4 4
0.13814707E-00	1.3072715E	0.14997443E	0.26961587E-02	4 4
0.13814707E-00	1.1784755E	0.95897350E	0.26961587E-02	4 4
0.36868524E-01	0.60880501E	-0.11398818E	0.26961587E-02	4 4
0.82931516E-01	1.29129222E	-0.112335042E	0.26961587E-02	4 4
0.88003264E-01	1.3301913E	0.16227492E	0.26961587E-02	4 5
0.16806377E-00	1.29983307E	0.14619173E	0.26961587E-02	5 5
0.3667153E-01	1.1721651E	0.93757255E	0.26961587E-02	5 5
0.36339643E-01	0.60442174E	-0.11851547E	0.26961587E-02	5 5
0.82931516E-01	1.29129222E	-0.112335042E	0.26961587E-02	5 5
0.85464005E-01	1.3108601E	0.15183441E	0.26961587E-02	6 6
0.16326858E-00	1.2811548E	0.13717461E	0.26961587E-02	6 6
0.13279344E-00	1.1554152E	0.8257473E	0.26961587E-02	6 6
0.34705687E-01	0.59067702E	-0.13251153E	0.26961587E-02	6 6
0.82931516E-01	1.29129222E	-0.112335042E	0.26961587E-02	6 7
0.79392660E-01	1.2634408E	0.12901696E	0.26961587E-02	6 7
0.15164423E-00	1.2347035E	0.11683237E	0.26961587E-02	6 7
0.12291981E-00	1.1116311E	0.75020336E	0.26961587E-02	7 7
0.29956402E-01	0.54877533E	-0.17338986E	0.26961587E-02	7 7
0.82931516E-01	1.29129222E	-0.112335042E	0.26961587E-02	7 7
0.64003878E-01	1.1344043E	0.81710129E	0.26961587E-02	8 8
0.12185185E-00	1.1067915E	0.73649745E	0.26961587E-02	8 8
0.96674807E-01	0.98583994E	0.44238906E	0.26961587E-02	8 8
0.16887335E-01	0.41203140E	-0.29050958E	0.26961587E-02	8 8
0.82931516E-01	1.29129222E	-0.112335042E	0.26961587E-02	8 9
0.16349057E-01	0.57333851E	-0.14974566E	0.26961587E-02	9 9
0.29524913E-01	0.54480874E	-0.17712619E	0.26961587E-02	9 9
0.16849285E-01	0.41156694E	-0.29087009E	0.26961587E-02	9 9
0.48887759E-02	0.22169189E	-0.65901104E	0.26961587E-02	9 9
0.82931516E-01	1.29129222E	-0.112335042E	0.26961587E-02	10 10
0.41465758E-01	1.29129222E	-0.112335042E	0.26961587E-02	10 10
0.82931516E-01	1.29129222E	-0.112335042E	0.26961587E-02	10 10
0.82931516E-01	1.29129222E	-0.112335042E	0.26961587E-02	10 10
0.41465758E-01	1.29129222E	-0.112335042E	0.26961587E-02	10 10

TOTAL FORCE (FORCT) = 0.36379664E 01 TIME(20) = 0.26961587E-02
TOTAL FLUX (FLUXT) = 0.42601639E-05 TIME(20) = 0.26961587E-02

FORCE (I, J)	BL (I, J)	HL (I, J)	TIME (40)	(I, J)
0.3576449E-01	0.119922383E 01	0.10322275E 03	0.53923169E-02	1 2
0.66466665E-01	0.115602235E 01	0.88452752E 02	0.53923169E-02	2 3
0.47051156E-03	0.97263501E 00	0.41505075E 02	0.53923169E-02	2 3
0.48473524E-01	0.98722704E-01	-0.49088437E 03	0.53923169E-02	2 3
0.41465758E-01	-0.12912922E 01	-0.25000000E 03	0.53923169E-02	2 3
0.71327852E-01	0.11975518E 01	0.10261101E 02	0.53923169E-02	2 3
0.13255854E-00	0.115443927E 01	0.87930042E 02	0.53923169E-02	3 3
0.93818423E-03	0.97116685E 00	0.41206053E 02	0.53923169E-02	3 3
0.94599222E-03	0.97519964E-01	-0.47151751E 02	0.53923169E-02	3 3
0.82931516E-01	-0.12912922E 01	-0.25000000E 03	0.53923169E-02	3 3
0.70601285E-01	0.11914369E 01	0.10041833E 02	0.53923169E-02	4 4
0.13119754E-00	0.11484513E 01	0.86045487E 02	0.53923169E-02	4 4
0.92769228E-01	0.96572138E 00	0.40105393E 02	0.53923169E-02	4 4
0.85610344E-03	0.92771135E-01	-0.49400901E 02	0.53923169E-02	4 4
0.82931516E-01	-0.12912922E 01	-0.25000000E 03	0.53923169E-02	4 4
0.68914851E-01	0.11771211E 01	0.95435143E 02	0.53923169E-02	5 5
0.12801585E-00	0.11344402E 01	0.81720982E 02	0.53923169E-02	5 5
0.90253160E-01	0.95253503E 00	0.37493915E 02	0.53923169E-02	5 5
0.63924025E-03	0.80164447E-01	-0.50055926E 02	0.53923169E-02	5 5
0.82931516E-01	-0.12912922E 01	-0.25000000E 03	0.53923169E-02	5 5
0.65268479E-01	0.11455565E 01	0.72675309E 02	0.53923169E-02	6 6
0.12108719E-00	0.11033132E 01	0.85133845E 02	0.53923169E-02	6 6
0.84640805E-01	0.92244322E 00	0.31815028E 02	0.53923169E-02	6 6
0.23541576E-03	0.49648283E-01	-0.51653909E 02	0.53923169E-02	6 6
0.82931516E-01	-0.12912922E 01	-0.25000000E 03	0.53923169E-02	6 6
0.57332665E-01	0.10736576E 01	0.64710458E 02	0.53923169E-02	6 6
0.10592006E-00	0.10319030E 01	0.54445374E 02	0.53923169E-02	7 7
0.72158296E-01	0.85171223E 00	0.19750106E 02	0.53923169E-02	7 7
0.90730357E-04	-0.30201412E-01	-0.55466334E 02	0.53923169E-02	7 7
0.82931516E-01	-0.12912922E 01	-0.25000000E 03	0.53923169E-02	7 7
0.38818311E-01	0.88345791E 00	0.24948010E 02	0.53923169E-02	8 8
0.70459558E-01	0.84162710E 00	0.18165620E 02	0.53923169E-02	8 8
0.43156781E-01	0.65875600E 00	-0.60091966E 01	0.53923169E-02	8 8
0.41887528E-02	-0.20520698E-01	0.64908683E 02	0.53923169E-02	8 8
0.82931516E-01	-0.12912922E 01	-0.25000000E 03	0.53923169E-02	9 9
0.2548614E-05	0.71584135E-02	-0.53675393E 02	0.53923169E-02	9 9
0.16373291E-03	-0.40571190E-01	-0.55975630E 02	0.53923169E-02	9 9
0.42287968E-02	-0.20618552E-00	0.64967045E 02	0.53923169E-02	9 9
0.32571247E-01	-0.57222515E 00	0.92946576E 02	0.53923169E-02	9 9
0.82931516E-01	-0.12912922E 01	-0.25000000E 03	0.53923169E-02	10 10
0.41465758E-01	-0.12912922E 01	-0.25000000E 03	0.53923169E-02	10 10
0.82931516E-01	-0.12912922E 01	-0.25000000E 03	0.53923169E-02	10 10
0.82931516E-01	-0.12912922E 01	-0.25000000E 03	0.53923169E-02	10 10
0.41465758E-01	-0.12912922E 01	-0.25000000E 03	0.53923169E-02	10 10

TOTAL FORCE (FORCT) = 0.26857109E 01 TIME(40) = 0.53923169E-02
TOTAL FLUX (FLUXT) = 0.20819284E-05 TIME(40) = 0.53923169E-02

FORCE (I, J)	BI (I, J)	HI (I, J)	TIME (80)	(I, J)
0.23130637E-01	0.96443628E 00	0.37847550E 02	0.10784632E-01	2
0.40516707E-01	0.90257192E 00	0.28244255E 02	0.10784632E-01	2
0.00204983E-01	0.63445693E 00	-0.86852137E 01	0.10784632E-01	2
0.91549834E-02	-0.42903579E-01	-0.80315129E 03	0.10784632E-01	2
0.41465975E-01	-0.12912922E 01	-0.25000000E 02	0.10784632E-01	2
0.45986420E-01	0.96156699E 00	0.39274509E 02	0.10784632E-01	3
0.80525393E-01	0.89973824E 00	0.27744881E 02	0.10784632E-01	3
0.39692173E-01	0.63168737E 00	-0.87835838E 01	0.10784632E-01	3
0.18387320E-01	-0.42996110E-00	-0.80387130E 02	0.10784632E-01	3
0.82931516E-01	-0.12912922E 01	-0.25000000E 03	0.10784632E-01	3
0.45028677E-01	0.95150121E 00	0.37292323E 02	0.10784632E-01	4
0.78752437E-01	0.89977818E 00	0.26022827E 02	0.10784632E-01	4
0.38469558E-01	0.62188252E 00	-0.10029289E 02	0.10784632E-01	4
0.18668772E-01	-0.43321913E-00	-0.80648596E 02	0.10784632E-01	4
0.82931516E-01	-0.12912922E 01	-0.25000000E 03	0.10784632E-01	4
0.42922333E-01	0.92898010E 00	0.33010077E 02	0.10784632E-01	5
0.74945462E-01	0.86742610E 00	0.22282109E 02	0.10784632E-01	5
0.35767052E-01	0.59964100E 00	-0.12341783E 02	0.10784632E-01	5
0.19329875E-01	-0.44082303E-00	-0.81253704E 02	0.10784632E-01	5
0.82931516E-01	-0.12912922E 01	-0.25000000E 03	0.10784632E-01	5
0.38615819E-01	0.88114499E 00	0.24558090E 02	0.10784632E-01	6
0.66850295E-01	0.81978773E 00	0.14833172E 02	0.10784632E-01	6
0.30278623E-01	0.55171884E 00	-0.17060275E 02	0.10784632E-01	6
0.20778389E-01	-0.45704155E-00	-0.82585043E 02	0.10784632E-01	6
0.82931516E-01	-0.12912922E 01	-0.25000000E 03	0.10784632E-01	6
0.29810452E-01	0.77419266E 00	0.25000000E 02	0.10784632E-01	7
0.50548927E-01	0.71286242E 00	0.83166630E 01	0.10784632E-01	7
0.19521182E-01	0.44299906E-00	-0.35128806E-00	0.10784632E-01	7
0.23971021E-01	-0.49089999E-00	-0.26594459E 02	0.10784632E-01	7
0.82931516E-01	-0.12912922E 01	-0.25000000E 03	0.10784632E-01	7
0.1871676E-01	0.48856322E-00	0.25000000E 02	0.10784632E-01	8
0.18036178E-01	0.442581602E-00	-0.27970458E 02	0.10784632E-01	8
0.21488351E-02	0.14697765E-00	-0.45475545E 02	0.10784632E-01	8
0.31623027E-01	-0.55383428E 00	-0.92130306E 02	0.10784632E-01	8
0.82931516E-01	-0.12912922E 01	-0.25000000E 03	0.10784632E-01	8
0.1332951E-01	-0.47734925E-00	-0.84287893E 02	0.10784632E-01	9
0.24492793E-01	-0.47621388E-00	-0.85914662E 02	0.10784632E-01	9
0.31690815E-01	-0.56443828E 00	-0.92138714E 02	0.10784632E-01	9
0.54914397E-01	-0.74300692E 00	-0.11218879E 03	0.10784632E-01	9
0.82931516E-01	-0.12912922E 01	-0.25000000E 03	0.10784632E-01	10
0.41465758E-01	-0.12912922E 01	-0.25000000E 03	0.10784632E-01	10
0.82931516E-01	-0.12912922E 01	-0.25000000E 03	0.10784632E-01	10
0.82931516E-01	-0.12912922E 01	-0.25000000E 03	0.10784632E-01	10
0.41465758E-01	-0.12912922E 01	-0.25000000E 03	0.10784632E-01	10

TOTAL FORCE (FORGT) = 0.20513958E 01 TIME(80) = 0.10784632E-01
TOTAL FLUX (FLUXT) = -0.34615710E-06 TIME(80) = 0.10784632E-01

FORCE(I,J)	BI(I,J)	HI(I,J)	TIME(100)	(I,J)
0.18559017E-01	0.86388759E-00	0.21705081E-02	0.13480790E-01	2
0.131346868E-01	0.793389273E-00	0.110669226E-02	0.13480790E-01	3
0.11854441E-01	0.48820844E-00	-0.128113222E-02	0.13480790E-01	4
0.138633061E-01	0.52795187E-00	-0.887543225E-02	0.13480790E-01	5
0.414655758E-01	0.12912922E-01	0.25000000E-02	0.13480790E-01	6
0.36834265E-01	0.86057902E-00	0.210602032E-02	0.13480790E-01	23
0.62176703E-01	0.790612335E-00	0.230922266E-02	0.13480790E-01	3
0.273393401E-01	0.48494940E-00	-0.230922266E-02	0.13480790E-01	4
0.27804430E-01	0.52869688E-00	-0.888222591E-02	0.13480790E-01	5
0.82931516E-01	0.12912922E-01	-0.25000000E-02	0.13480790E-01	6
0.35851645E-01	0.84902268E-00	0.193224502E-02	0.13480790E-01	23
0.60385093E-01	0.77913842E-00	0.89978344E-02	0.13480790E-01	4
0.22301017E-01	0.473349139E-00	-0.24069635E-02	0.13480790E-01	4
0.28084927E-01	0.53135699E-00	-0.89066953E-02	0.13480790E-01	4
0.82931516E-01	0.12912922E-01	-0.25000000E-02	0.13480790E-01	5
0.33713473E-01	0.82331590E-00	0.15366369E-02	0.13480790E-01	5
0.56485146E-01	0.75355834E-00	0.55395380E-02	0.13480790E-01	5
0.19942318E-01	0.44775204E-00	-0.26207998E-02	0.13480790E-01	5
0.28727540E-01	0.53740162E-00	-0.89625838E-02	0.13480790E-01	5
0.82931516E-01	0.12912922E-01	-0.25000000E-02	0.13480790E-01	6
0.29406694E-01	0.76893184E-00	0.75987871E-02	0.13480790E-01	6
0.48643907E-01	0.69930118E-00	-0.12978579E-02	0.13480790E-01	6
0.15343007E-01	0.39273987E-00	-0.30529381E-02	0.13480790E-01	6
0.30095617E-01	0.55004900E-00	-0.90811701E-02	0.13480790E-01	6
0.82931516E-01	0.12912922E-01	-0.25000000E-02	0.13480790E-01	7
0.20840845E-01	0.64775837E-00	0.14556853E-02	0.13480790E-01	7
0.33184131E-01	0.57758346E-00	-0.39210112E-02	0.13480790E-01	7
0.71629496E-02	0.26834635E-00	0.3332411E-02	0.13480790E-01	7
0.33020057E-01	0.57615411E-00	-0.93332411E-02	0.13480790E-01	7
0.82931516E-01	0.12912922E-01	-0.25000000E-02	0.13480790E-01	8
0.50948890E-02	0.32006052E-00	0.35769892E-02	0.13480790E-01	8
0.61390605E-02	0.24842809E-00	-0.40476813E-02	0.13480790E-01	8
0.44717512E-03	0.67048423E-01	-0.57302518E-02	0.13480790E-01	8
0.39822282E-01	0.63272183E-00	-0.99157947E-02	0.13480790E-01	8
0.82931516E-01	0.12912922E-01	-0.25000000E-02	0.13480790E-01	9
0.15932756E-01	0.56599189E-00	0.92339200E-02	0.13480790E-01	9
0.33509656E-01	0.58040980E-00	-0.93752974E-02	0.13480790E-01	9
0.39885331E-01	0.63322252E-00	-0.99211878E-02	0.13480790E-01	9
0.60279059E-01	0.77845405E-00	-0.11693345E-02	0.13480790E-01	9
0.82931516E-01	0.12912922E-01	-0.25000000E-02	0.13480790E-01	10
0.41465758E-01	0.12912922E-01	-0.25000000E-02	0.13480790E-01	10
0.82931516E-01	0.12912922E-01	-0.25000000E-02	0.13480790E-01	10
0.82931516E-01	0.12912922E-01	-0.25000000E-02	0.13480790E-01	10
0.41465758E-01	0.12912922E-01	-0.25000000E-02	0.13480790E-01	10

TOTAL FORCE (FORCI) = 0.18838430E 01 TIME(100) = 0.13480790E-01
TOTAL FLUX (FLUXT) = -0.11199690E-05 TIME(100) = 0.13480790E-01

MAGNETIC FIELD INTENSITY BI(I, J) AT TIME (5) = 0.67403973E-03

2	0.242900222E 03	0.23367802E 03	0.19316754E 03	0.778000727E 02	-0.250000000E 03
3	0.242891179E 03	0.23366942E 03	0.19316351E 03	0.778000281E 02	-0.250000000E 03
4	0.24281161E 03	0.23360194E 03	0.19312953E 03	0.77795920E 02	-0.250000000E 03
5	0.24232931E 03	0.23318644E 03	0.19290102E 03	0.77759462E 02	-0.250000000E 03
6	0.23980681E 03	0.23095629E 03	0.19154983E 03	0.77476157E 02	-0.250000000E 03
7	0.22880605E 03	0.22095668E 03	0.18481231E 03	0.75534222E 02	-0.250000000E 03
8	0.19001767E 03	0.18461831E 03	0.15738061E 03	0.64510448E 02	-0.250000000E 03
9	0.77157515E 02	0.75497995E 02	0.64508279E 02	0.12613811E 02	-0.250000000E 03
10	-0.25000000E 03	-0.25000000E 03	-0.25000000E 03	-0.250000000E 02	-0.250000000E 03

TOTAL FORCE (FORCT) = 0.52970427E 01 TIME(5) = 0.67403973E-03
 TOTAL FLUX (FLUXT) = 0.65340440E-05 TIME(5) = 0.67403973E-03

FLUX DENSITY MATRIX BI(I, J) AT TIME(5) = 0.67403973E-03

2	0.14442192E 01	0.14336488E 01	0.13803514E 01	0.11212639E 01	-0.12912922E 01
3	0.14442076E 01	0.14336387E 01	0.13803454E 01	0.11212623E 01	-0.12912922E 01
4	0.14441179E 01	0.14335594E 01	0.13802953E 01	0.11212474E 01	-0.12912922E 01
5	0.14435775E 01	0.14330709E 01	0.13779958E 01	0.11211227E 01	-0.12912922E 01
6	0.14407289E 01	0.14304301E 01	0.13779548E 01	0.11201519E 01	-0.12912922E 01
7	0.14278539E 01	0.14181910E 01	0.13677137E 01	0.11134300E 01	-0.12912922E 01
8	0.13756634E 01	0.13674177E 01	0.13212968E 01	0.10728835E 01	-0.12912922E 01
9	0.1190571E 01	0.11133033E 01	0.10728758E 01	0.80466644E 01	-0.12912922E 01
10	-0.12912922E 01	-0.12912922E 01	-0.12912922E 01	-0.12912922E 01	-0.12912922E 01

TOTAL FORCE (FORCT) = 0.52970427E 01 TIME(5) = 0.67403973E-03
 TOTAL FLUX (FLUXT) = 0.65340440E-05 TIME(5) = 0.67403973E-03

THEORETICAL FORCE(I, J) AT TIME(5) = 0.67403973E-03

2	0.51868869E-01	0.10222475E-00	0.94765411E-01	0.62529676E-01	0.41465758E-01
3	0.10373606E-00	0.20444663E-00	0.18952918E-00	0.12505901E-00	0.82931516E-01
4	0.10372317E-00	0.204424C3E-00	0.18951543E-00	0.12505569E-00	0.82931516E-01
5	0.10364556E-00	0.20428472E-00	0.18942284E-00	0.12502786E-00	0.82931516E-01
6	0.10323692E-00	0.20353252E-00	0.18887325E-00	0.12481143E-00	0.82931516E-01
7	0.10140003E-00	0.20006448E-00	0.18667759E-00	0.12331796E-00	0.82931516E-01
8	0.94122820E-01	0.18599571E-00	0.17366058E-00	0.11450005E-00	0.82931516E-01
9	0.62283787E-01	0.12328993E-00	0.11449826E-00	0.64406883E-01	0.82931516E-01
10	0.41465758E-01	0.82931516E-01	0.82931516E-01	0.82931516E-01	0.41465758E-01

TOTAL FORCE (FORCT) = 0.52970427E 01 TIME(5) = 0.67403973E-03
 TOTAL FLUX (FLUXT) = 0.65340440E-05 TIME(5) = 0.67403973E-03

MAGNETIC FIELD INTENSITY HI(I, J) AT TIME (I0) = 0.13480794E-02

2	0.21910437E 03	0.20339263E 03	0.14936348E 03	0.31698651E 02	-0.25000000E 03
3	0.21898415E 03	0.20329663E 03	0.14932033E 03	0.31693427E 02	-0.25000000E 03
4	0.21836215E 03	0.20279021E 03	0.14907693E 03	0.31659298E 02	-0.25000000E 03
5	0.21609063E 03	0.20089148E 03	0.14807831E 03	0.31486394E 02	-0.25000000E 03
6	0.20889440E 03	0.19470092E 03	0.14448617E 03	0.30688532E 02	-0.25000000E 03
7	0.18899771E 03	0.17705328E 03	0.13312110E 03	0.27370962E 02	-0.25000000E 03
8	0.14028606E 03	0.13239589E 03	0.10096276E 03	0.14972406E 02	-0.25000000E 03
9	0.29744595E 02	0.27213221E 02	0.14911997E 02	-0.30434300E 02	-0.25000000E 03
10	-0.25000000E 03	-0.25000000E 03	-0.25000000E 03	-0.25000000E 03	-0.25000000E 03

TOTAL FORCE (FORCT) = 0.45787790E 01 TIME(I0) = 0.13480794E-02
 TOTAL FLUX (FLUXT) = 0.56813462E-05 TIME(I0) = 0.13480794E-02

FLUX DENSITY MATRIX BI(I, J) AT TIME(I0) = 0.13480794E-02

2	0.14158496E 01	0.13949812E 01	0.13060828E 01	0.92186183E 00	-0.12912922E 01
3	0.14155963E 01	0.13948479E 01	0.13059986E 01	0.92183325E 00	-0.12912922E 01
4	0.14147047E 01	0.13941433E 01	0.13055235E 01	0.92164657E 00	-0.12912922E 01
5	0.14119884E 01	0.13914831E 01	0.13035659E 01	0.92069965E 00	-0.12912922E 01
6	0.14025009E 01	0.13826019E 01	0.12966411E 01	0.91630560E 00	-0.12912922E 01
7	0.13741264E 01	0.13553961E 01	0.12725488E 01	0.89759270E 00	-0.12912922E 01
8	0.12878161E 01	0.12709592E 01	0.11929654E 01	0.82036305E 00	-0.12912922E 01
9	0.91105458E 00	0.89668476E 00	0.82028227E 00	0.39399585E-00	-0.12912922E 01
10	-0.12912922E 01	-0.12912922E 01	-0.12912922E 01	-0.12912922E 01	-0.12912922E 01

TOTAL FORCE (FORCT) = 0.45787790E 01 TIME(I0) = 0.13480794E-02
 TOTAL FLUX (FLUXT) = 0.56813462E-05 TIME(I0) = 0.13480794E-02

THEORETICAL FORCE(I, J) AT TIME(I0) = 0.13480794E-02

2	0.49851105E-01	0.96784820E-01	0.84842205E-01	0.42267073E-01	0.41465758E-01
3	0.99680693E-01	0.19353264E-00	0.16966254E-00	0.84528904E-01	0.82931516E-01
4	0.99569172E-01	0.19333717E-00	0.16953912E-00	0.84494670E-01	0.82931516E-01
5	0.99159152E-01	0.19260004E-00	0.16903107E-00	0.84321138E-01	0.82931516E-01
6	0.97831076E-01	0.19014932E-00	0.16718078E-00	0.83518212E-01	0.82931516E-01
7	0.93912602E-01	0.18273975E-00	0.16103289E-00	0.80141807E-01	0.82931516E-01
8	0.82685616E-01	0.16068070E-00	0.14156511E-00	0.66944157E-01	0.82931516E-01
9	0.41231864E-01	0.79979756E-01	0.66930976E-01	0.15441298E-01	0.82931516E-01
10	0.41465759E-01	0.82931516E-01	0.82931516E-01	0.82931516E-01	0.41465758E-01

TOTAL FOPCE (FORCT) = 0.45787790E 01 TIME(I0) = 0.13480794E-02
 TOTAL FLUX (FLUXT) = 0.56813462E-05 TIME(I0) = 0.13480794E-02

MAGNETIC FIELD INTENSITY H(I,J) AT TIME (20) = 0.26961587E-02

2	0.16882860E 03	0.15174057E 03	0.96829750E 02	-0.11235042E 02	-0.25000000E 03
3	0.16840447E 03	0.15138979E 03	0.96650832E 02	-0.11263760E 02	-0.25000000E 03
4	0.16671370E 03	0.14997443E 03	0.95897850E 02	-0.11398818E 02	-0.25000000E 03
5	0.16227492E 03	0.14619173E 03	0.93757255E 02	-0.11851547E 02	-0.25000000E 03
6	0.15183441E 03	0.13711461E 03	0.88257473E 02	-0.13251153E 02	-0.25000000E 03
7	0.12901696E 03	0.11683237E 03	0.75020336E 02	-0.17338986E 02	-0.25000000E 03
8	0.81710129E 02	0.73649745E 02	0.442338906E 02	-0.29050958E 02	-0.25000000E 03
9	-0.14974566E 02	-0.17712619E 02	-0.29087009E 02	-0.65901104E 02	-0.25000000E 03
10	-0.25000000E 03	-0.25000000E 03	-0.25000000E 03	-0.25000000E 03	-0.25000000E 03

TOTAL FORCE (FORCI) = 0.36379664E 01 TIME(20) = 0.26961587E-02
 TOTAL FLUX (FLUXT) = 0.42601639E-05 TIME(20) = 0.26961587E-02

FLUX DENSITY MATRIX B(I,J) AT TIME(20) = 0.26961587E-02

2	0.13416639E 01	0.13106801E 01	0.11811876E 01	0.61038322E 00	-0.12912922E 01
3	0.13409360E 01	0.13100064E 01	0.11806685E 01	0.61010677E 00	-0.12912922E 01
4	0.13380148E 01	0.13072715E 01	0.11784755E 01	0.60880501E 00	-0.12912922E 01
5	0.13301913E 01	0.12998307E 01	0.11721651E 01	0.60442174E 00	-0.12912922E 01
6	0.13108601E 01	0.12811548E 01	0.11554152E 01	0.59067702E 00	-0.12912922E 01
7	0.12634408E 01	0.12347035E 01	0.1116311E 01	0.5487533E 00	-0.12912922E 01
8	0.11344043E 01	0.11067915E 01	0.985833994E 00	0.41203140E-00	-0.12912922E 01
9	0.57333851E 00	0.54480874E 00	0.41156694E-00	-0.22169189E-00	-0.12912922E 01
10	-0.12912922E 01	-0.12912922E 01	-0.12912922E 01	-0.12912922E 01	-0.12912922E 01

TOTAL FOPCE (FORCI) = 0.36379664E 01 TIME(20) = 0.26961587E-02
 TOTAL FLUX (FLUXT) = 0.42601639E-05 TIME(20) = 0.26961587E-02

THEORETICAL FORCE(I,J) AT TIME(20) = 0.26961587E-02

2	0.44763909E-01	0.85440540E-01	0.69391819E-01	0.18530010E-01	0.41465758E-01
3	0.89430710E-01	0.17070544E-00	0.13866169E-00	0.37026459E-01	0.82931516E-01
4	0.89041477E-01	0.16999342E-00	0.13814707E-00	0.36868624E-01	0.82931516E-01
5	0.88003264E-01	0.16806377E-00	0.13667153E-00	0.36339643E-01	0.82931516E-01
6	0.85464005E-01	0.16326878E-00	0.13279344E-00	0.34705687E-01	0.82931516E-01
7	0.79392660E-01	0.15164423E-00	0.12291981E-00	0.29956402E-01	0.82931516E-01
8	0.64003878E-01	0.12185185E-00	0.96674807E-01	0.16887335E-01	0.82931516E-01
9	0.16349057E-01	0.29524913E-01	0.16849285E-01	0.48887759E-02	0.82931516E-01
10	0.41465758E-01	0.82931516E-01	0.82931516E-01	0.82931516E-01	0.41465758E-01

TOTAL FOPCE (FORCI) = 0.36379664E 01 TIME(20) = 0.26961587E-02
 TOTAL FLUX (FLUXT) = 0.42601639E-05 TIME(20) = 0.26961587E-02

MAGNETIC FIELD INTENSITY H(I,J) AT TIME (40) = 0.53923169E-02

2	0.10322275E 03	0.88452752E 02	0.41505075E 02	-0.49088437E 02	0.25000000E 03
3	0.10261101E 03	0.879300042E 02	0.41206053E 02	-0.49151751E 02	0.25000000E 03
4	0.10041833E 03	0.86045487E 02	0.40105373E 02	-0.49400901E 02	0.25000000E 03
5	0.95435143E 02	0.81720982E 02	0.37493915E 02	-0.50055926E 02	0.25000000E 03
6	0.85138415E 02	0.72675309E 02	0.31805028E 02	-0.51653909E 02	0.25000000E 03
7	0.64710458E 02	0.54445374E 02	0.19750106E 02	-0.55466334E 02	0.25000000E 03
8	0.24948010E 02	0.18165620E 02	-0.60091966E 01	-0.64908683E 02	0.25000000E 03
9	-0.53675393E 02	0.55975630E 02	-0.64967045E 02	-0.92946576E 02	0.25000000E 03
10	-0.25000000E 03	-0.25000000E 03	-0.25000000E 03	-0.25000000E 03	0.25000000E 03

TOTAL FORCE (FORCT) = 0.26857109E 01 TIME(40) = 0.53923169E-02
 TOTAL FLUX (FLUXT) = 0.20819284E-05 TIME(40) = 0.53923169E-02

FLUX DENSITY MATRIX B(I,J) AT TIME(40) = 0.53923169E-02

2	0.11992383E 01	0.11560235E 01	0.97263501E 00	0.98722704E-01	0.12912922E 01
3	0.11975518E 01	0.11543927E 01	0.97116685E 00	0.97519964E-01	0.12912922E 01
4	0.11914369E 01	0.11484513E 01	0.96572138E 00	0.92771135E-01	0.12912922E 01
5	0.11771211E 01	0.11344402E 01	0.95253503E 00	0.80164447E-01	0.12912922E 01
6	0.11455565E 01	0.11033133E 01	0.92244322E 00	0.48648283E-01	0.12912922E 01
7	0.10736576E 01	0.10319030E 01	0.85171223E 00	0.30201413E-01	0.12912922E 01
8	0.98345791E 00	0.84162710E 00	0.65875600E 00	-0.20520698E-00	0.12912922E 01
9	0.71584135E-02	0.40571190E-01	-0.20618552E-00	-0.57222515E 00	0.12912922E 01
10	-0.12912922E 01	-0.12912922E 01	-0.12912922E 01	-0.12912922E 01	0.12912922E 01

TOTAL FORCE (FORCT) = 0.26857109E 01 TIME(40) = 0.53923169E-02
 TOTAL FLUX (FLUXT) = 0.20819284E-05 TIME(40) = 0.53923169E-02

THEORETICAL FORCE(I,J) AT TIME(40) = 0.53923169E-02

2	0.35764449E-01	0.66466666E-01	0.47051156E-01	0.48473524E-03	0.41465758E-01
3	0.71327352E-01	0.13255854E-00	0.93818438E-01	0.94599222E-03	0.82931516E-01
4	0.70601285E-01	0.13119754E-00	0.92769232E-01	0.85610344E-03	0.82931516E-01
5	0.68914851E-01	0.12201585E-00	0.90253160E-01	0.63924025E-03	0.82931516E-01
6	0.65268493E-01	0.12108719E-00	0.84640805E-01	0.23541576E-03	0.82931516E-01
7	0.57332665E-01	0.10592006E-00	0.72158296E-01	0.90730857E-04	0.82931516E-01
8	0.38819811E-01	0.70459558E-01	0.43156781E-01	0.41887528E-02	0.82931516E-01
9	0.25486140E-05	0.16373292E-03	0.42287958E-02	0.32571247E-01	0.82931516E-01
10	0.41465758E-01	0.82931516E-01	0.82931516E-01	0.82931516E-01	0.41465758E-01

TOTAL FORCE (FORCT) = 0.26857109E 01 TIME(40) = 0.53923169E-02
 TOTAL FLUX (FLUXT) = 0.20819284E-05 TIME(40) = 0.53923169E-02

MAGNETIC FIELD INTENSITY H(I,J) AT TIME (60) = 0.80884747E-02

2	0.65058924E 02	0.52145334E 02	0.11096895E 02	-0.68250201E 02	-0.25000000E 03
3	0.644451001E 02	0.51620746E 02	0.10785875E 02	-0.683222758E 02	-0.25000000E 03
4	0.62327000E 02	0.49780778E 02	0.96800489E 01	-0.69593042E 02	-0.25000000E 03
5	0.57672172E 02	0.45720726E 02	0.71830125E 01	-0.69249194E 02	-0.25000000E 03
6	0.48366828E 02	0.37526801E 02	0.19886278E 01	-0.70727625E 02	-0.25000000E 03
7	0.30335285E 02	0.21432260E 02	-0.86481190E 01	-0.74027571E 02	-0.25000000E 03
8	-0.422992123E 01	-0.10143340E 02	-0.30944101E 02	-0.81852385E 02	-0.25000000E 03
9	-0.72617624E 02	-0.74528828E 02	-0.81913723E 02	-0.10507315E 03	-0.25000000E 03
10	-0.25000000E 03	-0.25000000E 03	-0.25000000E 03	-0.25000000E 03	-0.25000000E 03

TOTAL FORCE (FORCT) = 0.22928122E 01 TIME(60) = 0.80884747E-02
 TOTAL FLUX (FLUXT) = 0.62974814E-06 TIME(60) = 0.80884747E-02

FLUX DENSITY MATRIX B(I,J) AT TIME(60) = 0.80884747E-02

2	0.10750026E 01	0.10219438E 01	0.79410351E 00	-0.25944519E-00	-0.12912922E 01
3	0.10726532E 01	0.10196393E 01	0.79191195E 00	-0.26058387E-00	-0.12912922E 01
4	0.106433353E 01	0.10114572E 01	0.78404745E 00	-0.26481148E-00	-0.12912922E 01
5	0.10454895E 01	0.99283726E 00	0.76586188E 00	-0.27498307E-00	-0.12912922E 01
6	0.10050625E 01	0.95270347E 00	0.72602072E 00	-0.29743733E-00	-0.12912922E 01
7	0.91435279E 00	0.86220591E 00	0.63480040E 00	-0.34535298E-00	-0.12912922E 01
8	0.67379107E 00	0.62080295E 00	0.38723662E-00	-0.44812144E-00	-0.12912922E 01
9	-0.32524107E-00	-0.35237791E-00	-0.44887252E-00	-0.68533394E 00	-0.12912922E 01
10	-0.12912922E 01	-0.12912922E 01	-0.12912922E 01	-0.12912922E 01	-0.12912922E 01

TOTAL FOPCE (FORCT) = 0.22928122E 01 TIME(60) = 0.80884747E-02
 TOTAL FLUX (FLUXT) = 0.62974814E-06 TIME(60) = 0.80884747E-02

THEORETICAL FORCE(I,J) AT TIME(60) = 0.80884747E-02

2	0.28738201E-01	0.51942707E-01	0.31363516E-01	0.33478174E-02	0.41465758E-01
3	0.57225447E-01	0.10341742E-00	0.62381281E-01	0.67545367E-02	0.82931516E-01
4	0.56341376E-01	0.10176432E-00	0.61148413E-01	0.69754802E-02	0.82931516E-01
5	0.54363807E-01	0.98052040E-01	0.58344700E-01	0.75216380E-02	0.82931516E-01
6	0.50240808E-01	0.90285081E-01	0.52432254E-01	0.88001782E-02	0.82931516E-01
7	0.41581303E-01	0.73947329E-01	0.40084353E-01	0.11863885E-01	0.82931516E-01
8	0.22579847E-01	0.38336109E-01	0.14916033E-01	0.19975237E-01	0.82931516E-01
9	0.52611569E-02	0.12351447E-01	0.20042253E-01	0.46720730E-01	0.82931516E-01
10	0.41465758E-01	0.82931516E-01	0.82931516E-01	0.82931516E-01	0.41465758E-01

TOTAL FOPCE (FORCT) = 0.22928122E 01 TIME(60) = 0.80884747E-02
 TOTAL FLUX (FLUXT) = 0.62974814E-06 TIME(60) = 0.80884747E-02

THEORETICAL FORCF(I,J) AT TIME(80) = 0.10784632E-01

2	0.23130637E-01	0.40516707E-01	0.20020494E-01	0.91549833E-02	0.41465758E-01
3	0.459866420E-01	0.80525393E-01	0.39692173E-01	0.18387320E-01	0.82931516E-01
4	0.45028677E-01	0.78752437E-01	0.38469558E-01	0.18667772E-01	0.82931516E-01
5	0.42922333E-01	0.74865462E-01	0.35767052E-01	0.19329875E-01	0.82931516E-01
6	0.38615819E-01	0.66850295E-01	0.302778623E-01	0.20778389E-01	0.82931516E-01
7	0.27810452E-01	0.50548927E-01	0.19521132E-01	0.23971021E-01	0.82931516E-01
8	0.11871676E-01	0.18036178E-01	0.21488351E-02	0.31623027E-01	0.82931516E-01
9	0.11332951E-01	0.24492793E-01	0.31690815E-01	0.54914397E-01	0.82931516E-01
10	0.41465758E-01	0.82931516E-01	0.82931516E-01	0.82931516E-01	0.41465758E-01

TOTAL FORCE (FORCT) = 0.20513958E 01 TIME(80) = 0.10784632E-01
 TOTAL FLUX (FLUXT) = -0.34615710E-06 TIME(80) = 0.10784632E-01

FLUX DENSITY MATRIX BI(I,J) AT TIME(80) = 0.10784632E-01

2	0.96443628E 00	0.90257192E 00	0.63445693E 00	-0.42903579E-00	-0.12912922E 01
3	0.96156699E 00	0.89973824E 00	0.63168737E 00	-0.42994110E-00	-0.12912922E 01
4	0.95150121E 00	0.88977818E 00	0.62188252E 00	-0.43321913E-00	-0.12912922E 01
5	0.92898010E 00	0.86742610E 00	0.599644100E 00	-0.44082303E-00	-0.12912922E 01
6	0.88114499E 00	0.81978773E 00	0.55171884E 00	-0.45704155E-00	-0.12912922E 01
7	0.77419264E 00	0.71286242E 00	0.44299906E-00	-0.49089999E-00	-0.12912922E 01
8	0.48856322E-00	0.42581602E-00	0.14697765E-00	-0.563836428E 00	-0.12912922E 01
9	-0.47734925E-00	-0.49621388E-00	-0.56443828E 00	-0.743000692E 00	-0.12912922E 01
10	-0.12912922E 01	-0.12912922E 01	-0.12912922E 01	-0.12912922E 01	-0.12912922E 01

TOTAL FORCE (FORCT) = 0.20513958E 01 TIME(80) = 0.10784632E-01
 TOTAL FLUX (FLUXT) = -0.34615710E-06 TIME(80) = 0.10784632E-01

MAGNETIC FIELD INTENSITY HI(I,J) AT TIME (80) = 0.10784632E-01

2	0.39847560E 02	0.28241255E 02	-0.86852137E 01	-0.80315129E 02	-0.25000000E 03
3	0.39274509E 02	0.27744881E 02	-0.89835838E 01	-0.80387130E 02	-0.25000000E 03
4	0.37292323E 02	0.26022827E 02	-0.10029282E 02	-0.80648596E 02	-0.25000000E 03
5	0.33010077E 02	0.22282109E 02	-0.12341713E 02	-0.81259704E 02	-0.25000000E 03
6	0.24558090E 02	0.14838172E 02	-0.17060275E 02	-0.82585043E 02	-0.25000000E 03
7	0.83166630E 01	0.35128806E-00	-0.26594459E 02	-0.85451940E 02	-0.25000000E 03
8	-0.22780658E 02	-0.27970458E 02	-0.46475545E 02	-0.92130306E 02	-0.25000000E 03
9	-0.84287893E 02	-0.85914662E 02	-0.92188714E 02	-0.11218879E 03	-0.25000000E 03
10	-0.25000000E 03	-0.25000000E 03	-0.25000000E 03	-0.25000000E 03	-0.25000000E 03

TOTAL FORCE (FORCT) = 0.20513958E 01 TIME(80) = 0.10784632E-01
 TOTAL FLUX (FLUXT) = -0.34615710E-06 TIME(80) = 0.10784632E-01

MAGNETIC FIELD INTENSITY HI(I,J) AT TIME (100) = 0.13480790E-01

2	0.217055081E 02	0.110669266E 02	0.228113222E 02	-0.887543255E 02	-0.250000000E 03
3	0.211691900E 02	0.106020332E 02	-0.230922266E 02	-0.88822591E 02	-0.250000000E 03
4	0.193245020E 02	0.899783344E 01	-0.240697358E 02	-0.890669538E 02	-0.250000000E 03
5	0.15366369E 02	0.55395380E 01	-0.26207998E 02	-0.89625838E 02	-0.250000000E 03
6	0.75987871E 01	-0.12979579E 01	-0.30529381E 02	-0.90811701E 02	-0.250000000E 03
7	-0.72812592E 01	-0.14556853E 02	-0.39210112E 02	-0.93332241E 02	-0.250000000E 03
8	-0.35769892E 02	-0.40476813E 02	-0.57302518E 02	-0.99157947E 02	-0.250000000E 03
9	-0.32339200E 02	-0.93752974E 02	-0.99211878E 02	-0.116933345E 03	-0.250000000E 03
10	-0.25000000E 03	-0.25000000E 03	-0.25000000E 03	-0.25000000E 03	-0.25000000E 03

TOTAL FORCE (FORCT) = 0.18838430E 01 TIME(100) = 0.13480790E-01
 TOTAL FLUX (FLUXT) = -0.11199690E-05 TIME(100) = 0.13480790E-01

FLUX DENSITY MATRIX BI(I,J) AT TIME(100) = 0.13480790E-01

2	0.86388759E 00	0.79389273E 00	0.48820844E-00	-0.52795187E 00	-0.12912922E 01
3	0.86057902E 00	0.79061235E 00	0.48494940E-00	-0.52869688E 00	-0.12912922E 01
4	0.84902268E 00	0.77913842E 00	0.47349139E-00	-0.53135699E 00	-0.12912922E 01
5	0.82331590E 00	0.75355834E 00	0.44775204E-00	-0.53740162E 00	-0.12912922E 01
6	0.76893184E 00	0.69930118E 00	0.39273937E-00	-0.55004900E 00	-0.12912922E 01
7	0.64732505E-00	0.57758377E 00	0.26834635E-00	-0.57615411E 00	-0.12912922E 01
8	0.32004052E-00	0.24842809E-00	0.67048422E-01	-0.63272183E 00	-0.12912922E 01
9	-0.56599189E 00	-0.58040980E 00	-0.67332225E 00	-0.77845405E 00	-0.12912922E 01
10	-0.12912922E 01	-0.12912922E 01	-0.12912922E 01	-0.12912922E 01	-0.12912922E 01

TOTAL FORCE (FORCT) = 0.18838430E 01 TIME(100) = 0.13480790E-01
 TOTAL FLUX (FLUXT) = -0.11199690E-05 TIME(100) = 0.13480790E-01

THEORETICAL FORCE(I,J) AT TIME(100) = 0.13480790E-01

2	0.18559017E-01	0.31346868E-01	0.11854441E-01	0.13863061E-01	0.41465758E-01
3	0.36834265E-01	0.62176703E-01	0.23393401E-01	0.27804430E-01	0.82931516E-01
4	0.35851645E-01	0.60385093E-01	0.22301017E-01	0.28084927E-01	0.82931516E-01
5	0.33713473E-01	0.56485146E-01	0.19942318E-01	0.28727540E-01	0.82931516E-01
6	0.29406694E-01	0.48643972E-01	0.15343007E-01	0.30095617E-01	0.82931516E-01
7	0.20840845E-01	0.3184131E-01	0.71629496E-02	0.33020057E-01	0.82931516E-01
8	0.50948890E-02	0.61390605E-02	0.44717512E-03	0.39822282E-01	0.82931516E-01
9	0.15932756E-01	0.33509656E-01	0.39885331E-01	0.60279059E-01	0.82931516E-01
10	0.41465758E-01	0.82931516E-01	0.82931516E-01	0.82931516E-01	0.41465758E-01

TOTAL FORCE (FORCT) = 0.18838430E 01 TIME(100) = 0.13480790E-01
 TOTAL FLUX (FLUXT) = -0.11199690E-05 TIME(100) = 0.13480790E-01

(KZ = 3 PREVIOUS DATA FOR DATA SET NO. 3)

HS	=	0.11573666E 06	HSO	=	-0.09999999E 08
T	=	0.37366199E-05	TKM	=	0.99999999E-02
K	=	8	KO	=	8

STOP BECAUSE K IS GREATER THAN OR = TO KO

THIS IS A DECAY RUN

5.3 Discussion of Output Data and Results

For the sample problem considered, a magnetic field intensity of 250 amp-t/m was sufficient to saturate the core. Assuming uniform flux density, initial values of total force and flux can be calculated using equations (2-2) and (2-4); hence

$$(2-2) \quad \phi = BA = (1.452)(0.004)(0.008) \\ = 46.464 \times 10^{-6} \text{ webers}$$

$$(2-4) \quad F = \frac{B^2A}{2\mu_0} = \frac{(1.452)^2(0.004)(0.008)}{(2)(4\pi \times 10^{-7})} = 26.843 \text{ newtons}$$

Total values computed by summing $B(I,J)h^2$ and $FORCE(I,J)$ compare closely with those given above; thus one may conclude computations of total force and flux specified by subroutines FORCX and FLUX are valid approximations.

Distribution patterns are indicated by the print-out of Tapes 9 and 15. Tape 15 Output Data indicates values of field intensity and flux density for each point in the lattice of Fig. 12 in a convenient format. (Note physical position of $H(I,J)$ in Fig. 12 and matrix output.) Examining these values as time increases, i.e., for $TIME(0)$, $TIME(5)$, $TIME(10)$, . . . , $TIME(100)$, we note that points near the boundary approach boundary conditions much more rapidly than those at the center of the core. Given sufficient time, all points will reach the same values as that on the outer boundary. Let us think of the field intensity at each grid point as a vector pointing in the z direction (see Fig. 7), whose magnitude represents the

magnitude of field intensity at that point. If a membrane was stretched over the tips of the vectors, the surface generated as time increases represents the distribution pattern. We may visualize this distribution pattern by citing the following example. Suppose a soap bubble has been partially formed by exerting pressure behind a membrane across a tube of rectangular cross section. If we release the pressure sustaining the bubble, the membrane will return to its initial state. The changing shape of the bubble during the period of time in which it is decaying to its initial state, is analogous to the shape of the surface representing the flux decay distribution patterns. We should also note that force depends only on the magnitude of flux density, thus force is always positive. Furthermore, it will return to a value corresponding to $H = -250 \text{ amp-t/m}$ when all transients have decayed.

It is hoped that the solution given for the sample problem was useful in illustrating how one might obtain a solution for other equations of this type. An explicit expression for H and B was not found; however, the thesis does present details of how the Modified Euler Method of numerical integration may be used to establish a numerical solution which approximates the actual solution.

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APPENDICES

APPENDIX A

Terminology and Formulae

I. Terminology

Most symbols are defined in the context of the thesis. Definition of variables in the program are listed in section 4.2B and in the "comments" section of the program listing. (Appendix C, Cards MG90000 - MG90106.) Those variables not defined or frequently used are given in this section for convenience. Overbarred symbols indicate vector quantities.

1. $\bar{i}, \bar{j}, \bar{k}$ - Unit vectors in cartesian co-ordinates
2. $\bar{\nabla} = \left(\frac{\partial}{\partial x}, \frac{\partial}{\partial y}, \frac{\partial}{\partial z} \right)$ - Del-operator¹
3. $\bar{E} = (E_x, E_y, E_z)$ - Electric field intensity
4. $\bar{B} = (B_x, B_y, B_z)$ - Flux density
5. $\bar{H} = (H_x, H_y, H_z)$ - Magnetic field intensity
6. $\bar{D} = (D_x, D_y, D_z)$ - Displacement current density
7. $\bar{i} = (i_x, i_y, i_z)$ - Conduction current density
8. $\bar{dS} = (dS_x, dS_y, dS_z)$ - Differential surface
9. $\bar{dl} = (dl_x, dl_y, dl_z)$ - Differential length
10. $\sigma =$ - Material conductivity
11. $\mu_0 =$ - Permeability of free space
12. t - Time in seconds

¹Partheses denote vector components, e.g., $\bar{H} = (H_x, H_y, H_z)$
 $= H_x \bar{i} + H_y \bar{j} + H_z \bar{k}$.

13. h - Grid size in sections 3.1 and 3.2
14. D_i - $(dH/dt)_i$
15. P_i - $(\partial H/\partial t)_i$
16. I - Specifies row of matrix shown by Fig. 12
17. J - Specifies column of matrix shown by Fig. 12.
18. $B(I,J)$ - Flux density at grid point I,J
19. $H(I,J)$ - Magnetic field intensity at grid point I,J
20. $P_i(I,J)$ - $\partial H/\partial t$ at grid point I,J
21. $FORCE(I,J)$ - Force at each grid point
22. $FLUXT$ - Total flux
23. $FORCT$ - Total force.

II. Formulae

1. Curl of \bar{H}

$$\bar{\nabla} \times \bar{H} = \begin{vmatrix} \bar{i} & \bar{j} & \bar{k} \\ \partial/\partial x & \partial/\partial y & \partial/\partial z \\ H_x & H_y & H_z \end{vmatrix}$$

2. Divergence of \bar{H}

$$\bar{\nabla} \cdot \bar{H} = \frac{\partial H_x}{\partial x} + \frac{\partial H_y}{\partial y} + \frac{\partial H_z}{\partial z}$$

3. Gradient of ϕ

$$\bar{\nabla}(\phi) = \frac{\partial \phi}{\partial x} \bar{i} + \frac{\partial \phi}{\partial y} \bar{j} + \frac{\partial \phi}{\partial z} \bar{k}$$

4. Other useful relations

$$a. \quad \bar{\nabla} \cdot \bar{\nabla} = \bar{\nabla}^2$$

$$b. \quad \bar{\nabla} \times (\bar{\nabla} \cdot \bar{H}) = 0$$

$$c. \quad \bar{\nabla} \cdot (\bar{\nabla} \times \bar{H}) = 0$$

$$d. \quad \bar{\nabla} \times (\bar{\nabla} \times \bar{H}) = (\bar{\nabla} \cdot \bar{H}) - \bar{\nabla}^2 \bar{H}$$

$$e. \quad \bar{A} \times (\bar{B} \times \bar{C}) - (\bar{A} \times \bar{B}) \times \bar{C} \\ = \bar{A}(\bar{B} \cdot \bar{C}) - (\bar{A} \cdot \bar{B})\bar{C}$$

APPENDIX B

Units and Conversion Factors

Several systems of units used interchangeably are the MKS, CGS, and Mixed English. Conversion factors are given below.

	<u>MKS</u>		<u>CGS</u>		<u>MIXED ENGLISH</u>
Length	meter	$\times 10^2$	= cm	$\times 0.3937$	= inch
Time	second	$\times 1.0$	= second	$\times 1.0$	= second
Force	newton	$\times 10^3$	= dyne	$\times 2248.0 \times 10^{-6}$	= pound
Voltage	volt	$\times 1.0$	= volt	$\times 1.0$	= volt
Current	ampere	$\times 1.0$	= ampere	$\times 1.0$	= ampere
Resistance	ohm	$\times 1.0$	= ohm	$\times 1.0$	= ohm
Capacitance	farad	$\times 1.0$	= farad	$\times 1.0$	= farad
Flux	weber	$\times 10^8$	= maxwell	$\times 1.0$	= maxwell
Flux density	weber/m ²	$\times 10^4$	= gauss	$\times 6.45$	= lines/in ²
MMF	amp-turn	$\times .4\pi$	= gilbert	$\times 1/.4\pi$	= amp-turn
Magnetic intensity	amp-t/m	$\times .004\pi$	= oersted	$\times 2.02$	= amp-t/in

B-H Conversion Factors

B (lines/cm ²)	Gauss	$\times 6.450$	= B (lines/in ²)	Maxwells/in ²
B (lines/in ²)		$\times 0.155$	= B (lines/cm ²)	Gauss
H (oersteds)		$\times 2.020$	= H (amp-t/in)	
H (amp-t/in)		$\times 0.495$	= H (oersteds)	

APPENDIX C

Program Listing with Data Sets

The program is written in FORTRAN IV language for use with the IBM 7090/7094 IBSYS operating system. Processing is to be with the IJOB FORTRAN IV compiler. All cards with a "\$" sign in column 1 are system control cards used for all programs to be executed and are not unique to this program. All cards with a "C" in columns are comment cards and are not processed.

The program consists of a main program named EMEX and six subprograms FLUX, FORCX, FLUX, PPLT, INSE, and LOGI. The subprograms PPLT, INSE, and LOGI are provided to obtain graphical results while FLUD, FORCX, and FLUX are subroutines called by EMEX to perform specific calculations. One should also note that organization of the program is sequenced by numbers in columns 76-80 of the listing. These also aided in locating various sections within the program in the discussion of section 4.0.

The data input is on tape unit 5 and data output is on tape units 6, 7, 9, 15, and 16. Definition of symbols used is given in Appendix A, in the "comments" section of the program listing and in section 4.2B. If a larger matrix than 35 x 35 is desired, additional storage is required and the dimension statement must be altered accordingly.


```

C C C C C C
15 = VALUE USED FOR SPECIFYING BUILDUP OR DECAY
    USE 0 OR 1 FOR BUILDUP
    USE 2 FOR DECAY
CX = DIVISION WIDTHS
C1 = CONSTANT IN MODIFIED FROELICH APPROXIMATION
C2 = CONSTANT IN MODIFIED FROELICH APPROXIMATION
DT = TIME INCREMENT IN SECONDS
DIMENSION HO(35,35),DHO(35,35),H1(35,35),DHI(35,35),
1B1(35,35),FORCE(35,35),B(201),H(201),FFLUX( 800),
2TIME( 800)
COMMON I,J,UO,DX,T2,FLUXT,FORCE,BR,NX,NY,I5,B1,C1,C2
REVIND 9
REVIND15
REVIND16
48 READ (5,1) EO,HSO,BR,HMAX,CUO
    TO,TKW,COND,AX,3Y
    1 FORMAT (5E16.8)
    2 READ (5,2)CDT,KZ,NI,NJ,K,KO,KI,KG
    213 READ (5,213) B2,H2,B3,H3
    214 FORMAT (4E16.8)
    278 READ (5,278) H10,DHX,I3,I4,I5
    279 FORMAT (2E16.8,3I5)
    C1=(B2-BR)*H3*H2/((B2-BR)*H3-((B3-BR)*H2))
    HC=- (C2*BR)/(C1+BR)
    H(1)=ABS(H10)
    DO 624 I=1,13
    621 IF (H(I)-HC) 621,622,622
    HD=ABS(H(I)-HC)
    HD2=HC+HD
    B(I)=- (C1*HD2/(C2+HD2))+BR)
    GO TO 623
    622 B(I)= (C1*H(I)/(C2+H(I)))+3R
    623 H(I+1)=H(I)-DHX
    624 CONTINUE
    HC=C2*BR/(C1+BR)
    I6=2*13
    DO 627 I=13,16
    625 IF (H(I)-HC) 625,626,626
    HD1=-H(I)
    B(I)=- (C1*HD1/(C2+HD1))+BR)
    GO TO 627
    626 HX=H(I)-2.0*HC
    B(I)=(C1*HX/(C2+HX))+BR
    627 H(I+1)=H(I)+DHX
    WRITE ( 6,215)
    WRITE ( 9,215)
    MG90094
    MG90096
    MG90098
    MG90100
    MG90102
    MG90104
    MG90106
    MG90108
    MG90110
    MG90112
    MG90114
    MG90116
    MG90118
    MG90120
    MG90122
    MG90124
    MG90126
    MG90128
    MG90130
    MG90132
    MG90134
    MG90136
    MG90138
    MG90140
    MG90142
    MG90144
    MG90146
    MG90148
    MG90150
    MG90152
    MG90154
    MG90156
    MG90158
    MG90160
    MG90162
    MG90164
    MG90166
    MG90168
    MG90170
    MG90172
    MG90174
    MG90176
    MG90178
    MG90180
    MG90182
    MG90184
    MG90186
    MG90188

```



```

215 WRITE (15,215)
    FORMAT (1H1,8X,2H8R,14X,2HB2,14X,2HH2,14X,2HB3,14X,2HH3)
WRITE (6,216) BR,B2,H2,B3,H3
WRITE (9,216) BR,B2,H2,B3,H3
WRITE (15,216) BR,B2,H2,B3,H3
216 FORMAT (5E16.8)
WRITE (6,280)
WRITE (9,280)
WRITE (15,280)
280 FORMAT (1H0,7X,3HH10,13X,3HDHX,9X,2HI3,3X,2HI4)
WRITE (6,281) H10,DX,13,14
WRITE (9,281) H10,DX,13,14
WRITE (15,281) H10,DX,13,14
281 FORMAT (1H,22E16.8,2I5)
WRITE (6,222) C1,C2
WRITE (9,222) C1,C2
WRITE (15,222) C1,C2
222 FORMAT (1H0,5HC1 = E16.8,6X,5HC2 = E16.8/)
WRITE (6,217)
WRITE (9,217)
WRITE (15,217)
217 FORMAT (1H0,8X,4HB(1),12X,4HH(1),11X,1HI/)
DO 219 I=1,16
WRITE (6,218) B(I),H(I),I
WRITE (9,218) B(I),H(I),I
WRITE (15,218) B(I),H(I),I
218 FORMAT (1H,2E16.8,3X,I5)
219 CONTINUE
CALL PPLT(H,B,I6)
WRITE (6,610)
610 FORMAT (1H0,40X,45HF LUX DENSITY VS MAGNETIC FIELD INTENSITY )
NN=1
I1=MJ+1
I2=NI+1
NY=NJ+2
NXY=NZ+2
READ (5,1) (HO(I,J),J=2,NX),I=2,NY)
LI4=1
PI=3.14159
UO=6.0*PI*CUO
X=COND*CI*C2
T=TO
ANI=NI
CX=AX/(2.0*ANI)
ANJ=BY/(2.0*CX)
IF (I5-1) 600,600,602
600 WRITE (6,601)
    WRITE (9,601)

```

```

MG90190
MG90192
MG90194
MG90196
MG90198
MG90200
MG90202
MG90204
MG90206
MG90208
MG90210
MG90212
MG90214
MG90216
MG90218
MG90220
MG90222
MG90224
MG90226
MG90228
MG90230
MG90232
MG90234
MG90236
MG90238
MG90240
MG90242
MG90244
MG90246
MG90248
MG90250
MG90252
MG90254
MG90256
MG90258
MG90260
MG90262
MG90264
MG90266
MG90268
MG90270
MG90272
MG90274
MG90276
MG90278
MG90280
MG90282
MG90284

```

```

WRITE (15,601)
601 FORMAT (1H1,25HTHIS IS A BUILDUP RUN
)
602 WRITE (6,603)
WRITE (9,603)
WRITE (15,603)
603 FORMAT (1H1,25HTHIS IS A DECAY RUN
604 WRITE (6,399)
WRITE (9,399)
WRITE (15,399)
399 FORMAT (1H ,50HLATTICE COVERS ONLY 1/4 OF THE CROSS SECTION AREA )
WRITE (6,56) NI,NJ
WRITE (9,56) NI,NJ
WRITE (15,56) NI,NJ
56 FORMAT (1H ,28HLATTICE SIZE = NI BY NJ = 13,4H BY 13,
15X,20H(MAX = 48 BY 48)
WRITE (6,51) NI,NJ,ANJ,CDI
WRITE (9,51) NI,NJ,ANJ,CDI
WRITE (15,51) NI,NJ,ANJ,CDI
51 FORMAT (1H ,5HNJ = 15,4X,5HNJ = 15,4X,6HANJ = E16.8,
14X,CHCDI = E16.8)
ANJ1=ANJ*0.5
NJ1=ANJ1
IX=ADS (NJ1-NJ)
IF (IX-2) 58,58,59
59 WRITE (6,60)
WRITE (9,60)
WRITE (15,60)
60 FORMAT (1H0, 49HTHE VALUE OF NJ READ IN IS NOT CORRECT. RELOAD
61 STOP
58 CONTINUE
DX=CX**2
DT=(CDT*X*DX)/(C2*HMAX)**2)
DO 49 I=1,NY
DO 49 J=1,NX
DHO(I,J)=0.0
HI(I,J)=0.0
DI(I,J)=0.0
BI(I,J)=0.0
FORCE(I,J)=0.0
49 CONTINUE
WRITE (6,39)
WRITE (9,39)
WRITE (15,39)
39 FORMAT (1H0,7X,2HCX,14X,2HC1,14X,2HC2,14X,2HDI)
WRITE (6,40) CX,C1,C2,DI
WRITE (9,40) CX,C1,C2,DI
WRITE (15,40) CX,C1,C2,DI
MG90284
MG90288
MG90290
MG90292
MG90294
MG90296
MG90298
MG90300
MG90302
MG90304
MG90306
MG90308
MG90310
MG90312
MG90314
MG90316
MG90318
MG90320
MG90322
MG90324
MG90326
MG90328
MG90330
MG90332
MG90334
MG90336
MG90338
MG90340
MG90342
MG90344
MG90346
MG90348
MG90350
MG90352
MG90354
MG90356
MG90358
MG90360
MG90362
MG90364
MG90366
MG90368
MG90370
MG90372
MG90374
MG90376
MG90378
MG90380

```

```

40  FORMAT ( 1H , 4E16.8)
    WRITE ( 6, 400)
    WRITE ( 9, 400)
    WRITE (15, 400)
400  FORMAT ( 1H , 7X, 2HEO, 13X, 3HHSO, 14X, 2HBR, 13X, 4HHMAX, 12X, 3HCJO)
    WRITE ( 6, 401)
    WRITE ( 9, 401)
    WRITE (15, 401)
401  FORMAT ( 1H , 5E16.8)
    WRITE ( 6, 402)
    WRITE ( 9, 402)
    WRITE (15, 402)
402  FORMAT ( 1H , 7X, 2HTO, 13X, 3HTKM, 13X, 4HCOND, 13X, 2HAX, 14X, 2HBY)
    WRITE ( 6, 401)
    WRITE ( 9, 401)
    WRITE (15, 401)
403  FORMAT ( 1H , 7X, 3HCDT, 9X, 2HKZ, 3X, 2HNI, 3X, 2HNJ, 4X, 1HK, 3X, 2HKO, 3X,
12HKI, 3X, 2HKG)
    WRITE ( 6, 404)
    WRITE ( 9, 404)
    WRITE (15, 404)
404  FORMAT ( 1H , E16.8, 7F5)
    WRITE ( 6, 405)
    WRITE ( 9, 405)
    WRITE (15, 405)
405  FORMAT ( 1H , 7X, 2HB2, 14X, 2HH2, 14X, 2HB3, 14X, 2HH3)
    WRITE ( 6, 406)
    WRITE ( 9, 406)
    WRITE (15, 406)
406  FORMAT ( 1H , 4E16.8)
    WRITE ( 6, 407)
    WRITE ( 9, 407)
    WRITE (15, 407)
407  FORMAT ( 1H , 6X, 3HHIO, 13X, 3HDHX, 10X, 2HI3, 3X, 2HI4, 3X, 2HI5)
    WRITE ( 6, 408)
    WRITE ( 9, 408)
    WRITE (15, 408)
408  FORMAT ( 1H , 2E16.8, 3I5)
    WRITE ( 6, 500)
    WRITE ( 9, 500)
    WRITE (15, 500)
500  FORMAT ( 1H0, 50HEO
    WRITE ( 6, 501)
    WRITE ( 9, 501)
    WRITE (15, 501)

```

= MAXIMUM MATRIX ERROR

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MG90382
MG90384
MG90386
MG90388
MG90390
MG90392
MG90394
MG90396
MG90398
MG90400
MG90402
MG90404
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MG90408
MG90410
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MG90440
MG90442
MG90444
MG90446
MG90448
MG90450
MG90452
MG90454
MG90456
MG90458
MG90460
MG90462
MG90464
MG90466
MG90468
MG90470
MG90472
MG90474
MG90476

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501 FORMAT (1H,50HHSO) = MINIMUM SUM OF HI(I,J)) MG90478
 WRITE (6,502)) MG90480
 WRITE (15,502)) MG90482
 502 FORMAT (1H,50HBR) = RESIDUAL FLUX DENSITY) MG90484
 WRITE (6,503)) MG90486
 WRITE (9,503)) MG90488
 WRITE (15,503)) MG90490
 503 FORMAT (1H,50HMAX) = MAXIMUM FIELD INTENSITY APPLIED) MG90494
 WRITE (6,504)) MG90496
 WRITE (9,504)) MG90498
 WRITE (15,504)) MG90500
 504 FORMAT (1H,50HCUO) = CONSTANT USED TO ALLOW ALL UNITS) MG90502
 WRITE (6,505)) MG90504
 WRITE (9,505)) MG90506
 WRITE (15,505)) MG90508
 505 FORMAT (1H,50HIO) = INITIAL VALUE OF TIME IN SECONDS) MG90510
 WRITE (6,506)) MG90512
 WRITE (9,506)) MG90514
 WRITE (15,506)) MG90516
 506 FORMAT (1H,50HMKH) = SPECIFIES MAXIMUM TIME) MG90518
 WRITE (6,507)) MG90520
 WRITE (9,507)) MG90522
 WRITE (15,507)) MG90524
 507 FORMAT (1H,50HCOND) = MATERIAL CONDUCTIVITY) MG90526
 WRITE (6,508)) MG90528
 WRITE (9,508)) MG90530
 WRITE (15,508)) MG90532
 508 FORMAT (1H,50HAX) = CORE WIDTH) MG90534
 WRITE (6,509)) MG90536
 WRITE (9,509)) MG90538
 WRITE (15,509)) MG90540
 509 FORMAT (1H,50HBY) = CORE HEIGHT) MG90542
 WRITE (6,510)) MG90544
 WRITE (9,510)) MG90546
 WRITE (15,510)) MG90548
 510 FORMAT (1H,50HKZ) = NO OF SETS OF DATA) MG90550
 WRITE (6,511)) MG90552
 WRITE (9,511)) MG90554
 WRITE (15,511)) MG90556
 511 FORMAT (1H,50HCDT) = CONSTANT MULTIPLIER FOR THE TIME INCREMENT) MG90558
 WRITE (6,512)) MG90560
 WRITE (9,512)) MG90562
 WRITE (15,512)) MG90564
 512 FORMAT (1H,50HHI) = NO OF DIVISIONS IN THE X DIRECTION) MG90566
 WRITE (6,513)) MG90568
 WRITE (9,513)) MG90570
 WRITE (15,513)) MG90572

513 FORMAT (1H, 50HNJ)
 WRITE (6, 514)
 WRITE (9, 514)
 WRITE (15, 514)
 514 FORMAT (1H, 50HK)
 WRITE (6, 515)
 WRITE (9, 515)
 WRITE (15, 515)
 515 FORMAT (1H, 50HKO)
 WRITE (6, 634)
 WRITE (9, 634)
 WRITE (15, 634)
 634 FORMAT (1H, 60HKI)
 *ITS
 WRITE (6, 1580)
 WRITE (9, 1580)
 WRITE (15, 1580)
 1580 FORMAT (1H, 50HKG)
 WRITE (6, 516)
 WRITE (9, 516)
 WRITE (15, 516)
 516 FORMAT (1H, 50HB2)
 WRITE (6, 517)
 WRITE (9, 517)
 WRITE (15, 517)
 517 FORMAT (1H, 50HH2)
 WRITE (6, 518)
 WRITE (9, 518)
 WRITE (15, 518)
 518 FORMAT (1H, 50HB3)
 WRITE (6, 519)
 WRITE (9, 519)
 WRITE (15, 519)
 519 FORMAT (1H, 50HH3)
 WRITE (6, 520)
 WRITE (9, 520)
 WRITE (15, 520)
 520 FORMAT (1H, 50HHO)
 WRITE (6, 521)
 WRITE (9, 521)
 WRITE (15, 521)
 521 FORMAT (1H, 50HDX)
 WRITE (6, 522)
 WRITE (9, 522)
 WRITE (15, 522)
 522 FORMAT (1H, 50HI3)
 WRITE (6, 523)
 WRITE (9, 523)

= NO OF DIVISIONS IN THE Y DIRECTION)
 = NO OF INCREMENTS OR TIME SUBSCRIPT)
 = MAXIMUM NO OF INCREMENTS PER RUN)
 = SPECIFIES INCREMENTS FOR FORCT AND FLUXT PLO)
 = SPECIFIES INCREMENT FOR OUTPUT H1, B1, FORCE)
 = FLUX DENSITY FROM HYSTERESIS CURVE)
 = FIELD INTENSITY CORRESPONDING TO B2 VALUE)
 = FLUX DENSITY FROM HYSTERESIS CURVE)
 = FIELD INTENSITY CORRESPONDING TO B3 VALUE)
 = MAXIMUM VALUE OF FIELD INTENSITY FOR B VS H)
 = INCREMENTS OF FIELD INTENSITY IN B VS H)
 = MAXIMUM NO OF POINTS ON B VS H CURVE)

MG90574
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 MG90659
 MG90660

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WRITE ( 15, 523)
523 FORMAT ( 1H, 45HI4
150H USE 1 FOR OUTPUT TAPE 9 AND 15 = USED TO SPECIFY OUTPUT FORMAT
250H USE 3 FOR OUTPUT TAPE 9 ONLY (ALL DATA)
350H USE 4 FOR OUTPUT TAPE 15 ONLY (HI)
450H USE 5 FOR OUTPUT TAPE 15 ONLY (BI)
550H USE 6 FOR OUTPUT TAPE 15 ONLY (FORCE)
650H USE 7 FOR OUTPUT TAPE 15 ONLY (HI, BI, FORCE)
751H USE 6, 577) (ITERATIONS ON 9)
WRITE ( 9, 577)
WRITE ( 15, 577)
577 FORMAT ( 1H, 45HI5 = VALUE TO SPECIFY BUILDUP OR DECAY
250H USE 0 OR 1 FOR BUILDUP
250H USE 2 FOR DECAY
WRITE ( 6, 524)
WRITE ( 9, 524)
WRITE ( 15, 524)
524 FORMAT ( 1H, 50HCX = DIVISION WIDTHS
WRITE ( 6, 525)
WRITE ( 9, 525)
WRITE ( 15, 525)
525 FORMAT ( 1H, 50HC1 = CONSTANT IN MODIFIED FROELICH APPROXIMATION)
WRITE ( 6, 526)
WRITE ( 9, 526)
WRITE ( 15, 526)
526 FORMAT ( 1H, 50HC2 = CONSTANT IN MODIFIED FROELICH APPROXIMATION)
WRITE ( 6, 527)
WRITE ( 9, 527)
WRITE ( 15, 527)
527 FORMAT ( 1H, 50HDI = FINE INCREMENT IN SECONDS
41 FORMAT ( 1H, 41) K
114, 111) 7X, 7H(I, J), 9X, 7HBO(I, J), 9X, 7HHC(I, J), 9X, 5HTIME(
88 FORMAT ( 1H, 2X, 2H K, 2X, 19HTOTAL FORCE (FORCF), 11X, 10HTIME( K),
110X, 13HTOTAL FLUX (FLUXF))
DO 117 I=2, NY
DO 117 J=2, NX
CALL FLUD (110)
I2=BI(I, J)
CALL FORCX
117 CONTINUE
DO 119 I=2, NY
DO 119 J=2, NX
WRITE ( 9, 42) FORCE(I, J), BI(I, J), HO(I, J), T, I, J
42 FORMAT ( 1H, 4E16.8, 3X, 2I4)
119 CONTINUE

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/8X, MG90662
 /8X, MG90664
 /8X, MG90666
 /8X, MG90668
 /8X, MG90670
 /8X, MG90672
 /8X, MG90674
 /8X, MG90676
 /8X, MG90678
 MG90680
 MG90682
 MG90684
 MG90686
 /8X, MG90688
 /8X, MG90690
 MG90692
 MG90694
 MG90696
 MG90698
 MG90700
 MG90702
 MG90704
 MG90706
 MG90708
 MG90710
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 MG90750
 MG90752
 MG90754
 MG90756

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120 FORCT =0.0
    FLUXT=0.0
    DO 121 I=2,NY
    DO 121 J=2,NX
    FORCT=FORCT+FORCE(I,J)
    CALL FLUX
    CONTINUE
121 WRITE (9,200) FORCT,K,I
    200 FORMAT (1H0,22HTOTAL FORCE (FORCT) = E16.8,4X,5HTIME( I4,4H) = E
        116.8)
    WRITE (9,212) FLUXT,K,I
    212 FORMAT (1H ,22HTOTAL FLUX (FLUXT) = E16.8,4X,5HTIME(I4,4H) = E
        116.8/)
    WRITE ( 9,1600)
1600 FORMAT (1H1)
    89 WRITE ( 6,89) K,FORCT,I,FLUXT
    89 FORMAT (1H,14,2X,E16.8,11X,E16.8,8X,E16.8)
    290 FORMAT (1H1,42HMAGNETIC FIELD INTENSITY HO(I,J) AT TIME (I4,4H) =
        1E16.8)
    WRITE (15,271)
    DO 295 I=2,NY
    IF (NX-8) 491,492
    491 WRITE (15,272) I,(HO(I,J),J=2,NX)
    492 GO TO 295
    WRITE (15,272) I,(HO(I,J),J=2,8)
    WRITE (15,375) (HO(I,J),J=9,NX)
    295 CONTINUE
    WRITE (15,200) FORCT,K,I
    WRITE (15,212) FLUXT,K,I
    WRITE (15,270) K,I
    WRITE (15,271)
    DO 275 I=2,NY
    IF (NX-8) 489,490
    489 WRITE (15,272) I,(BI(I,J),J=2,NX)
    490 GO TO 275
    WRITE (15,272) I,(BI(I,J),J=2,8)
    275 CONTINUE
    WRITE (15,200) FORCT,K,I
    WRITE (15,212) FLUXT,K,I
    WRITE (15,316) K,I
    536 FORMAT (1H1,31HTHEORETICAL FORCE(I,J) AT TIME( I4, 4H) = E16.8)
    DO 537 I=2,NY
    IF (NX-8) 538,539
    538 WRITE (15,272) I,(FORCE(I,J),J=2,NX)

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MG90758
 MG90760
 MG90762
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 MG90772
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 MG90842
 MG90844
 MG90846
 MG90848

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539 GO TO 537
WRITE (15,272) I, (FORCE(I,J),J=2,8)
WRITE (15,375) (FORCE(I,J),J=9,NX)
537 CONTINUE
WRITE (15,200) FORCI,K,I
WRITE (15,212) FLUXI,K,I
IF (K-(K/KI)*KI) 631,630,631
630 KK=(K+KI)/KI
IFOPC (KK) = FORCI
TIME (KK) = FLUXI
KI=K
631 CONTINUE
I=I+DT
K=K+1
DO 15 J=2,I1
DO 15 J=2,I2
IF (I-2) 6,6,5
5 IF (J-2) 10,10,30
6 IF (J-2) 3,8,7
7 HO(I-1,J)=HO(I+1,J)
GO TO 30
8 HO(I;J-1)=HO(I,J+1)
HO(I-1,J)=HO(I+1,J)
GO TO 30
10 HO(I,J-1)=HO(I,J+1)
30 X1=(O2+HO(I,J))**2
X2=(HO(I,J+1))-2.0*HO(I,J)+HO(I,J-1))/DX
16 X3=(HO(I+1,J))-2.0*HO(I,J)+HO(I-1,J))/DX
DHO(I,J)=X1+(X2+X3)/X
DO 17 I=2,I1
DO 17 J=2,I2
17 H1(I,J)=HO(I,J)+DHO(I,J)*DT
I=NIJ+2
DO 9 J=2,NX
9 H1(I,J)=HO(I,J)
DO 90 I=2,NY
90 H1(I,J)=HO(I,J)
557 WRITE (9,558) NM,K,I
558 FORMAT (1H0,5HNI = I,5X,5HTIME(I,4H) = E16.8)
560 CONTINUE
18 DO 32 I=2,I1
DO 32 J=2,I2
IF (I-2) 20,20,19
19 IF (J-2) 23,23,24
20 IF (J-2) 22,22,21
21 H1(I-1,J)=H1(I+1,J)

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MG90850
MG90852
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MG90952


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22  GO 10 24
    H1(I,J-1)=H1(I,J+1)
    H1(I-1,J)=H1(I+1,J)
23  GO 10 24
    X4=(C2+H1(I,J))*2
24  X5=(H1(I,J+1)-2.0*H1(I,J)+H1(I,J-1))/DX
    X3=(H1(I+1,J)-2.0*H1(I,J)+H1(I-1,J))/DX
32  DH1(I,J)=X4*(X5+X6)/X
    NN=NN+1
    E=0.0
33  DO 34 I=2,I1
    DO 34 J=2,I2
        X7=(DH1(I,J)+DH1(I,J))/2.0
        T1=H0(I,J)+X7*DT
        E=E+ABS(T1-H1(I,J))
        H1(I,J)=T1
        CALL FLUD(H1)
        T2=B1(I,J)
        CALL FORCX
34  CONTINUE
    I=NY
    DO 201 J=2,NX
        CALL FLUD(H0)
        T2=B1(I,J)
        CALL FORCX
        CONTINUE
    J=NX
    DO 202 I=2,NY
        CALL FLUD(H0)
        T2=B1(I,J)
        CALL FORCX
        CONTINUE
    IF (NN-3) 204,204,211
201  IF (E=EI-E)
202  IF (DE-(0.01*E1)) 205,204,204
204  KN=1
    GO 10 206
205  KN=2
206  GO 10 (106,207),KN
207  WRITE (9,208)
    WRITE (6,208)
    WRITE (15,208)
208  I=I+1
    FORMAT (1H0,72HTHE ERROR DOES NOT CONVERGE FOR THIS SET OF DATA.)
    I=I+1
    WRITE (6,57) NN,E
    WRITE (9,57) NN,E
    WRITE (15,57) NN,E

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MG90954
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 MG91048

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GO 10 46
106 WRITE (9,224) NR,E
224 FORMAT (1H,5UNN = I4,4X,4HE = E16.8)
57 FORMAT (1H,5UNN = I4,4X,4HE = E16.8/)
E1=E
81 IF(E=EO) 81,81,18
FORC=0.0
FLUX=0.0
DO 83 I=2,NX
DO 81 J=2,NX
FORC=FORC+FORCE(I,J)
T2=B:(I,J)
CALL FLUX
83 CONTINUE
WRITE (6,89) K,FORCT,T,FLUX
14-2) 31,31,550
14-1) 1578,1578,1579
1578 IF(K-(K/KG)=KG) 550,1579,550
1579 CONTINUE
WRITE (9,123) K
123 FORMAT (1H),4X,10HFORCE(I,J),7X,7HBI(I,J),9X,7HH1(I,J),9X,5HTIME(
14,1H),7X,7H(I,J)
35 DO 80 I=2,NX
DO 80 J=2,NX
WRITE (9,42) FORCE(I,J),BI(I,J),HI(I,J),T,I,J
80 CONTINUE
WRITE (9,200) FORCT,K,T
WRITE (9,212) FLUX,K,T
WRITE (9,1600)
550 CONTINUE
IF (14-1) 1590,1590,1591
1590 IF (K-(K/KG)=KG) 555,1591,555
1591 CONTINUE
552 GO TO (552,555,552,553,554,552,555), I4
552 WRITE (15,296) K,T
296 FORMAT (1H1,42H MAGNETIC FIELD INTENSITY 12(I,J) AT TIME (I4,4H)
E16.8)
WRITE (15,271)
DO 297 I=2,NX
IF (NX-8) 389,389,390
WRITE (15,272) I,(HI(I,J),J=2,NX)
GO TO 297
WRITE (15,272) I,(HI(I,J),J=2,8)
WRITE (15,375) (HI(I,J),J=9,NX)
297 CONTINUE
WRITE (15,200) FORCT,K,T
WRITE (15,212) FLUX,K,T
GO TO (553,555,555,554,553,555), I4

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MG91050
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MG91126
MG91128
MG91130

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553 WRITE (15,270) K,T
270 FORMAT (1H1, 36HFLUX DENSITY MATRIX B1(I,J) AT TIME(14,4H) =
1E16.8)
271 WRITE (15,271)
FORMAT (1H0,13X,3HJ=2,15X,1H3,15X,1H4,15X,1H5,15X,1H6)
DO 273 I=2,NY
IF (NX-8) 379,379,380
WRITE (15,272) I,(B1(I,J),J=2,NX)
272 FORMAT (1H,14,1X,7E16.8)
GO TO 273
380 WRITE (15,272) I,(B1(I,J),J=2,8)
WRITE (15,375) (B1(I,J),J=9,NX)
375 FORMAT (1H,5X,7E16.8)
273 CONTINUE
WRITE (15,200) FORCT,K,T
WRITE (15,212) FLUXI,K,T
GO TO (554,555,555,555,554,555) K,T
554 WRITE (15,271)
DO 547 I=2,NY
IF (NX-8) 548,548,549
WRITE (15,272) I,(FORCE(I,J),J=2,NX)
548 GO TO 547
WRITE (15,272) I,(FORCE(I,J),J=2,8)
549 WRITE (15,375) (FORCE(I,J),J=9,NX)
547 CONTINUE
WRITE (15,200) FORCT,K,T
WRITE (15,212) FLUXI,K,T
IF (14-1) 555,555,635
555 CONTINUE
IF (K-(K/KI)*KI) 633,632,633
632 TFORC (KK) = FORCT
TFLUX (KK) = FLUXI
TIME (KK) = T
633 CONTINUE
IS=0,0
DO 43 I=2,NY
DO 43 J=2,NX
HS=HS+HI(I,J)
572 IF (IS-1) 572,572,571
571 IF (HS-HSO) 44,46,46
44 IF (HS-HSO) 46,46,44
276 IF (K-KO) 276,46,46
277 IF (T-TD) 277,46,46
K=K+1
HH=H

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MG911132
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DO 45 I=2, NY
DO 45 J=2, NX
45 HO(I, J)=HI(I, J)
CALL SSMATCH (I, KKK)
GO TO (46, 4), KKK
46 WRITE (6, 561) HS, HSO, I, TKN, K, KO
WRITE (9, 561) HS, HSO, I, TKN, K, KO
WRITE (15, 561) HS, HSO, I, TKN, K, KO
561 FORMAT (1H0, 6HHS = F16.8, 5X, 6HHSO = F16.8
2H , 6HK = F16.8, 5X, 6HKKH = F16.8
2H , 6HK = F15, 16X, 6HKO = F15
IF (15-1) 573, 574, 574
573 IF (HS-HSO) 569, 575, 575
574 IF (HS-HSO) 566, 566, 569
569 IF (T-TKN) 570, 567, 567
570 IF (K-KO) 565, 568, 568
575 WRITE (6, 576)
WRITE (9, 576)
WRITE (15, 576)
576 FORMAT (1H0, 60HSTOP BECAUSE HS IS GREATER THAN OR = TO HSO (BUILDUP
1)
GO TO 1565
566 WRITE (6, 562)
WRITE (9, 562)
WRITE (15, 562)
562 FORMAT (1H0, 50HSTOP BECAUSE HS IS LESS THAN OR = TO HSO (DECAY)
GO TO 1565
567 WRITE (6, 563)
WRITE (9, 563)
WRITE (15, 563)
563 FORMAT (1H0, 50HSTOP BECAUSE TIME (T) IS GREATER THAN OR = TO TKN
GO TO 1565
568 WRITE (6, 564)
WRITE (9, 564)
WRITE (15, 564)
564 FORMAT (1H0, 50HSTOP BECAUSE K IS GREATER THAN OR = TO KO
GO TO 1565
565 WRITE (6, 1566)
WRITE (9, 1566)
WRITE (15, 1566)
1566 FORMAT (1H0, 42HSTOP BECAUSE SENSE SWITCH 1 WAS TURNED ON
1565 CONTINUE
IF (15-1) 605, 605, 606
605 WRITE (6, 607)
WRITE (9, 607)
WRITE (15, 607)
607 FORMAT (1H0, 25HTHIS IS A BUILDUP RUN
GO TO 609

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MG91230
MG91232
MG91234
MG91236
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MG91299
MG91300
MG91301
MG91302
MG91304
MG91306
MG91308
MG91310
MG91312
MG91314

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606 WRITE ( 6,608)
WRITE ( 9,608)
608 WRITE ( 15,608)
609 FORMAT (1H0, 25HTHIS IS A DECAY RUN
CONTINUE
K2=KK-K1+1
CALL PPLT ( TIME(K1), TFORC(K1), K2)
WRITE ( 6,611)
611 FORMAT (1H0, 40X, 45HTOTAL FORCE VS TIME
CALL PPLT ( TIME(K1), TFLUX(K1), K2)
WRITE ( 6,612)
612 FORMAT (1H0, 40X, 45HTOTAL FLUX VS TIME
WRITE (16,1000) EO, HSO, BR, HMAX, CUO
WRITE (16,1000) T, TKH, COND, AX, BY
1000 FORMAT (5E16.8)
2000 WRITE (16,2000) CDT, KZ, NI, NJ, K, KO, KI, KG
WRITE ( E16.8, 7I5)
2130 WRITE (16,2130) B2, H2, B3, H3
2790 WRITE (16,2790) H10, DHX, I3, I4, I5
WRITE (2E16.8, 3I5)
1577 WRITE (16,1577) ((H1(I, J), J=2, NX), I=2, NY)
FORMAT (1H, 90HTHIS IS OUTPUT DATA TO BE USED FOR A CONTINUATION
1 RUN. CHANGE K TO K = 0
WRITE ( 6,1001) EO, HSO, BR, HMAX, CUO
1001 WRITE (1H, 5E16.8)
WRITE ( 6,1002) T, TKH, COND, AX, BY
1002 FORMAT (1H, 5E16.8)
2001 WRITE ( 6,2001) CDT, KZ, NI, NJ, K, KO, KI, KG
WRITE (1H, E16.8, 7I5)
2131 FORMAT (1H, 4E16.8)
2791 WRITE ( 6,2791) H10, DHX, I3, I4, I5
WRITE (1H, 2E16.8, 3I5)
1I4=6
I4=LI4
GO TO 31
635 WRITE ( 6,636) KZ, KZ
WRITE ( 9,636) KZ, KZ
636 WRITE (15,636) KZ, KZ
FORMAT(1H0, 6H(KZ = I4, 3X, 31HPREVIOUS DATA FOR DATA SET NO. I4, 1H))
KZ=KZ-1
285 IF (KZ) 47, 47, 48
47 CONTINUE
END FILE 9
END FILE 15

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MG91316
MG91318
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MG91398
MG91400
MG91402
MG91404

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END FILE 16
STOP
END
$IBFTC FLUX
SUBROUTINE FLUX (I0)
DIMENSION HO(35,35),FORCE(35,35),B1(35,35)
COMMON I,J,UO,DX,T2,FLUX1,FORCE,BR,IX,NY,I5,B1,C1,C2
IF (I5-1) 614,614,613
613 HC=(C2*BR)/(C1+BR)
GO TO 617
614 HC=(C2*BR)/(C1+BR)
IF (HO(I,J)-HC) 615,616,616
615 HOX=-HO(I,J)
B1(I,J)=-HC+HOX/(C2+H0X)*BR)
GO TO 620
616 X1=HO(I,J)-2.0*HC
B1(I,J)=(C1*X1/(C2+X1))+BR)
GO TO 620
617 IF (HO(I,J)-HC) 618,619,619
618 HD3=35*(HO(I,J)-HC)
HD4=HC+HD3
B1(I,J)=-HC+HD4/(C2+HD4)+BR)
GO TO 620
619 B1(I,J)=(C1*HO(I,J)/(C2+HO(I,J)))+BR)
620 RETURN
END
$IBFTC FORCX
SUBROUTINE FORCX
DIMENSION FORCE(35,35),F1(35,35)
COMMON I,J,UO,DX,T2,FLUX1,FORCE,DR,IX,NY,I5,B1,C1,C2
IF (I-2) 300,300,301
IF (J-2) 303,253,302
IF (J-IX) 252,253,253
IF (J-2) 303,303,304
IF (I-NY) 252,253,253
IF (I-NY) 305,306,306
IF (J-IX) 254,252,252
IF (J-IX) 252,253,253
252 FORCE(I,J)=((T2**2)-(DX/2.0))/(2.0*U0)
GO TO 256
253 FORCE(I,J)=((T2**2)-(DX/4.0))/(2.0*U0)
GO TO 256
254 FORCE(I,J)=((T2**2)-DX)/(2.0*U0)
256 CONTINUE
RETURN
END
$IBFTC FLUX
SUBROUTINE FLUX

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MG91406
MG91408
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MG91496
MG91498
MG91500

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TIME DEPENDENCE OF THE MAGNETIC
FIELD IN A RECTANGULAR TOROID

by

GLEN LEROY SHURTZ

B.S.M.E., Kansas State University, 1963
B.S.E.E., Kansas State University, 1964

AN ABSTRACT OF A MASTER'S THESIS

submitted in partial fulfillment of the
requirements for the degree

MASTER OF SCIENCE

Department of Electrical Engineering

KANSAS STATE UNIVERSITY
Manhattan, Kansas

1966

The purpose of this thesis was to investigate the time dependence of the magnetic field in a rectangular toroid when subjected to an applied magnetomotive force.

According to Lenz's law, eddy currents generated in the transient state induce an mmf in opposition to that applied. This results in a nonuniform time varying distribution of field intensity and flux within the core during a transition between states. An equation describing time dependence of this distribution pattern was derived by applying Maxwell's equations with the assumption that symmetrical boundary conditions were applied along the length of the toroid and no leakage could occur. Combining this equation with another which provided an approximation to the B-H relationship for a given ferromagnetic material, yielded the Hysteretic Diffusion Equation.

$$\frac{\partial H}{\partial t} = \frac{(C_2 + H)^2}{\sigma C_1 C_2} \left[\frac{\partial^2 H}{\partial x^2} + \frac{\partial^2 H}{\partial y^2} \right]$$

Its solution yields the time dependence of the field intensity as a function of position in the core.

The Modified Euler Method of numerical integration was used to evaluate time variance of the magnetic field intensity and flux distribution patterns from which the time dependence of total force and total flux was calculated.

Discussion of error criteria and selection of an appropriate time increment to insure convergence were discussed and a sample problem illustrating the method was given.

